

Quantum entanglement and the phases of matter

Twelfth Arnold Sommerfeld Lecture Series
January 31 - February 3, 2012

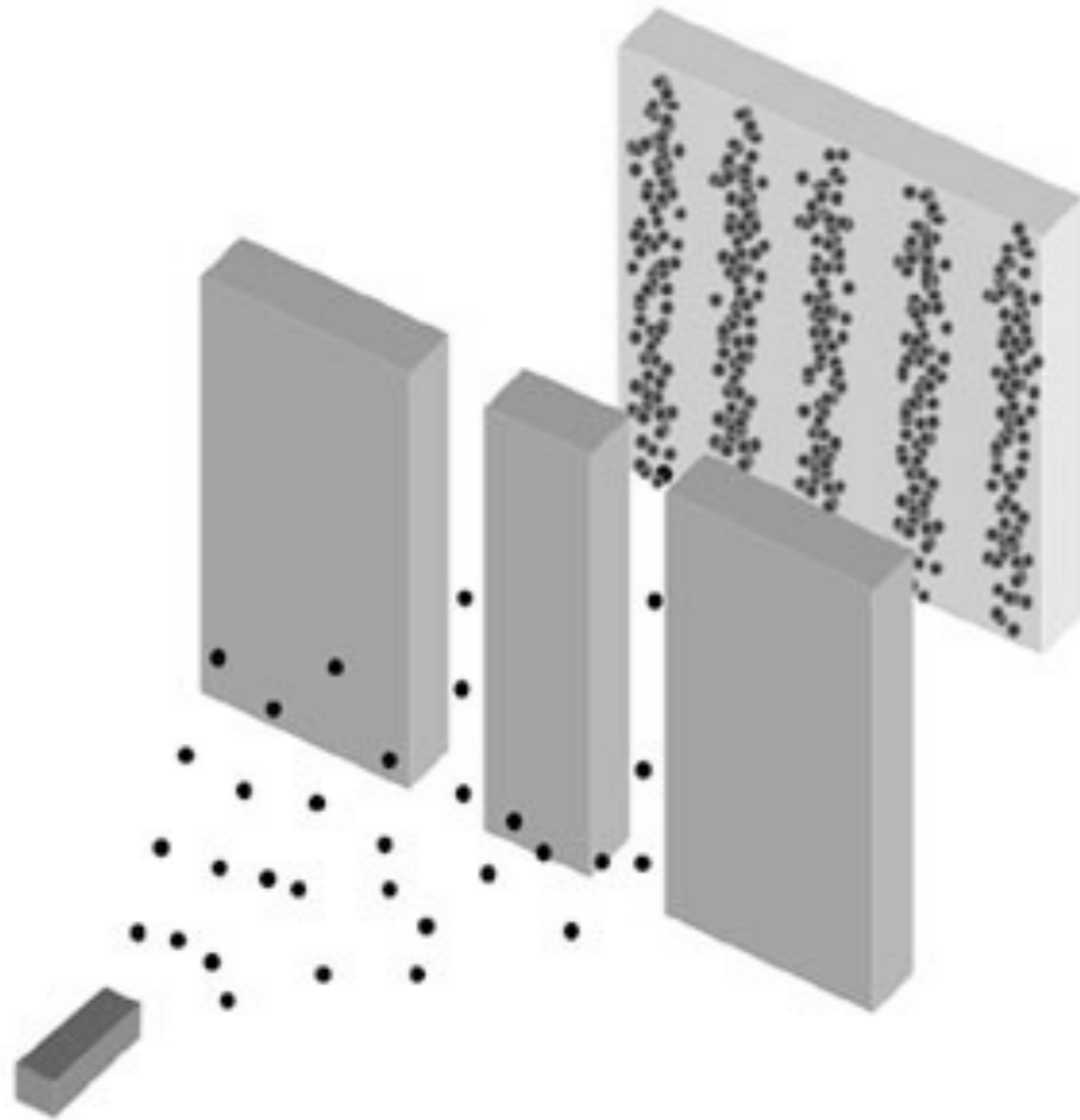
sachdev.physics.harvard.edu



**Quantum
superposition and
entanglement**

Quantum Superposition

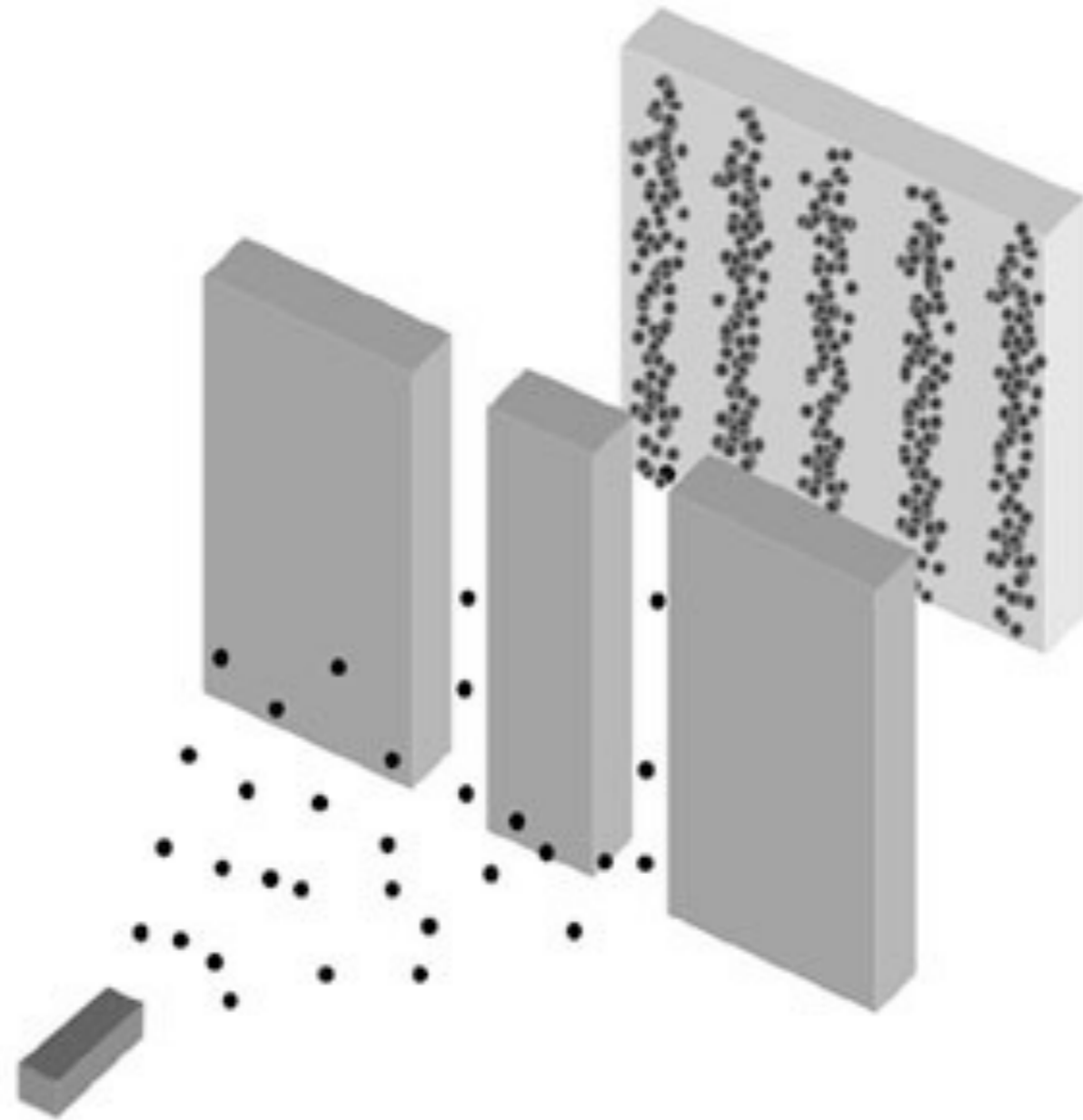
The double slit experiment



Interference of electrons

Quantum Superposition

The double slit experiment

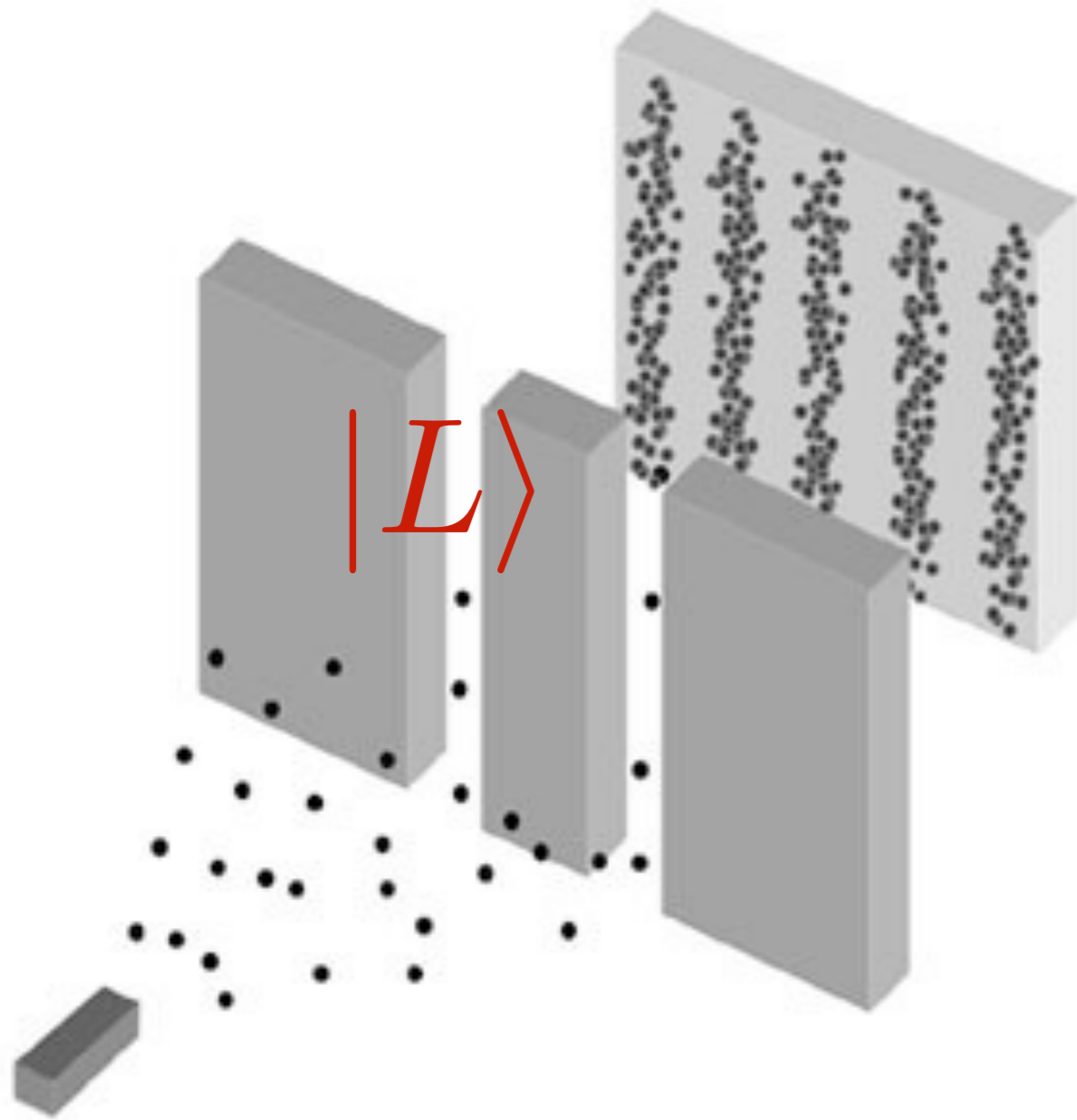


Each
electron
passes
through
both slits

Interference of electrons

Quantum Superposition

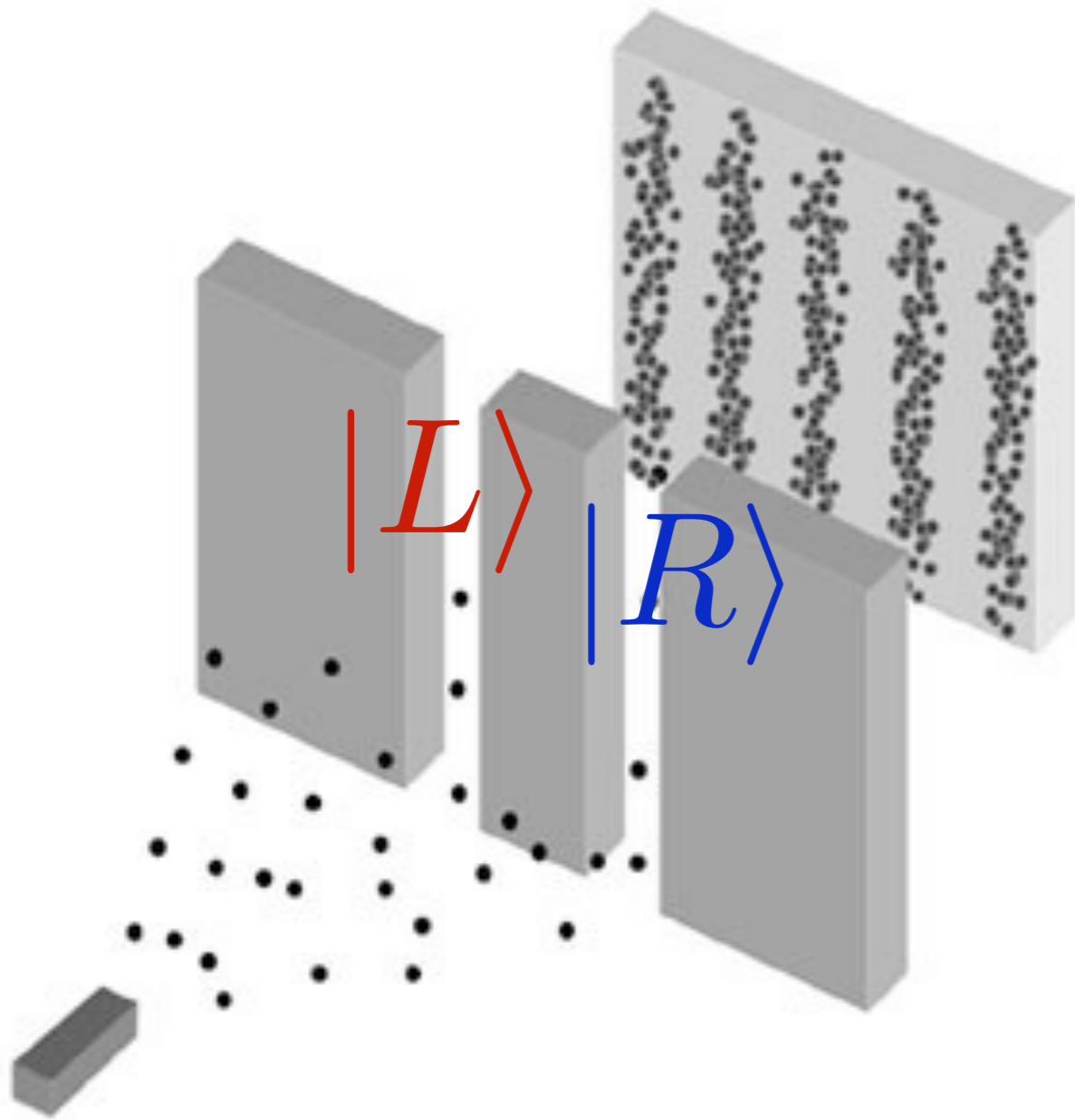
The double slit experiment



Let $|L\rangle$ represent the state with the electron in the left slit

Quantum Superposition

The double slit experiment

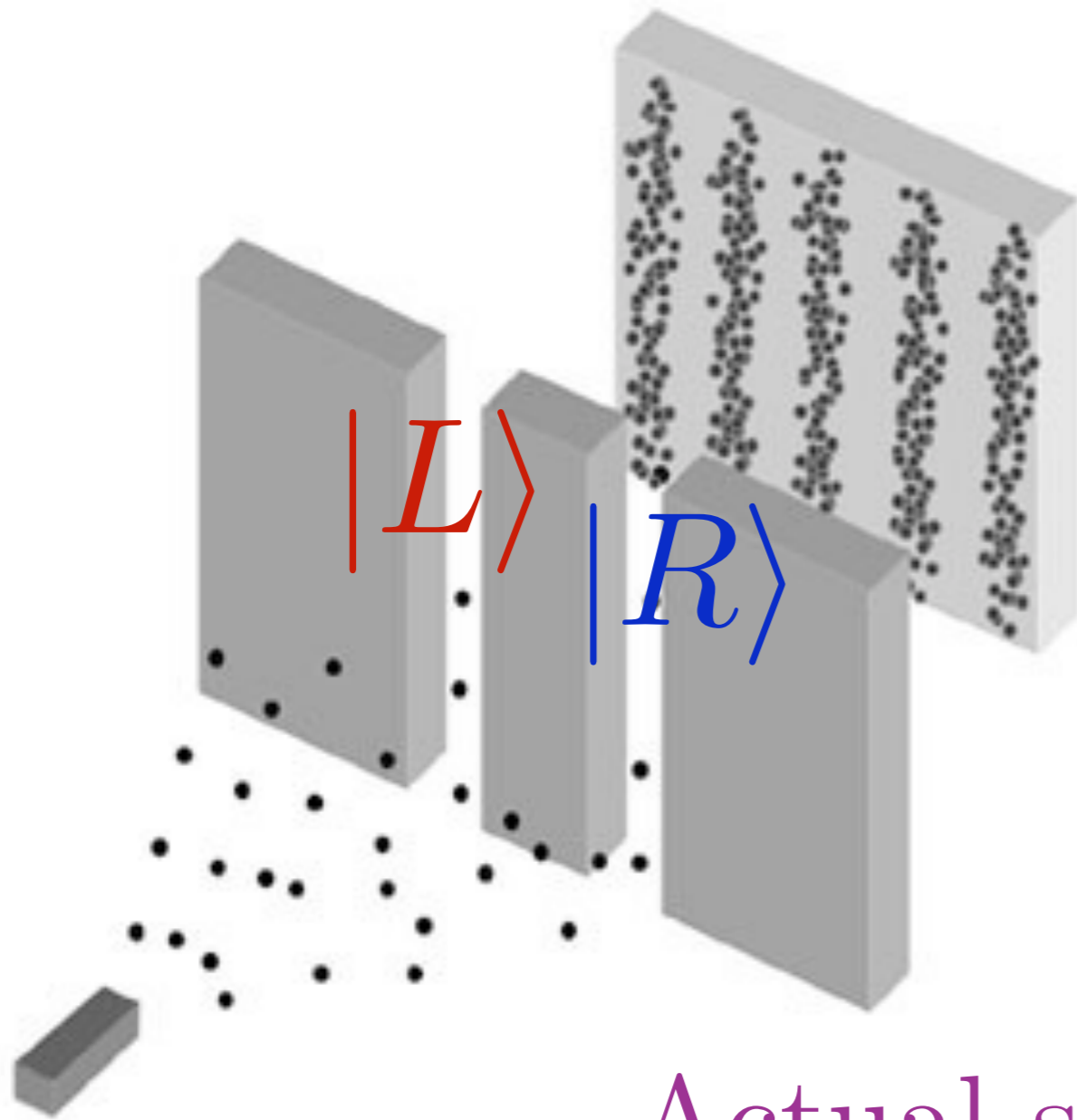


Let $|L\rangle$ represent the state with the electron in the left slit

And $|R\rangle$ represents the state with the electron in the right slit

Quantum Superposition

The double slit experiment



Let $|L\rangle$ represent the state with the electron in the left slit

And $|R\rangle$ represents the state with the electron in the right slit

Actual state of the electron is

$$|L\rangle + |R\rangle$$

Quantum Entanglement: quantum superposition with more than one particle

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Hydrogen atom:

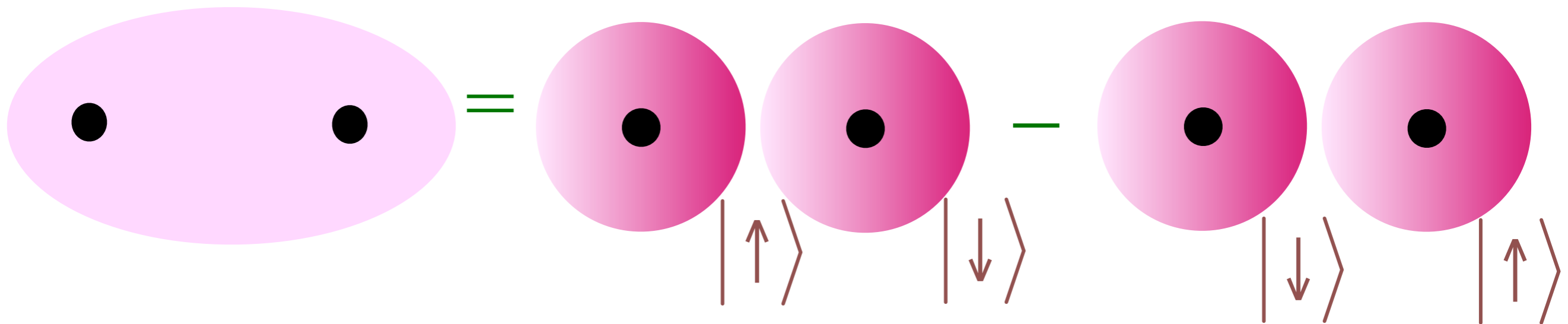


Quantum Entanglement: quantum superposition with more than one particle

Hydrogen atom:



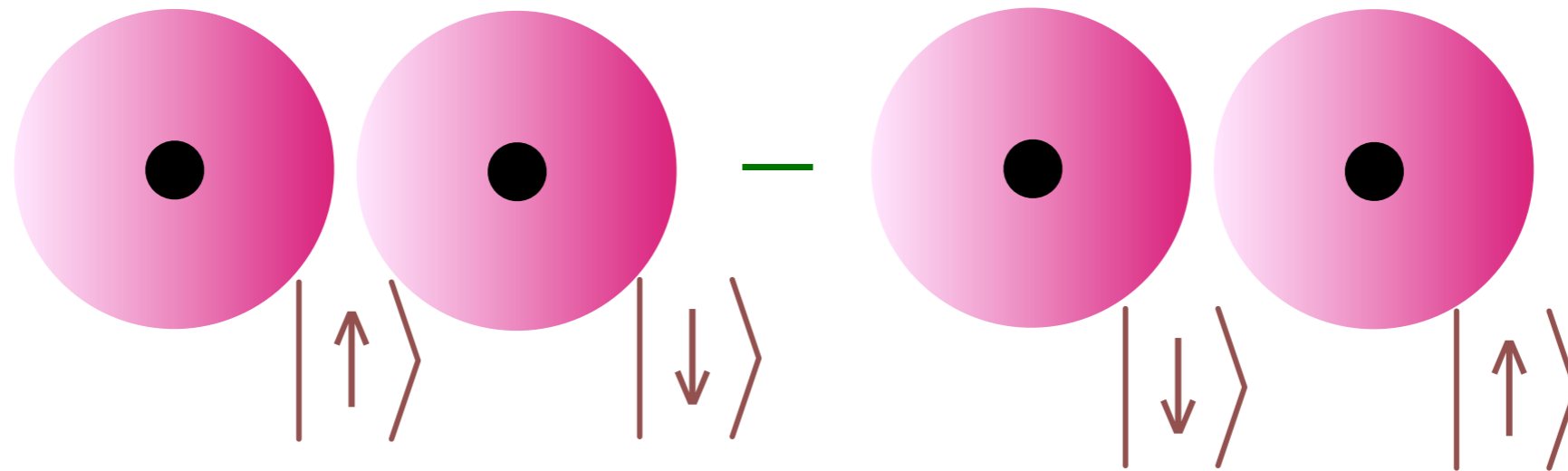
Hydrogen molecule:



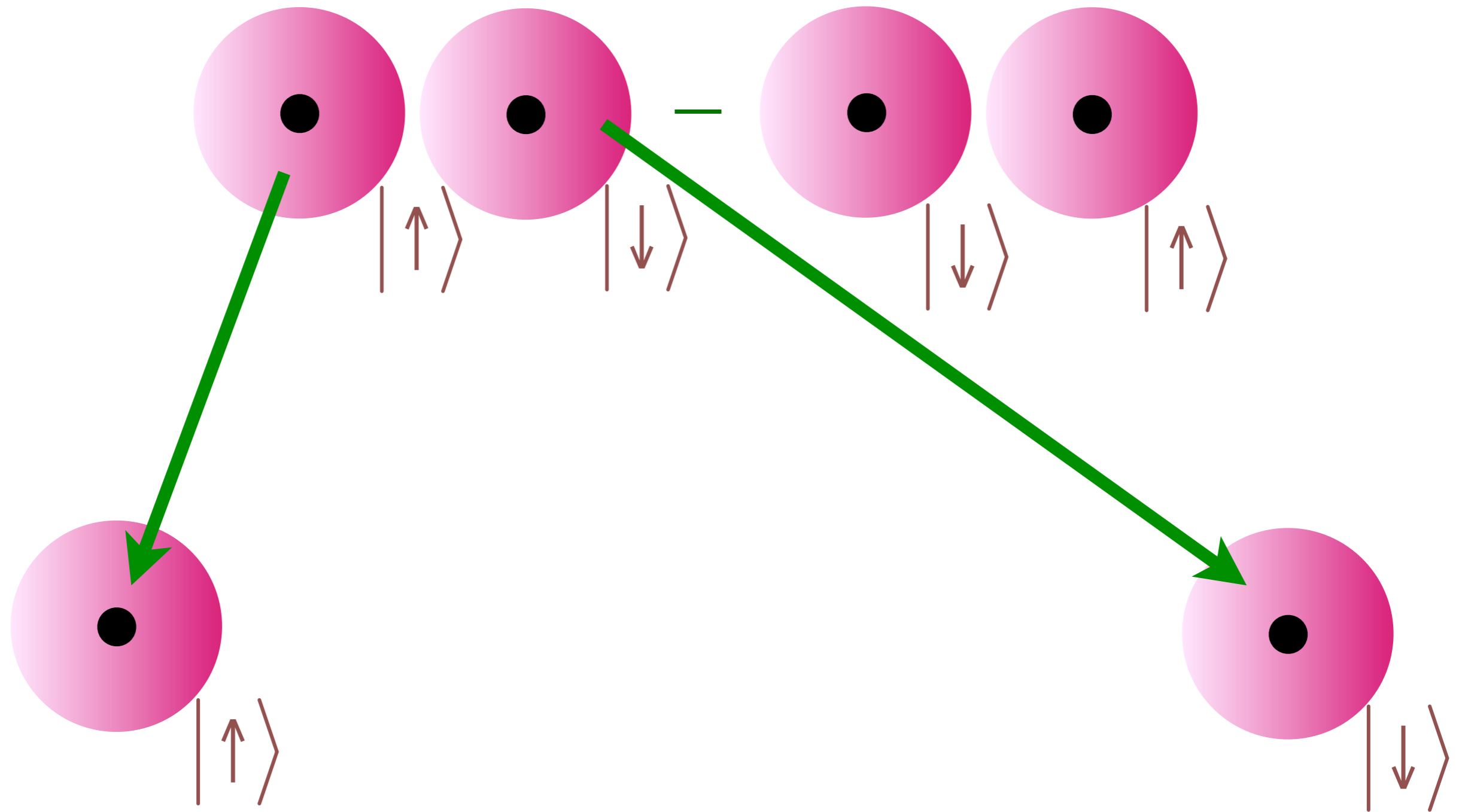
$$= \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Superposition of two electron states leads to non-local
correlations between spins

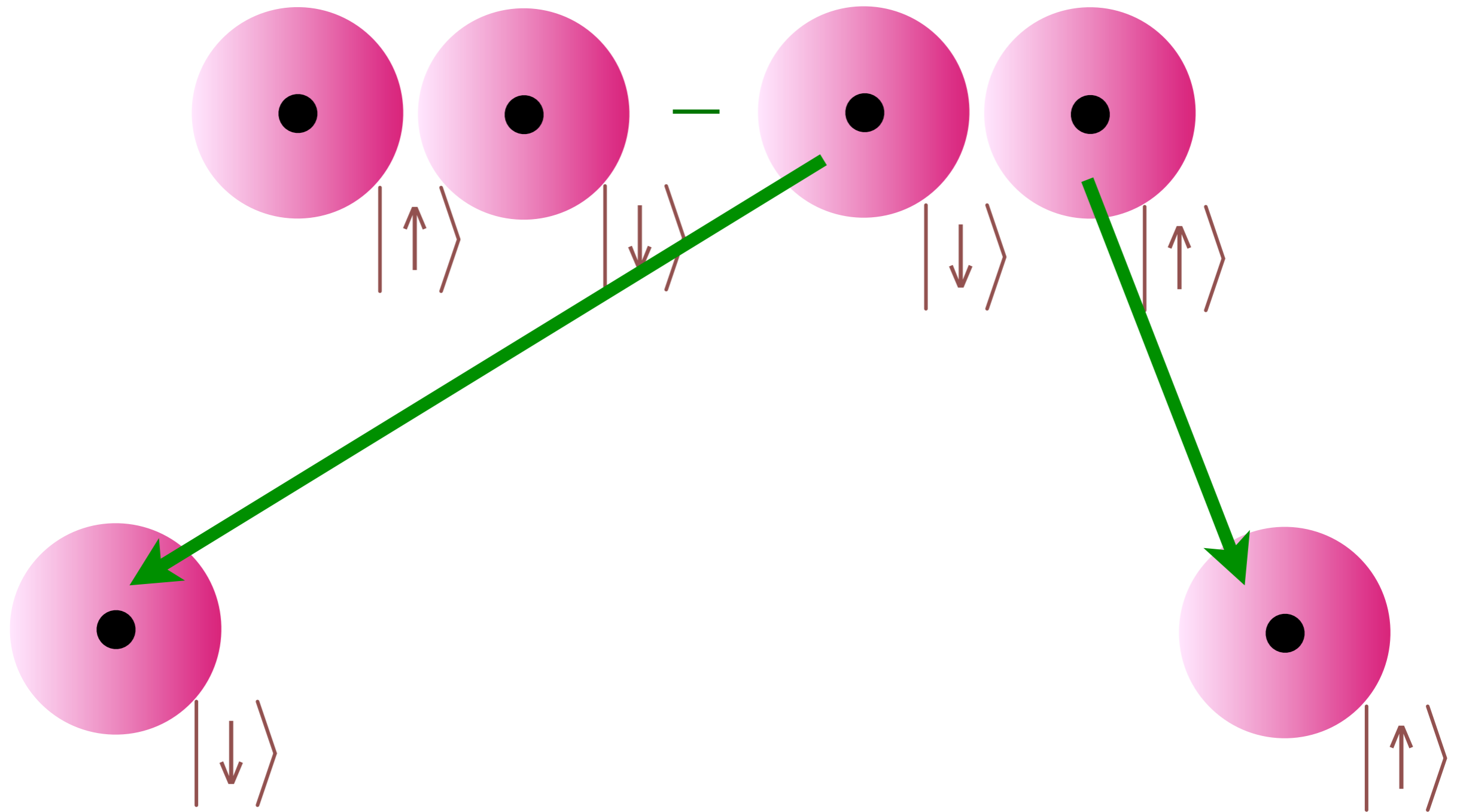
Quantum Entanglement: quantum superposition with more than one particle



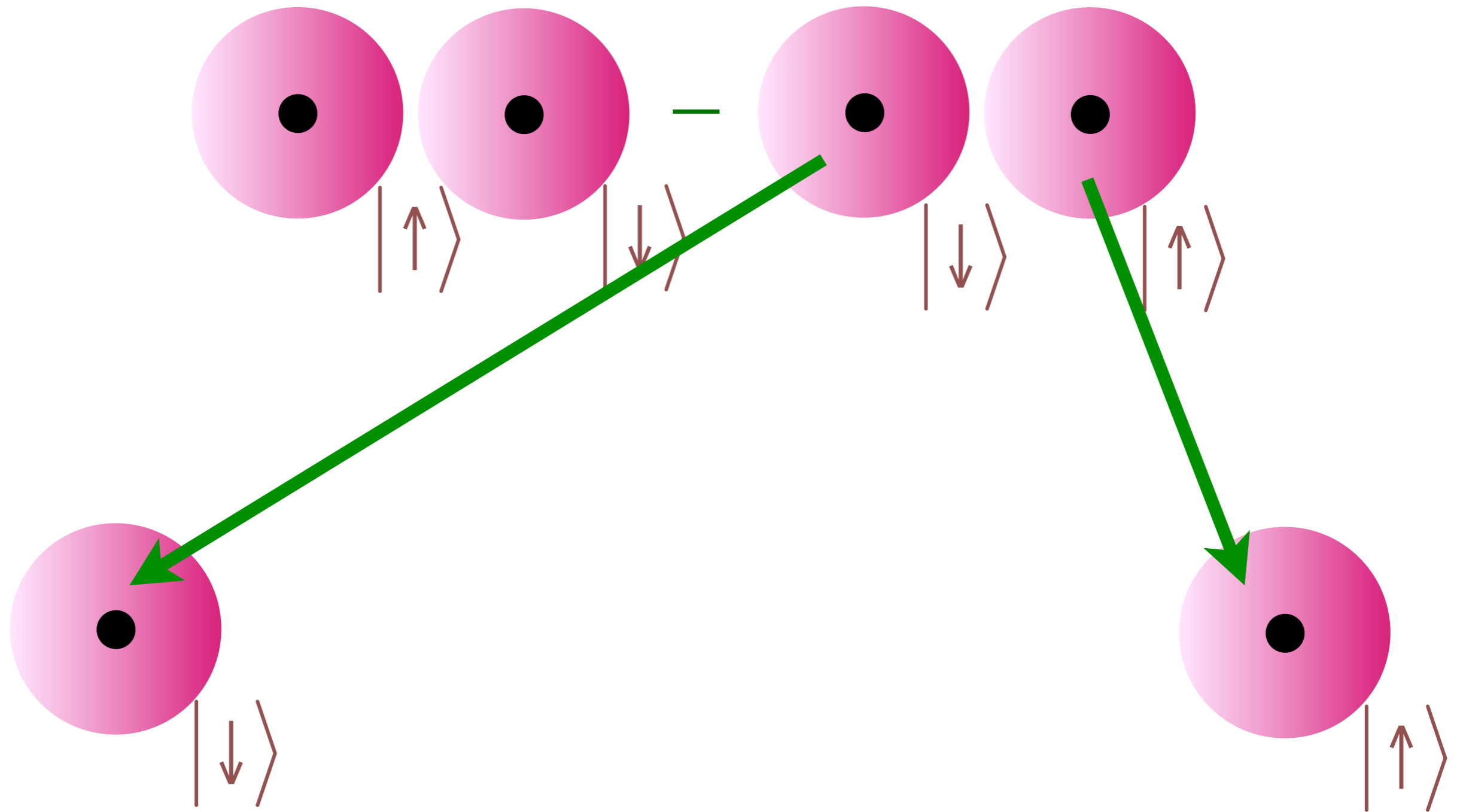
Quantum Entanglement: quantum superposition with more than one particle



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Quantum Entanglement: quantum superposition with more than one particle



Einstein-Podolsky-Rosen “paradox”: Non-local correlations between observations arbitrarily far apart

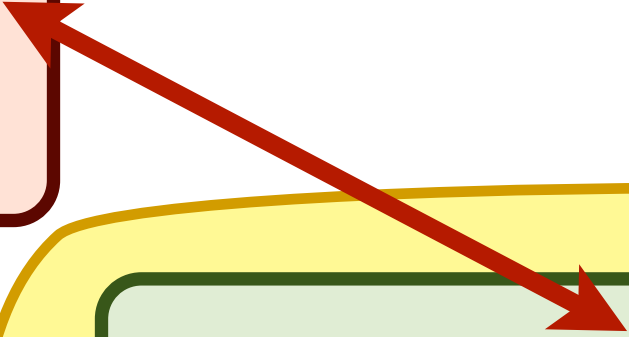
**Quantum
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**Quantum critical
points of electrons
in crystals**

**String theory
and black holes**

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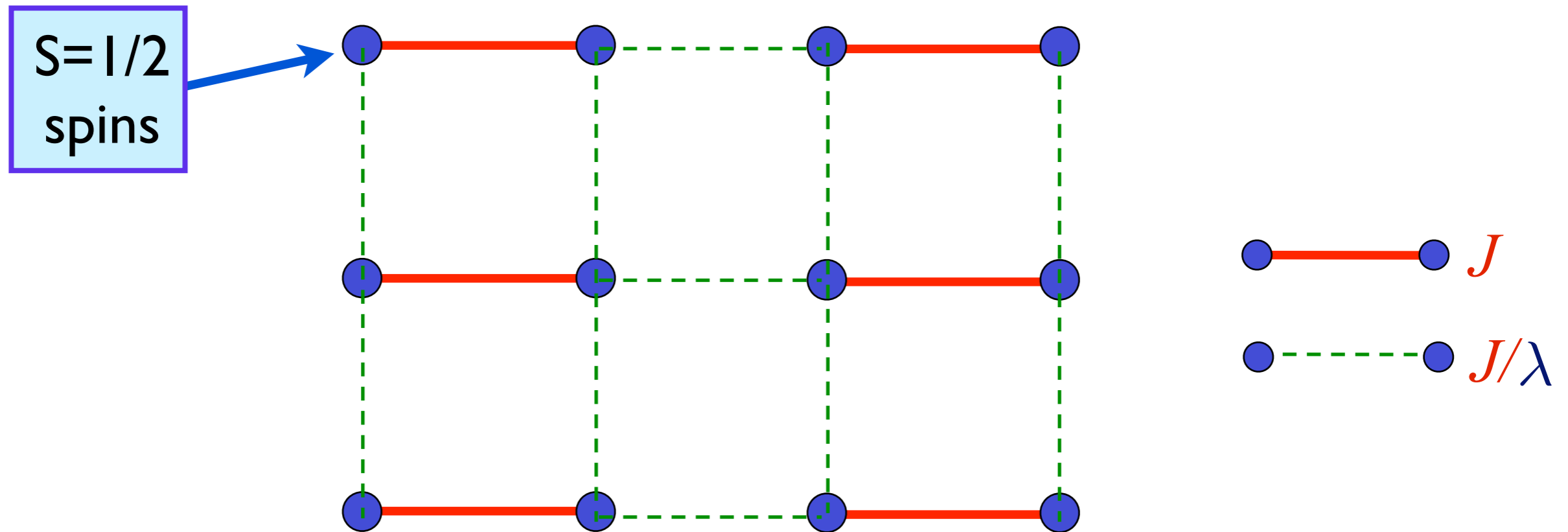


**Quantum critical
points of electrons
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**String theory
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Spinning electrons localized on a square lattice

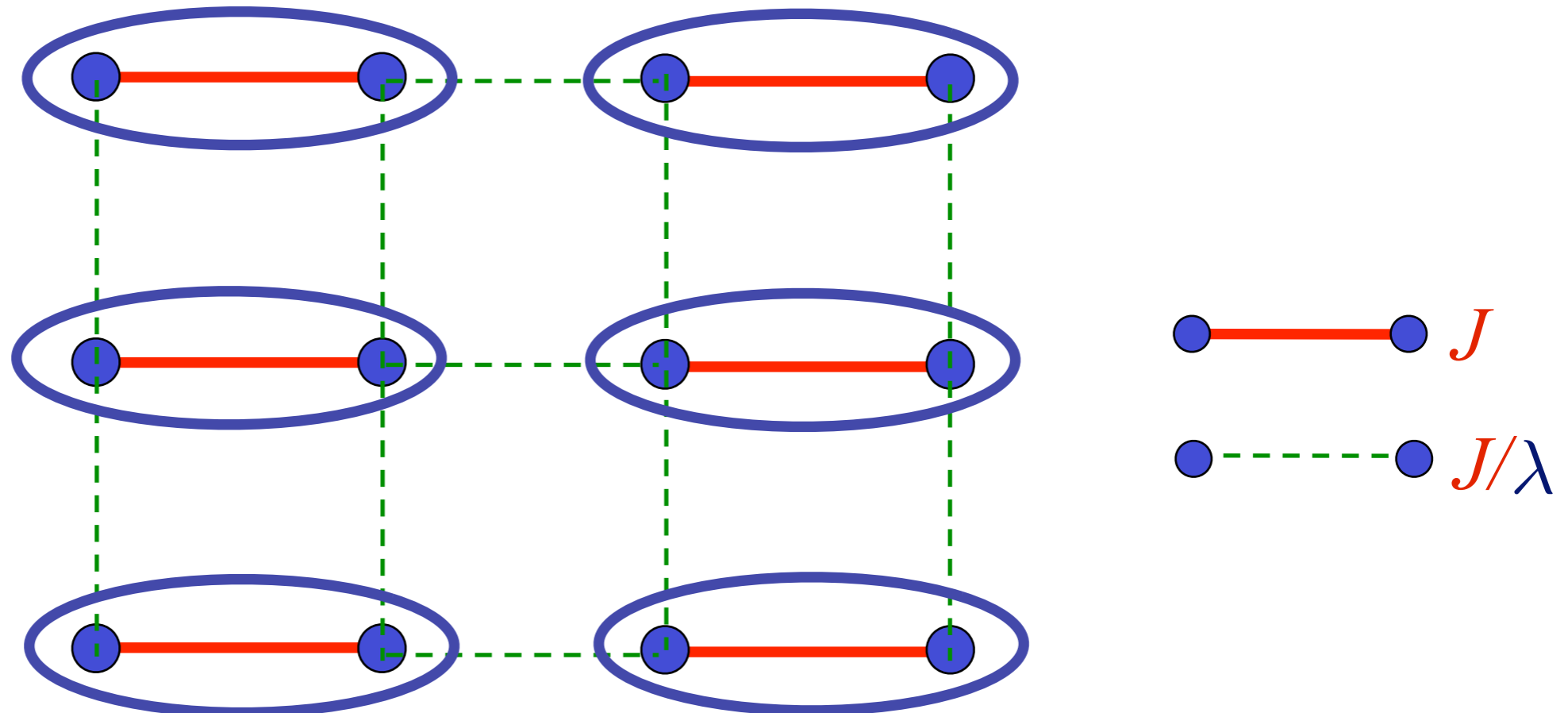
$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$



Examine ground state as a function of λ

Spinning electrons localized on a square lattice

$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

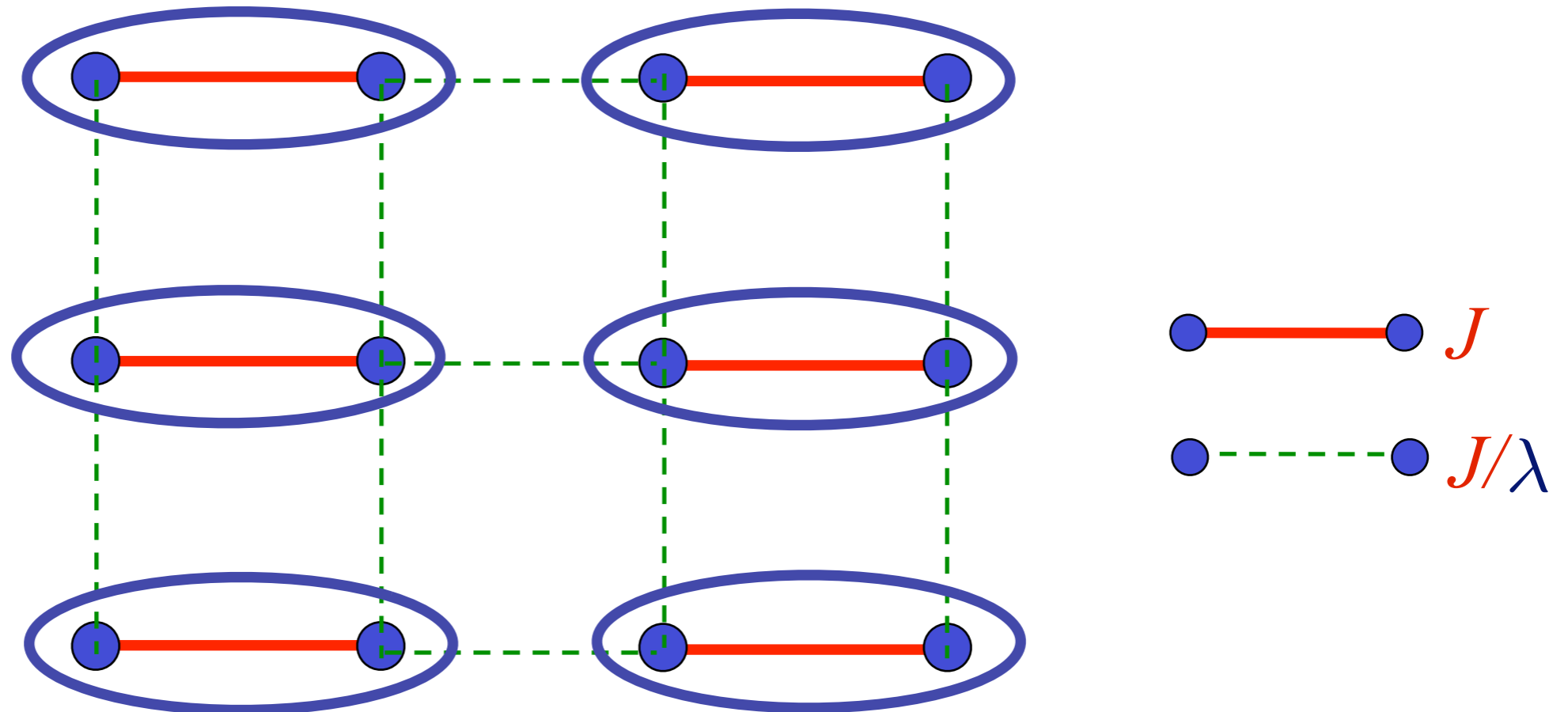


$$\text{Valence bond singlet} = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$

At large λ ground state is a “quantum paramagnet” with spins locked in valence bond singlets

Spinning electrons localized on a square lattice

$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

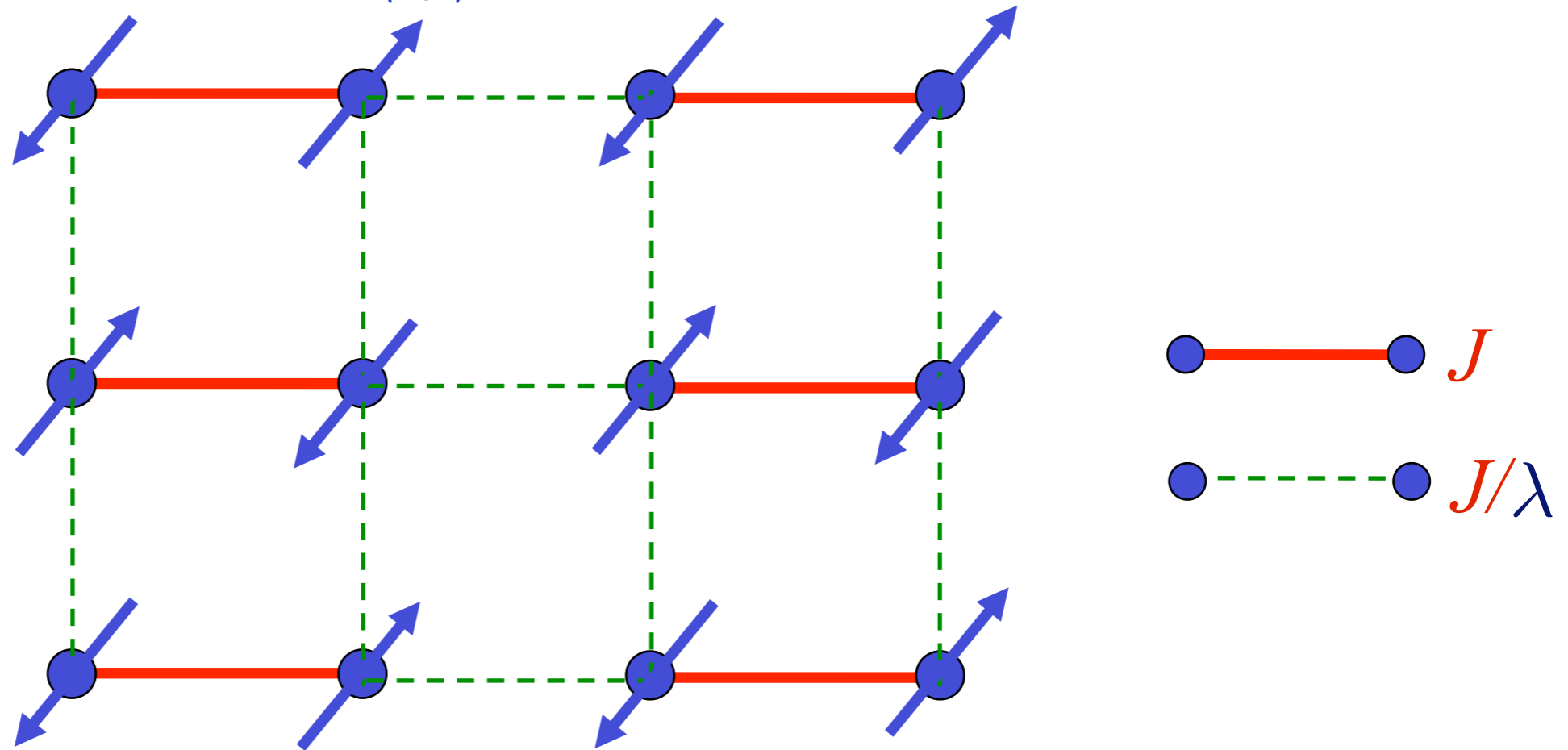


$$\text{[Pair of sites in a blue oval]} = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$

Nearest-neighbor spins are “entangled” with each other.
Can be separated into an Einstein-Podolsky-Rosen (EPR) pair.

Spinning electrons localized on a square lattice

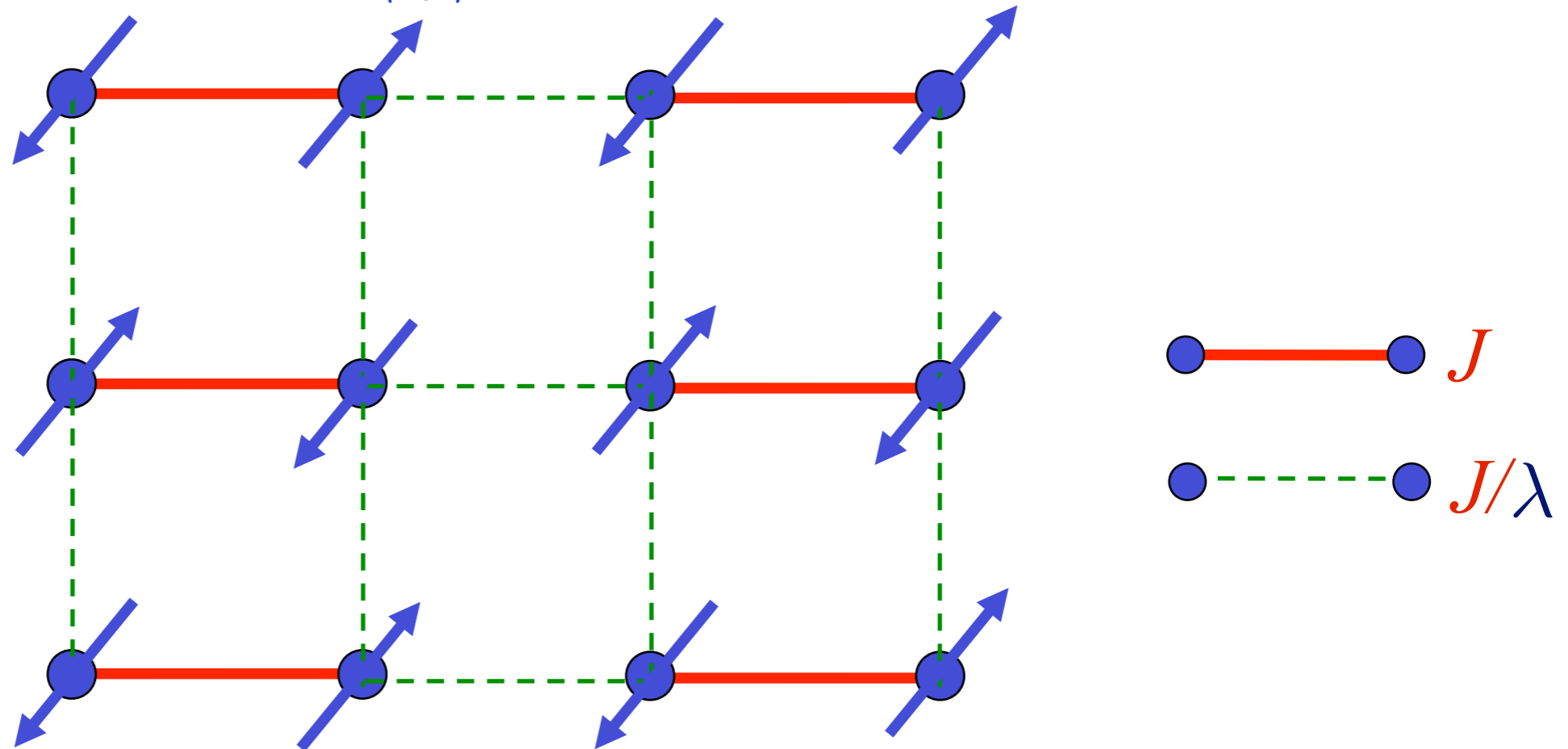
$$H = \sum_{\langle ij \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j$$



For $\lambda \approx 1$, the ground state has antiferromagnetic (“Néel”) order, and the spins align in a checkerboard pattern

Spinning electrons localized on a square lattice

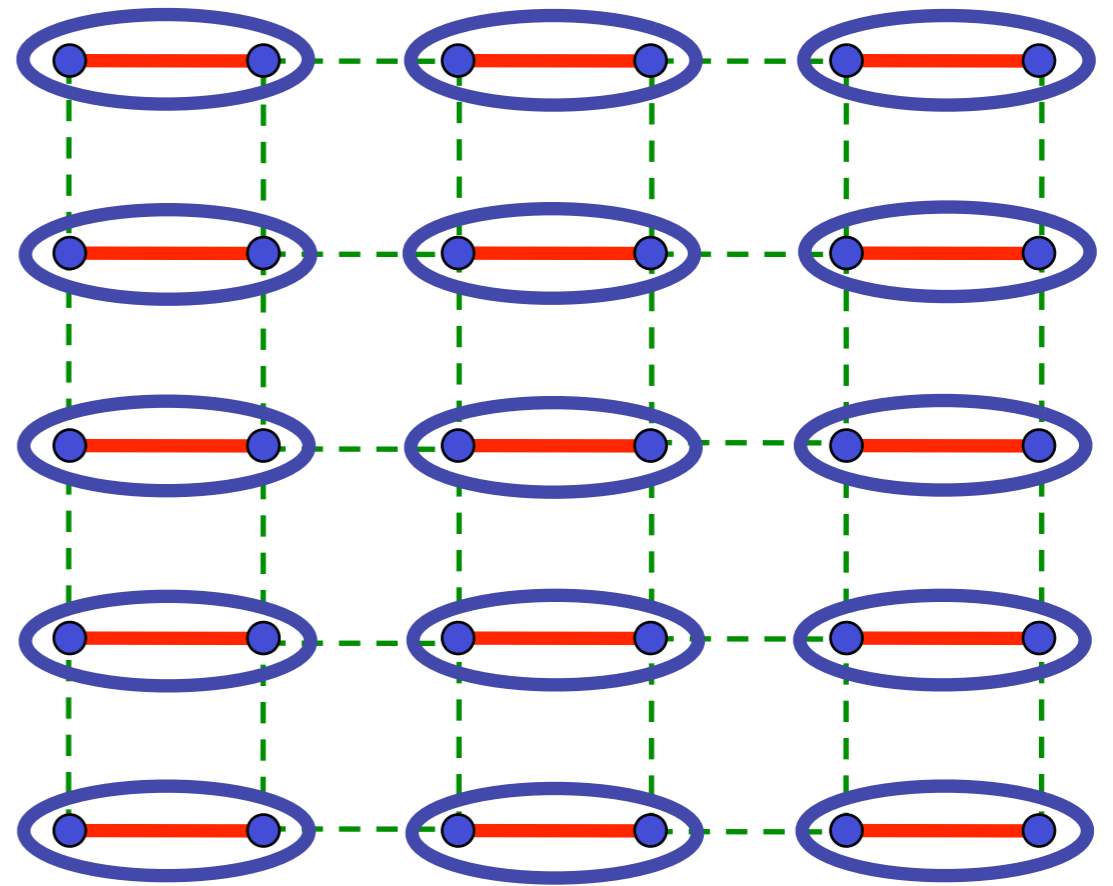
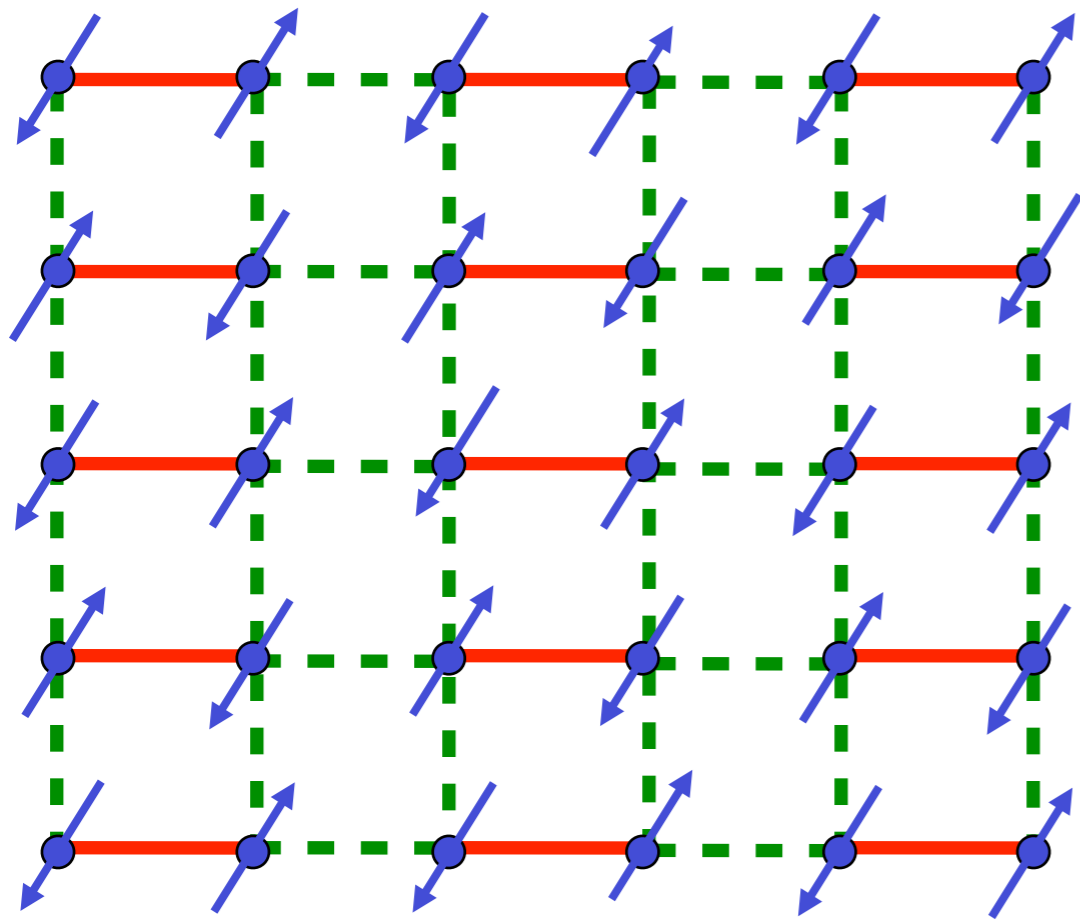
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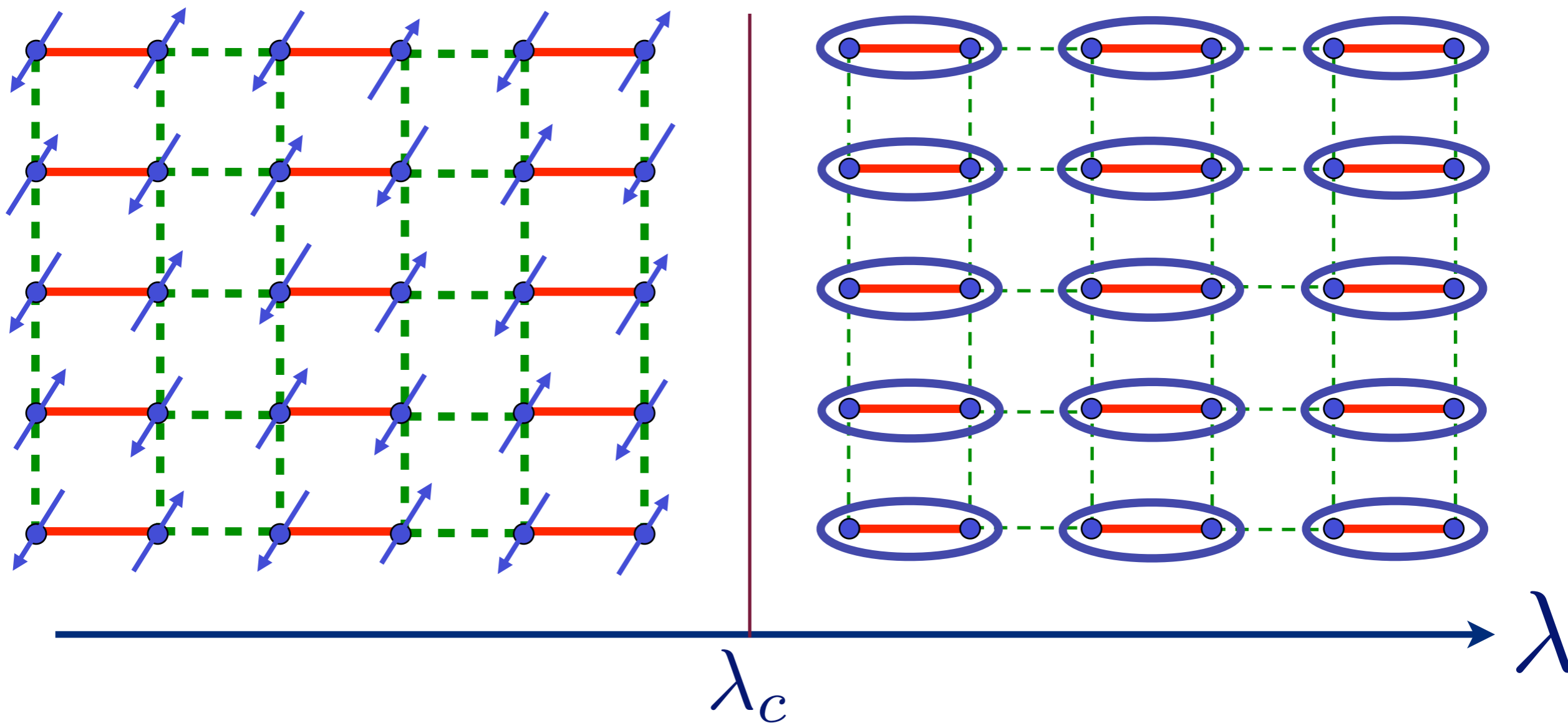
For $\lambda \approx 1$, the ground state has antiferromagnetic (“Néel”) order,
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No EPR pairs

$$\text{Diagram of two blue spheres connected by a red line, enclosed in a blue oval} = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$



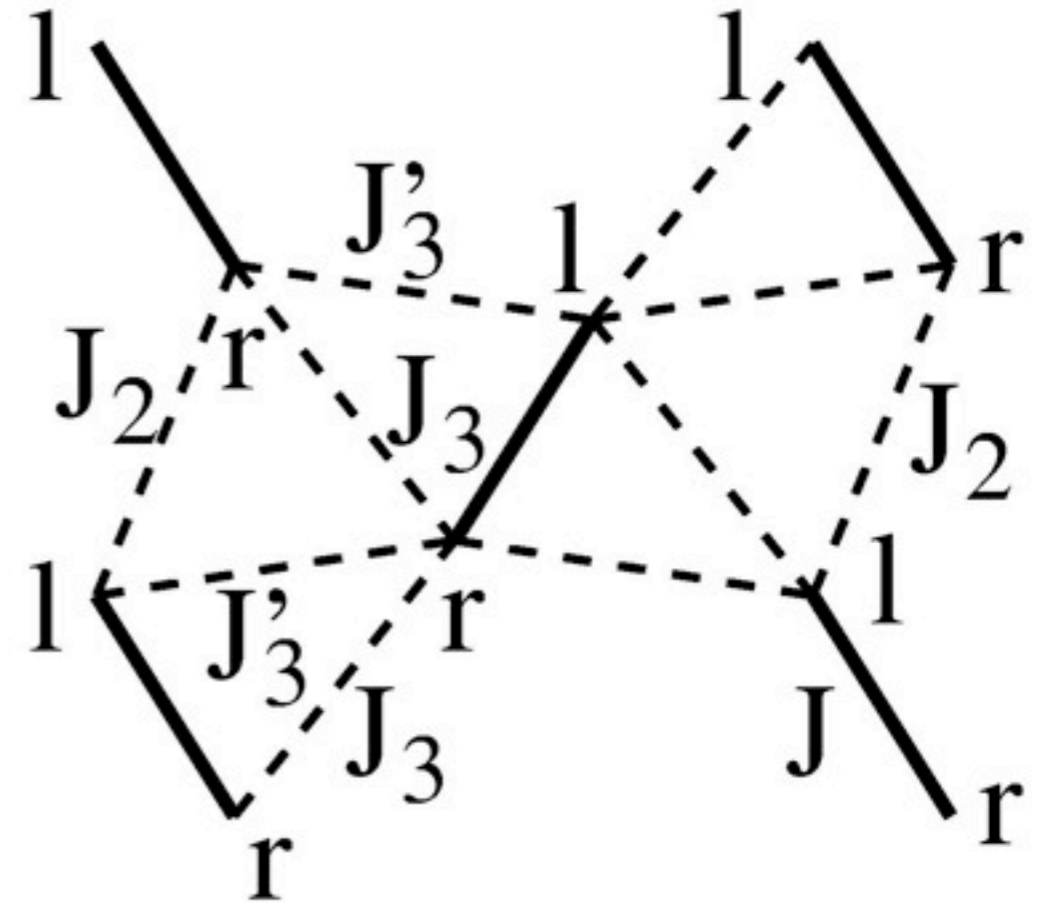
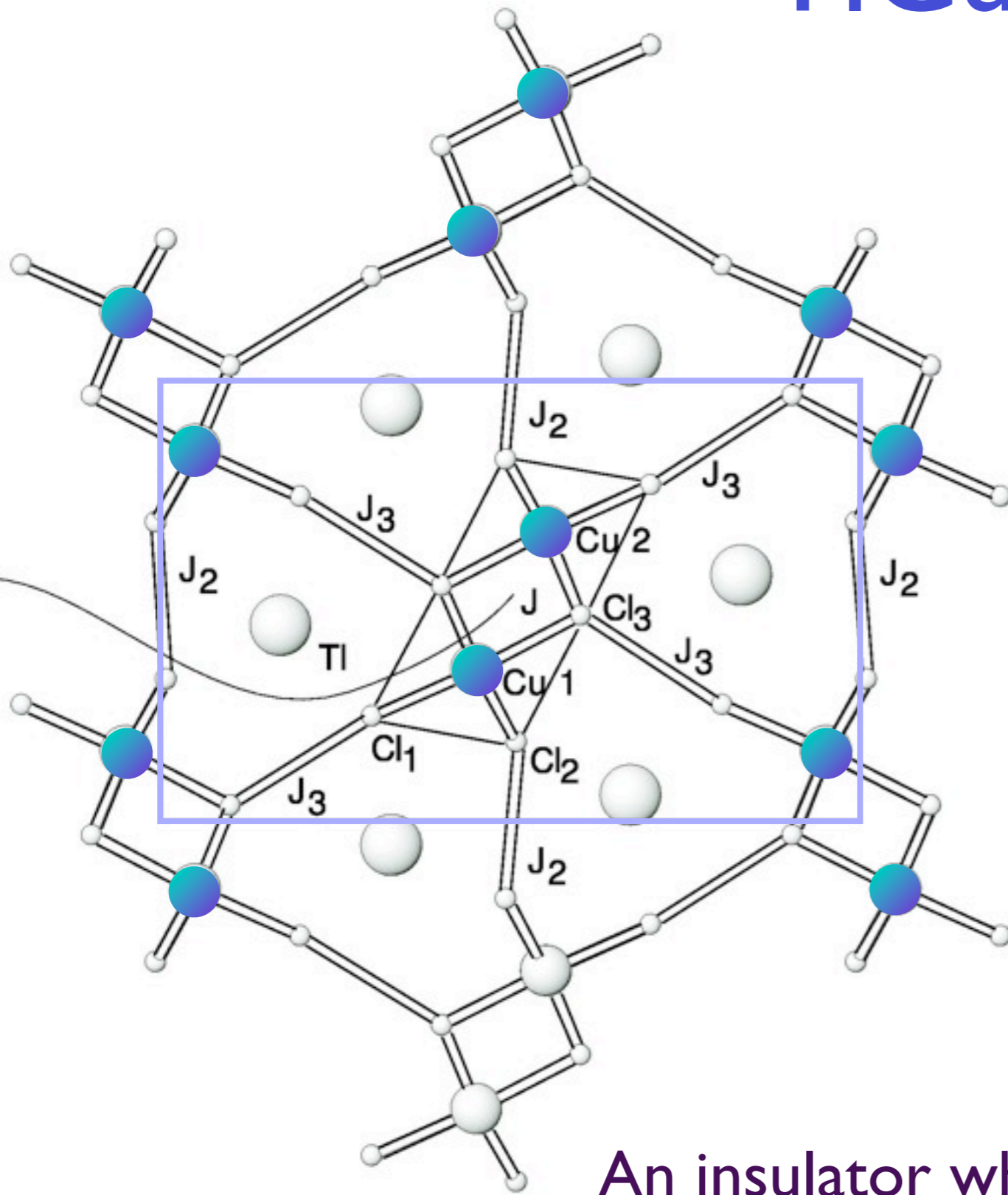
$$\text{Diagram of two blue spheres connected by a red line, enclosed in a blue oval} = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$



Pressure in TlCuCl_3

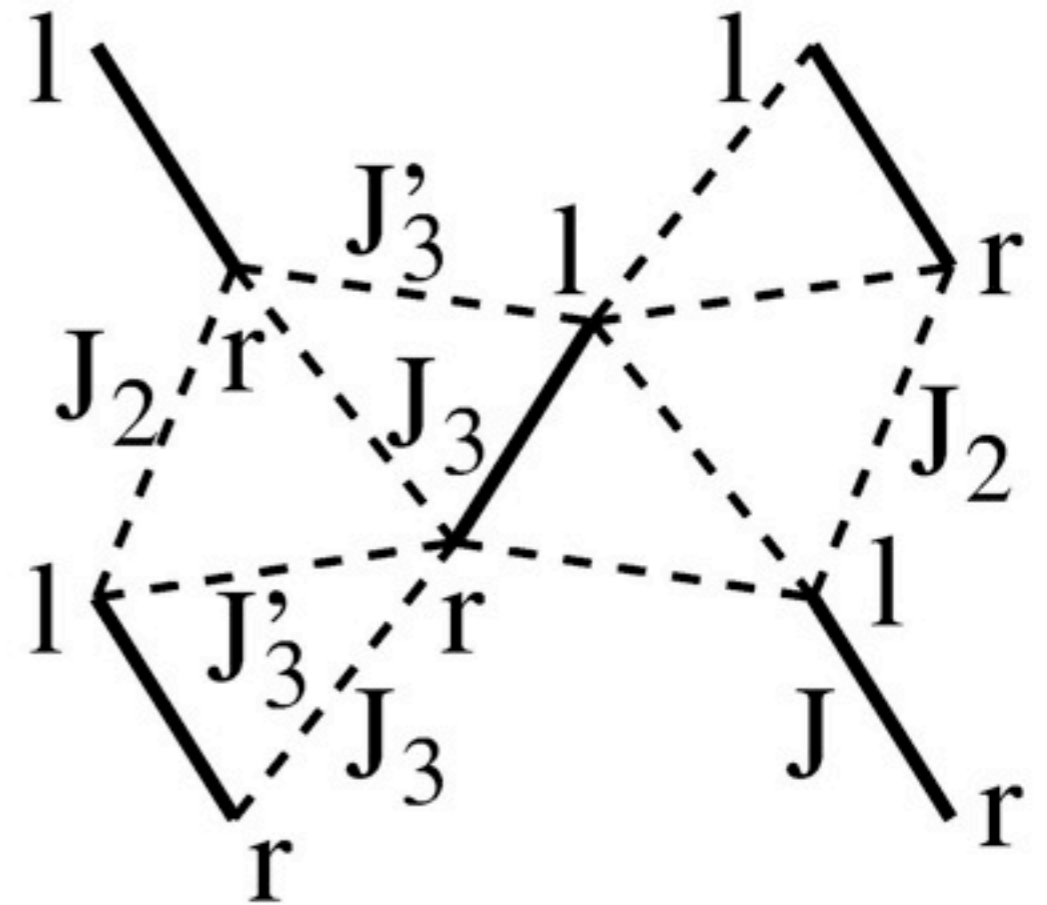
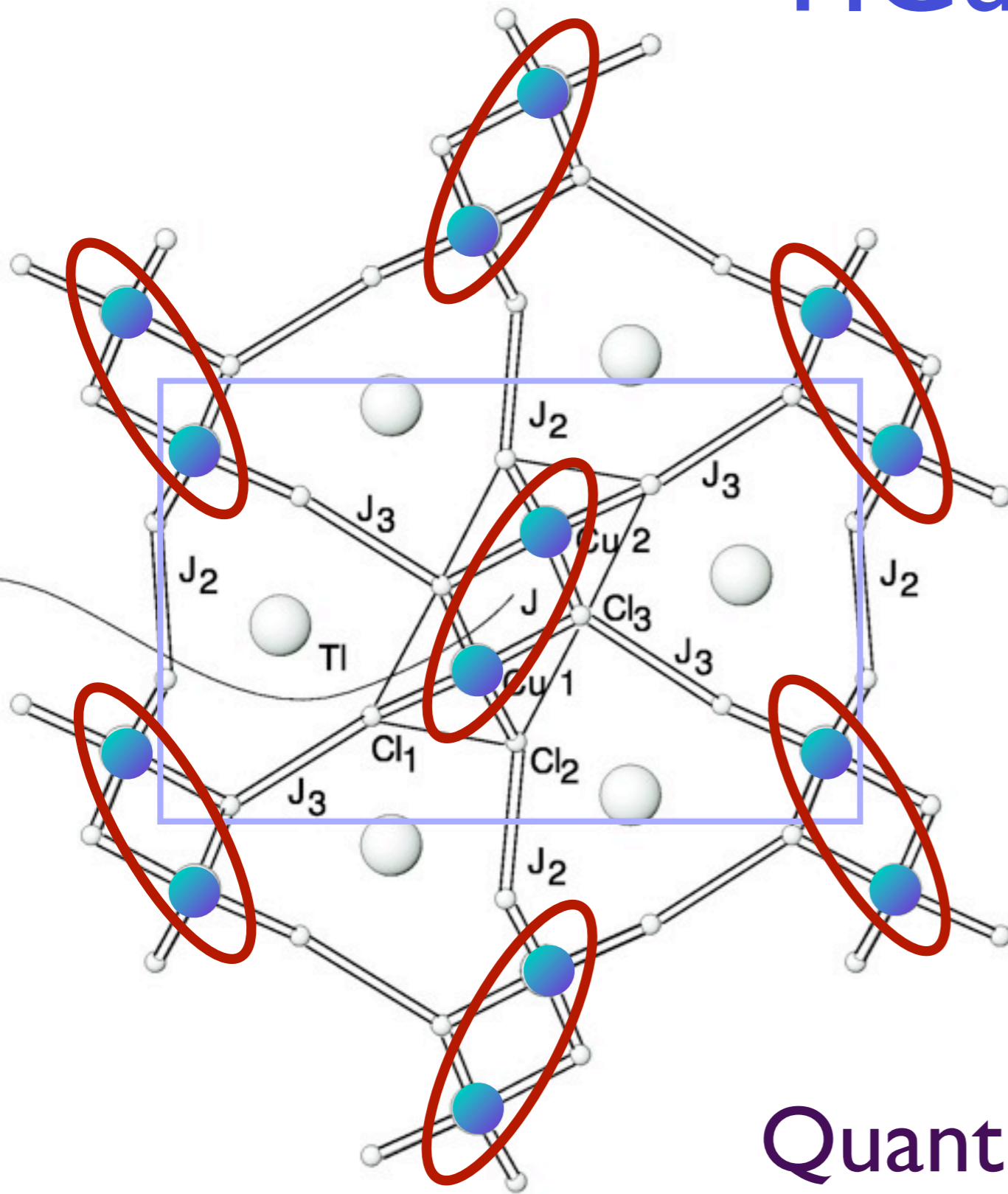
A. Oosawa, K. Kakurai, T. Osakabe, M. Nakamura, M. Takeda, and H. Tanaka,
Journal of the Physical Society of Japan, **73**, 1446 (2004).

TiCuCl₃



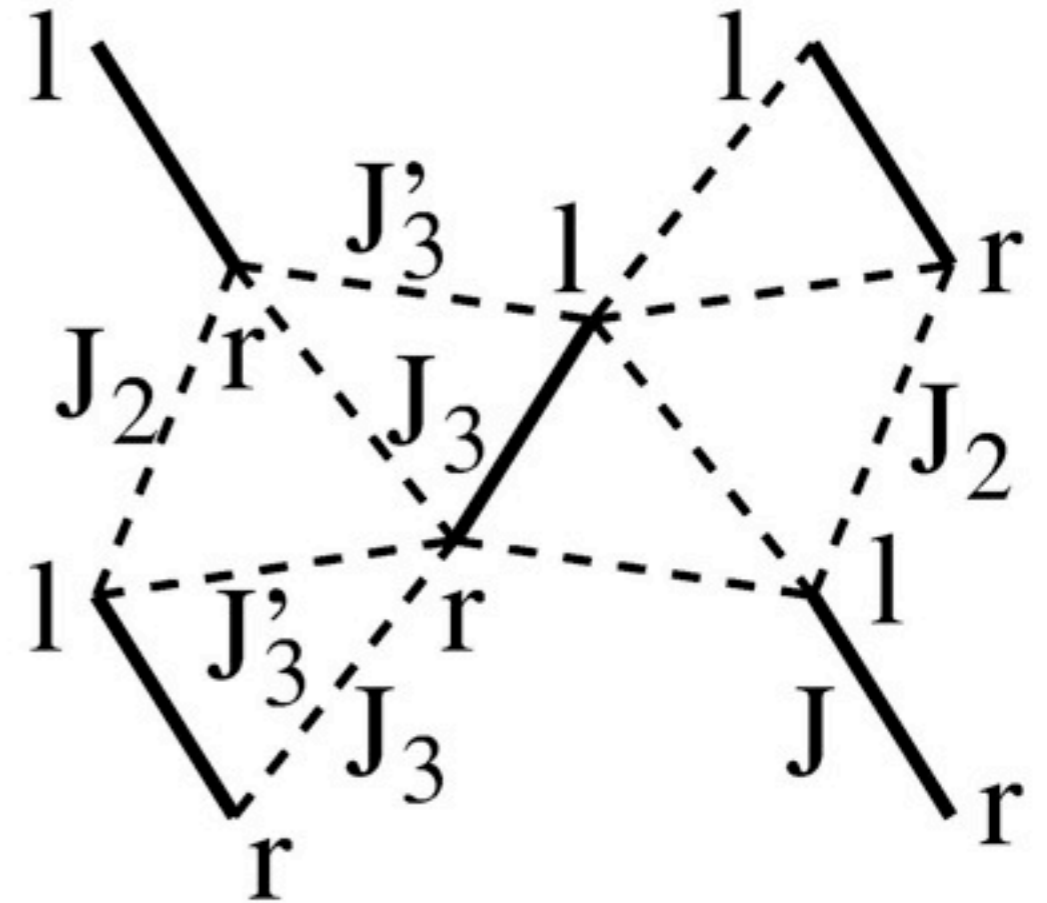
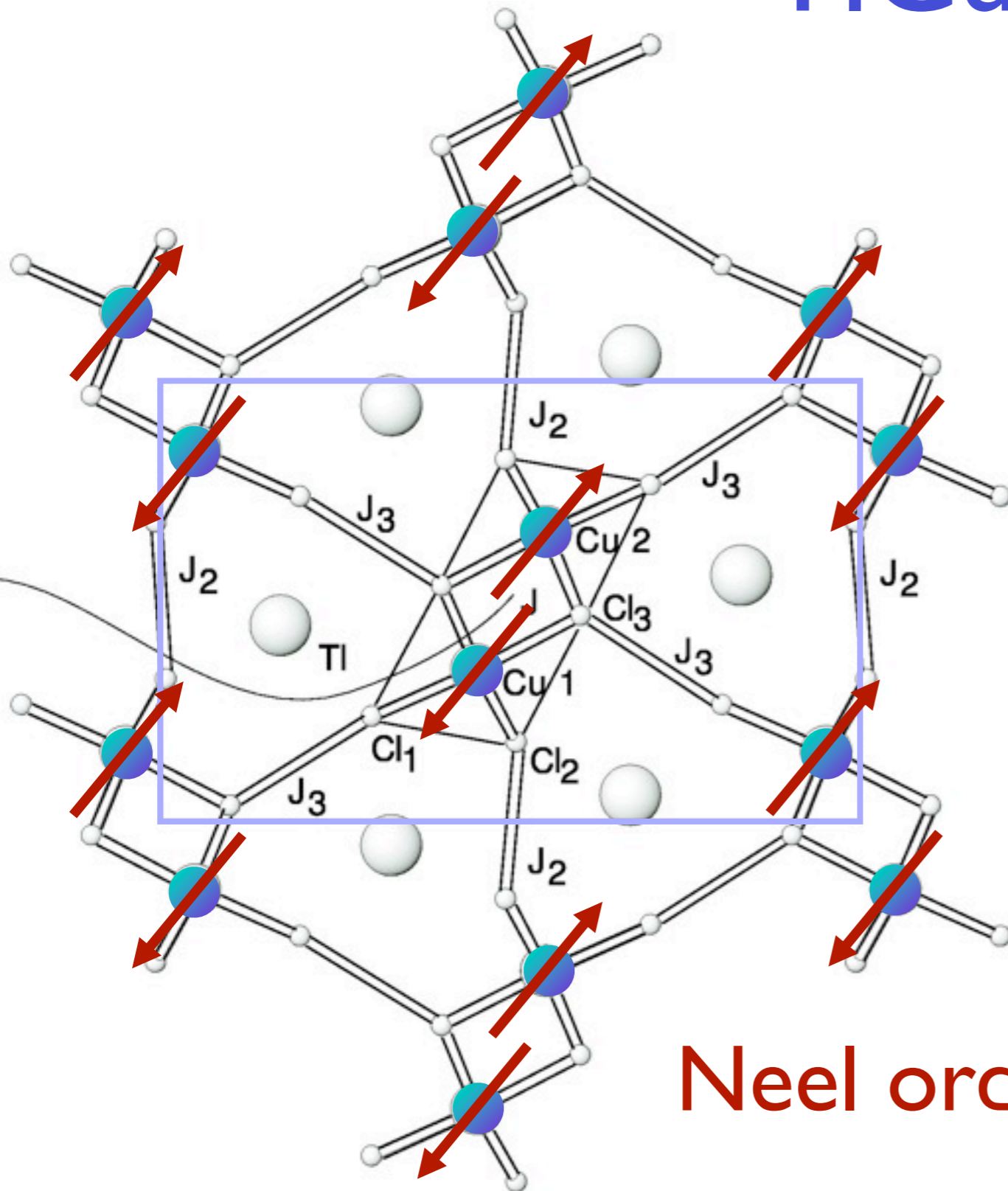
An insulator whose spin susceptibility vanishes exponentially as the temperature T tends to zero.

TlCuCl₃



Quantum paramagnet at
ambient pressure

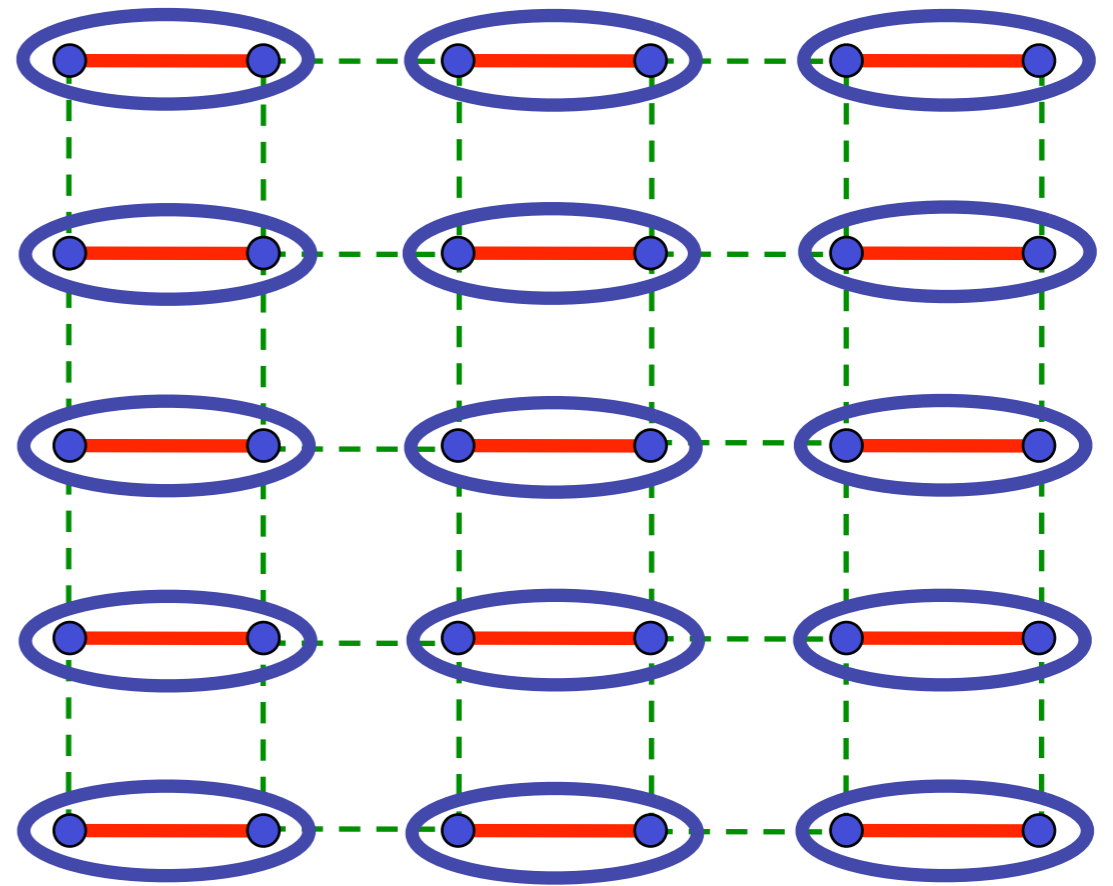
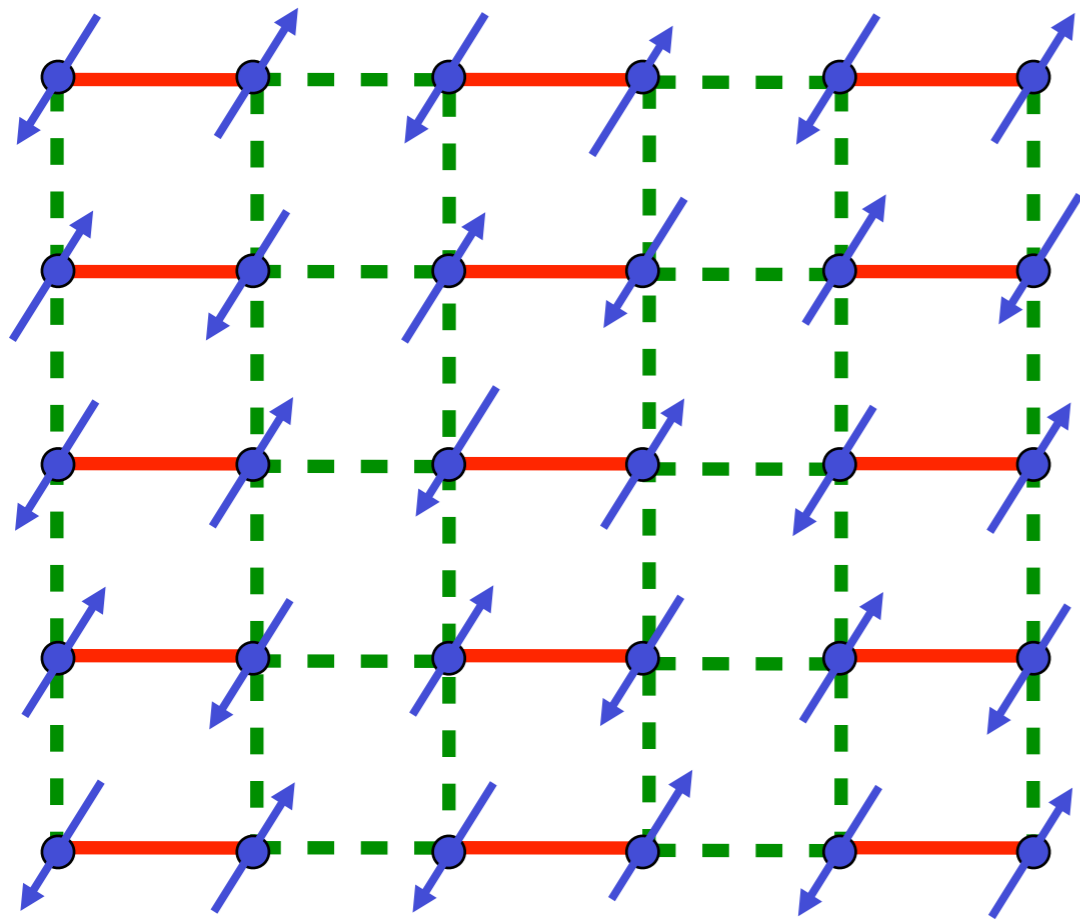
TlCuCl₃



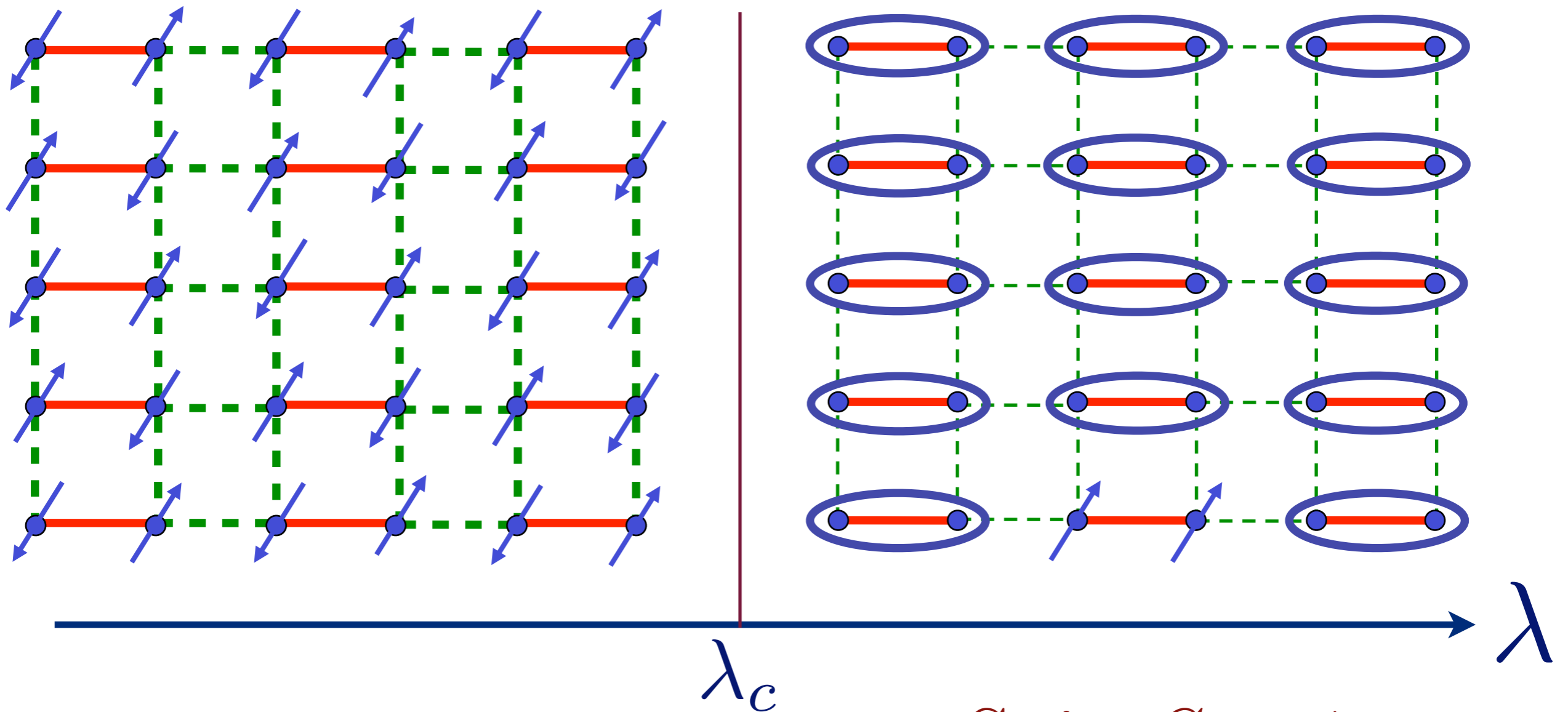
Neel order under pressure

A. Oosawa, K. Kakurai, T. Osakabe, M. Nakamura, M. Takeda, and H. Tanaka, *Journal of the Physical Society of Japan*, **73**, 1446 (2004).

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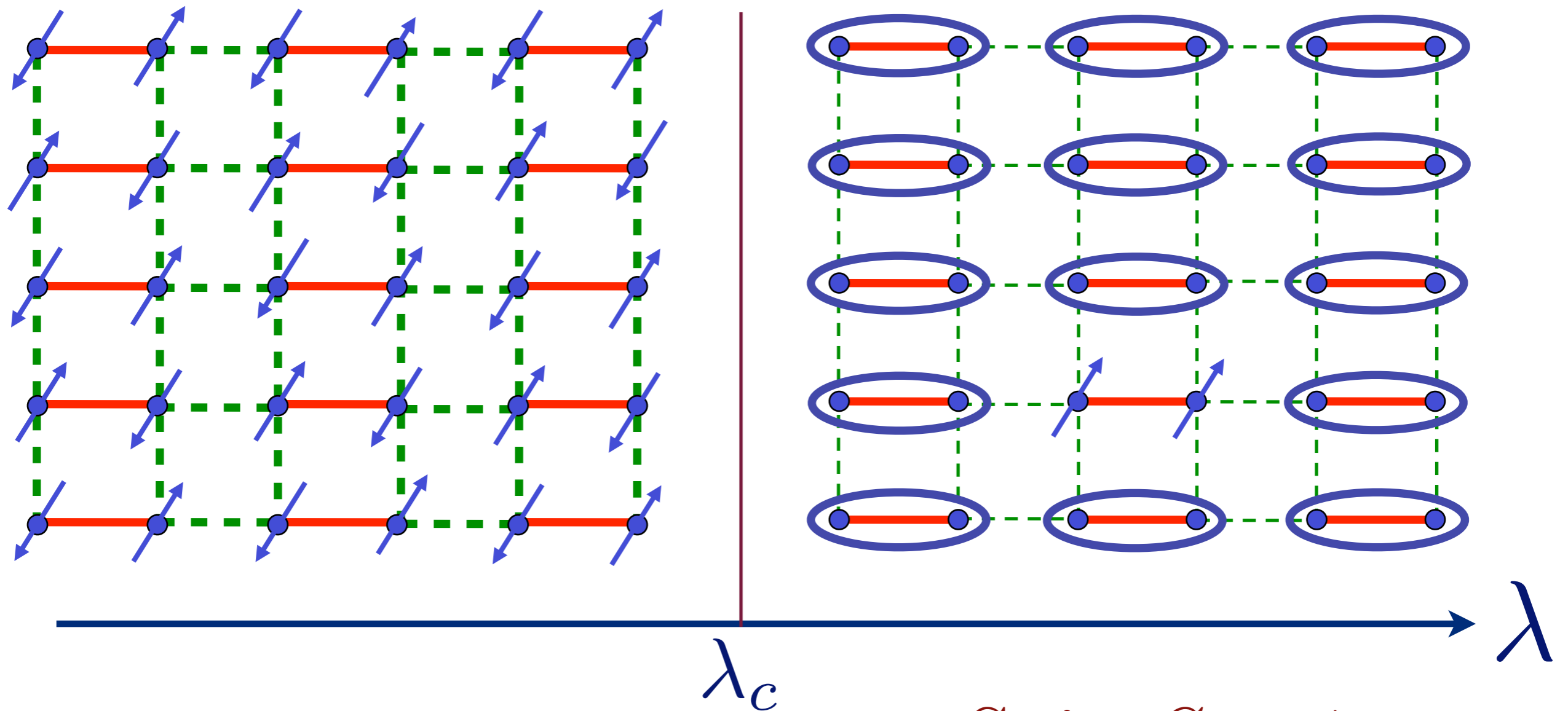


Excitation spectrum in the paramagnetic phase



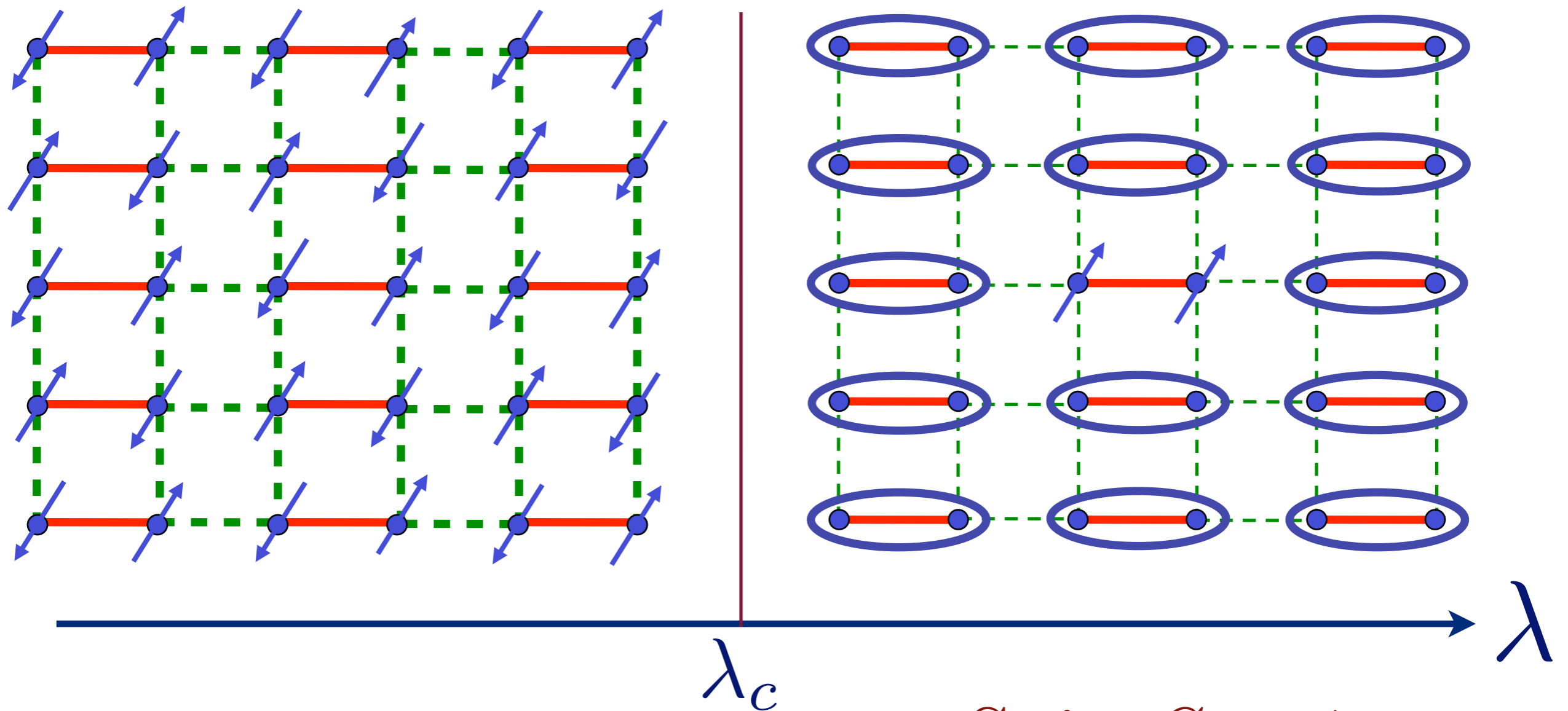
Spin $S = 1$
“triplon”

Excitation spectrum in the paramagnetic phase



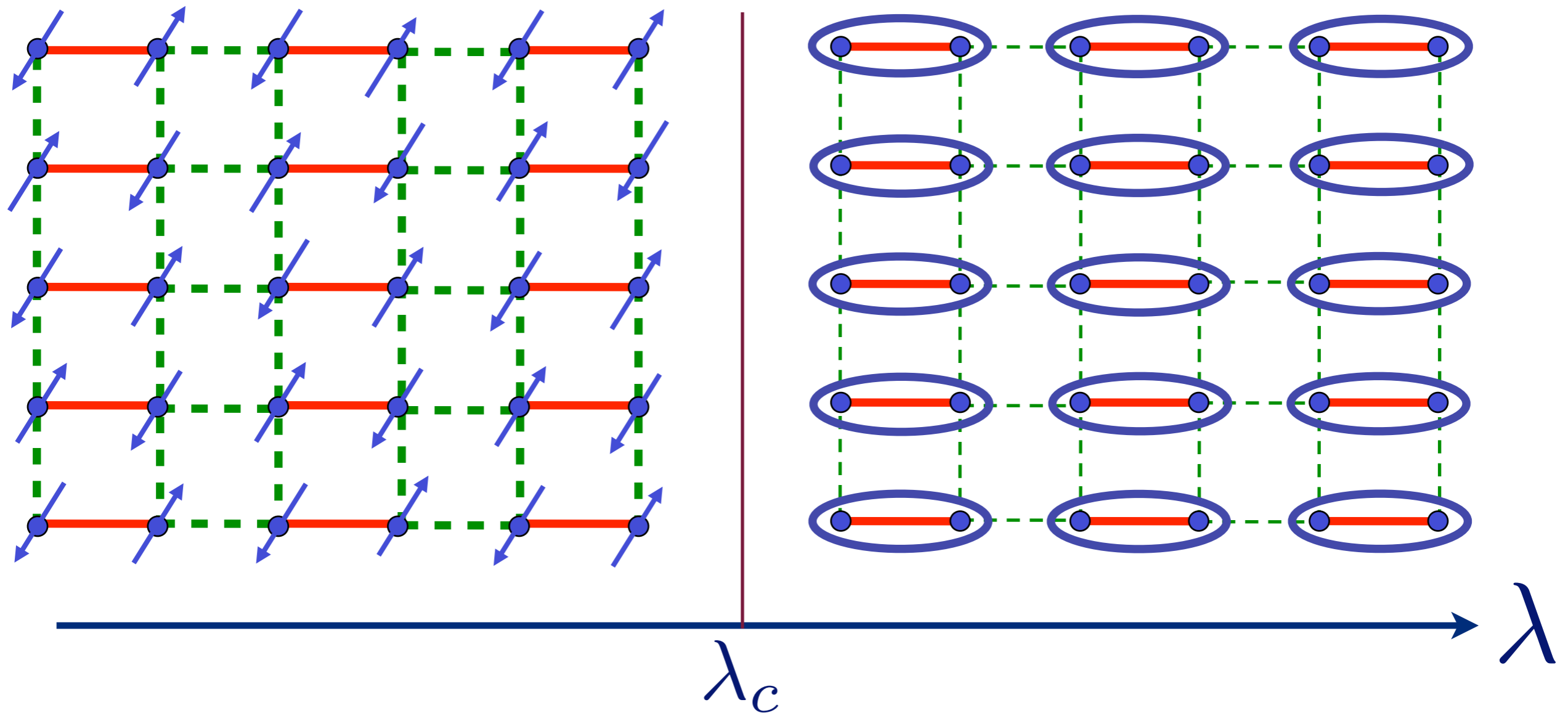
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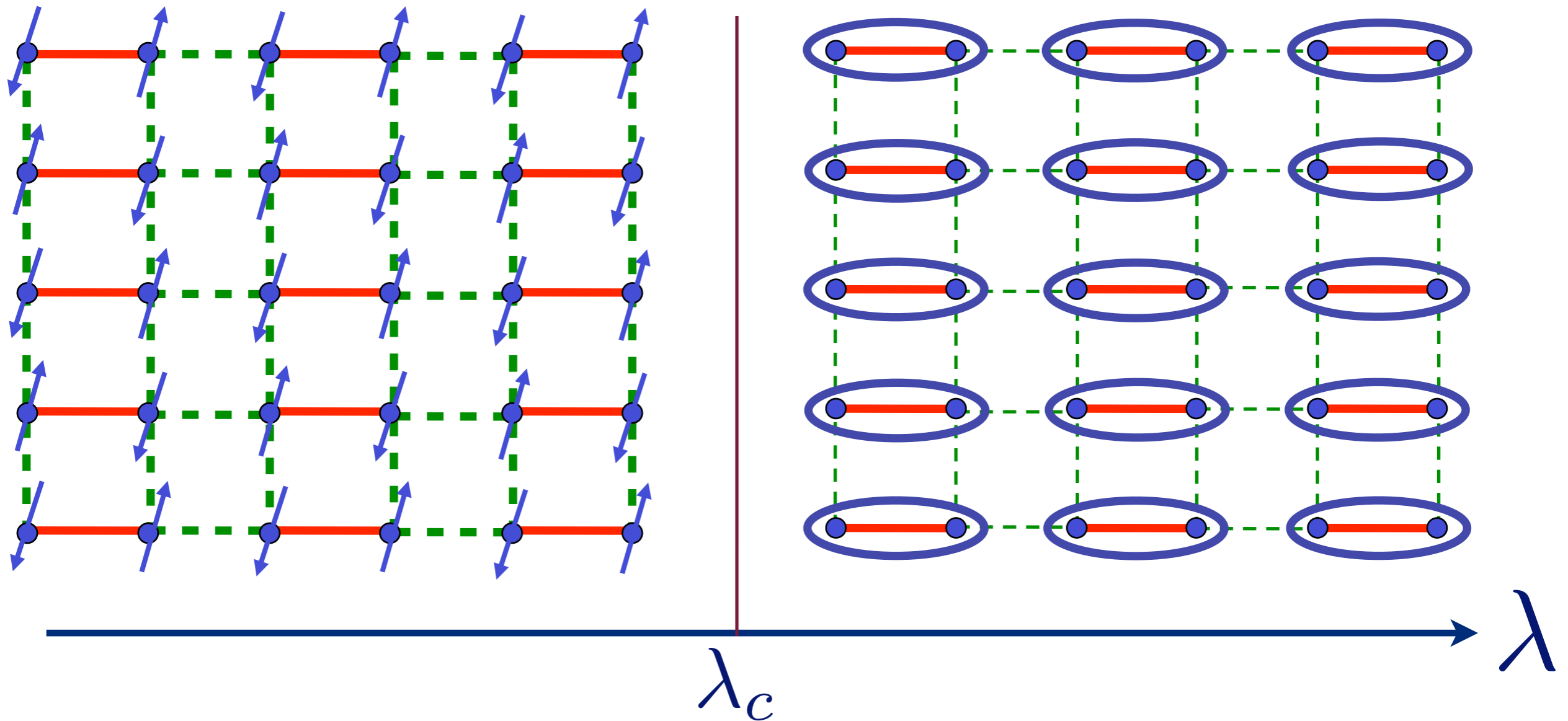
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Excitation spectrum in the Néel phase



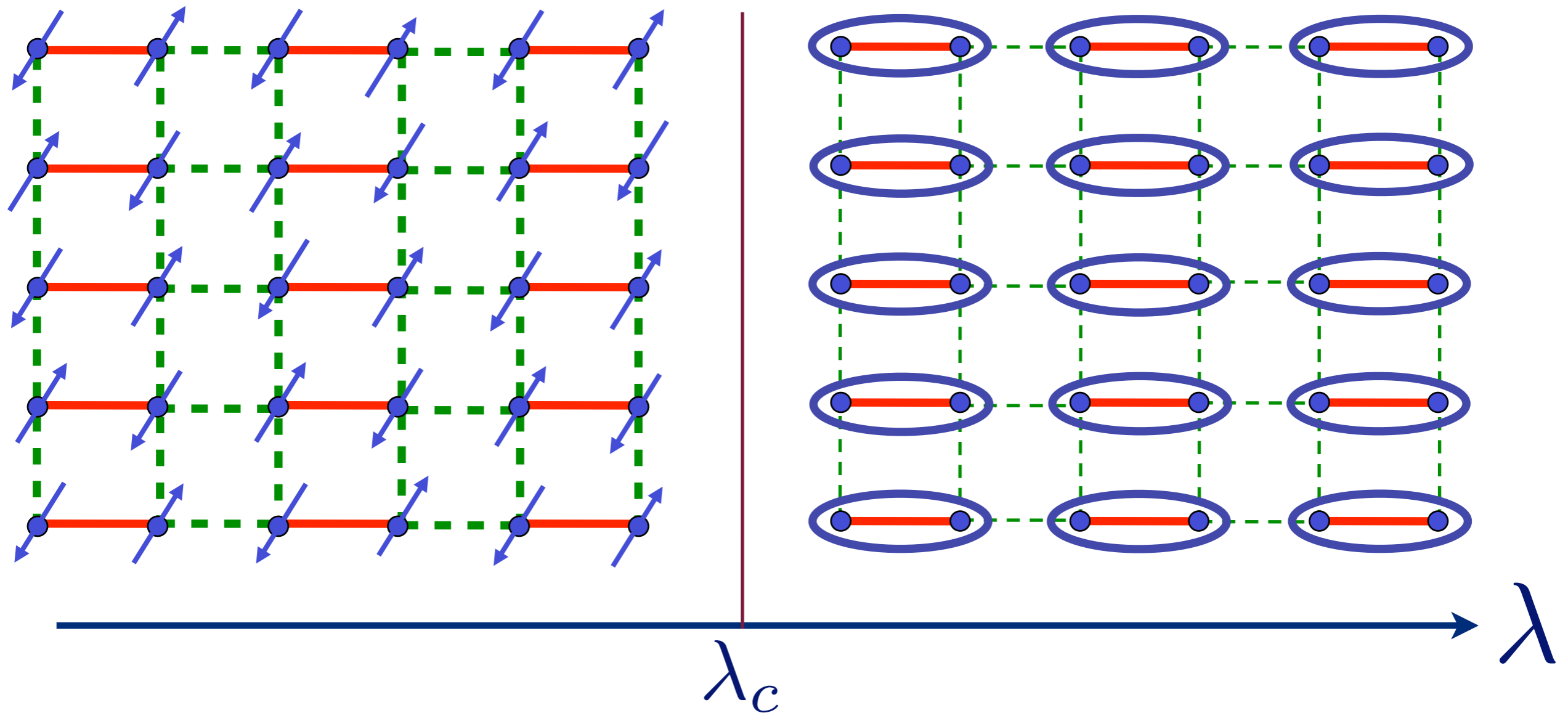
Spin waves

Excitation spectrum in the Néel phase



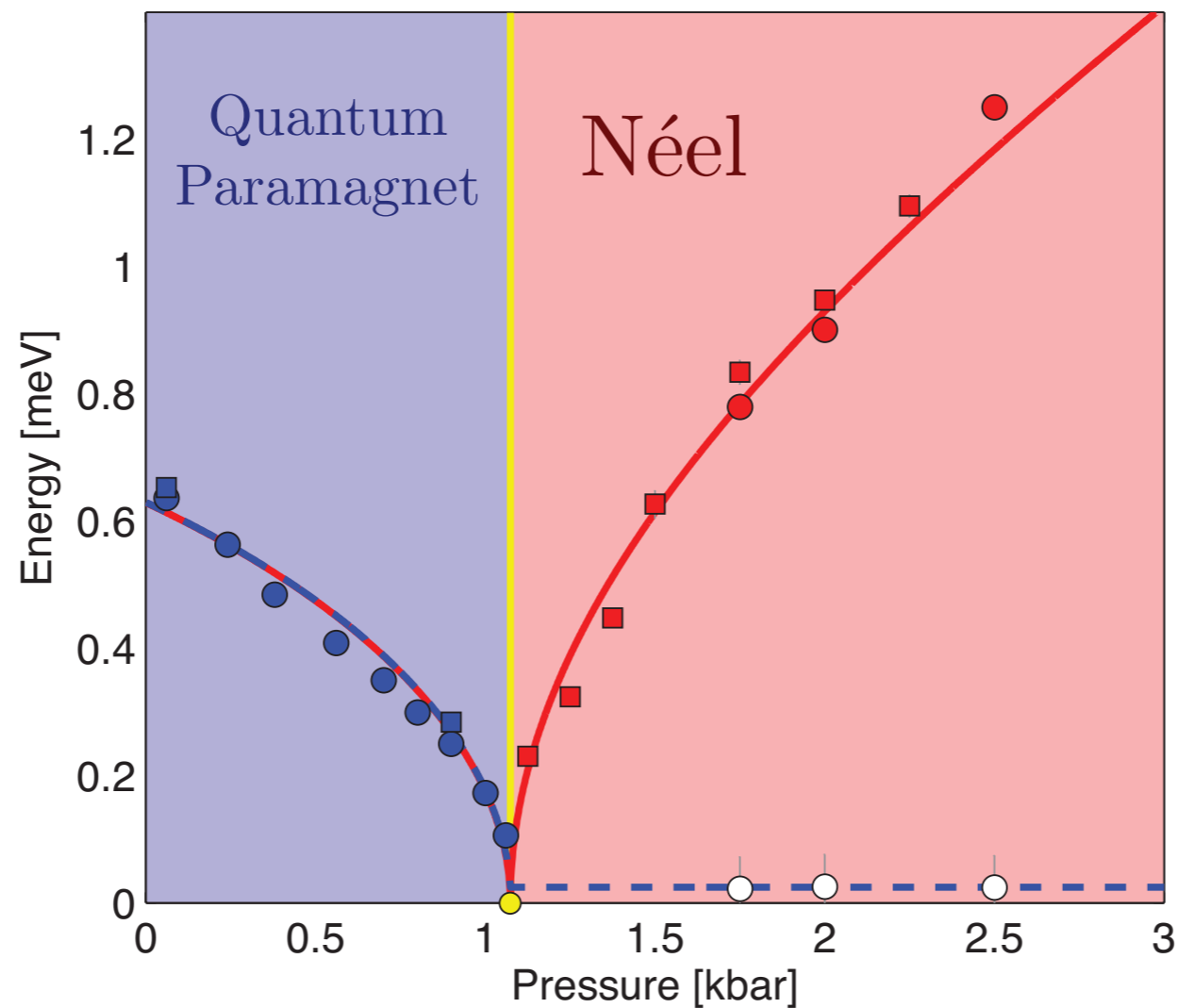
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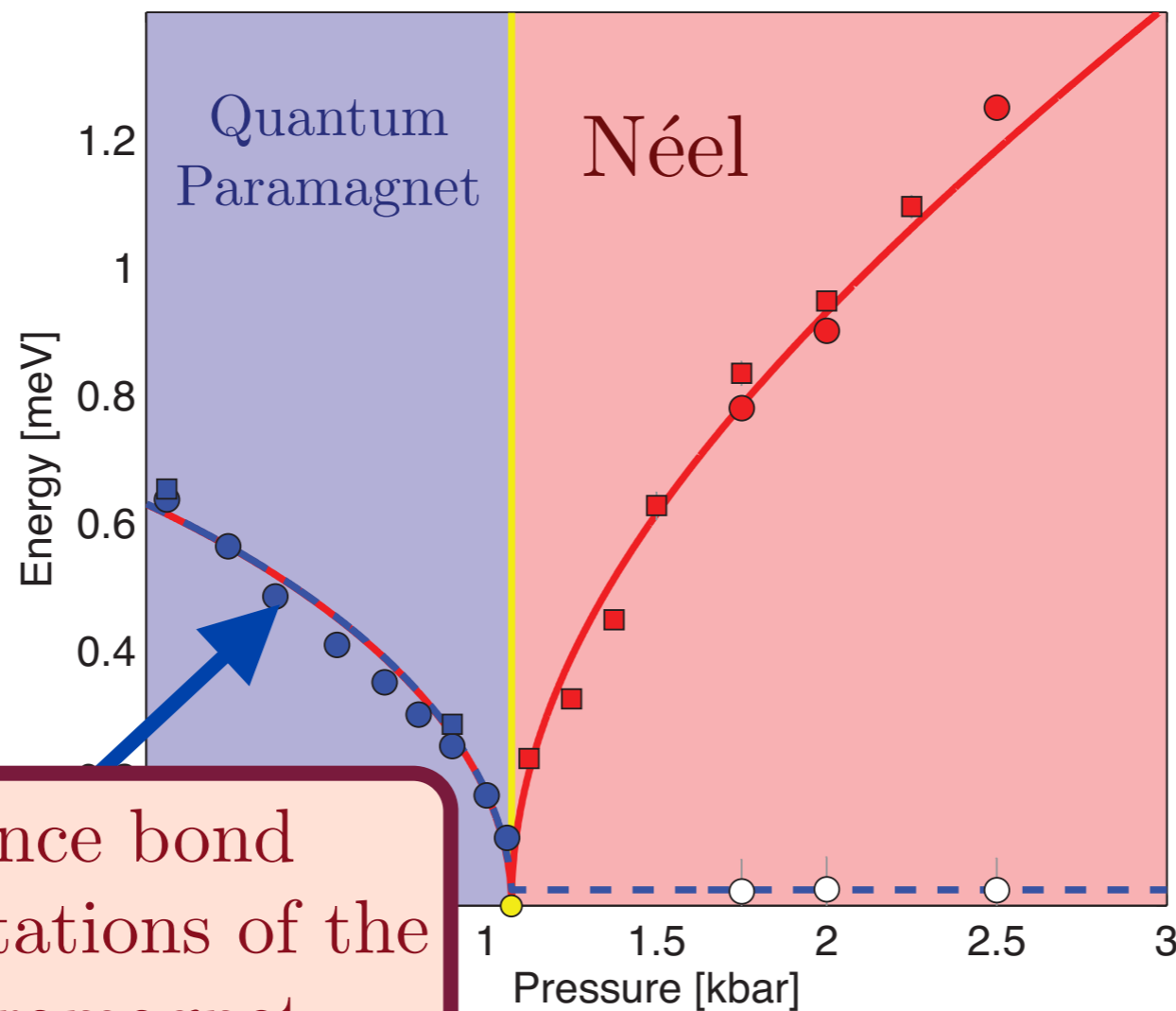
Spin waves

Excitations of TlCuCl_3 with varying pressure



Christian Ruedg, Bruce Normand, Masashige Matsumoto, Albert Furrer, Desmond McMorro, Karl Kramer, Hans-Ulrich Gudel, Severian Gvasaliya, Hannu Mutka, and Martin Boehm, *Phys. Rev. Lett.* **100**, 205701 (2008)

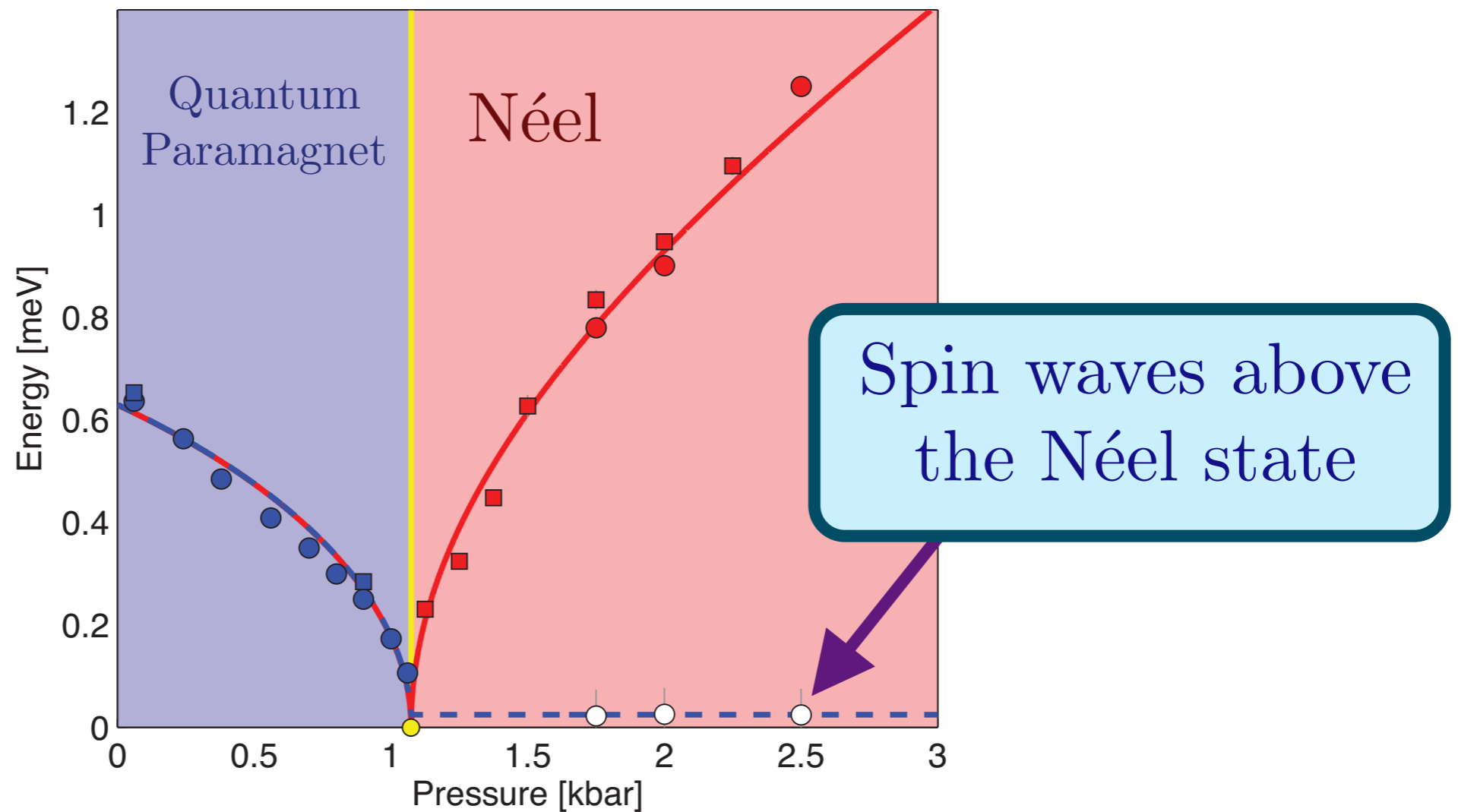
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Broken valence bond (“triplon”) excitations of the quantum paramagnet

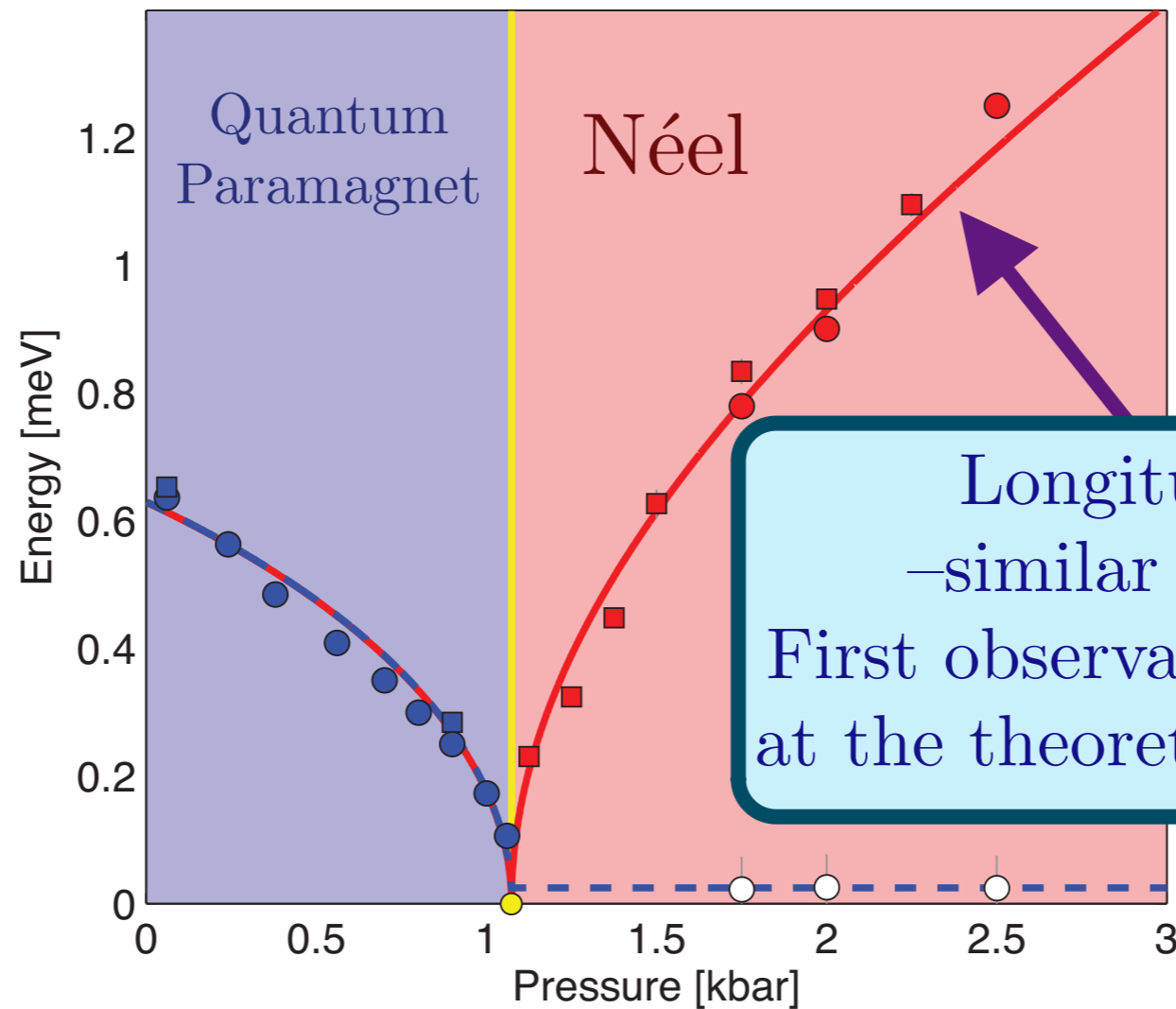
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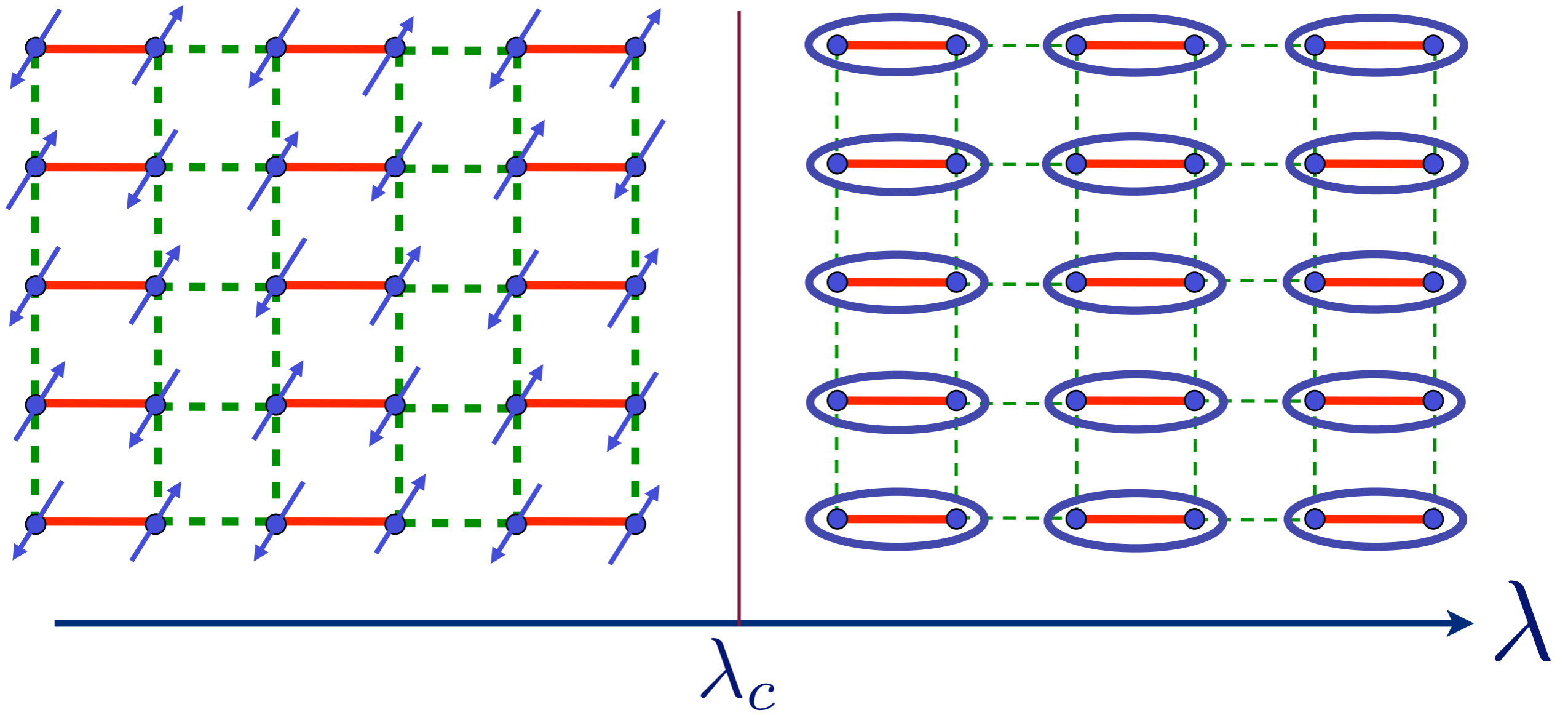


Longitudinal excitations
–similar to the Higgs boson
First observation of the Higgs boson
at the theoretically predicted energy!

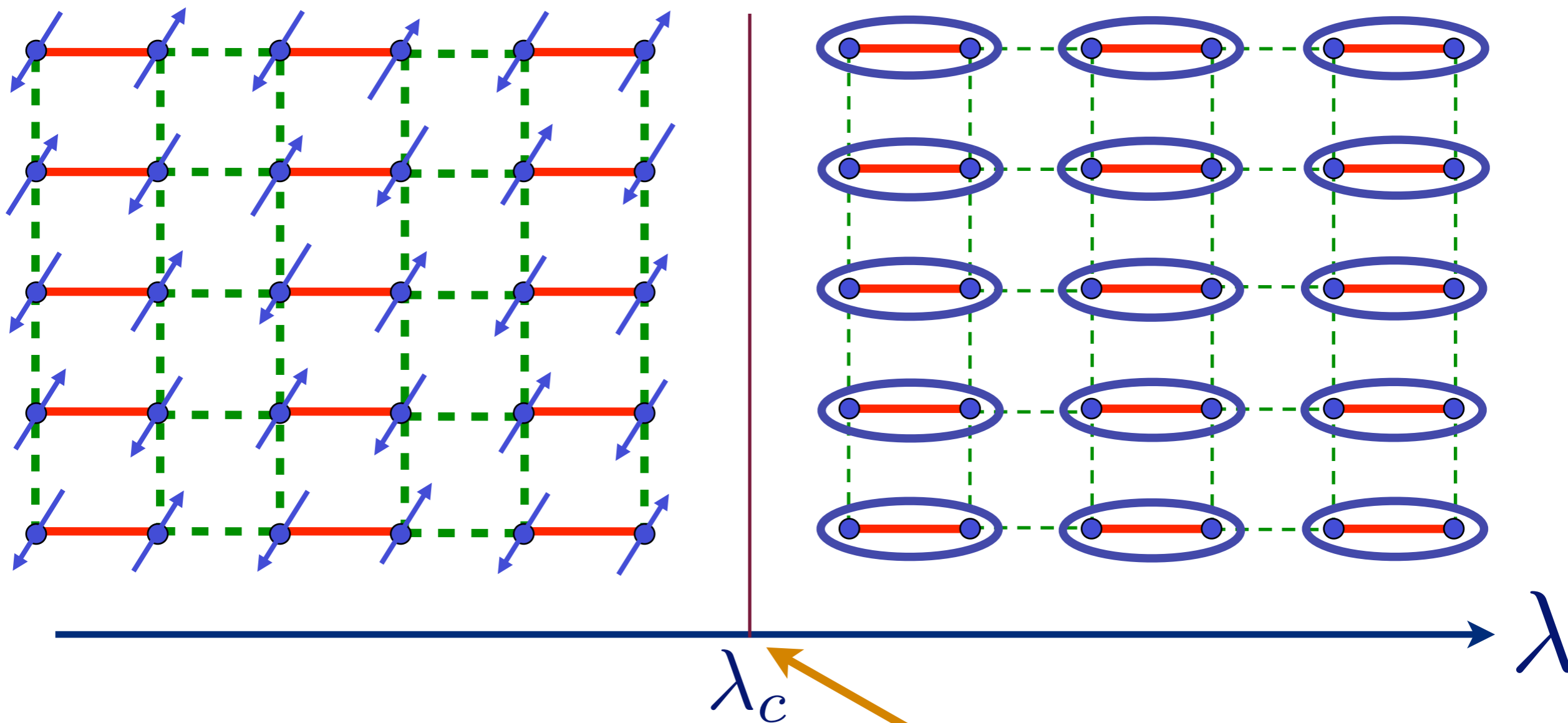
S. Sachdev,
Solvay conference,
arXiv:0901.4103

Christian Ruegg, Bruce Normand, Masashige Matsumoto, Albert Furrer, Desmond McMorrow, Karl Kramer, Hans-Ulrich Gudel, Severian Gvasaliya, Hannu Mutka, and Martin Boehm, *Phys. Rev. Lett.* **100**, 205701 (2008)

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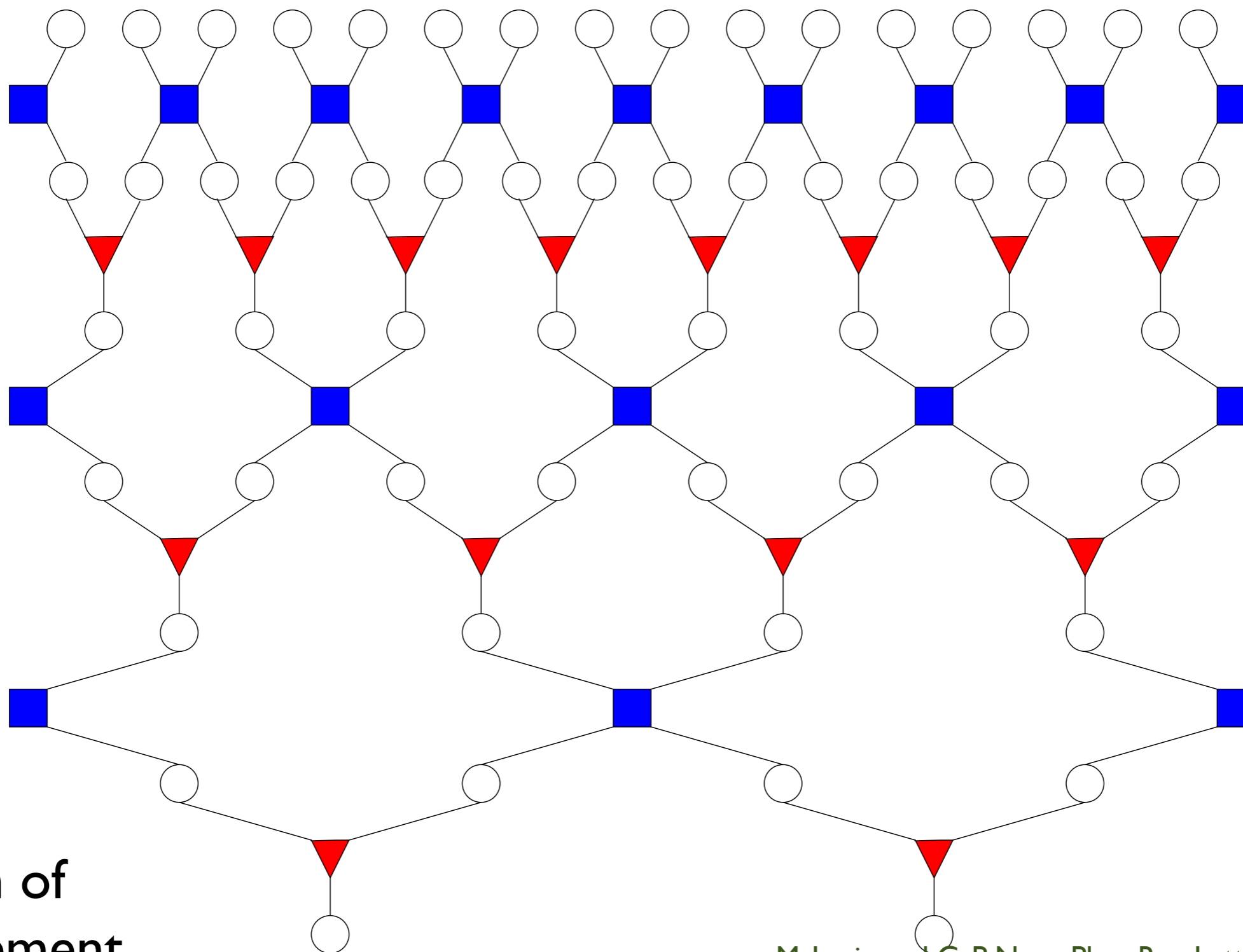
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Quantum critical point with non-local entanglement in spin wavefunction

Tensor network representation of entanglement at quantum critical point

D -dimensional
space



depth of
entanglement

M. Levin and C. P. Nave, Phys. Rev. Lett. 99, 120601 (2007)
F. Verstraete, M. M. Wolf, D. Perez-Garcia, and J. I. Cirac, Phys. Rev. Lett. 96, 220601 (2006)

Characteristics of quantum critical point

- Long-range entanglement

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**Quantum critical
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**String theory
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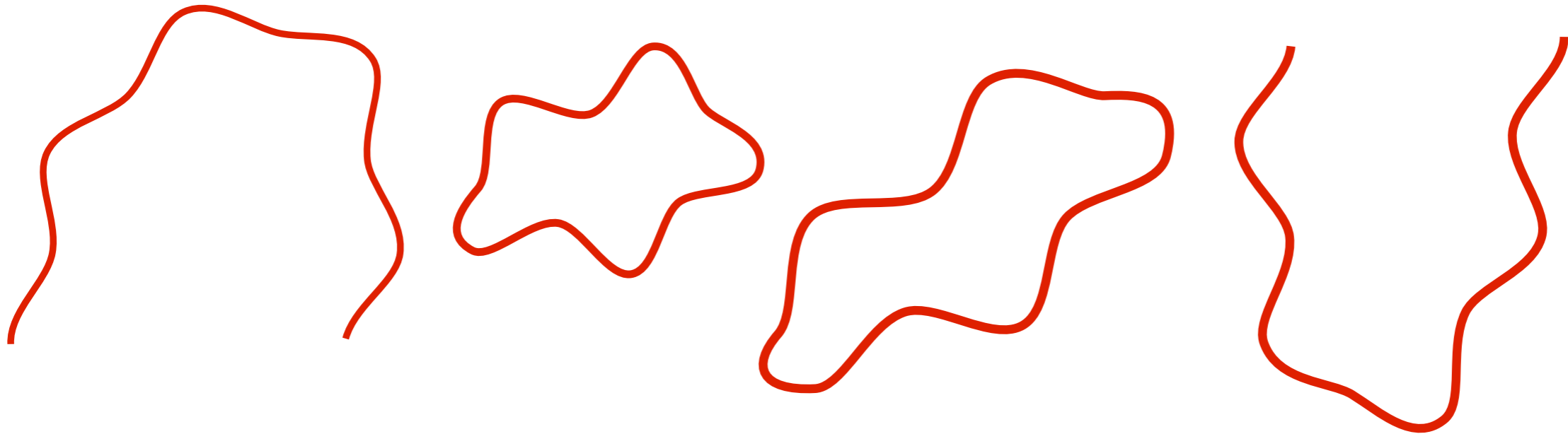
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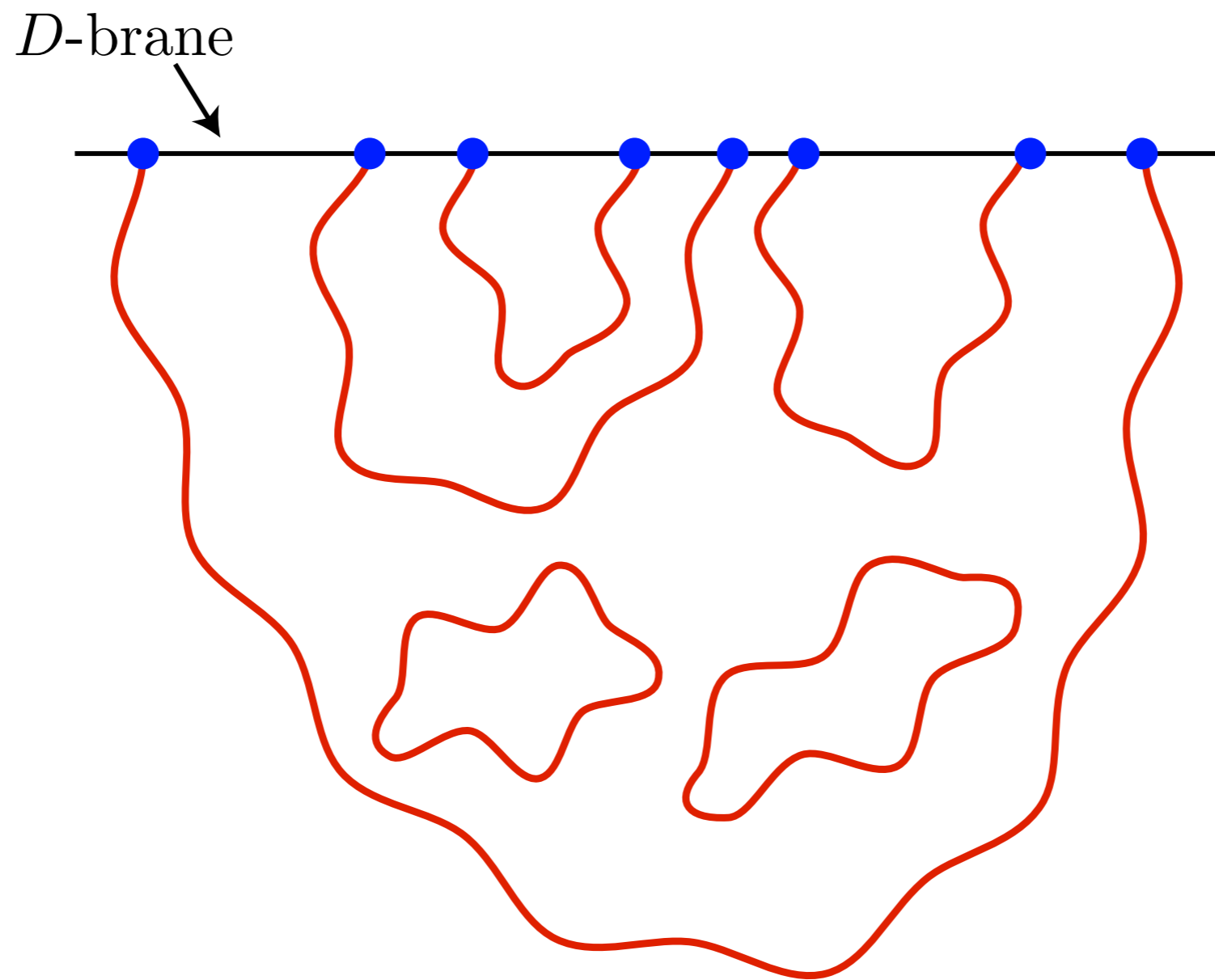
**String theory
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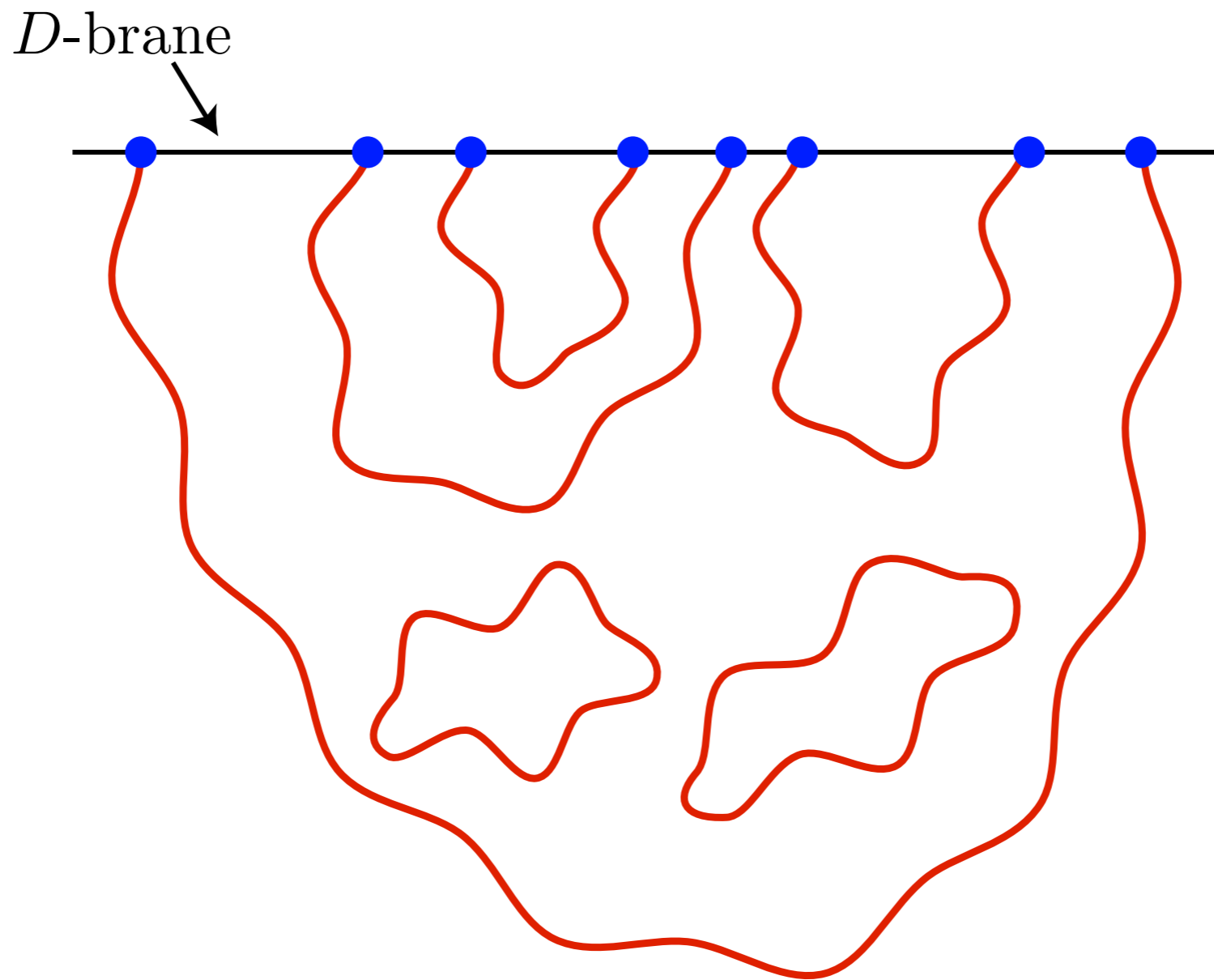
String theory



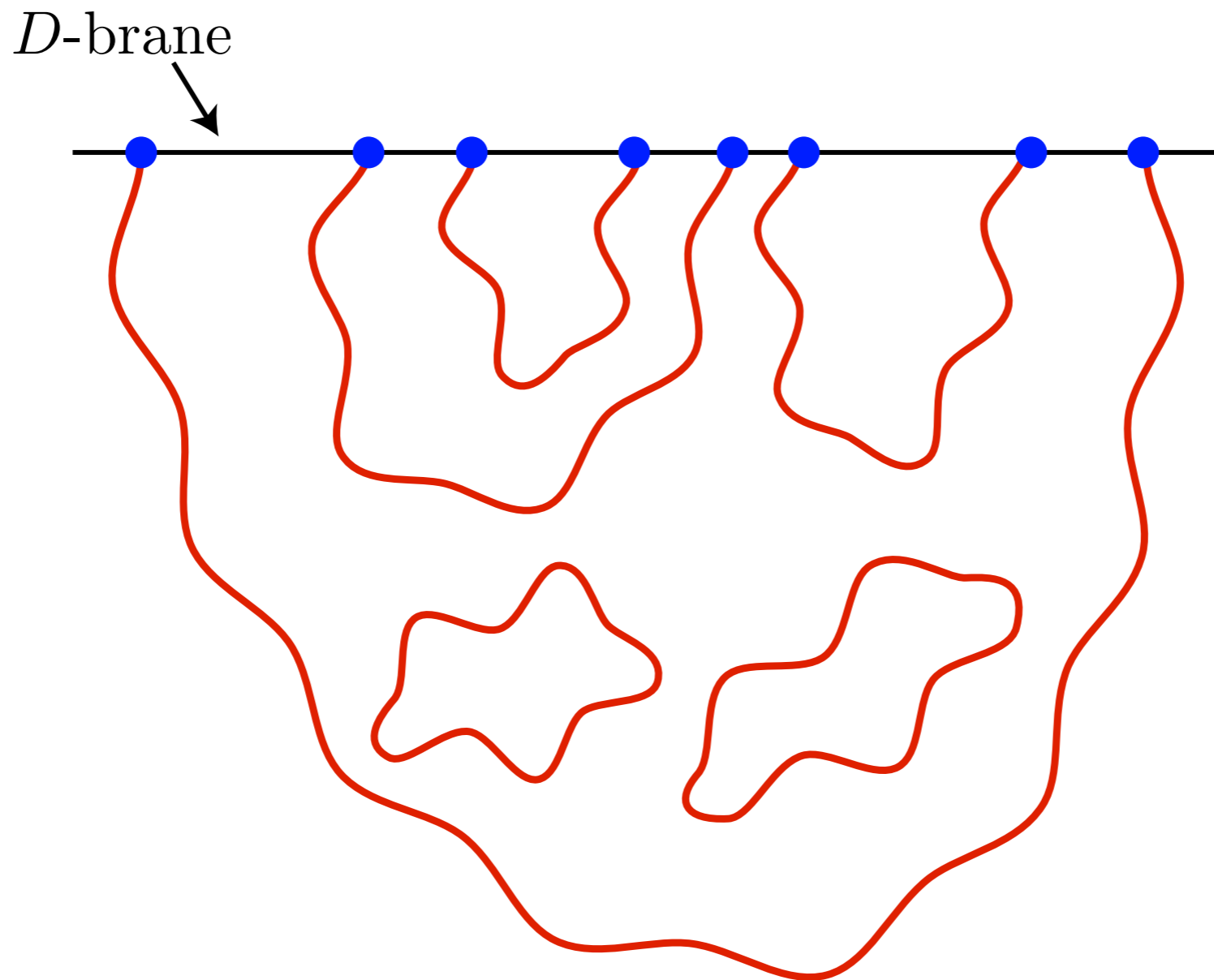
- Allows unification of the standard model of particle physics with gravity.
- Low-lying string modes correspond to gauge fields, gravitons, quarks ...



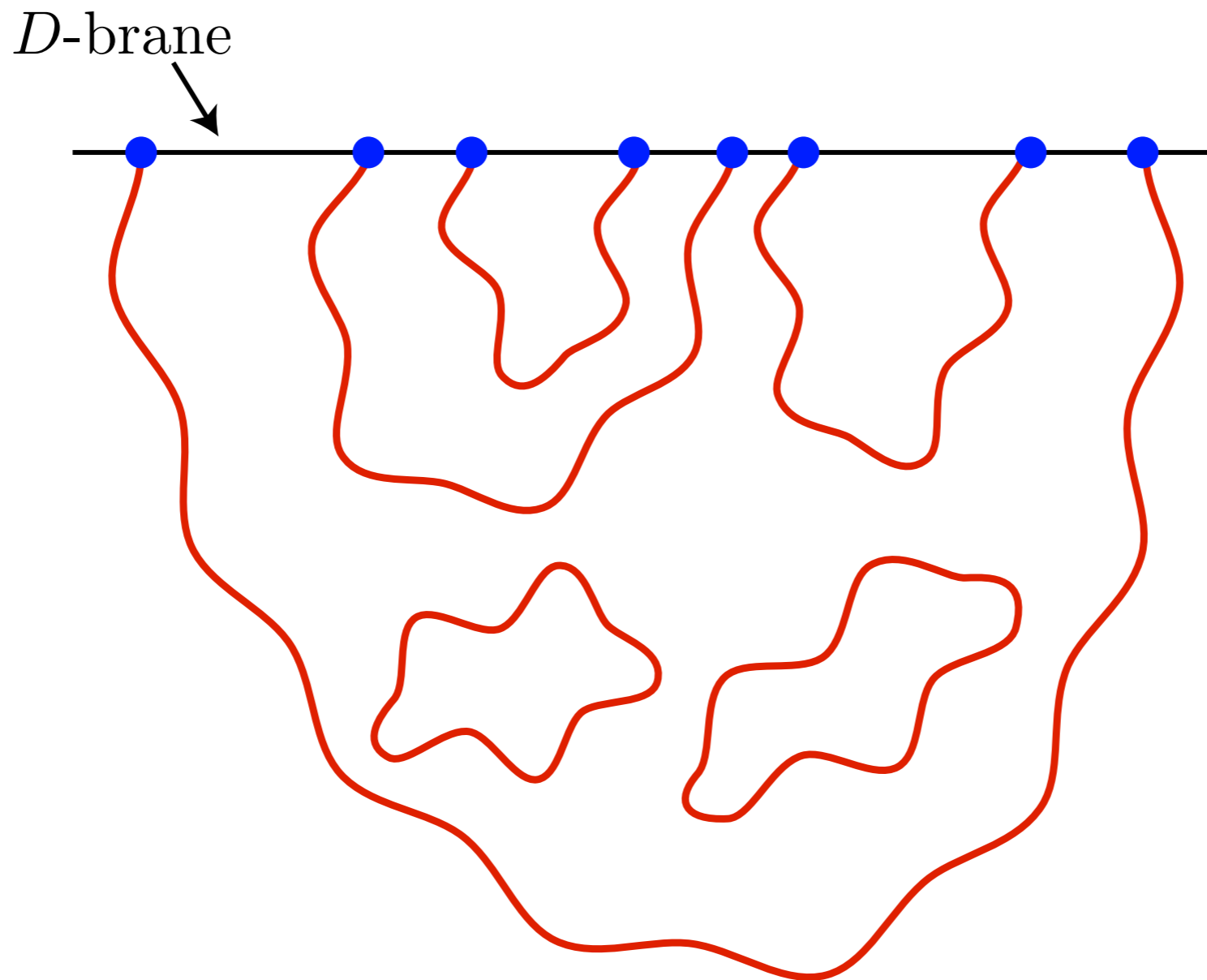
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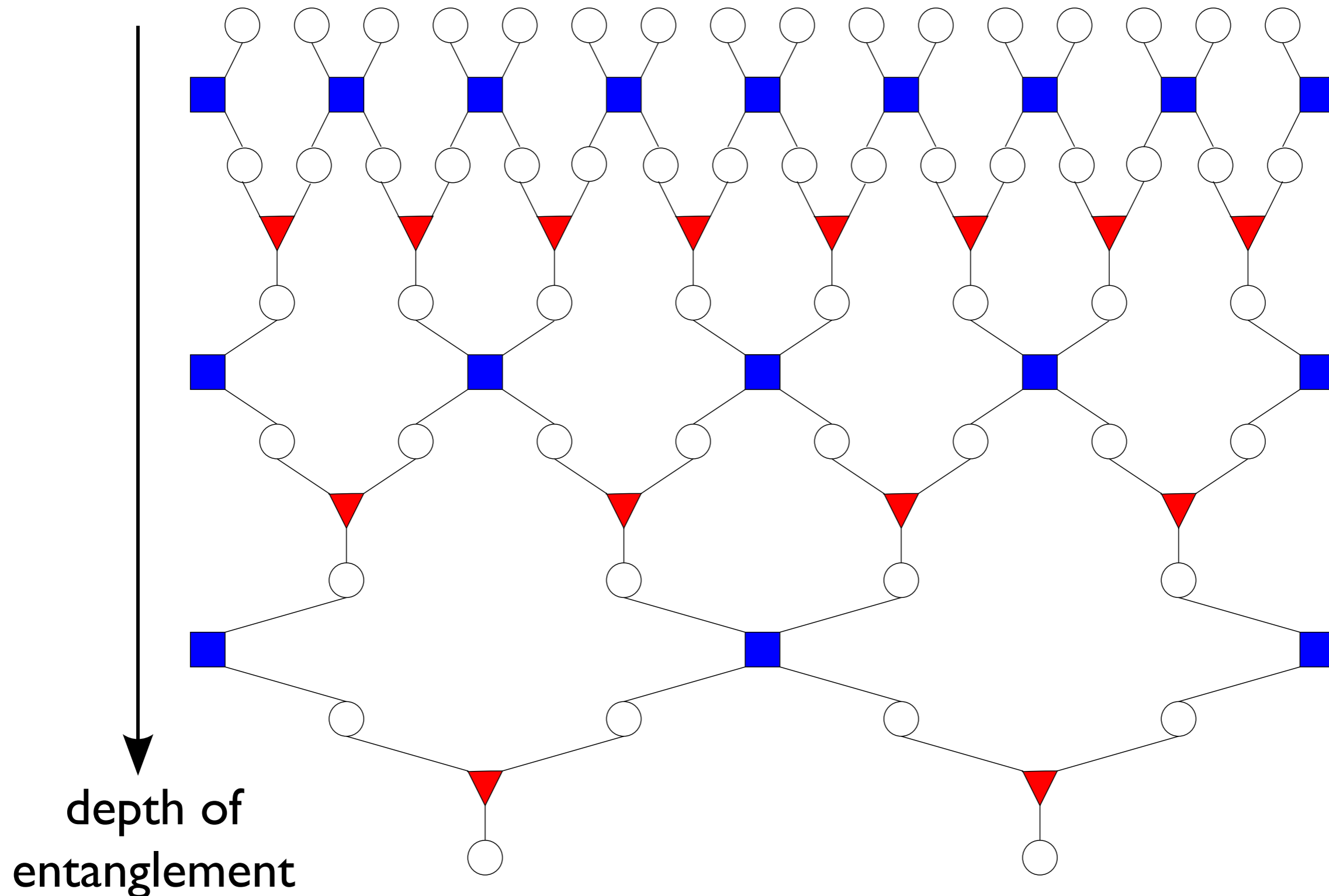
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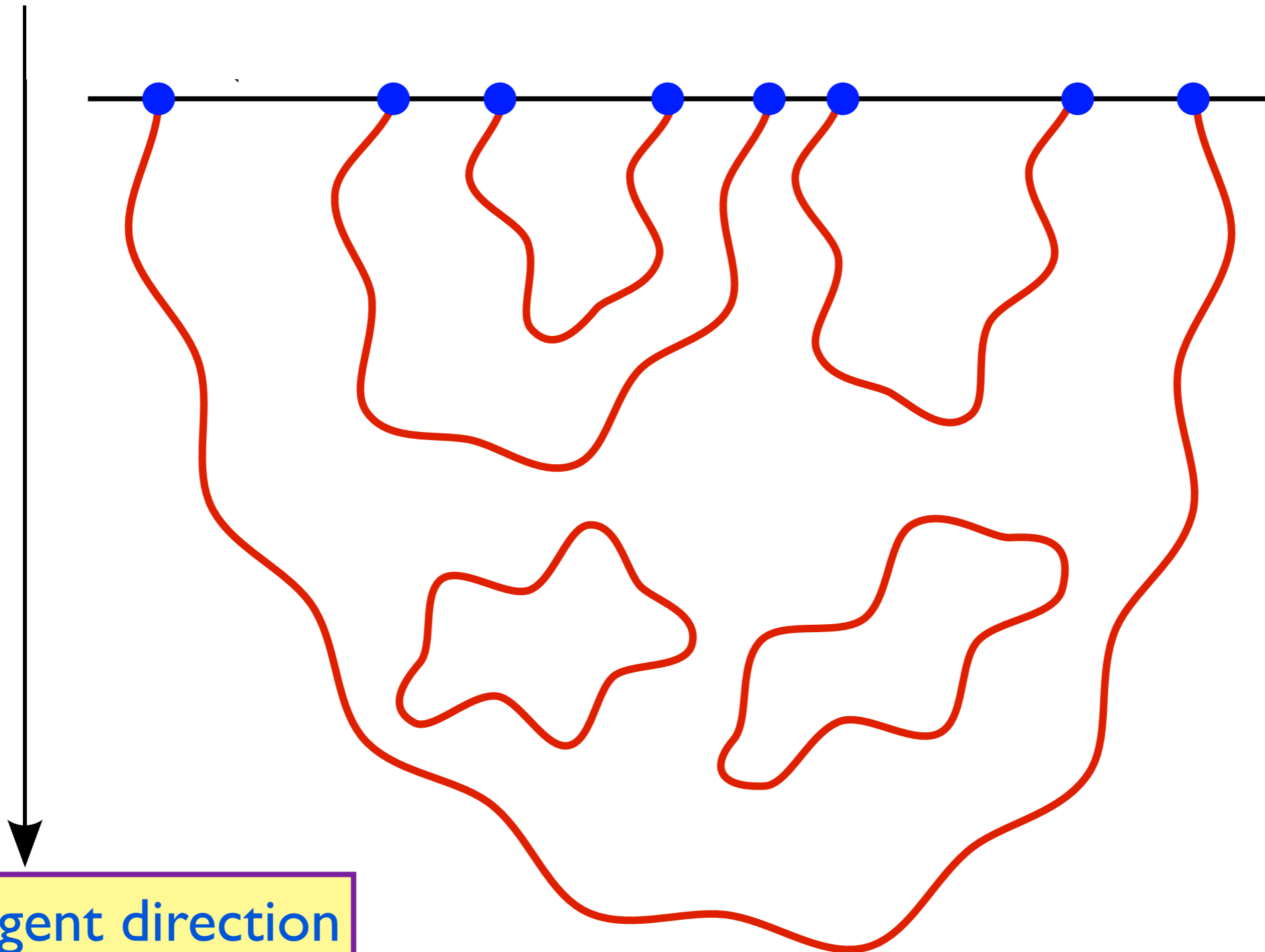
Tensor network representation of entanglement at quantum critical point

D -dimensional
space



String theory near
a D-brane

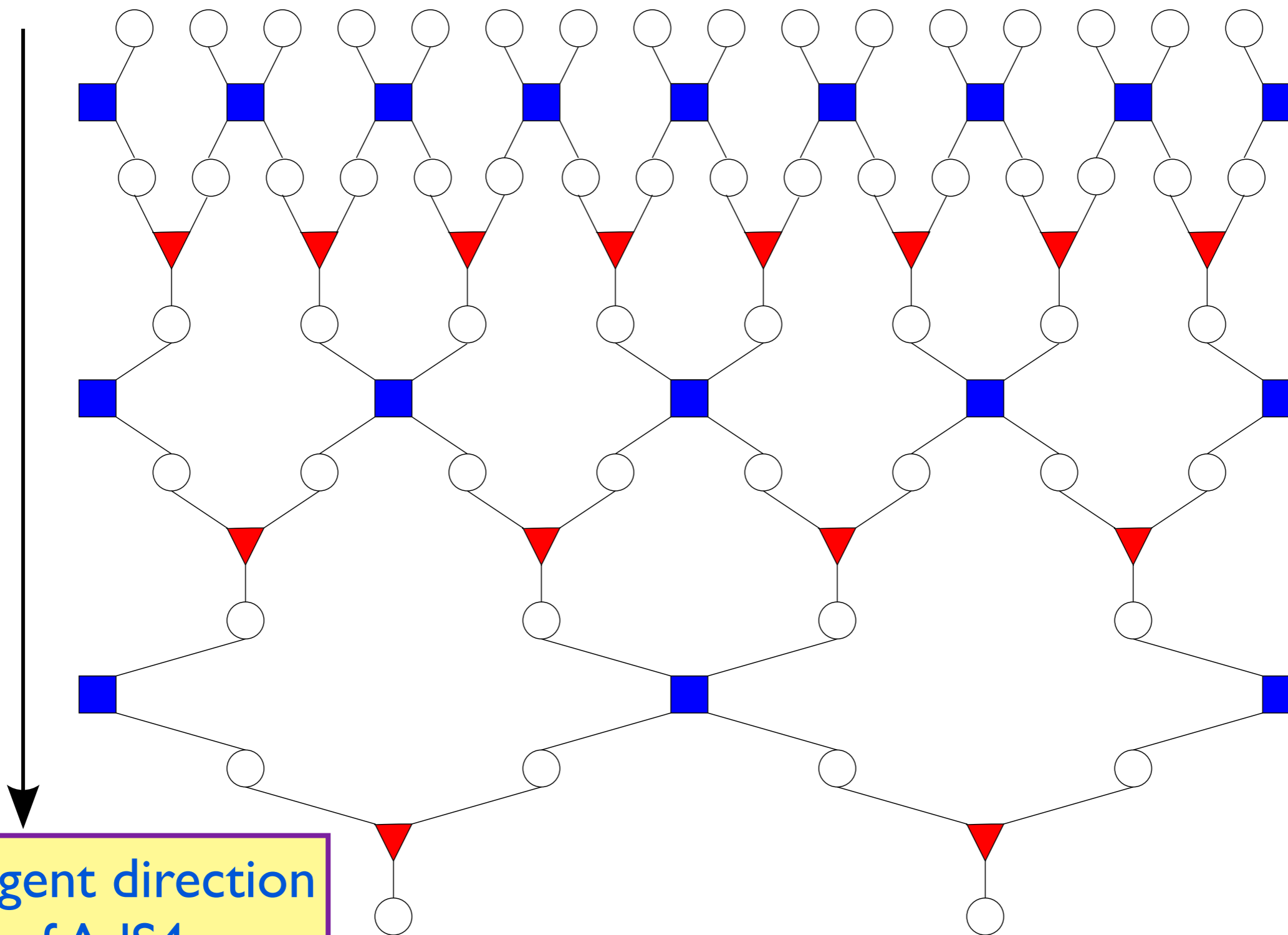
D -dimensional
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Emergent direction
of AdS4

Tensor network representation of entanglement at quantum critical point

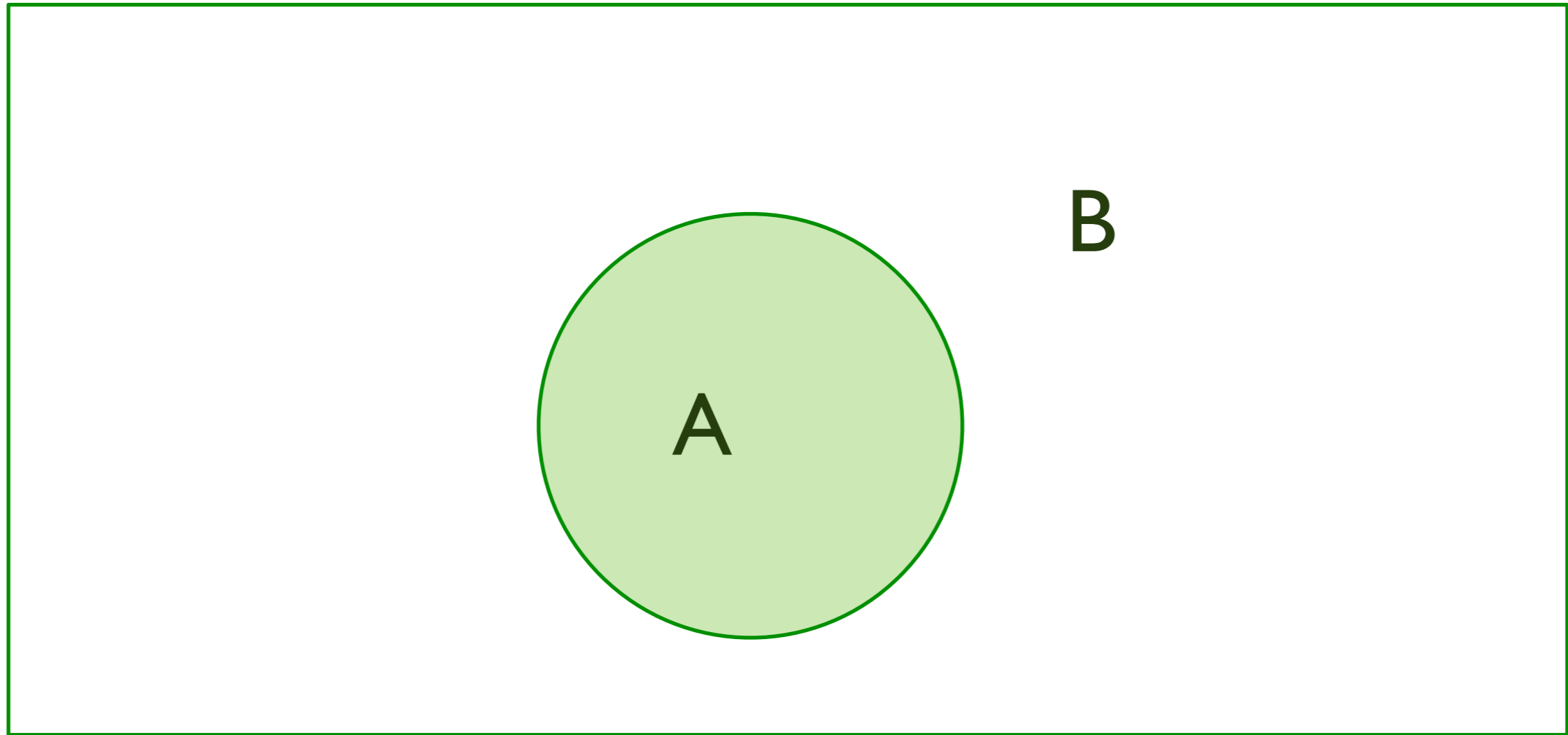
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Brian Swingle, arXiv:0905.1317

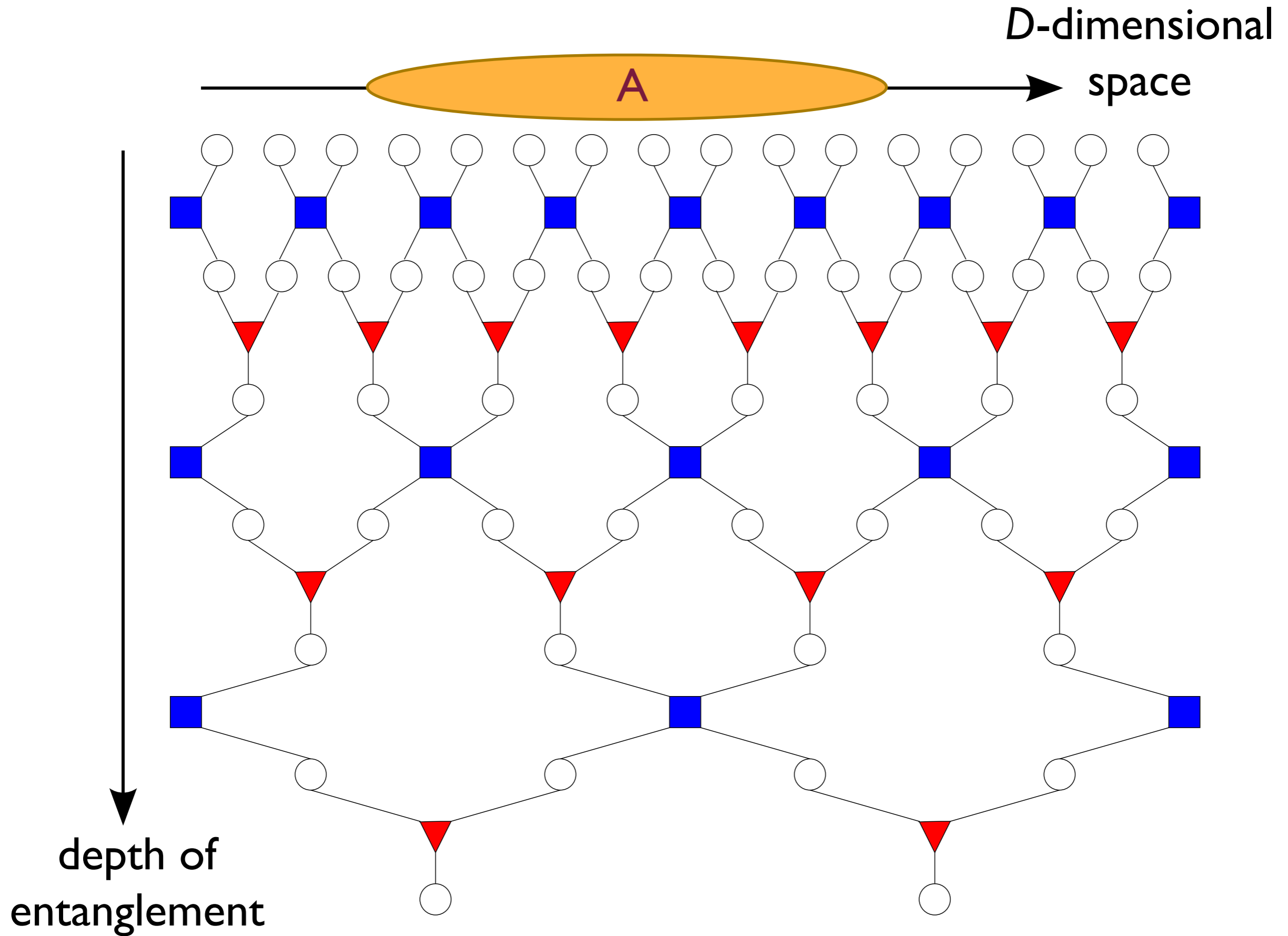
Entanglement entropy



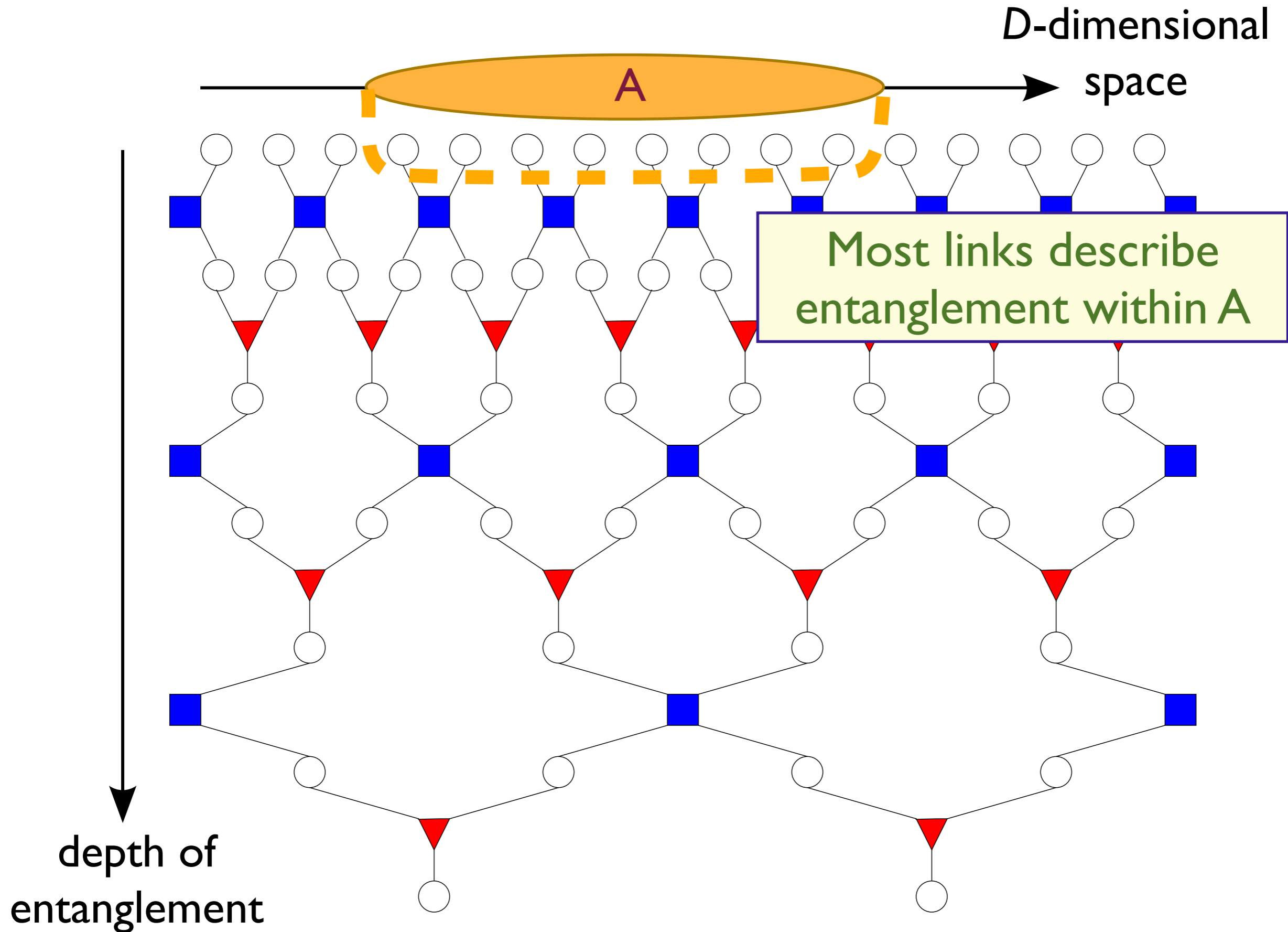
Measure strength of quantum entanglement of region A with region B .

$\rho_A = \text{Tr}_B \rho =$ density matrix of region A
Entanglement entropy $S_{EE} = -\text{Tr}(\rho_A \ln \rho_A)$

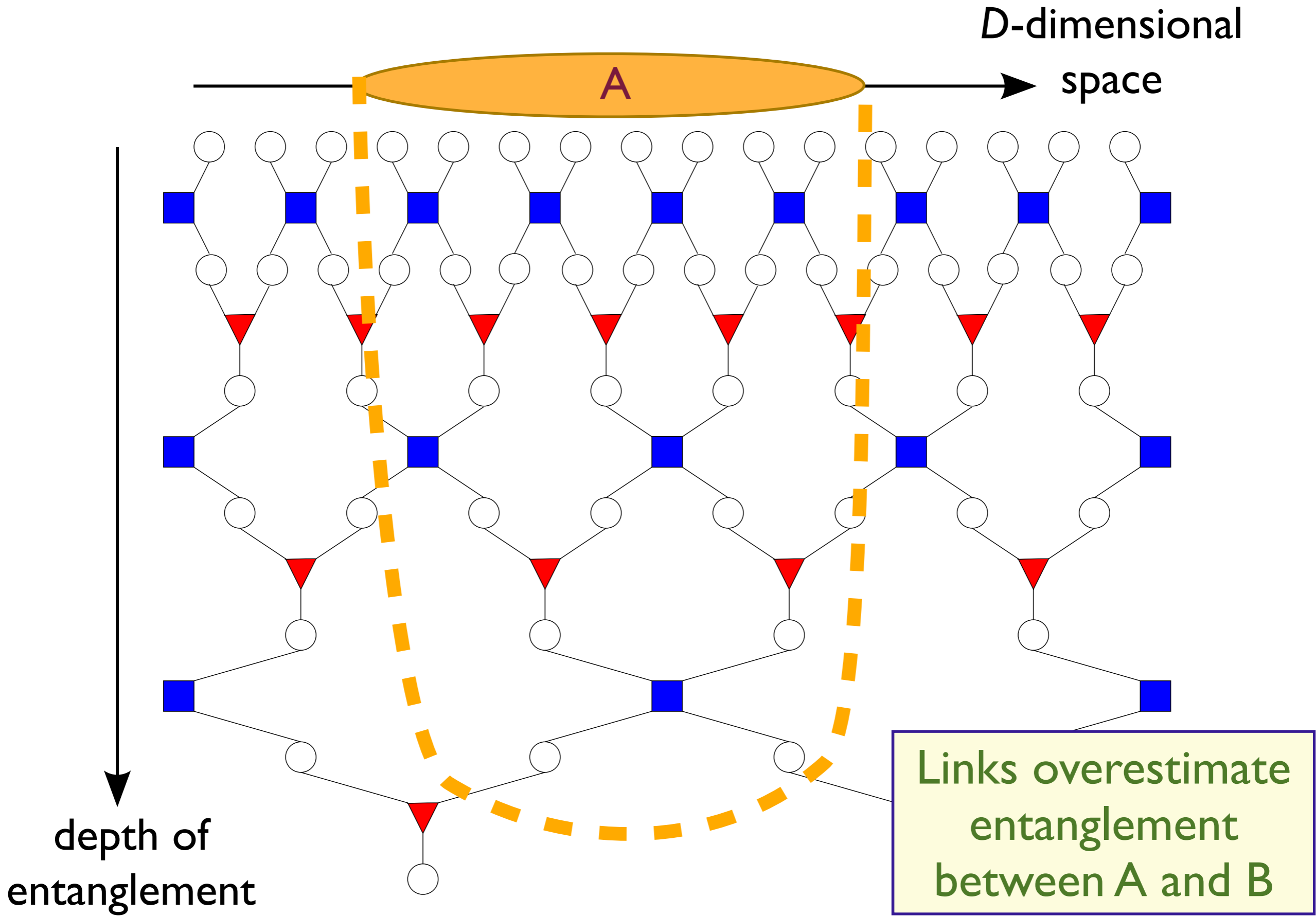
Entanglement entropy



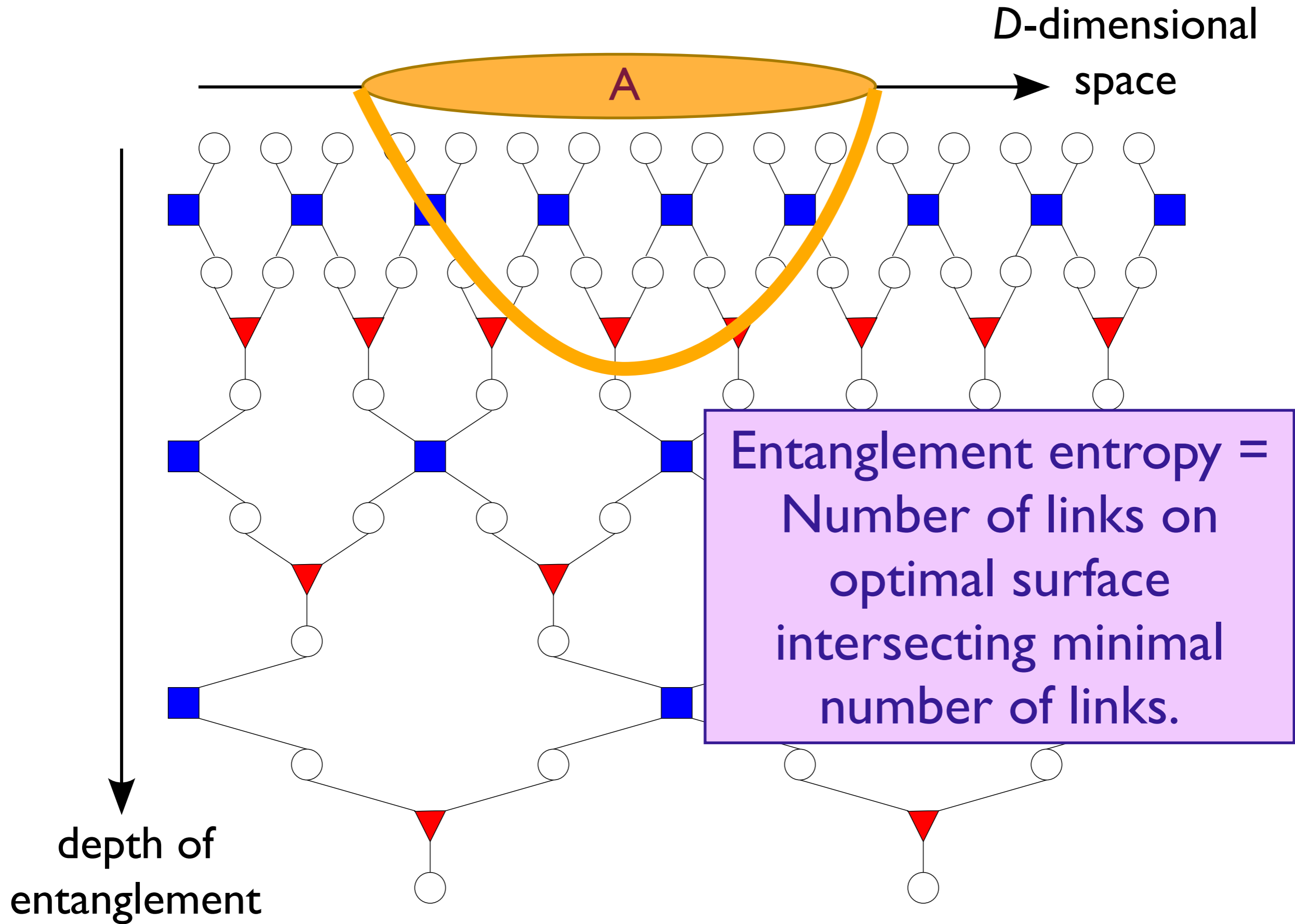
Entanglement entropy



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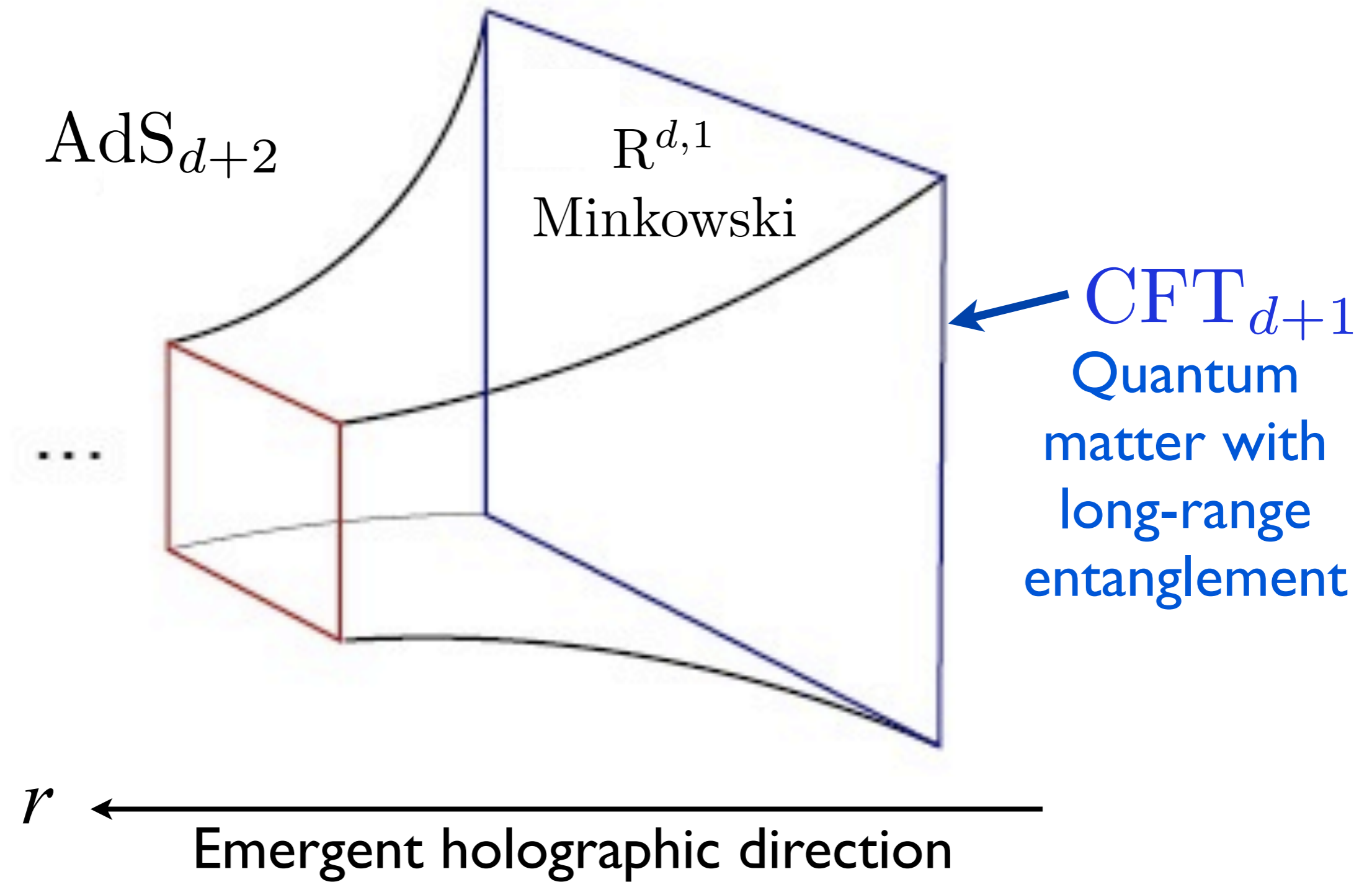


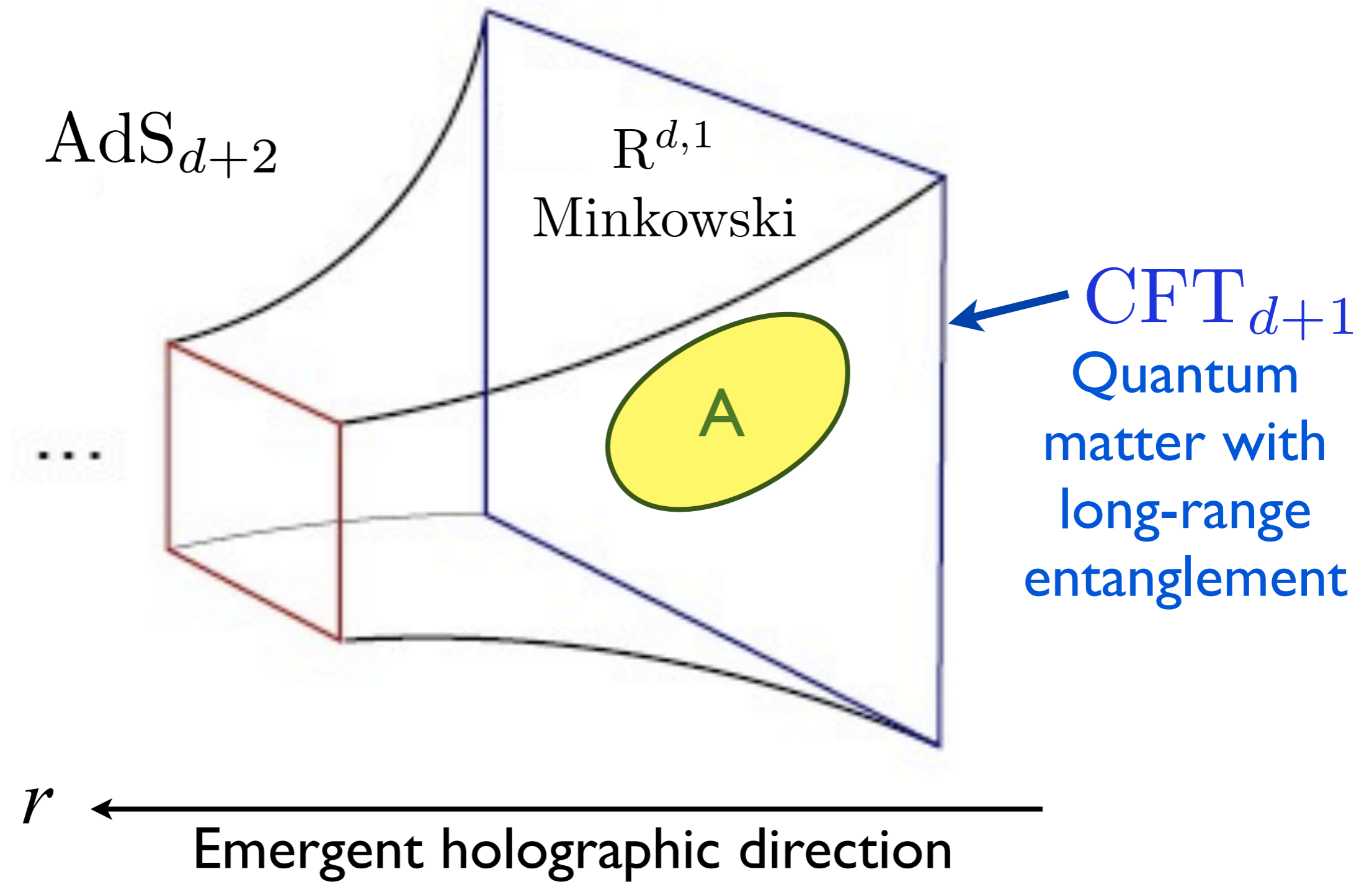
Entanglement entropy

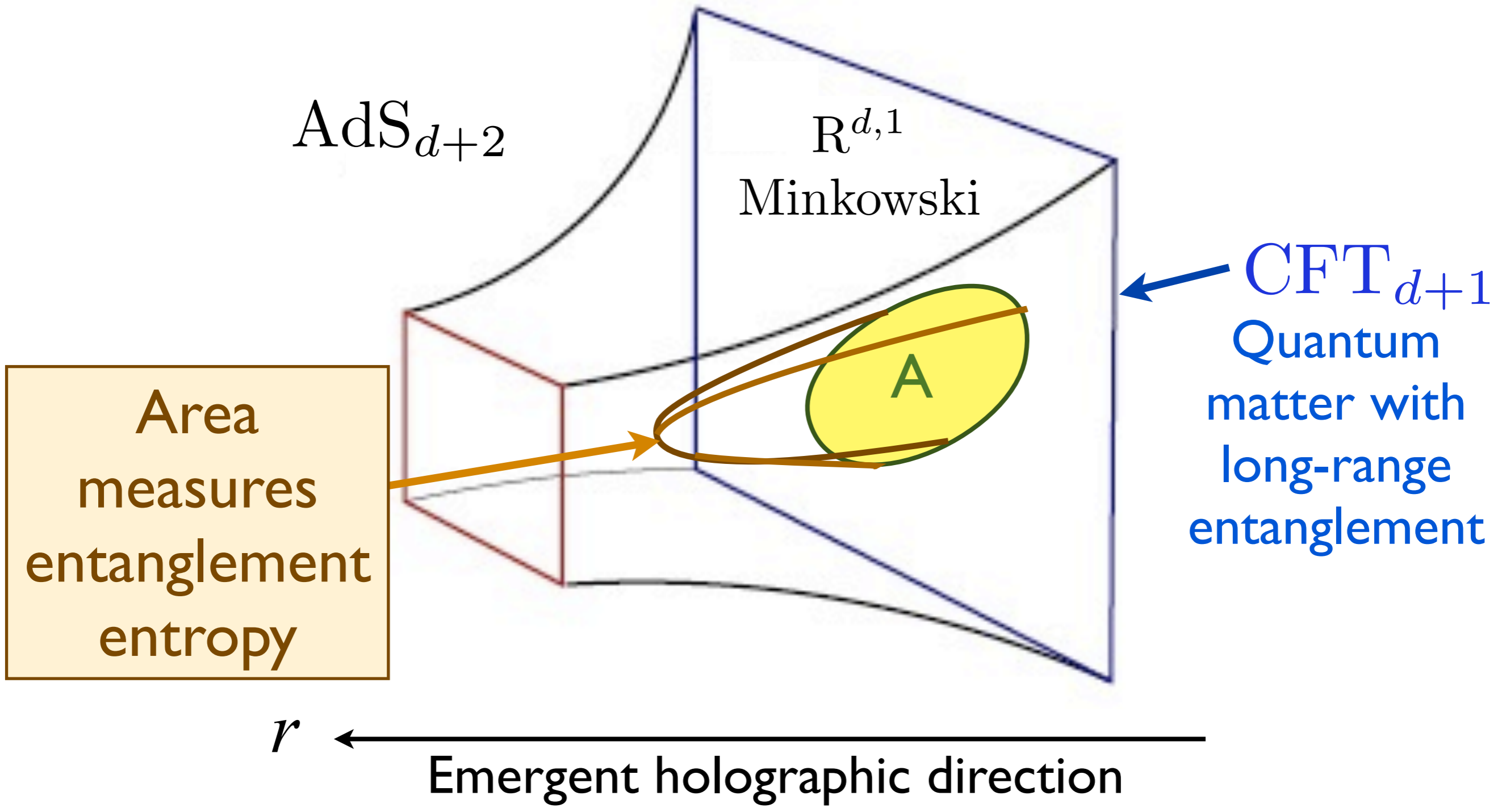
The entanglement entropy of a region A on the boundary equals the minimal area of a surface in the higher-dimensional space whose boundary co-incides with that of A .

This can be seen both the string and tensor-network pictures

S. Ryu and T. Takayanagi, Phys. Rev. Lett. 96, 18160 (2006).
Brian Swingle, arXiv:0905.1317







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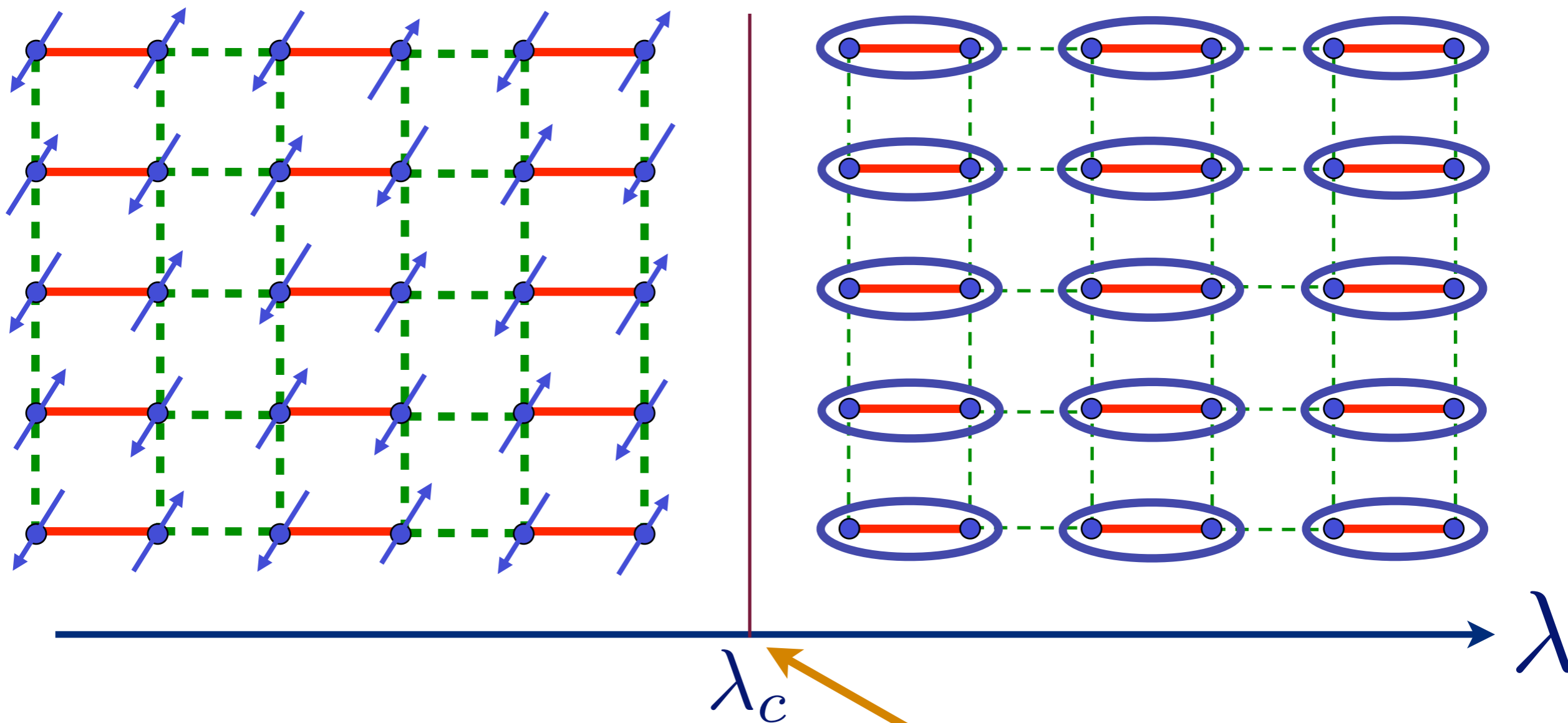
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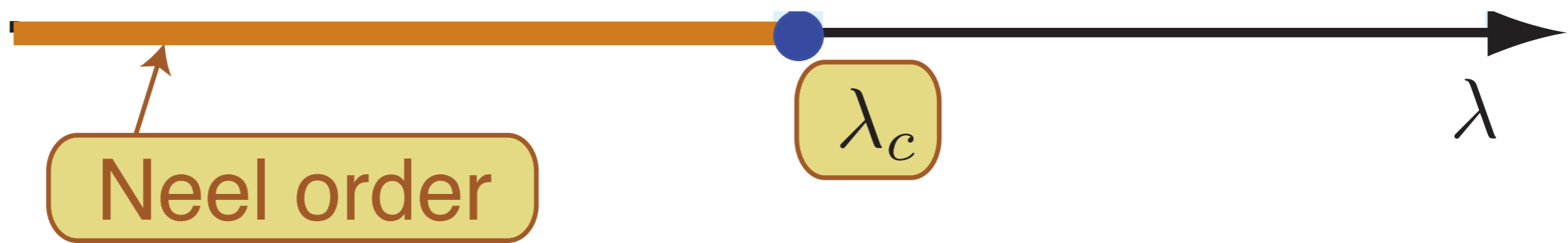
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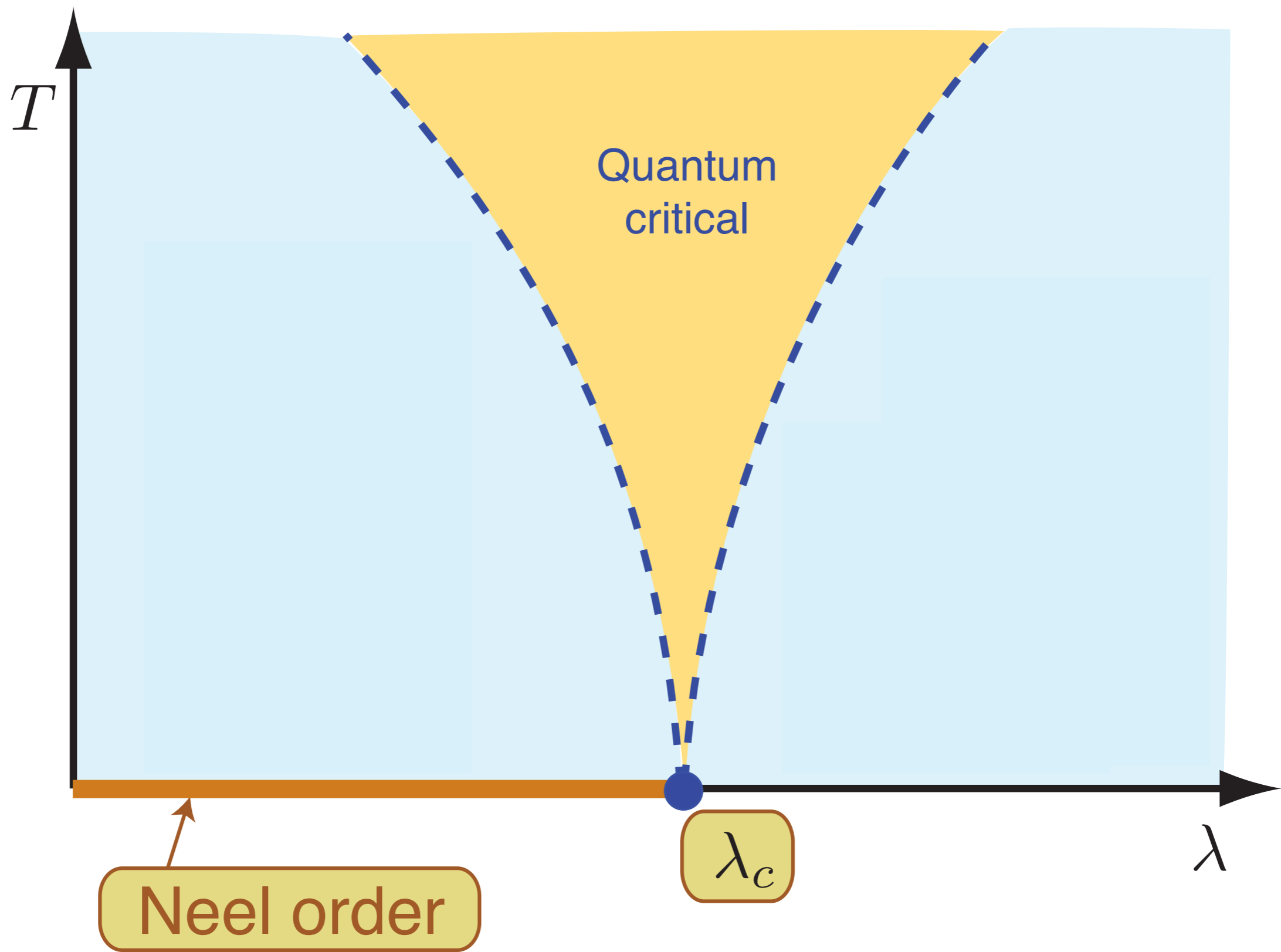
**String theory
and black holes**

$$\text{Diagram of two blue dots connected by a red line, enclosed in a blue oval} = \frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$

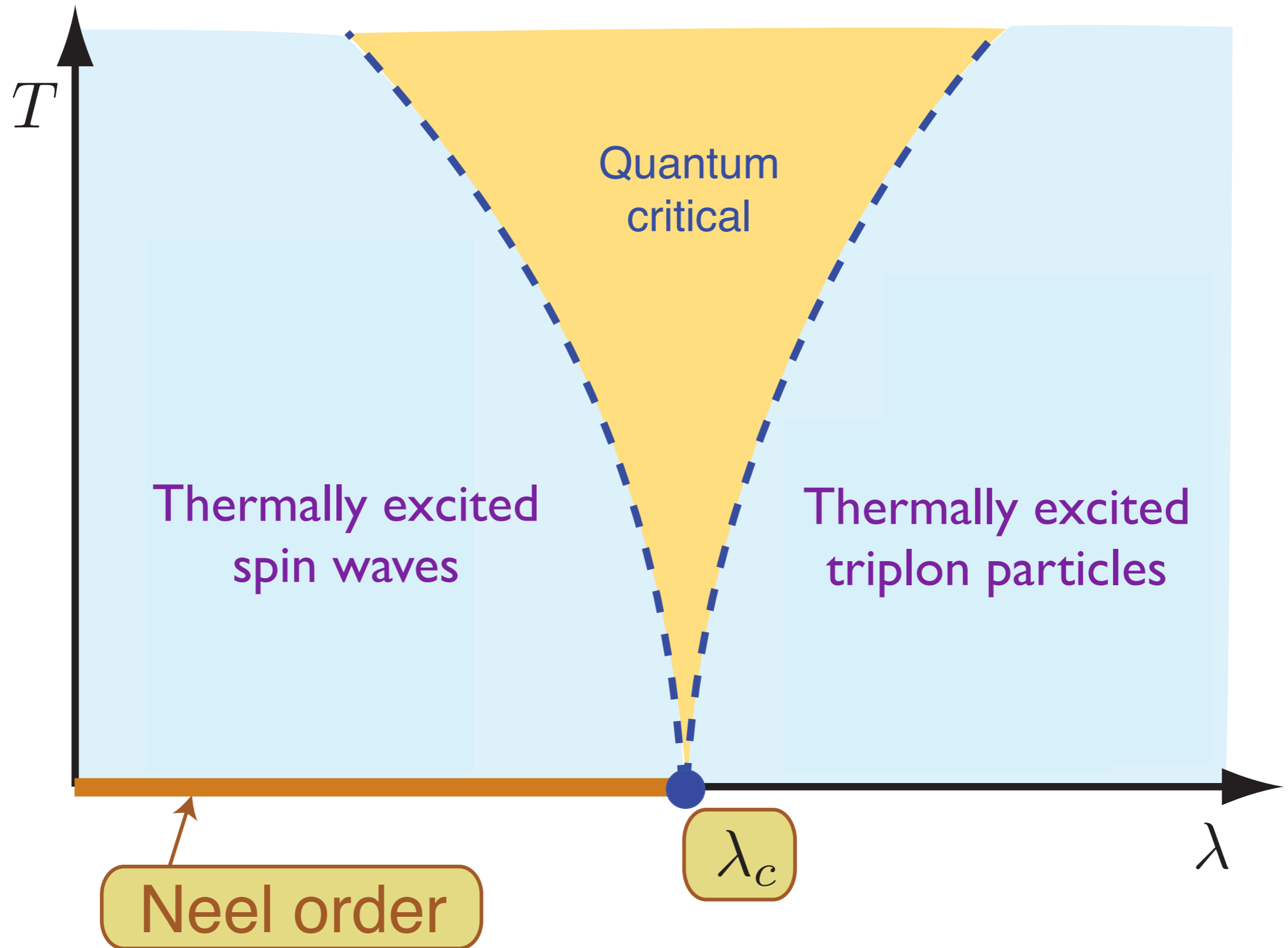


Quantum critical point with non-local entanglement in spin wavefunction

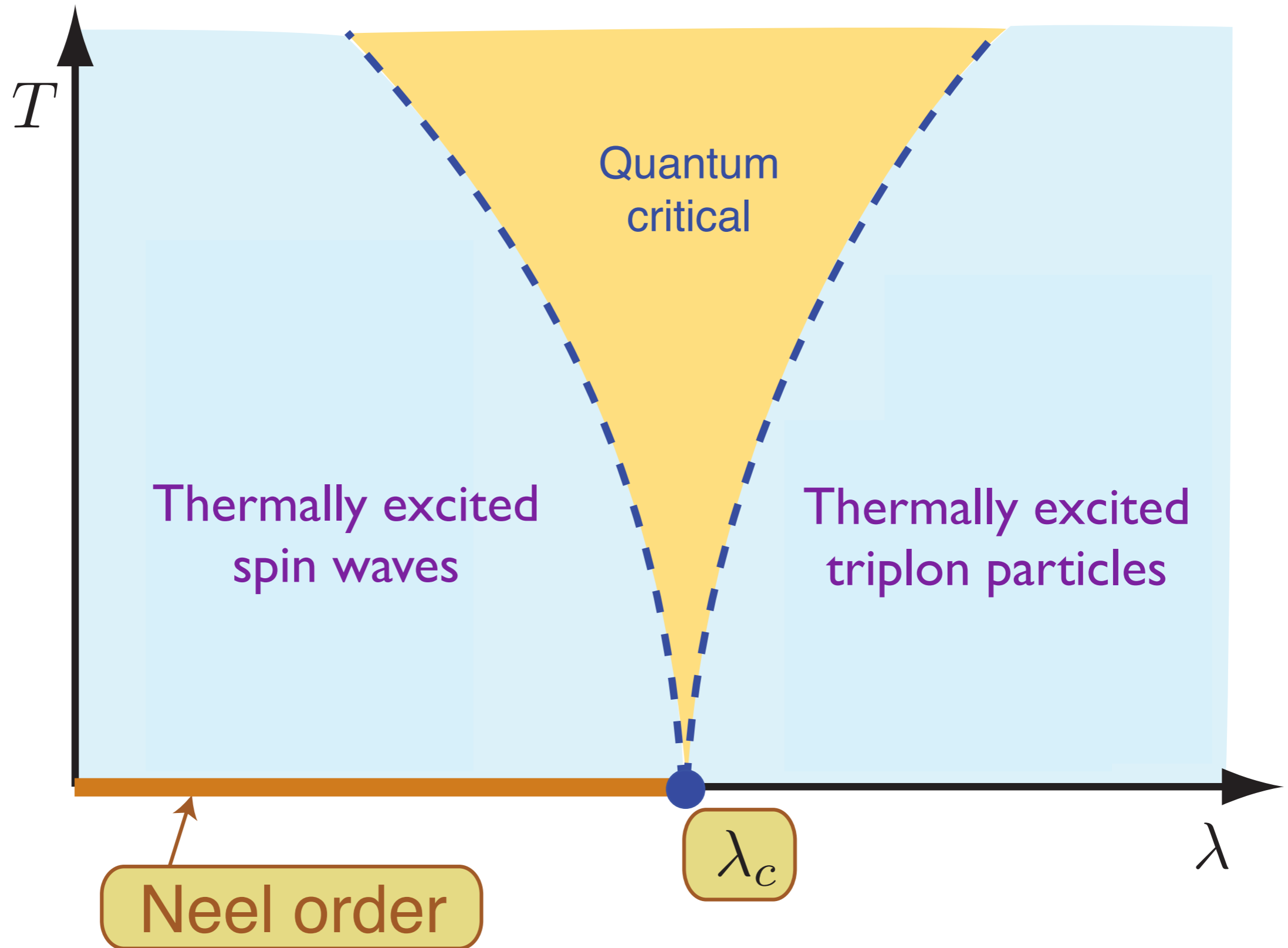




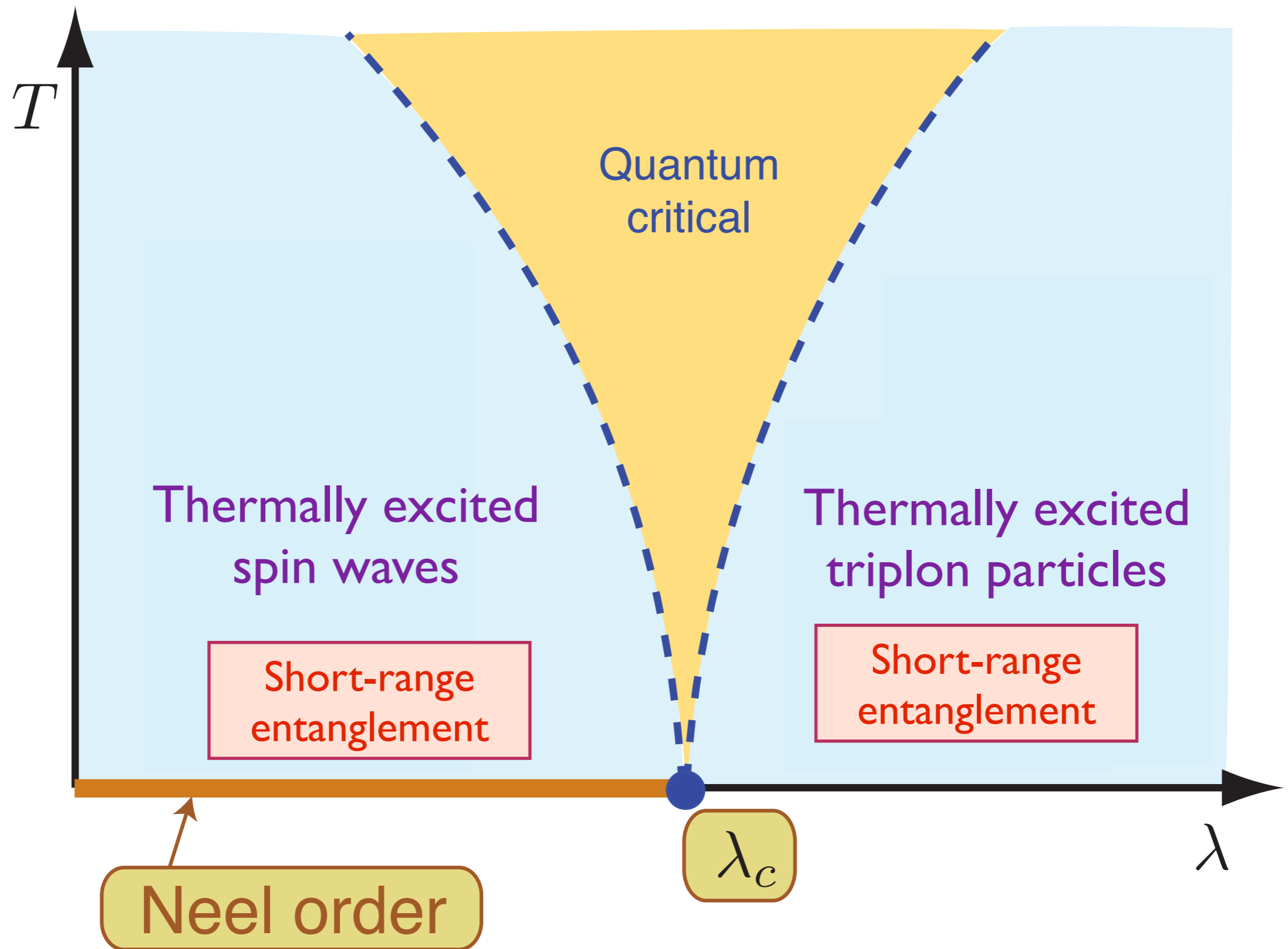
S. Sachdev and J. Ye, *Phys. Rev. Lett.* **69**, 2411 (1992).
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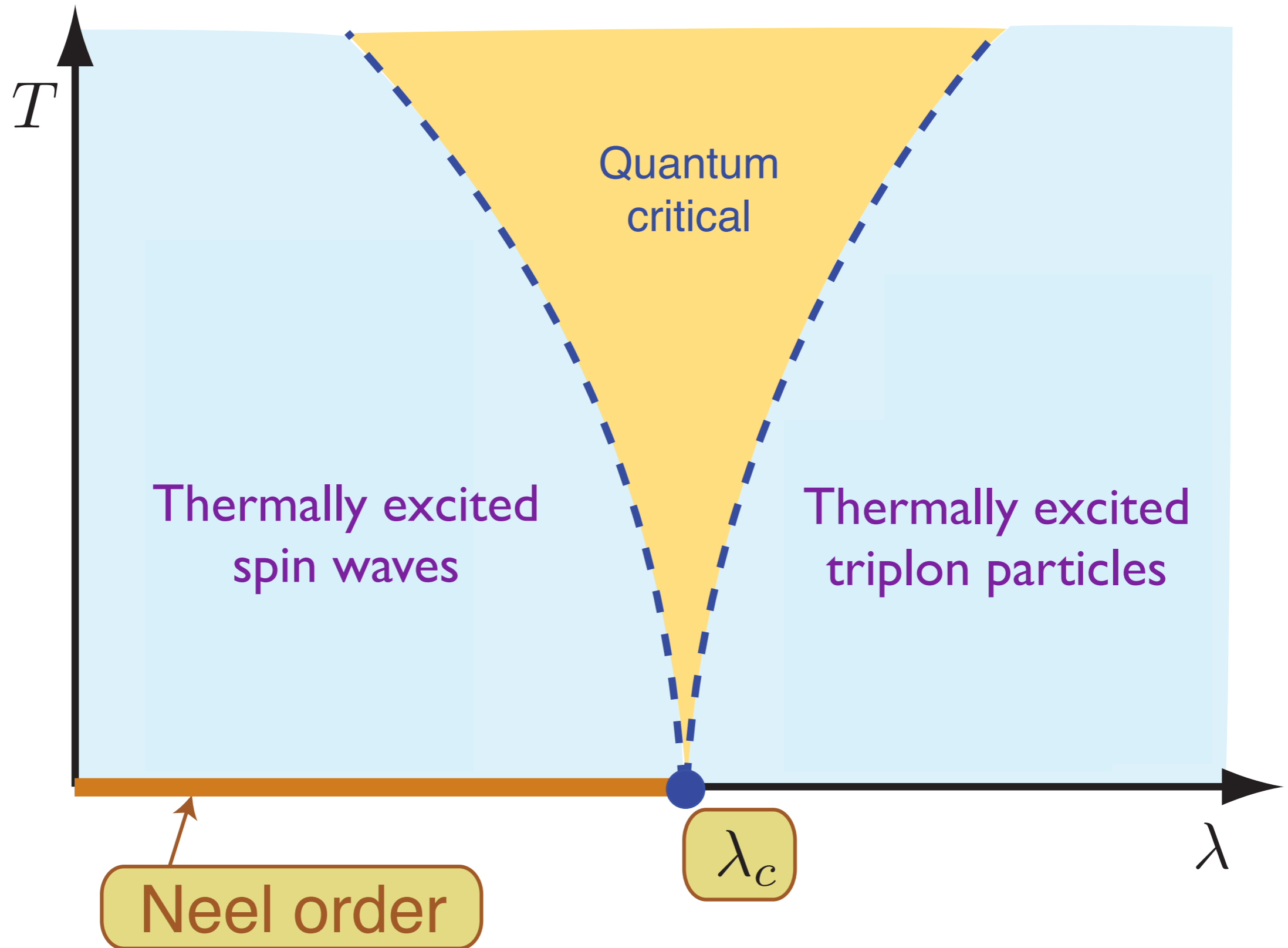
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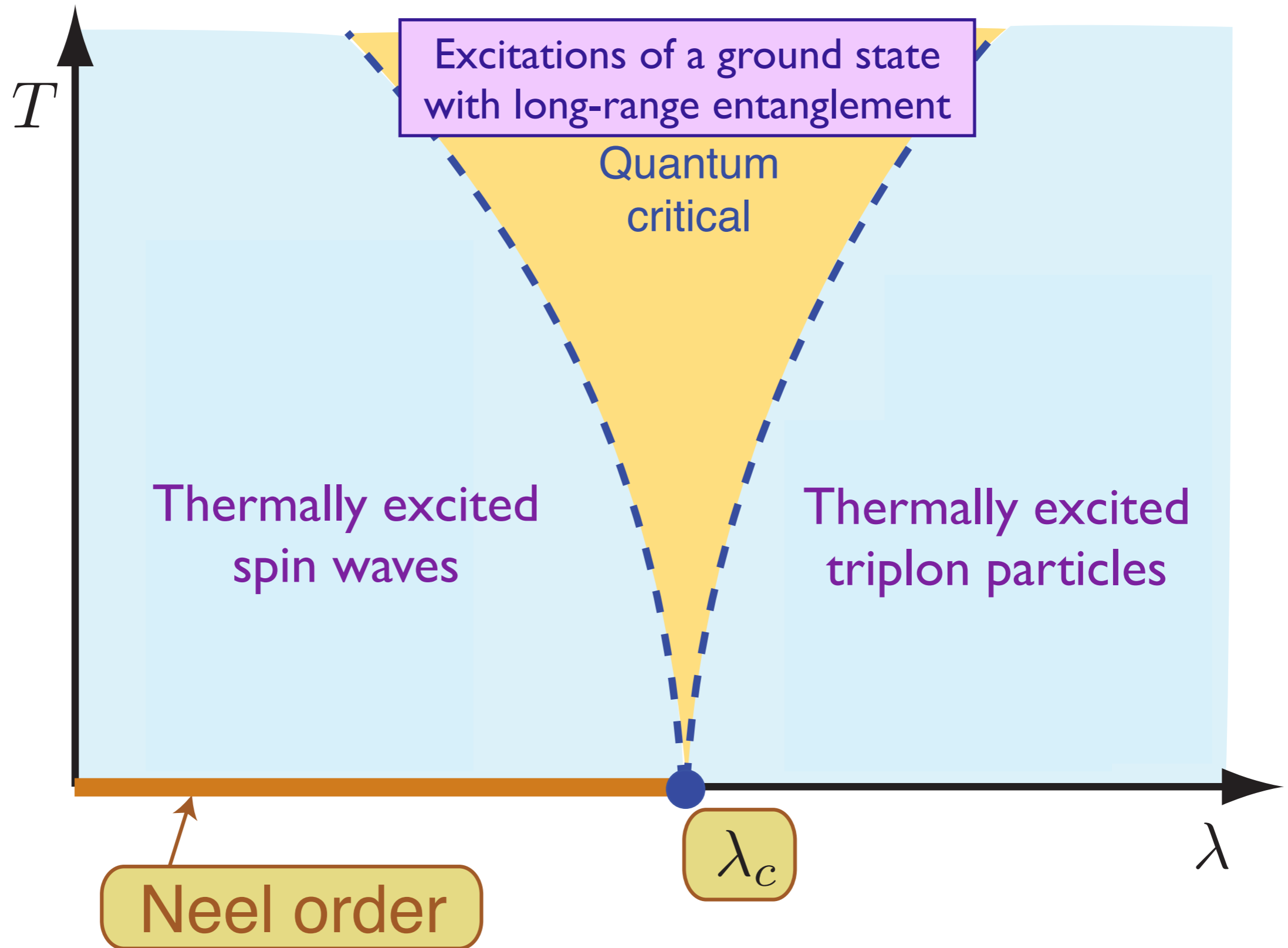
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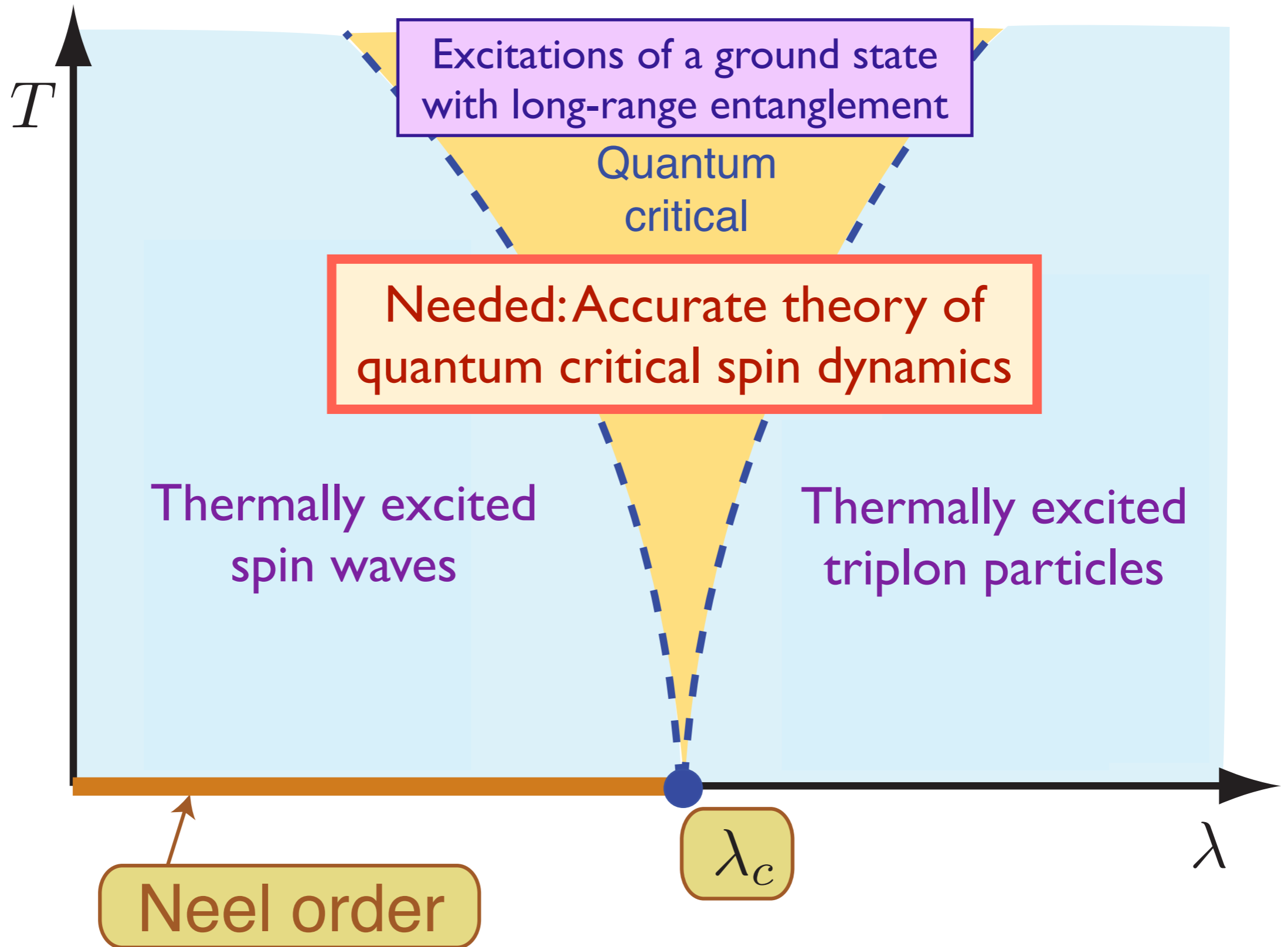
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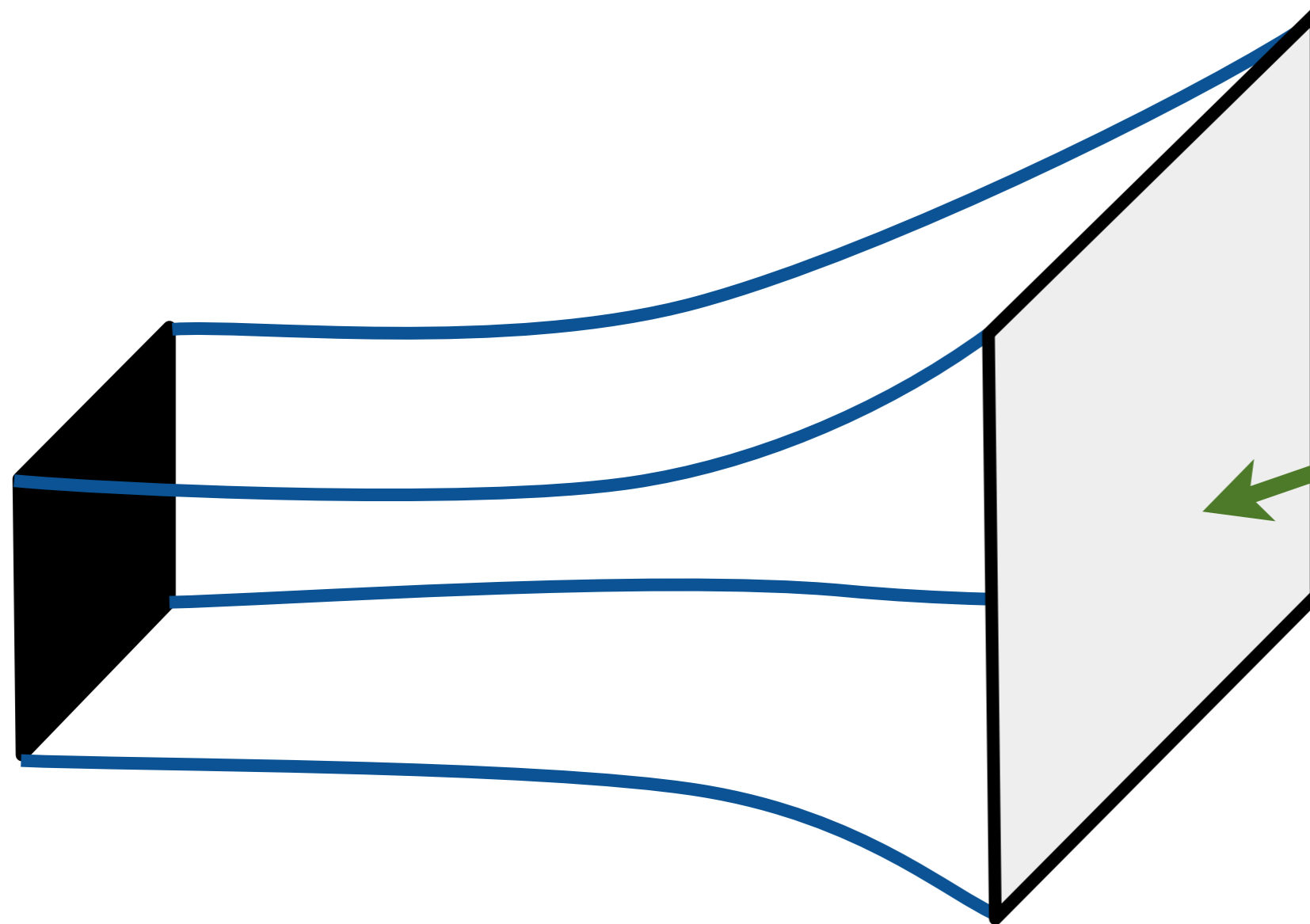
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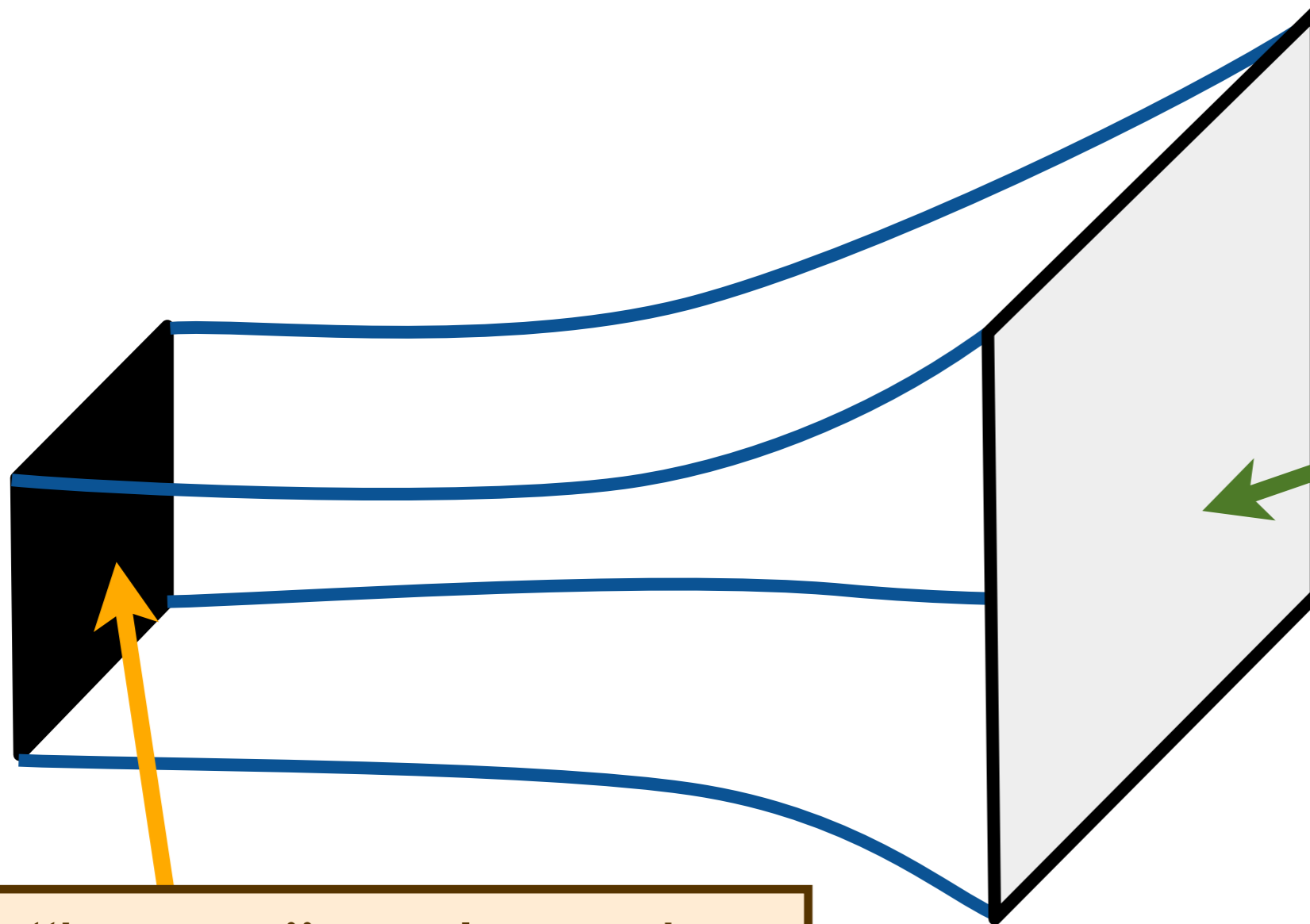


String theory at non-zero temperatures



A 2+1
dimensional
system at its
quantum
critical point

String theory at non-zero temperatures

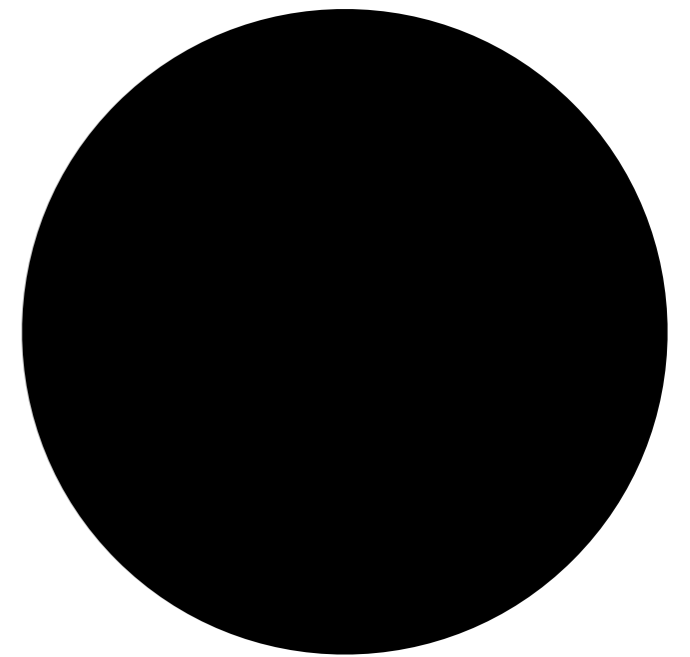


A "horizon", similar to the surface of a black hole !

A 2+1 dimensional system at its quantum critical point

Black Holes

Objects so massive that light is gravitationally bound to them.

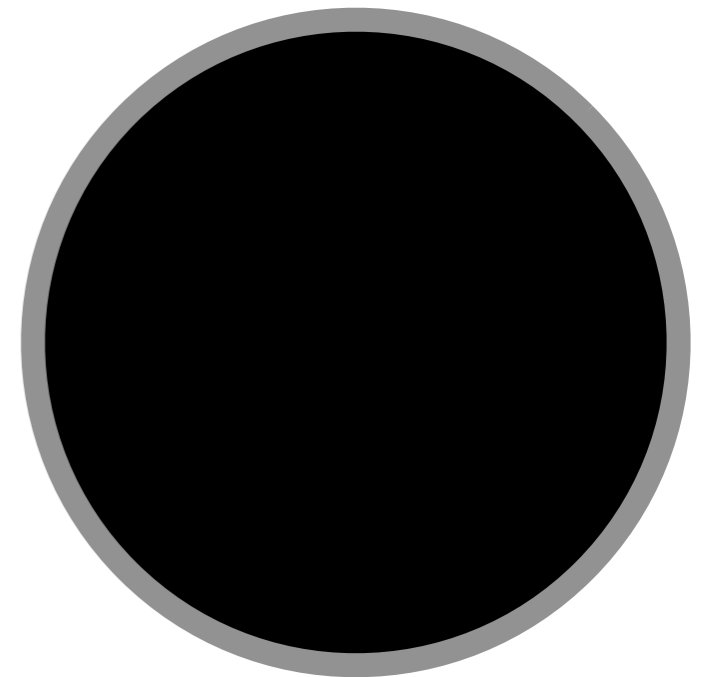


Black Holes

Objects so massive that light is gravitationally bound to them.

In Einstein's theory, the region inside the black hole **horizon** is disconnected from the rest of the universe.

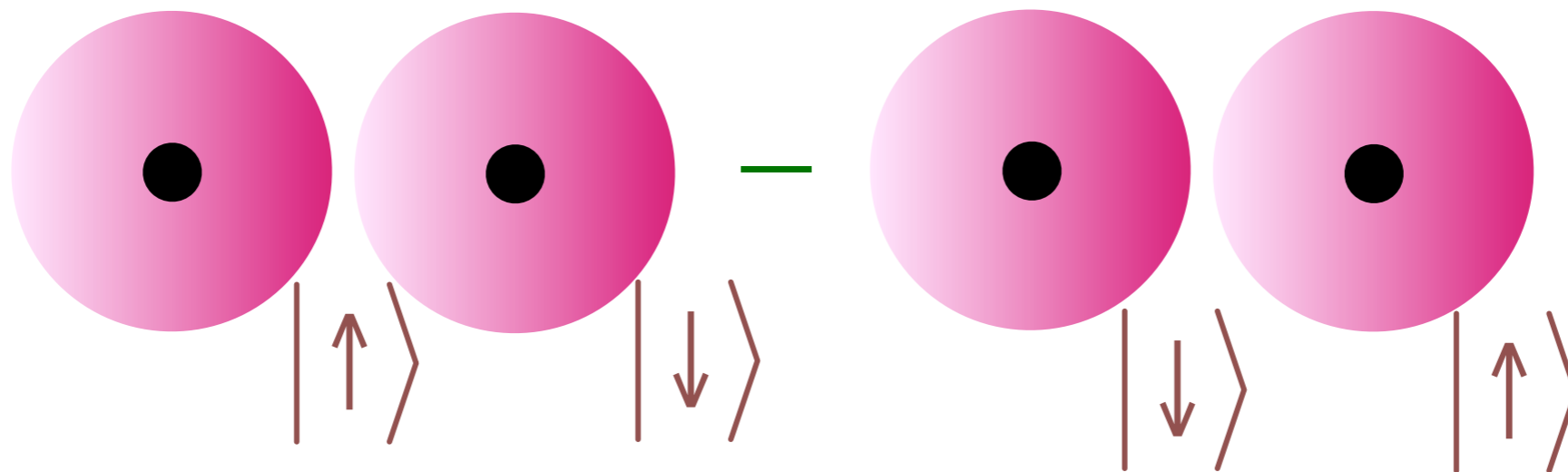
$$\text{Horizon radius } R = \frac{2GM}{c^2}$$



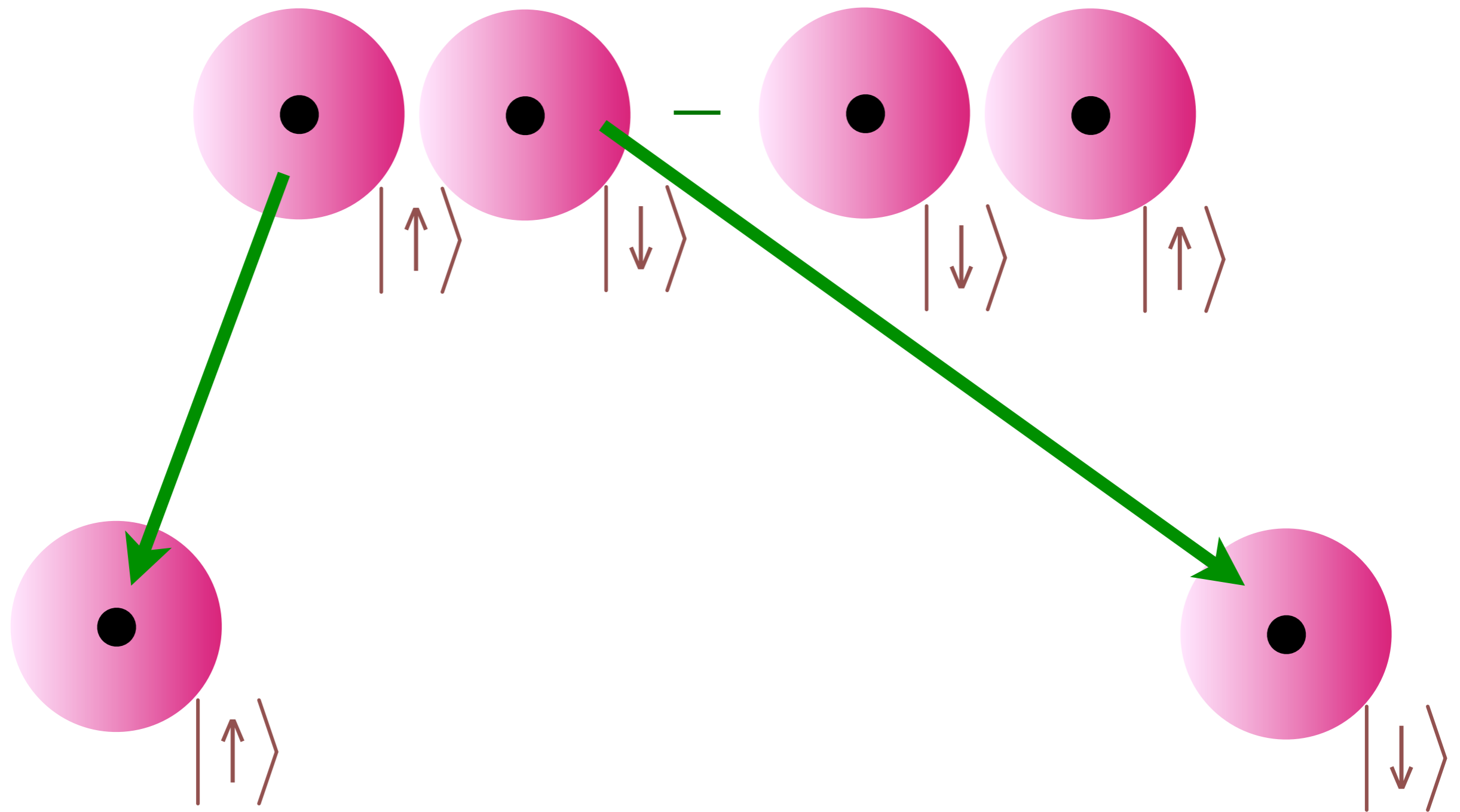
Black Holes + Quantum theory

Around 1974, Bekenstein and Hawking showed that the application of the quantum theory across a black hole horizon led to many astonishing conclusions

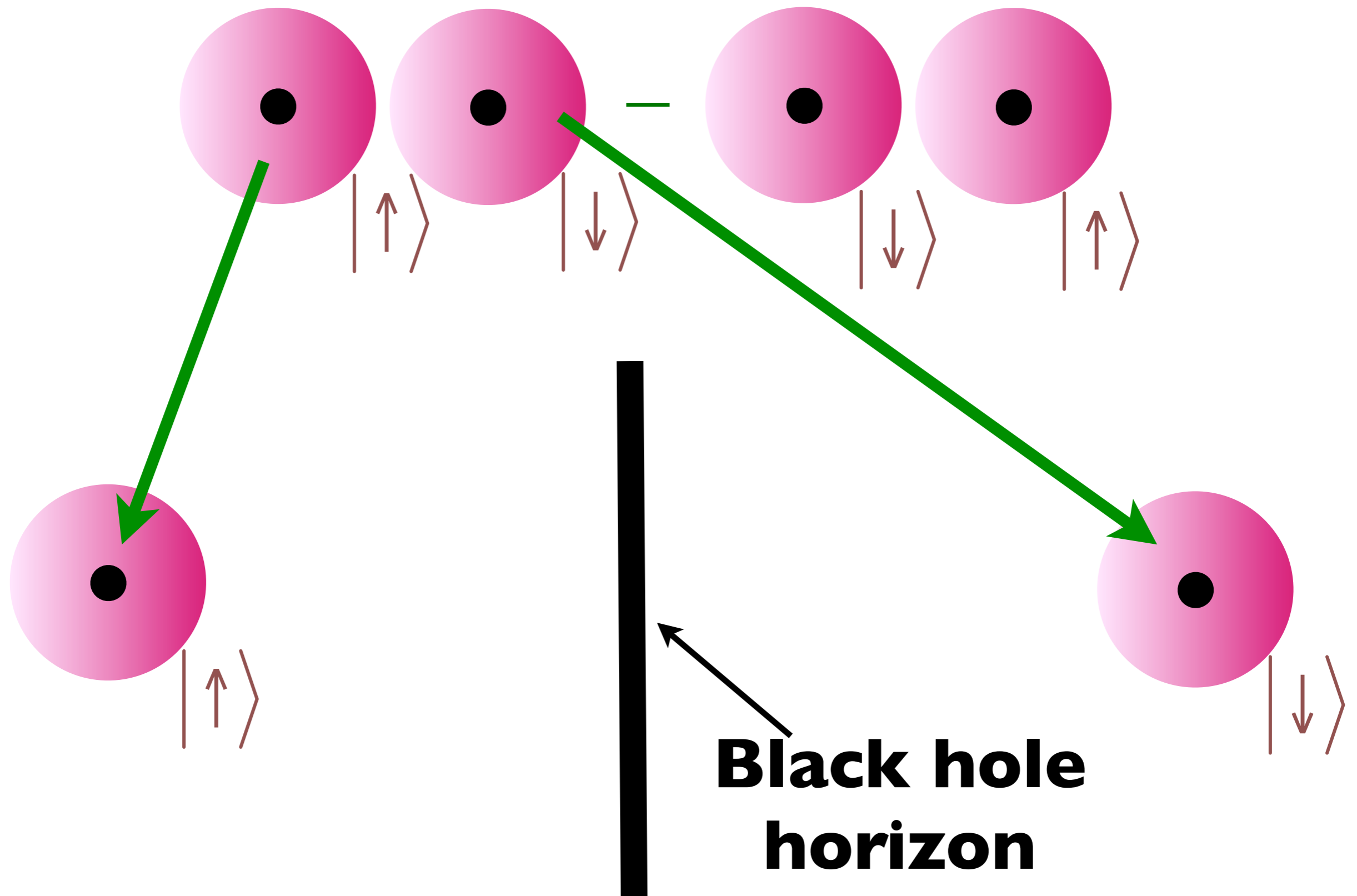
Quantum Entanglement across a black hole horizon



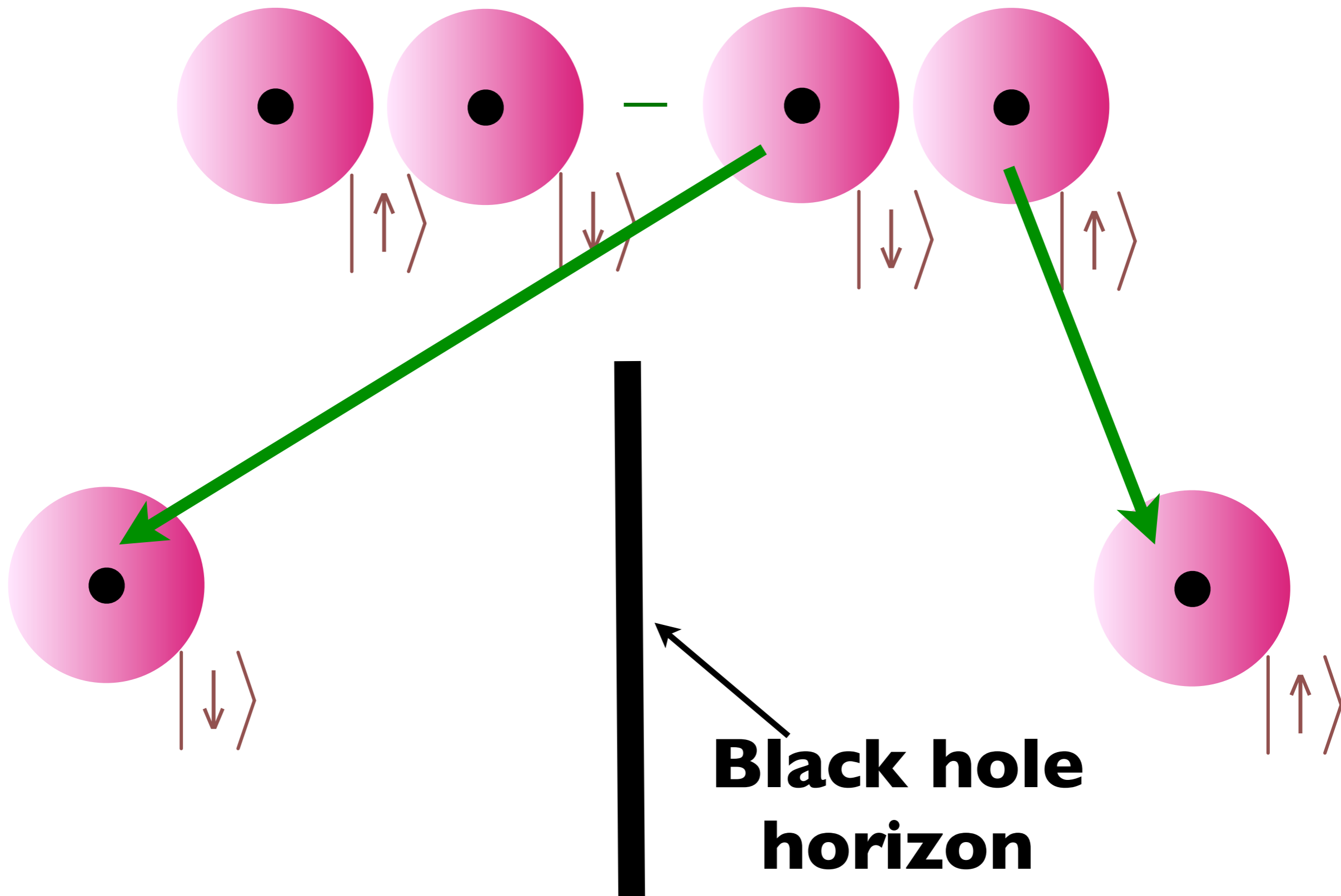
Quantum Entanglement across a black hole horizon



Quantum Entanglement across a black hole horizon

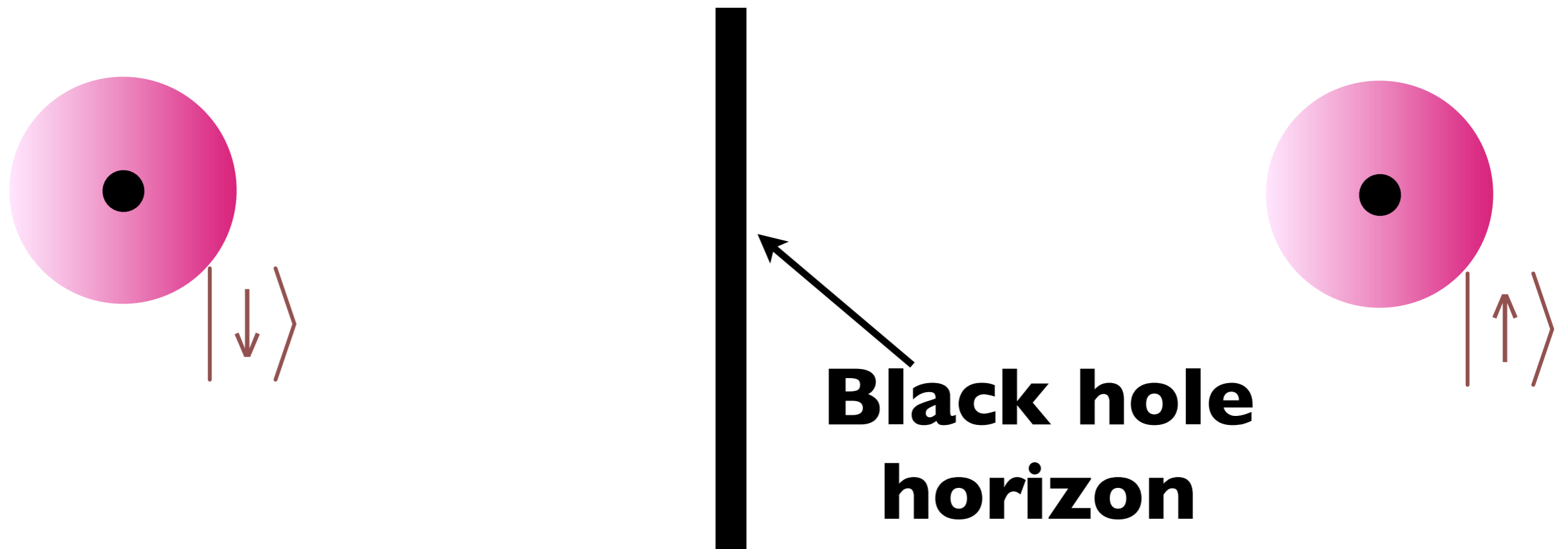


Quantum Entanglement across a black hole horizon



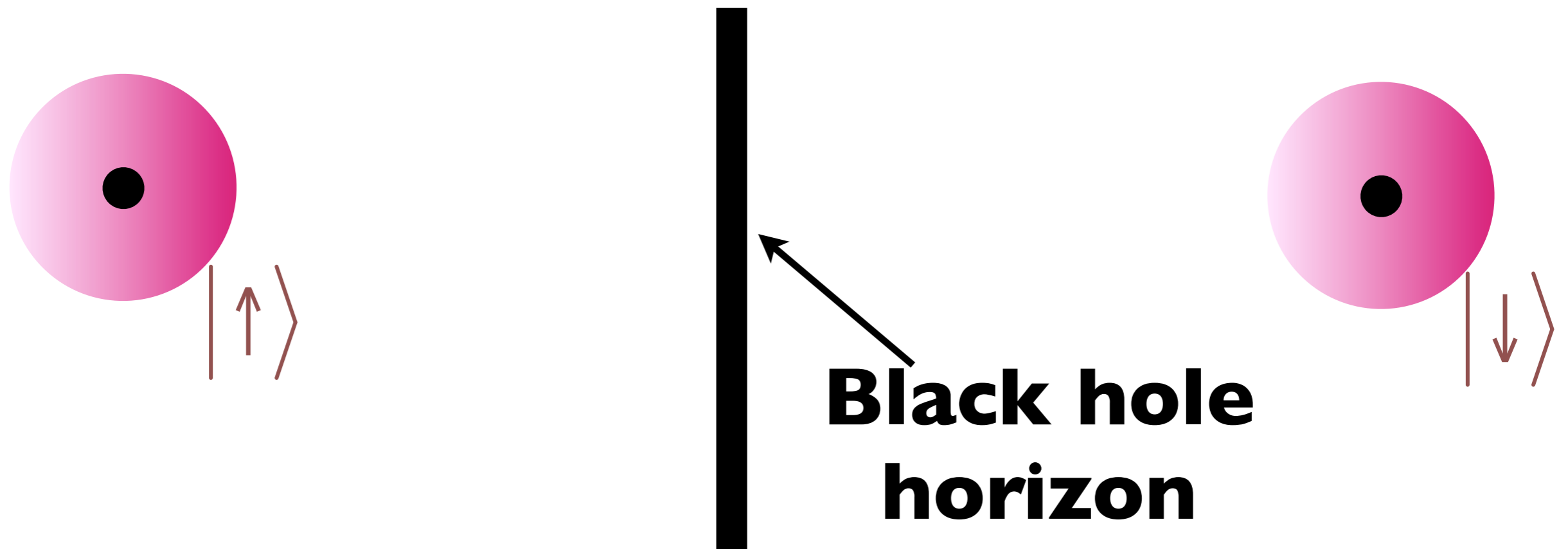
Quantum Entanglement across a black hole horizon

There is a non-local quantum entanglement between the inside and outside of a black hole



Quantum Entanglement across a black hole horizon

There is a non-local quantum entanglement between the inside and outside of a black hole

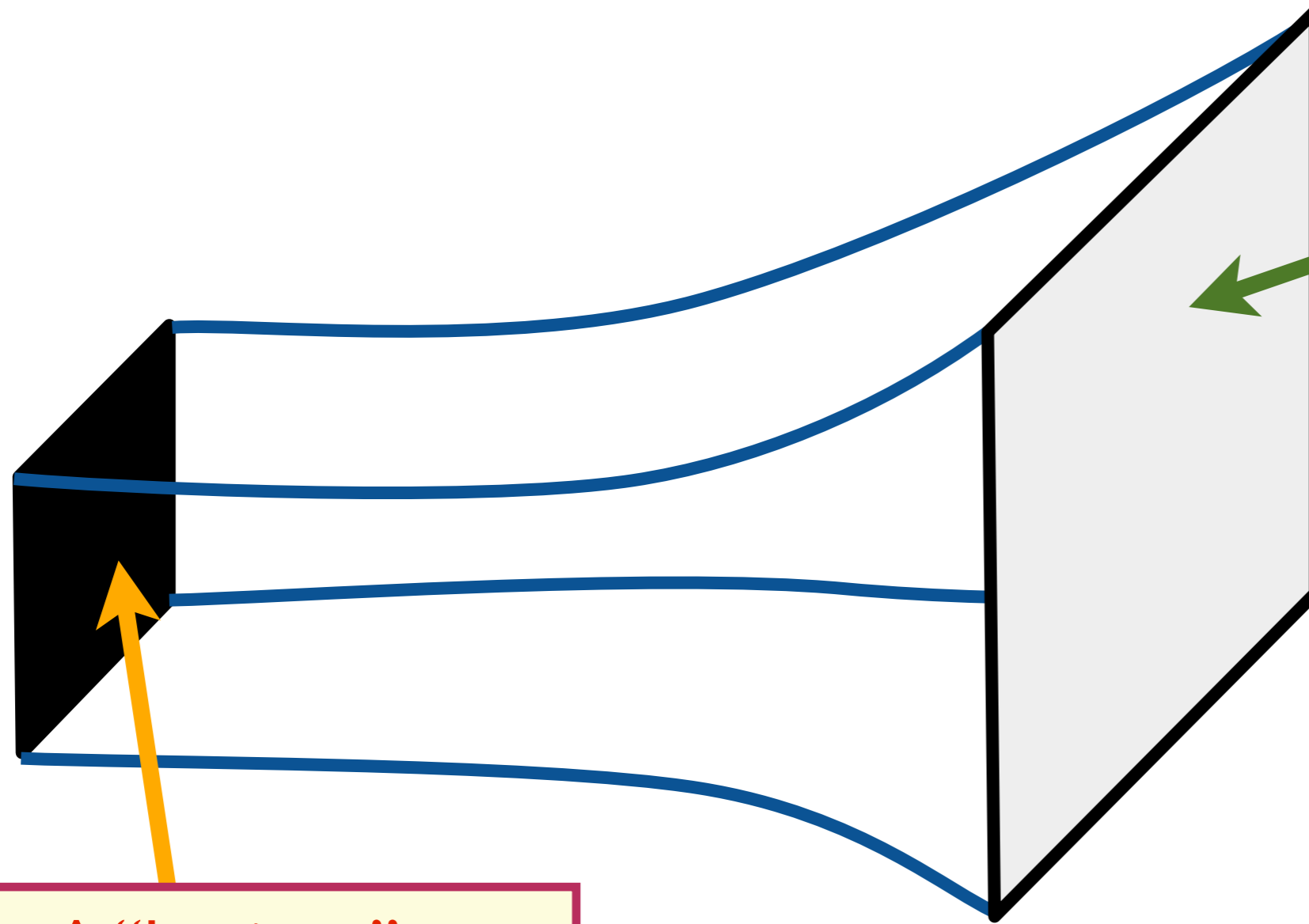


Quantum Entanglement across a black hole horizon

There is a non-local quantum entanglement between the inside and outside of a black hole

This entanglement leads to a black hole temperature (the Hawking temperature) and a black hole entropy (the Bekenstein entropy)

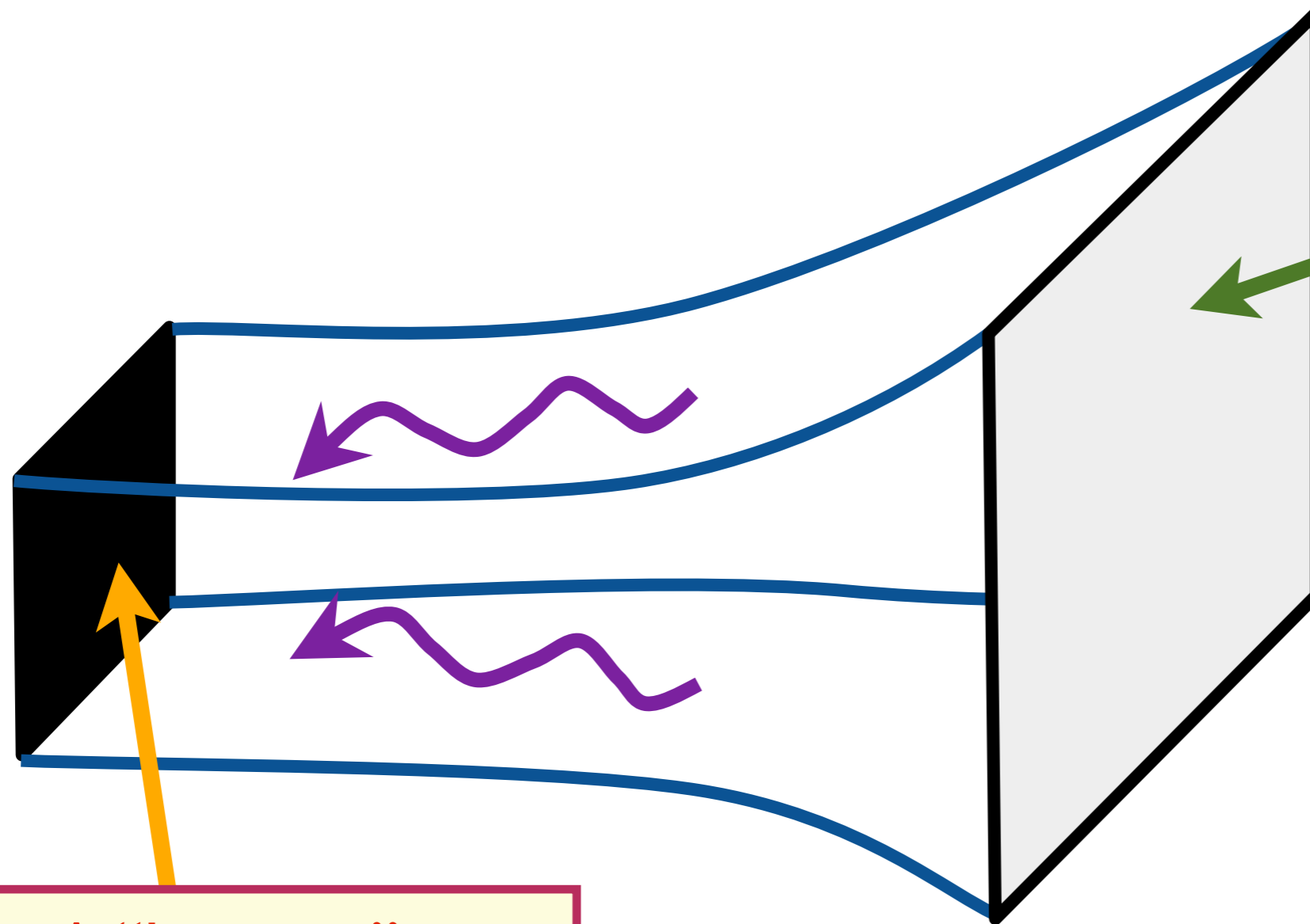
String theory at non-zero temperatures



A “horizon”,
whose temperature
and entropy equal
those of the quantum
critical point

A 2+1
dimensional
system at its
quantum
critical point

String theory at non-zero temperatures

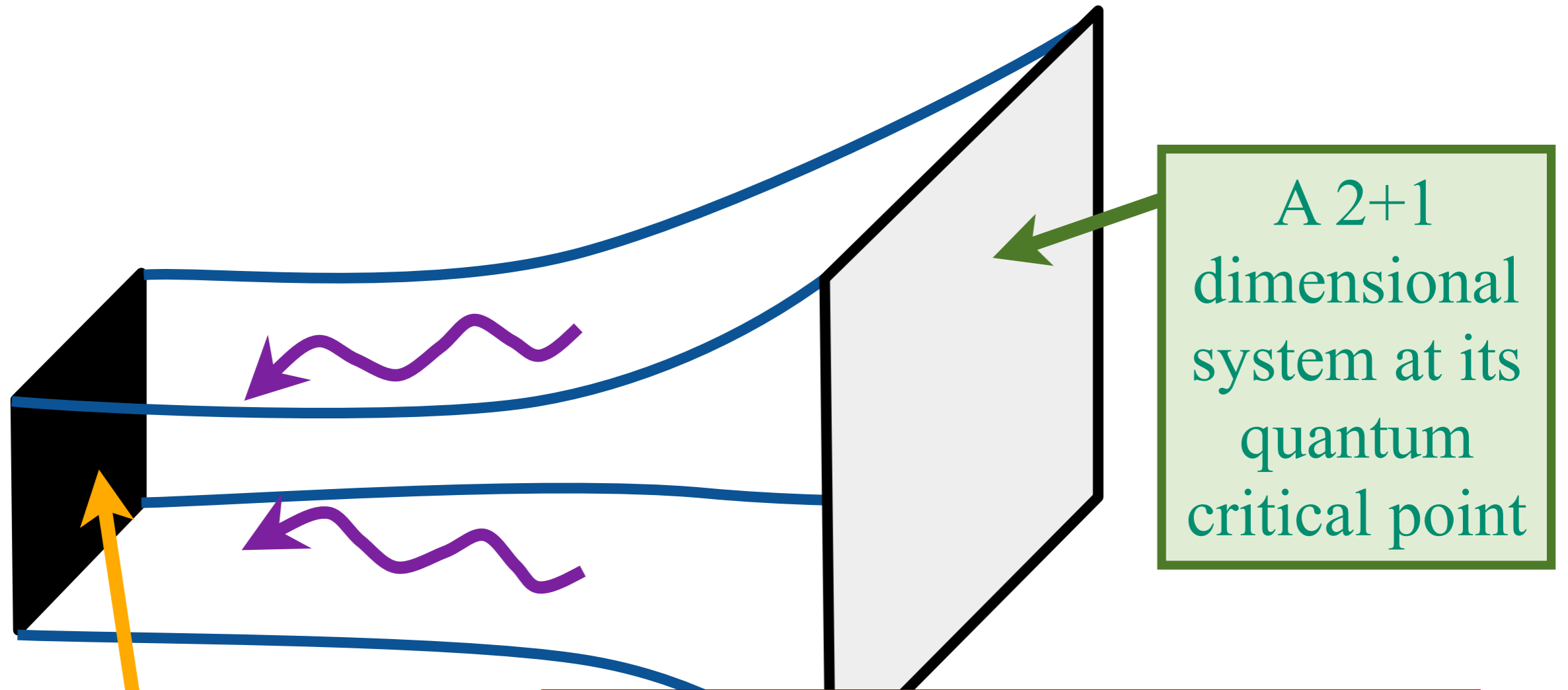


A 2+1
dimensional
system at its
quantum
critical point

A “horizon”,
whose temperature
and entropy equal
those of the quantum
critical point

Friction of quantum
criticality = waves
falling into black brane

String theory at non-zero temperatures



A “horizon”,
whose temperature
and entropy equal
those of the quantum
critical point

An (extended) Einstein-Maxwell
provides successful description of
dynamics of quantum critical
points at non-zero temperatures
(where no other methods apply)

A 2+1
dimensional
system at its
quantum
critical point

**Quantum
superposition and
entanglement**

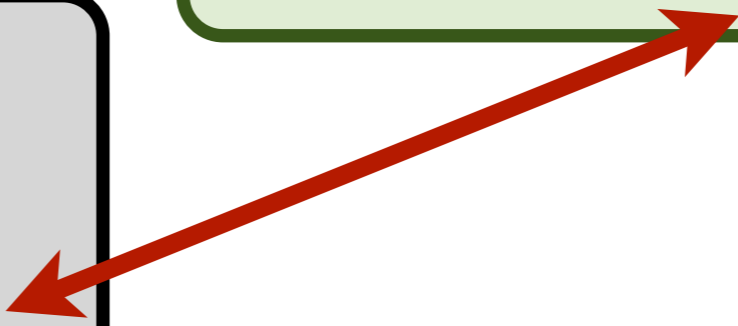
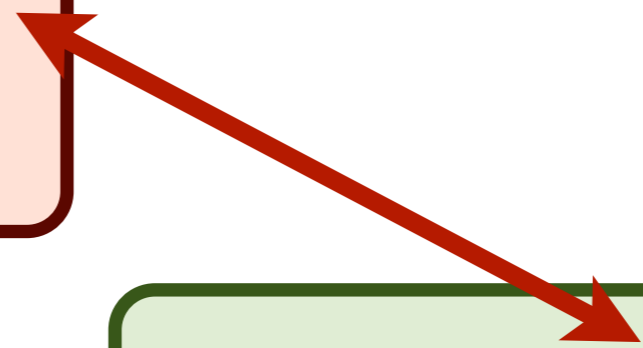
**Quantum critical
points of electrons
in crystals**

**String theory
and black holes**

**Quantum
superposition and
entanglement**

**Quantum critical
points of electrons
in crystals**

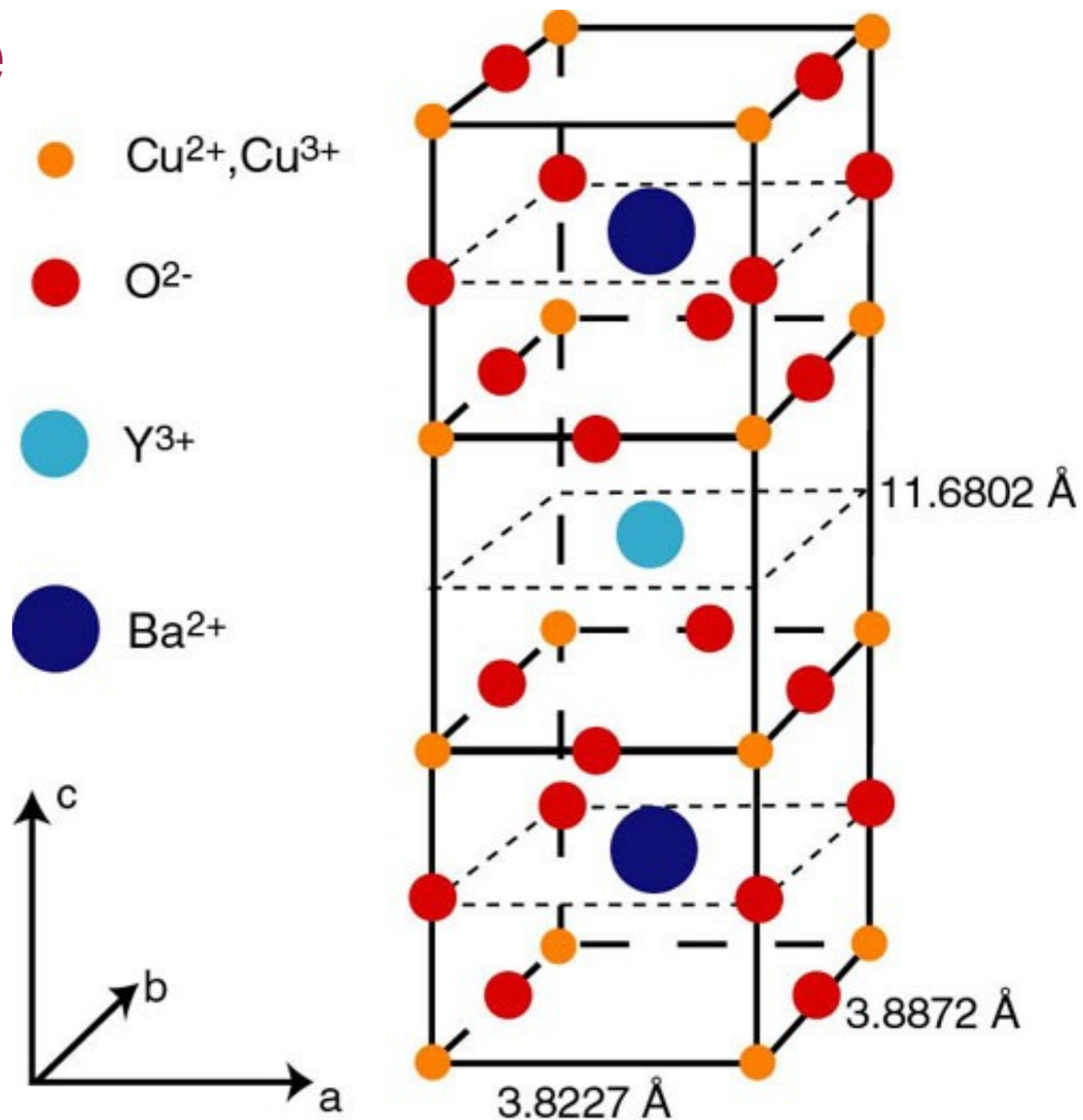
**String theory
and black holes**



**Metals, "strange metals", and
high temperature
superconductors**

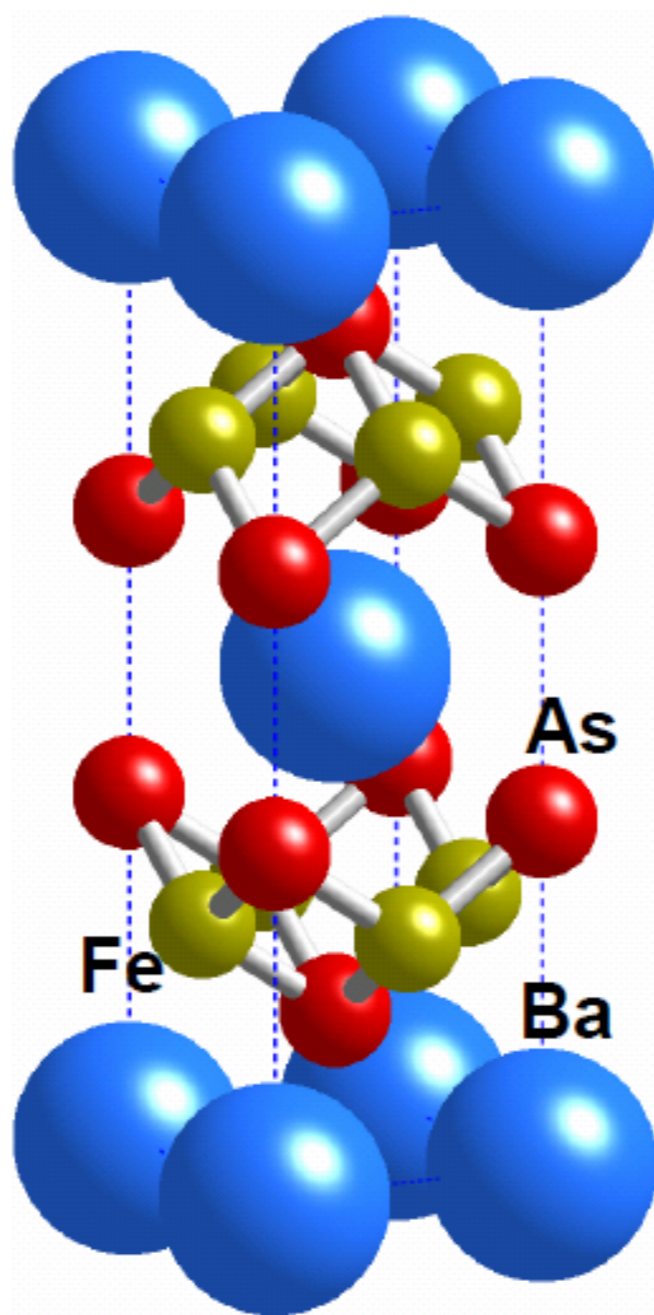
**Insights from gravitational
"duals"**

High temperature superconductors

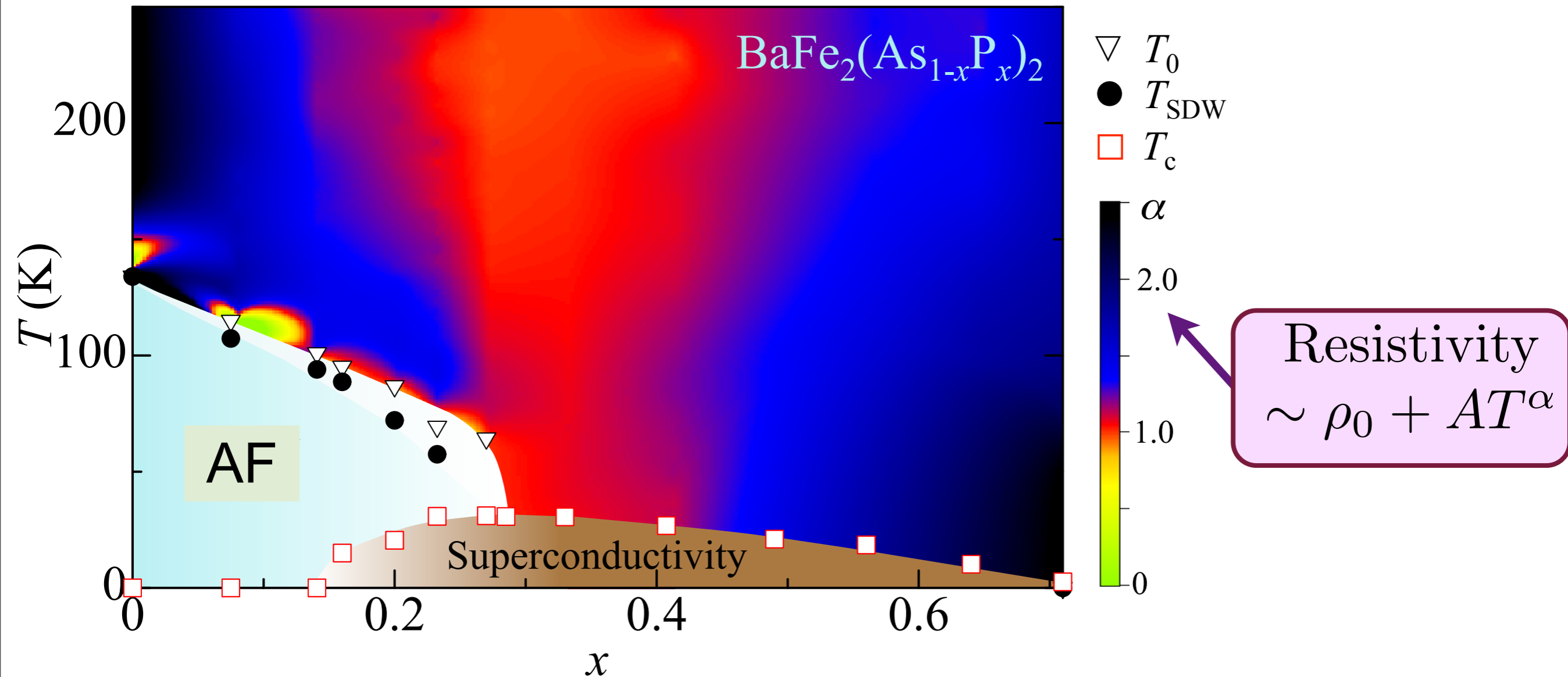


Iron pnictides:

a new class of high temperature superconductors

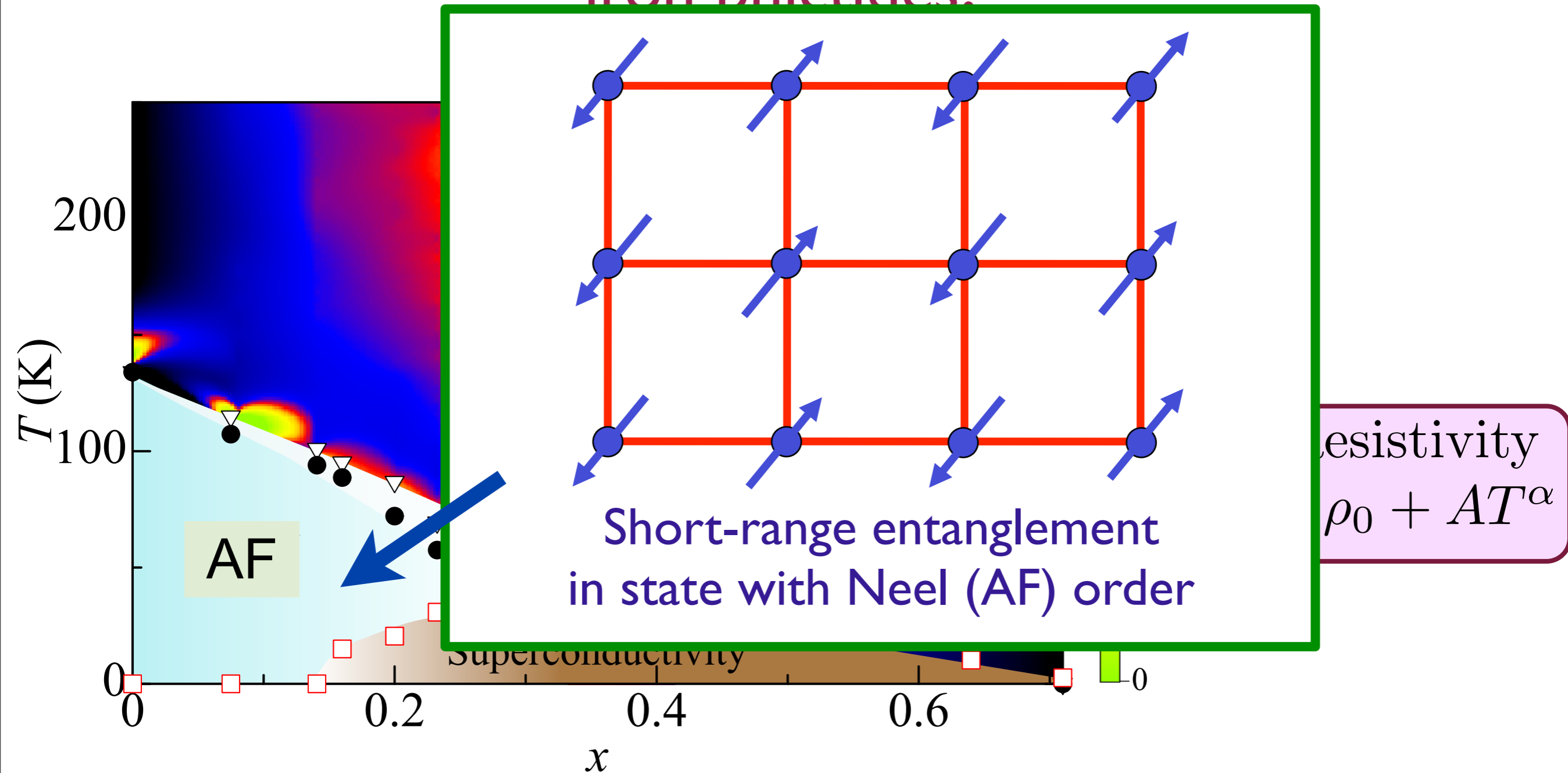


Temperature-doping phase diagram of the iron pnictides:



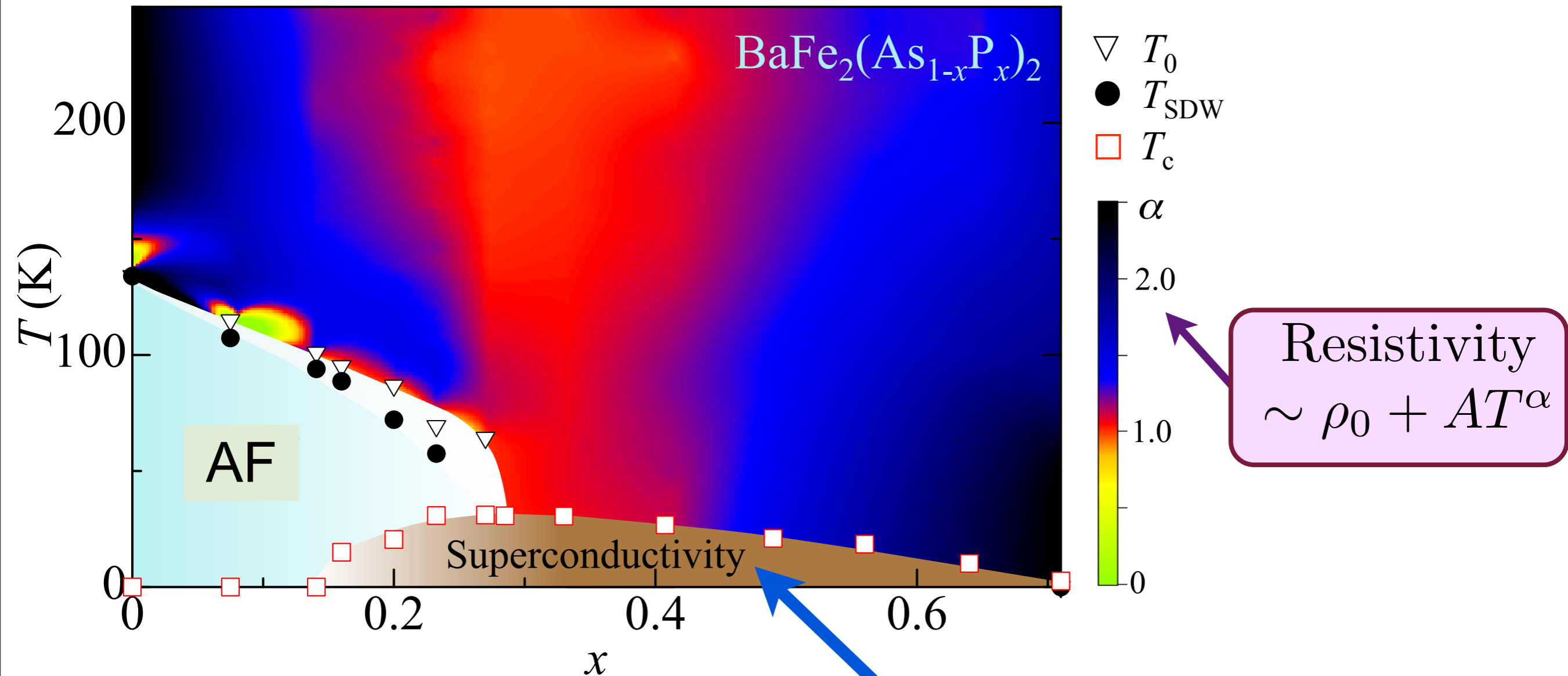
S. Kasahara, T. Shibauchi, K. Hashimoto, K. Ikada, S. Tonegawa, R. Okazaki, H. Shishido, H. Ikeda, H. Takeya, K. Hirata, T. Terashima, and Y. Matsuda, *Physical Review B* **81**, 184519 (2010)

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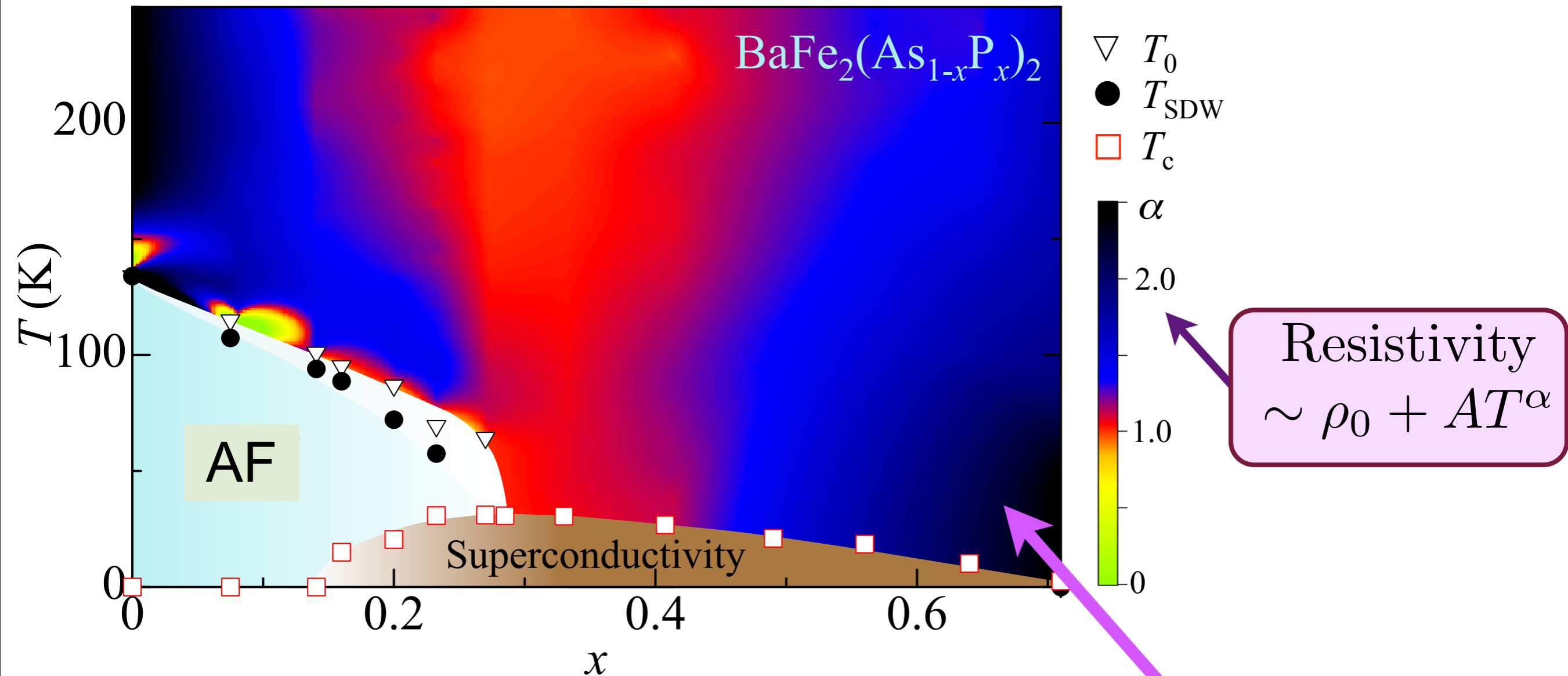
Temperature-doping phase diagram of the iron pnictides:



Superconductor
 Bose condensate of pairs of electrons
 Short-range entanglement

S. Kasahara, T. Shibauchi,
 H. Ikeda

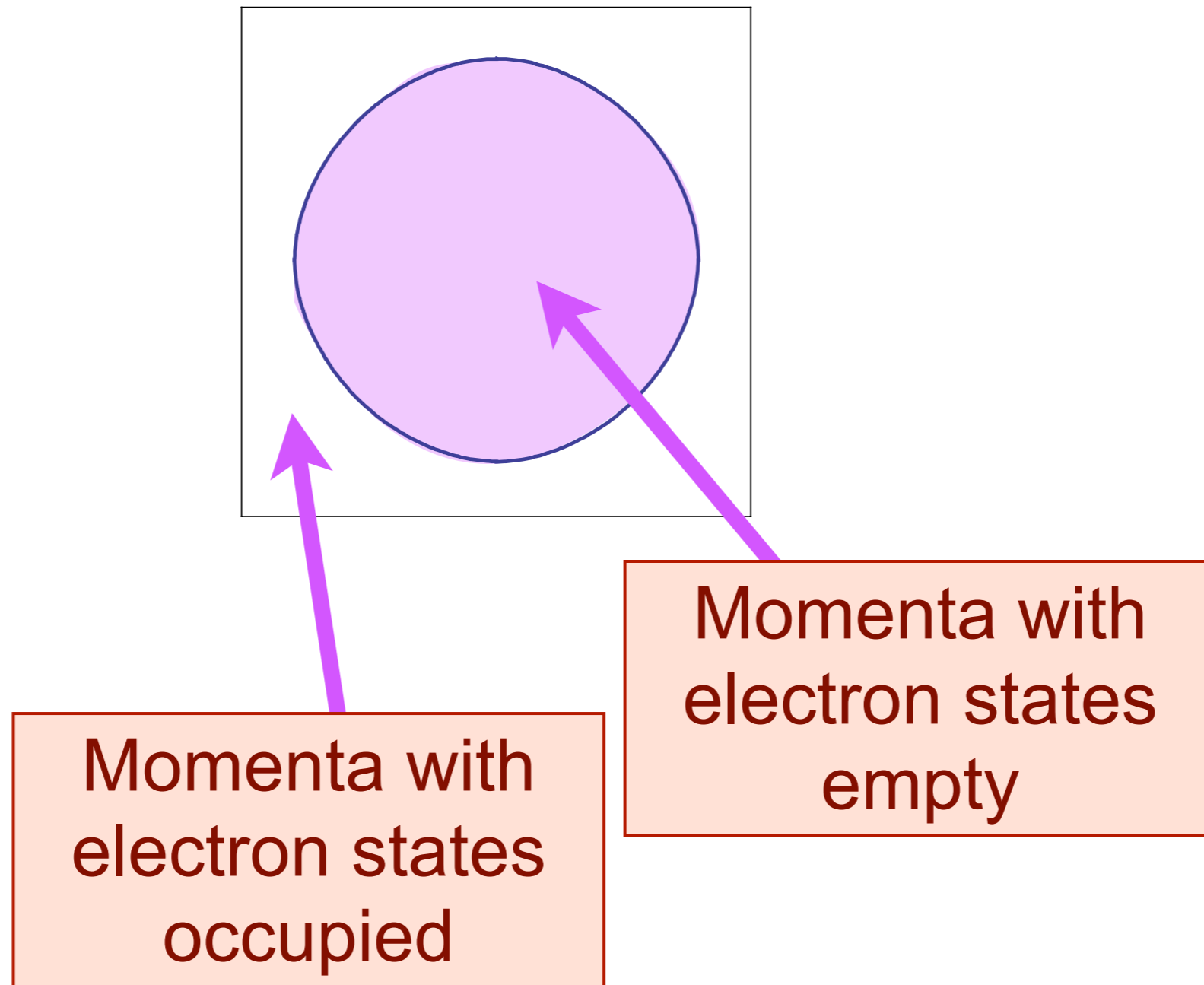
Temperature-doping phase diagram of the iron pnictides:



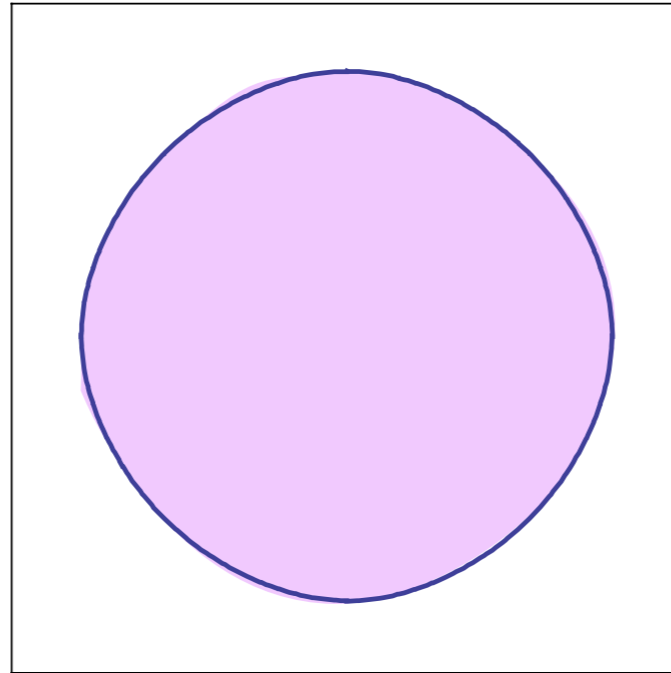
Ordinary metal
(Fermi liquid)

S. Kasahara, T. Shibauchi, K. Hashimoto, K. Ikada, S.
 H. Ikeda, H. Takeya, K. Hirata, T. Terasa
Physical Review B **81**, 184519 (2010)

Sommerfeld theory of ordinary metals



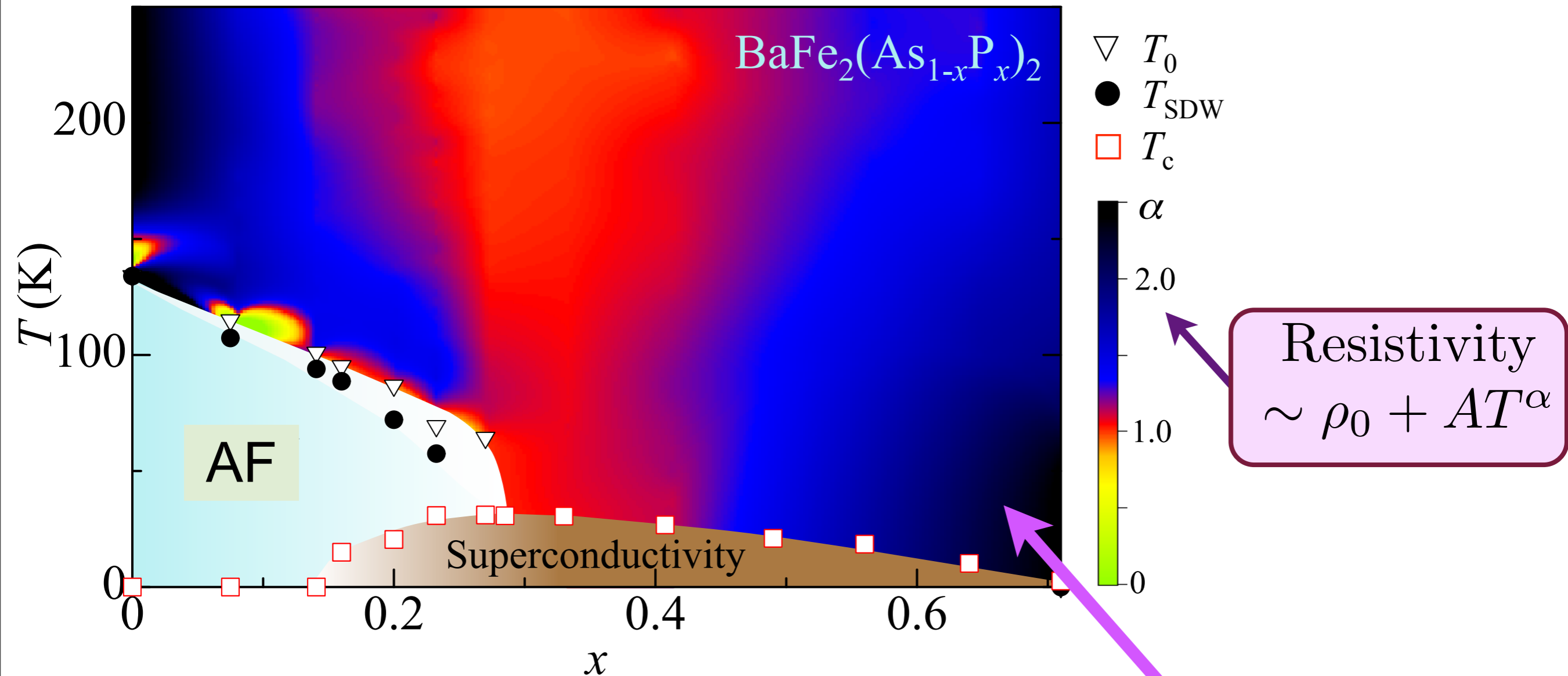
Sommerfeld theory of ordinary metals



Key feature of the Sommerfeld theory: the Fermi surface

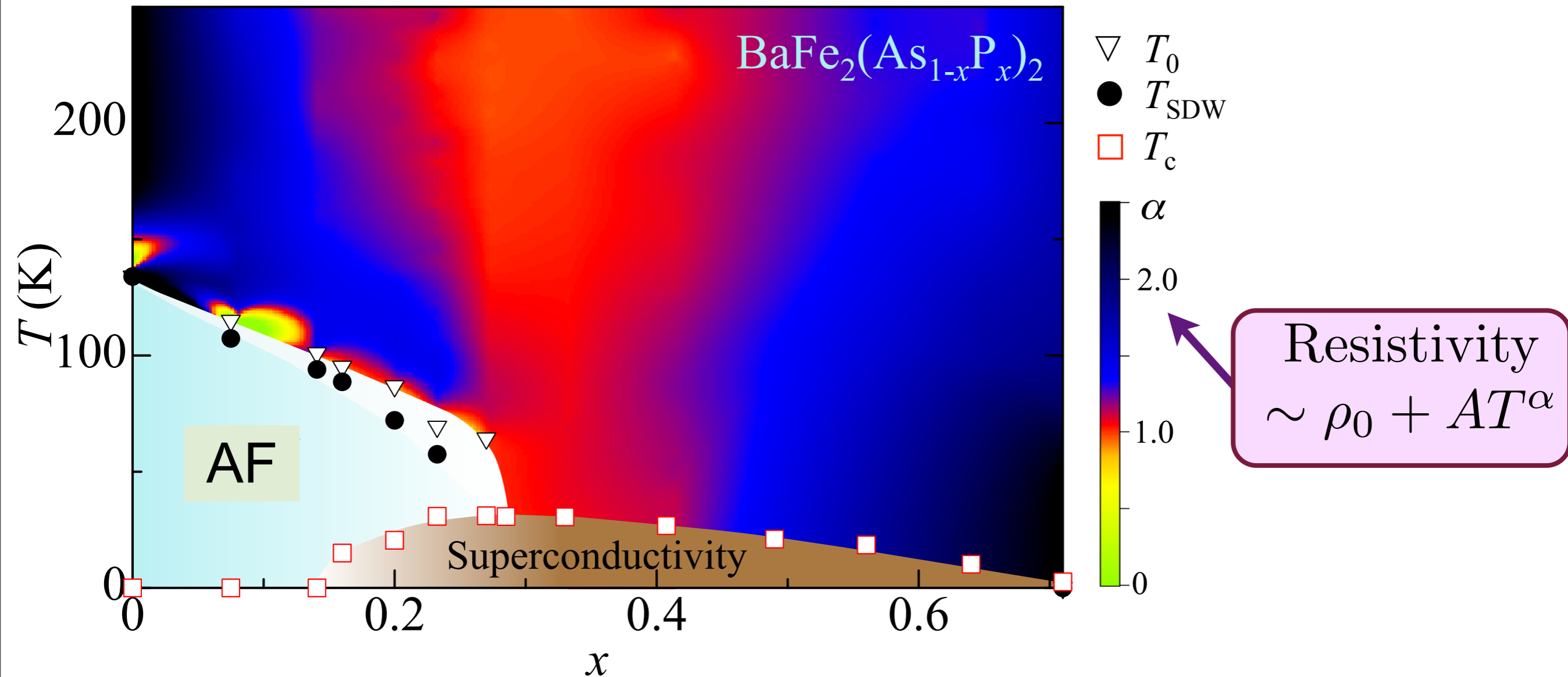
The Fermi surface separates regions of occupied and empty electron states, and is responsible for most of the familiar properties of ordinary metals, such as resistivity $\sim T^2$.

Temperature-doping phase diagram of the iron pnictides:



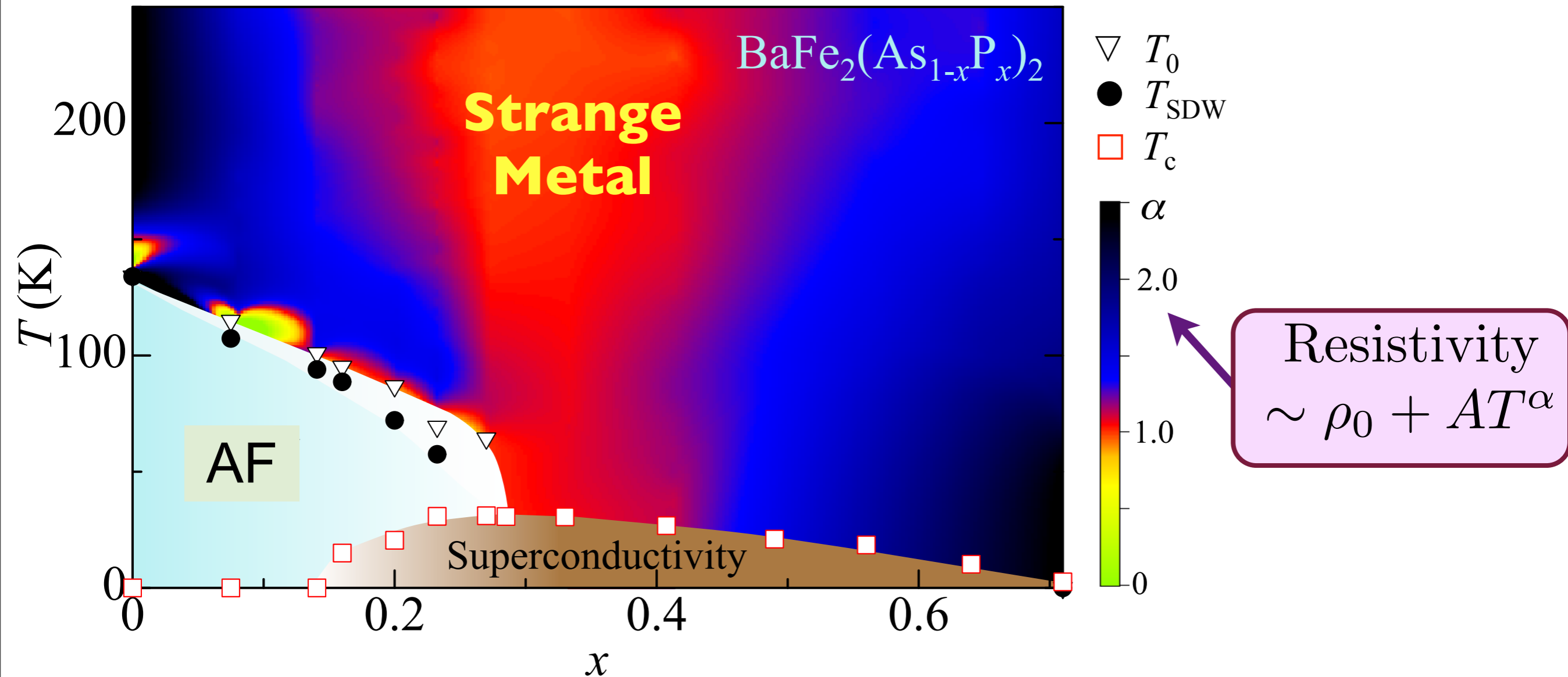
S. Kasahara, T. Shibauchi, K. Hashimoto, K. Ikada, S.
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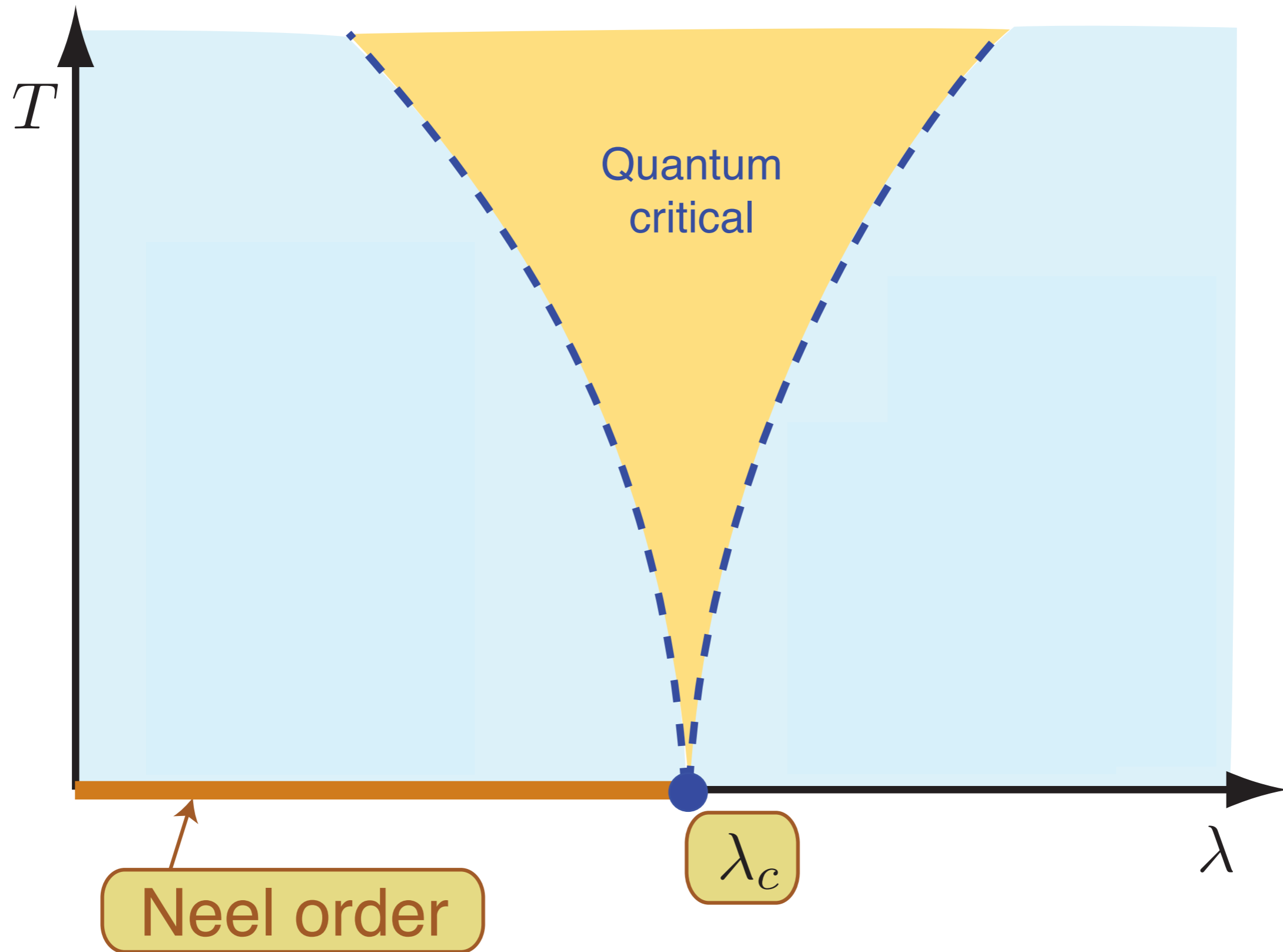
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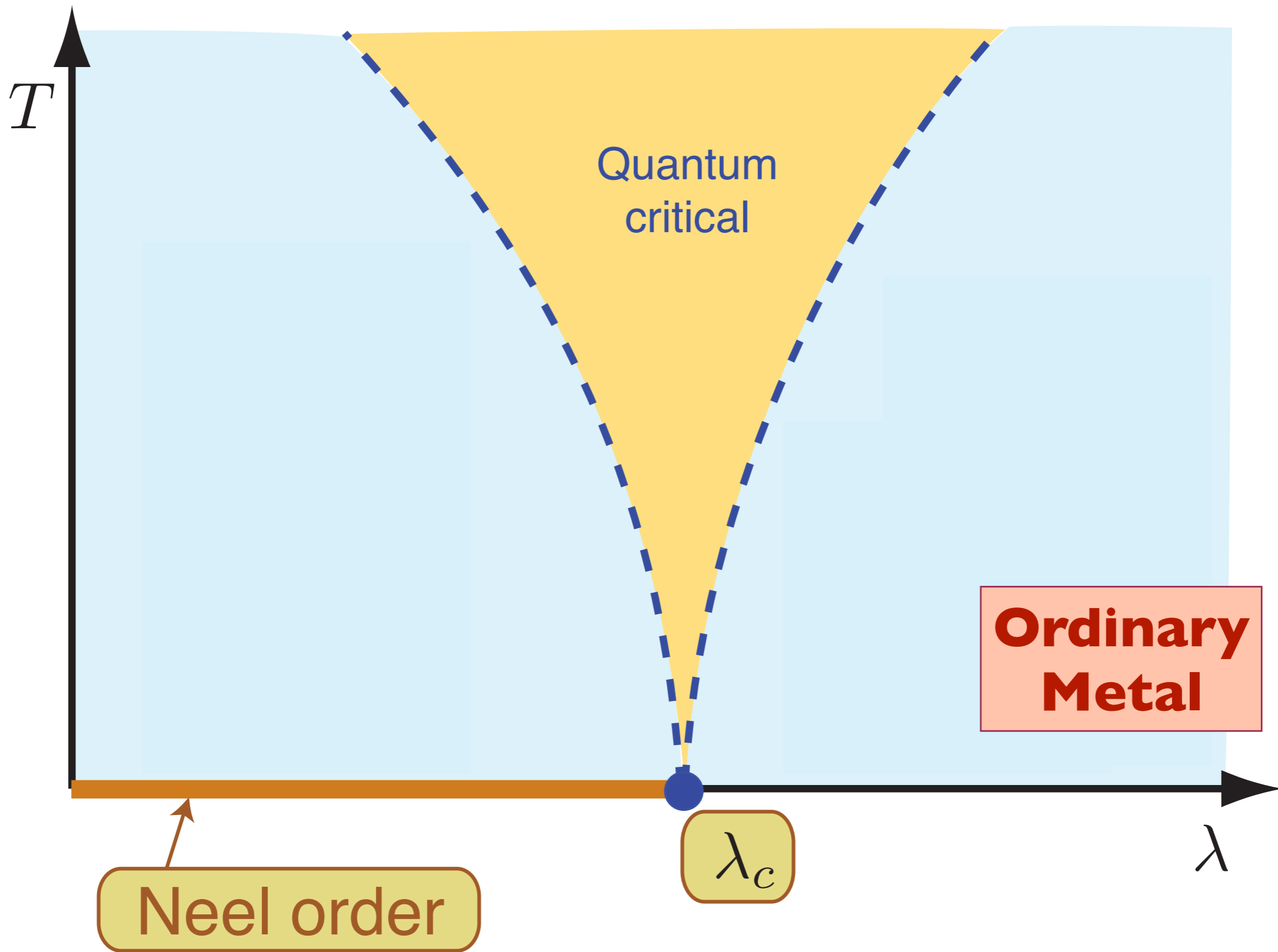
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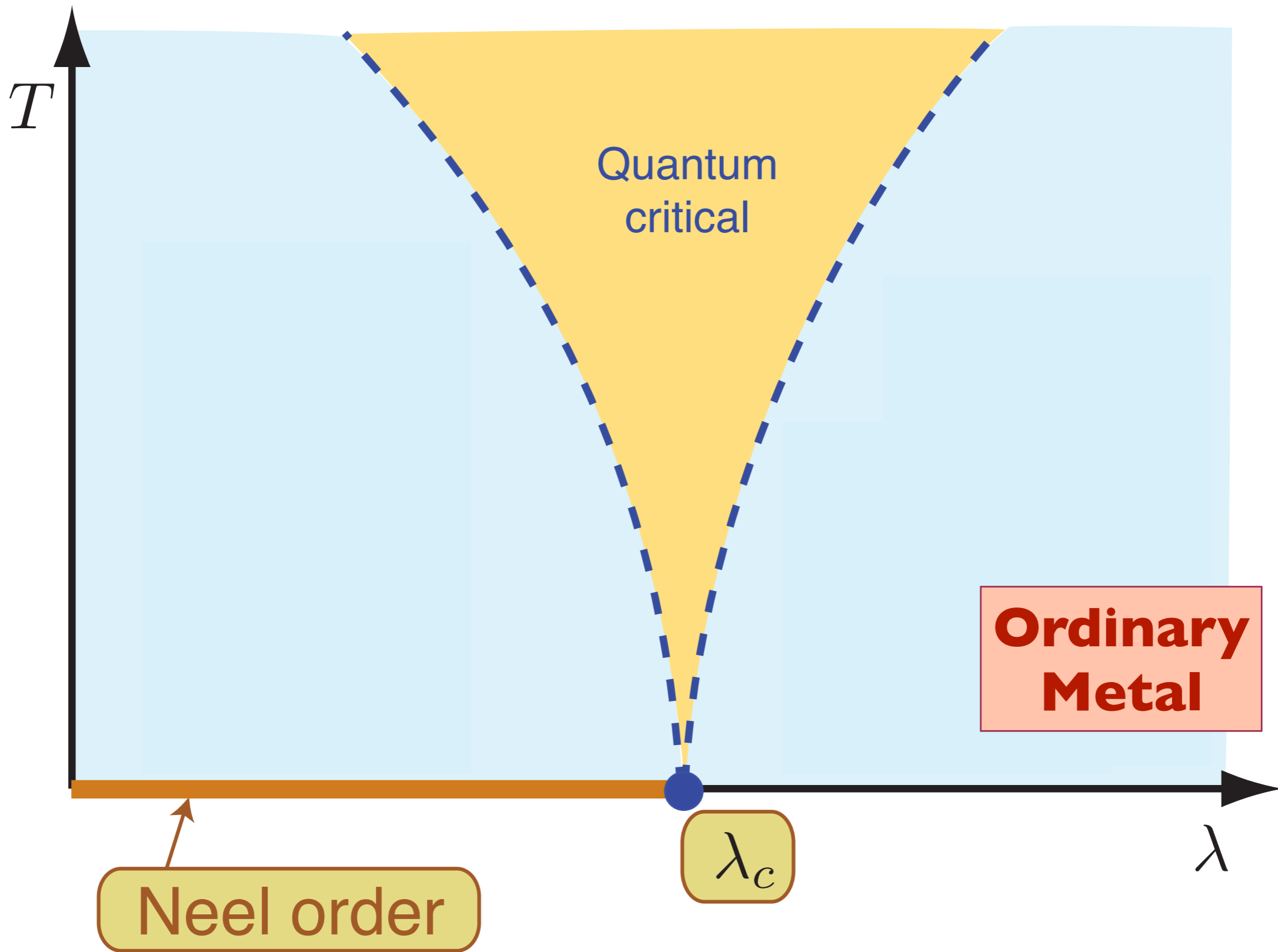


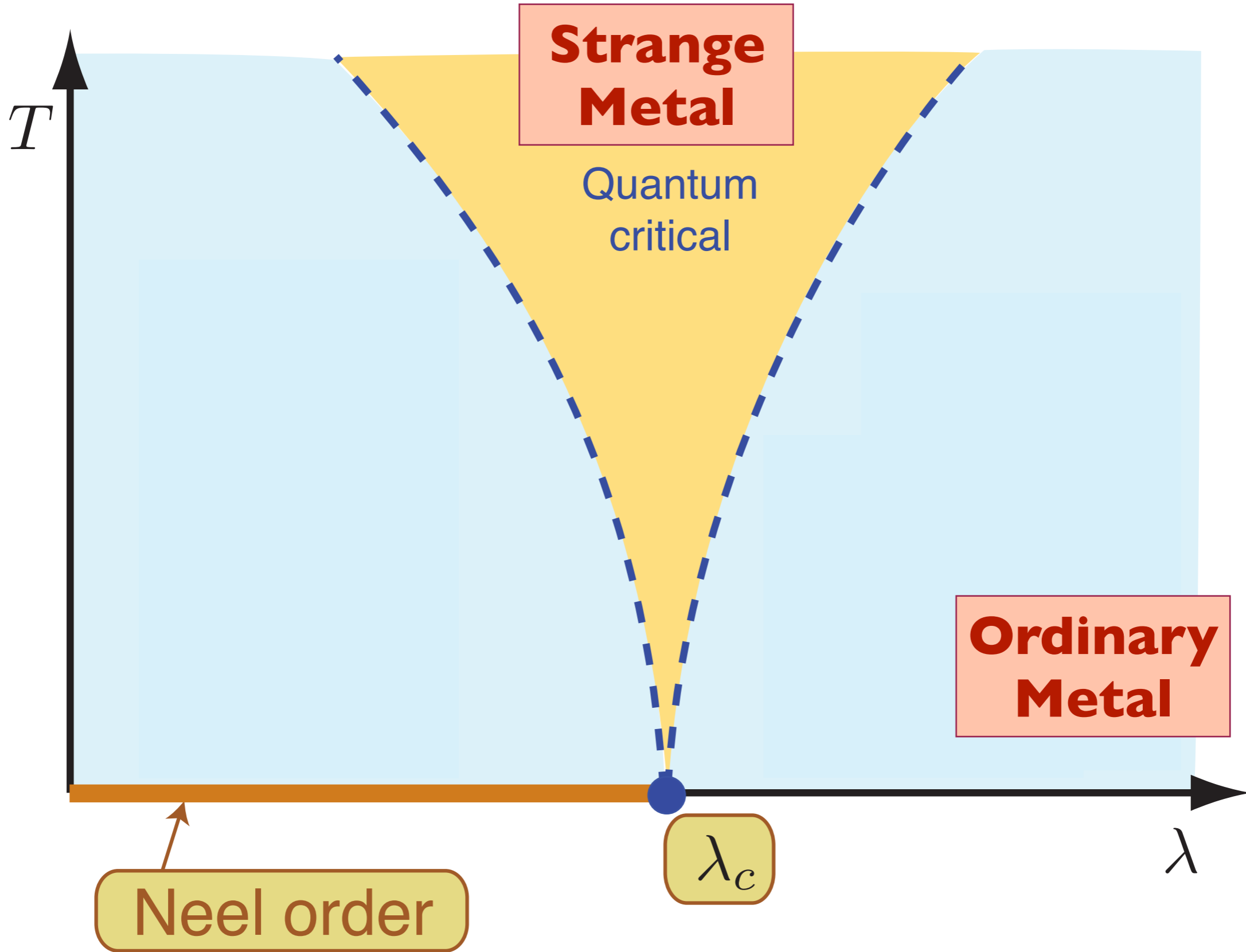
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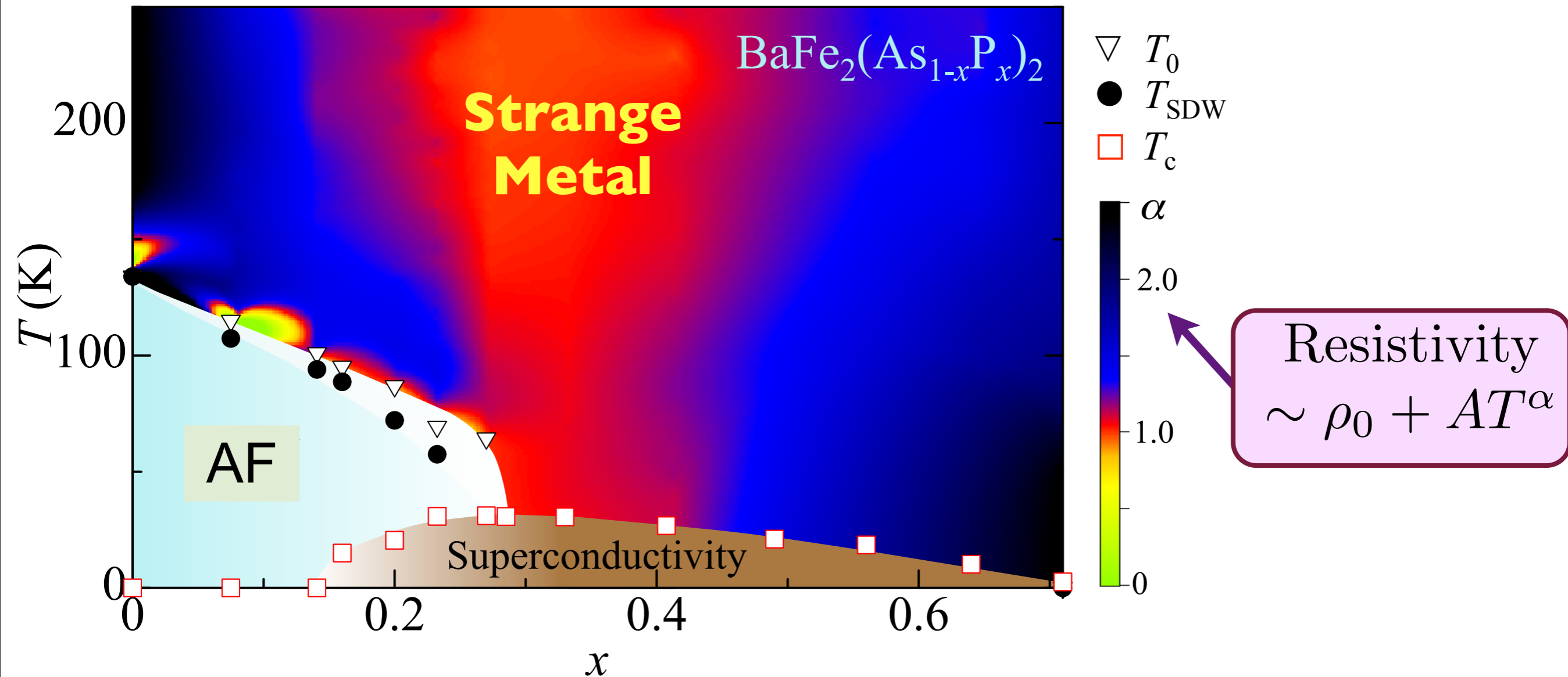






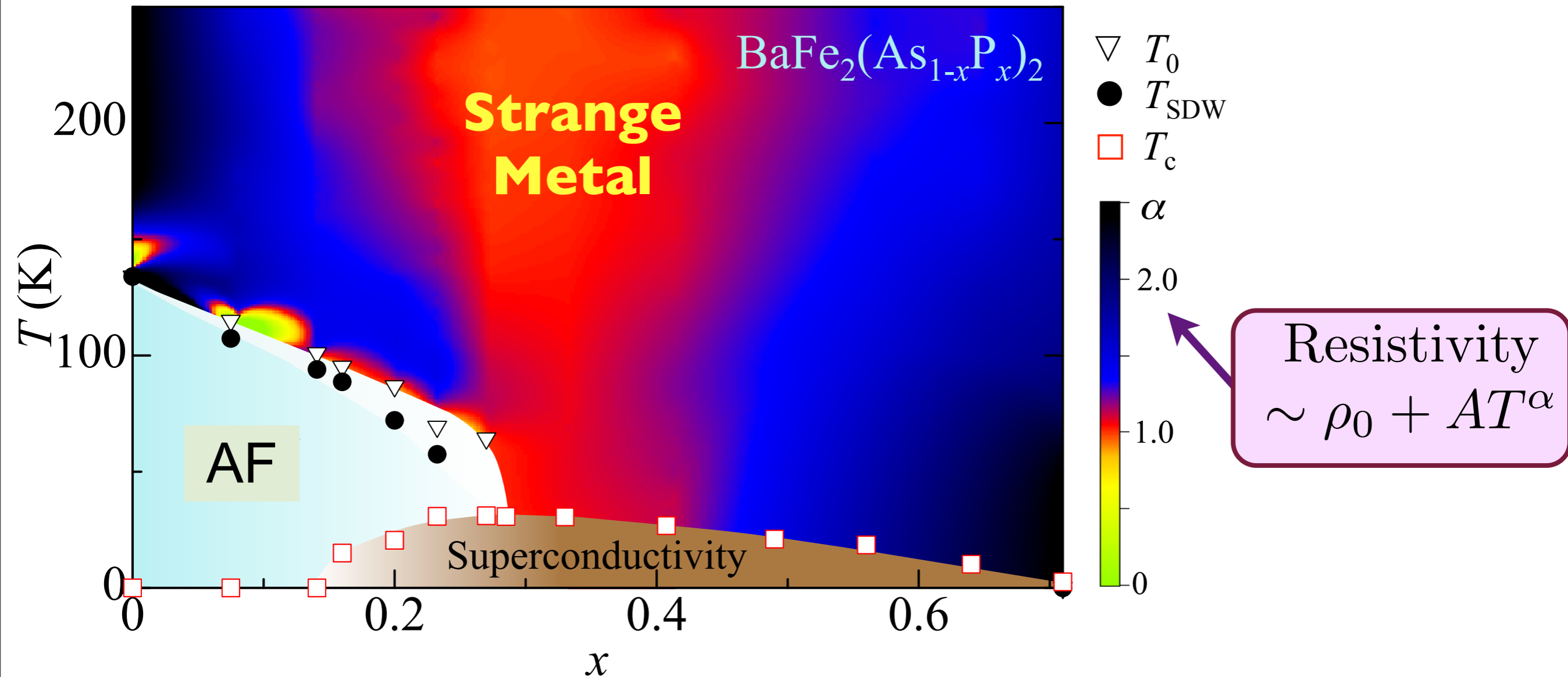


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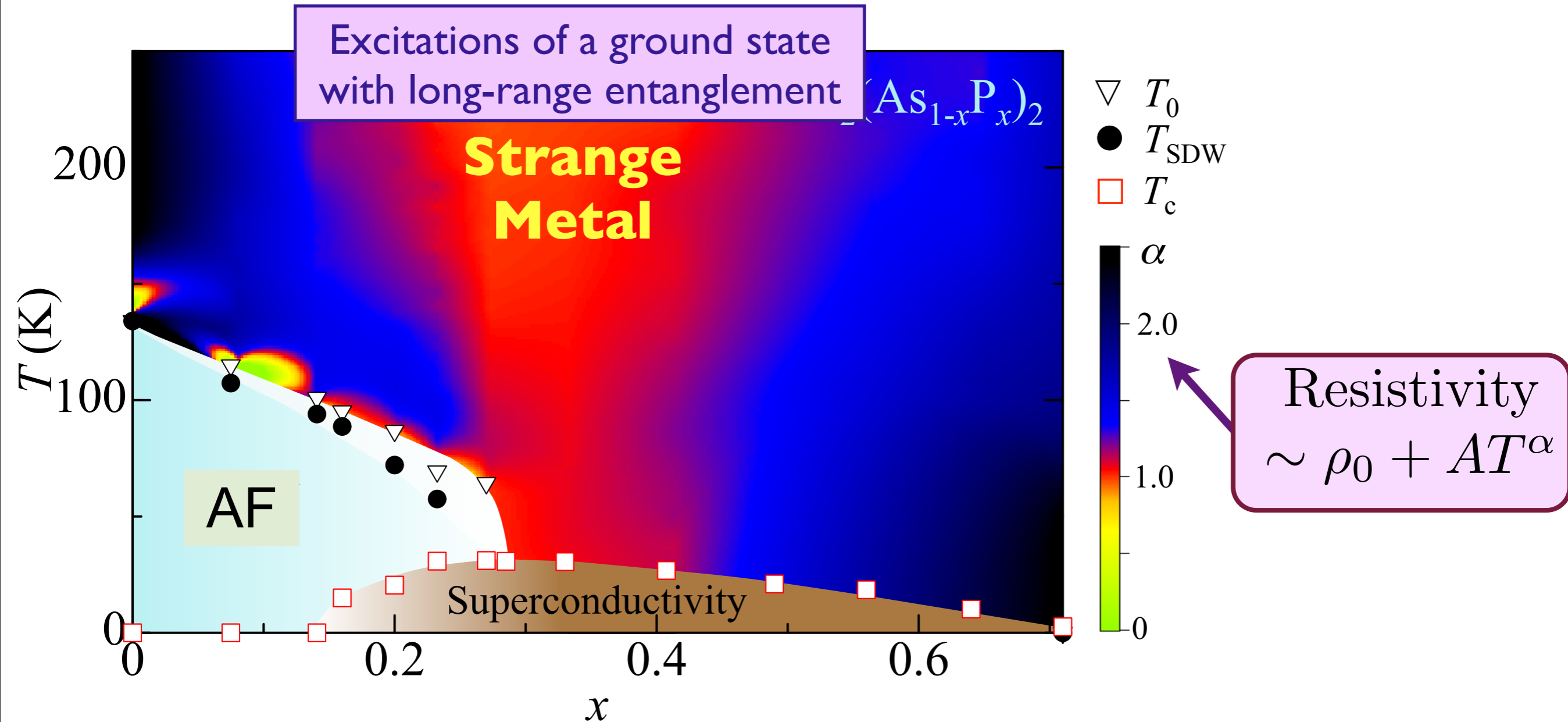
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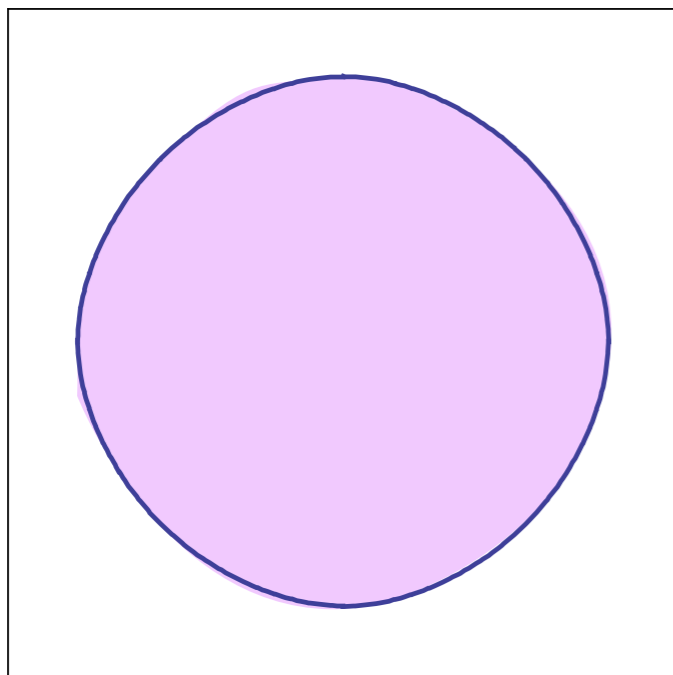
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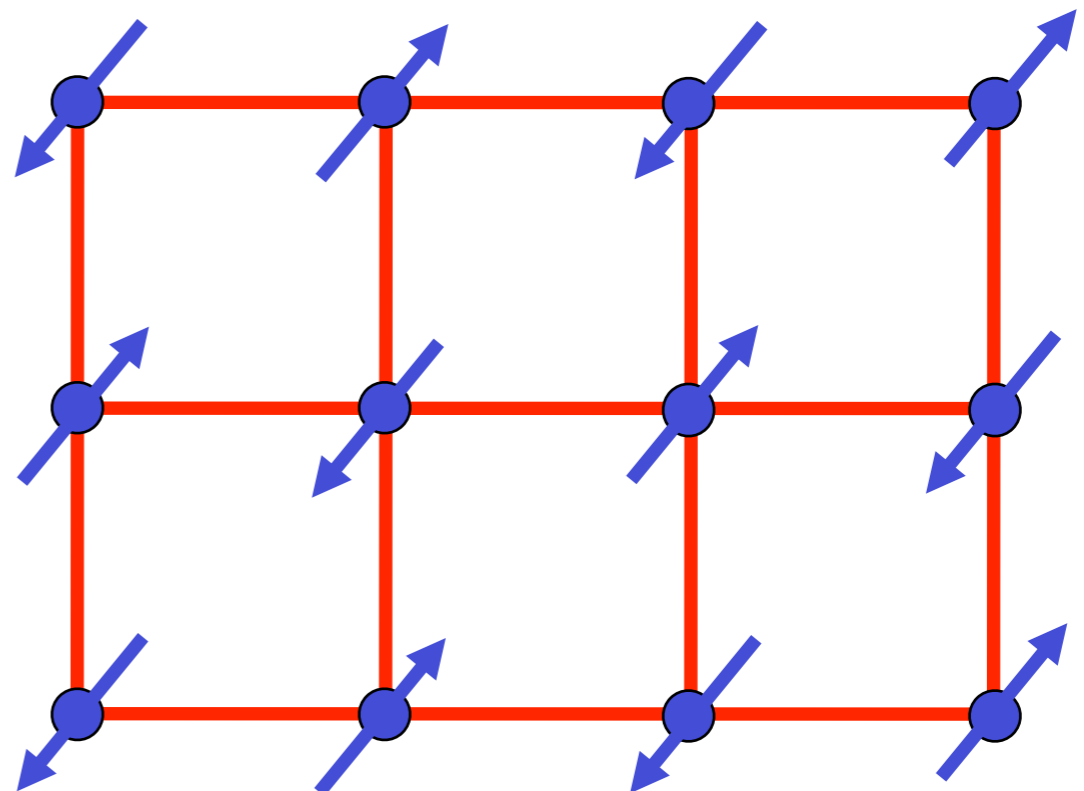
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Key (difficult) problem:

Describe quantum critical points and phases of systems with Fermi surfaces leading to metals with novel types of long-range entanglement



+



Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Can we obtain gravitational theories of superconductors and ordinary Sommerfeld metals ?

Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Can we obtain gravitational theories of superconductors and ordinary Sommerfeld metals ?

Yes

Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Do the “holographic” gravitational theories yield metals distinct from ordinary Sommerfeld metals ?

Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Do the “holographic” gravitational theories yield metals distinct from ordinary Sommerfeld metals ?

Yes, lots of them, with many “strange” properties !

Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Do any of the holographic “strange metals” have the correct type of long-range entanglement linked to Fermi surfaces ?

Challenge to string theory:

Describe quantum critical points and phases of metallic systems

Do any of the holographic “strange metals” have the correct type of long-range entanglement linked to Fermi surfaces ?

Yes, a very select subset has the proper logarithmic violation of the area law of entanglement entropy !!

These are now being studied intensively.....

Conclusions

Phases of matter with long-range quantum entanglement are prominent in numerous modern materials.

Conclusions

Simplest examples of long-range entanglement are at quantum-critical points of insulating antiferromagnets

Conclusions

More complex examples in metallic states are experimentally ubiquitous, but pose difficult strong-coupling problems to conventional methods of field theory

Conclusions

String theory and gravity in emergent dimensions offer a remarkable new approach to describing states with long-range quantum entanglement.

Conclusions

String theory and gravity in emergent dimensions offer a remarkable new approach to describing states with long-range quantum entanglement.

Much recent progress offers hope of a holographic description of “strange metals”