



Intertwining topological order and broken symmetries in insulators and metals



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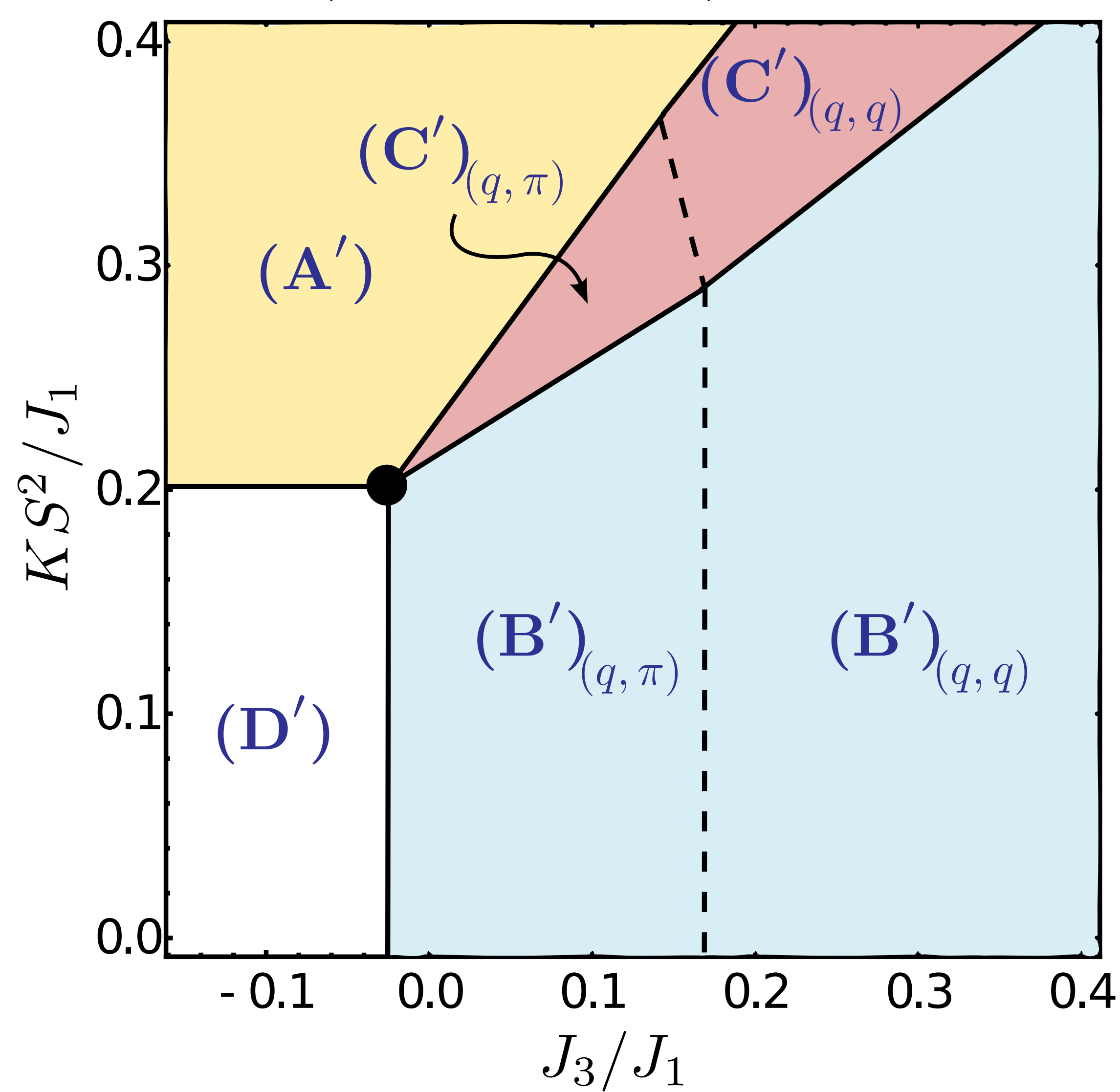
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arXiv:1703.00014

arXiv:1705.06289

Classical phase diagram of square lattice antiferromagnet with near-neighbor exchanges J_1, J_2, J_3, J_4 and ring-exchange K

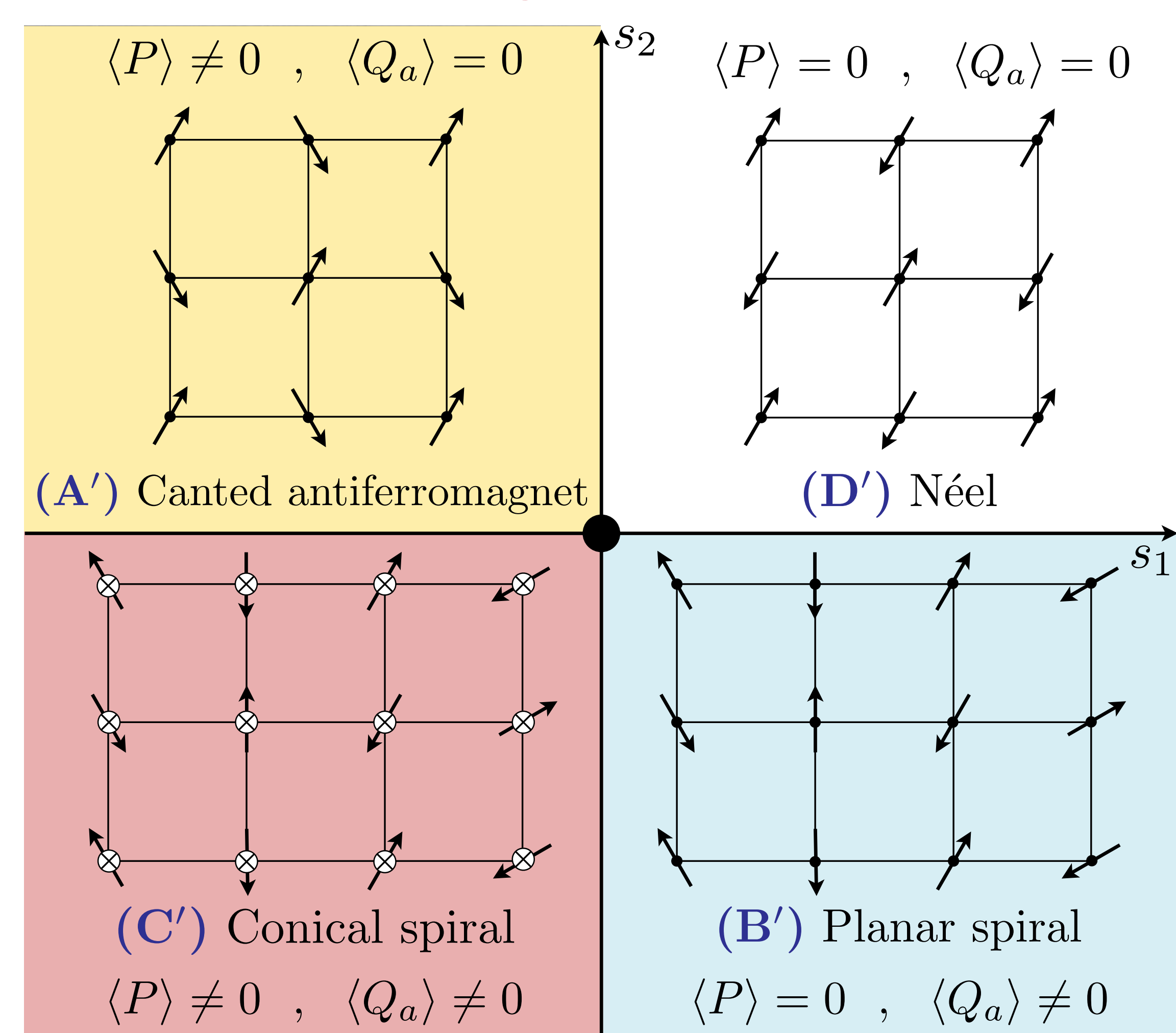
$$J_2/J_1 = 0.05, J_4/J_1 = -0.1$$



Insulators

CP^1 model of bosonic spinons z_α coupled to a $U(1)$ gauge field a_μ , and charge 2 Higgs fields $P \sim \varepsilon_{\alpha\beta} z_\alpha \partial_t z_\beta$ and $Q_a \sim \varepsilon_{\alpha\beta} z_\alpha \partial_a z_\beta$ ($a = x, y$).

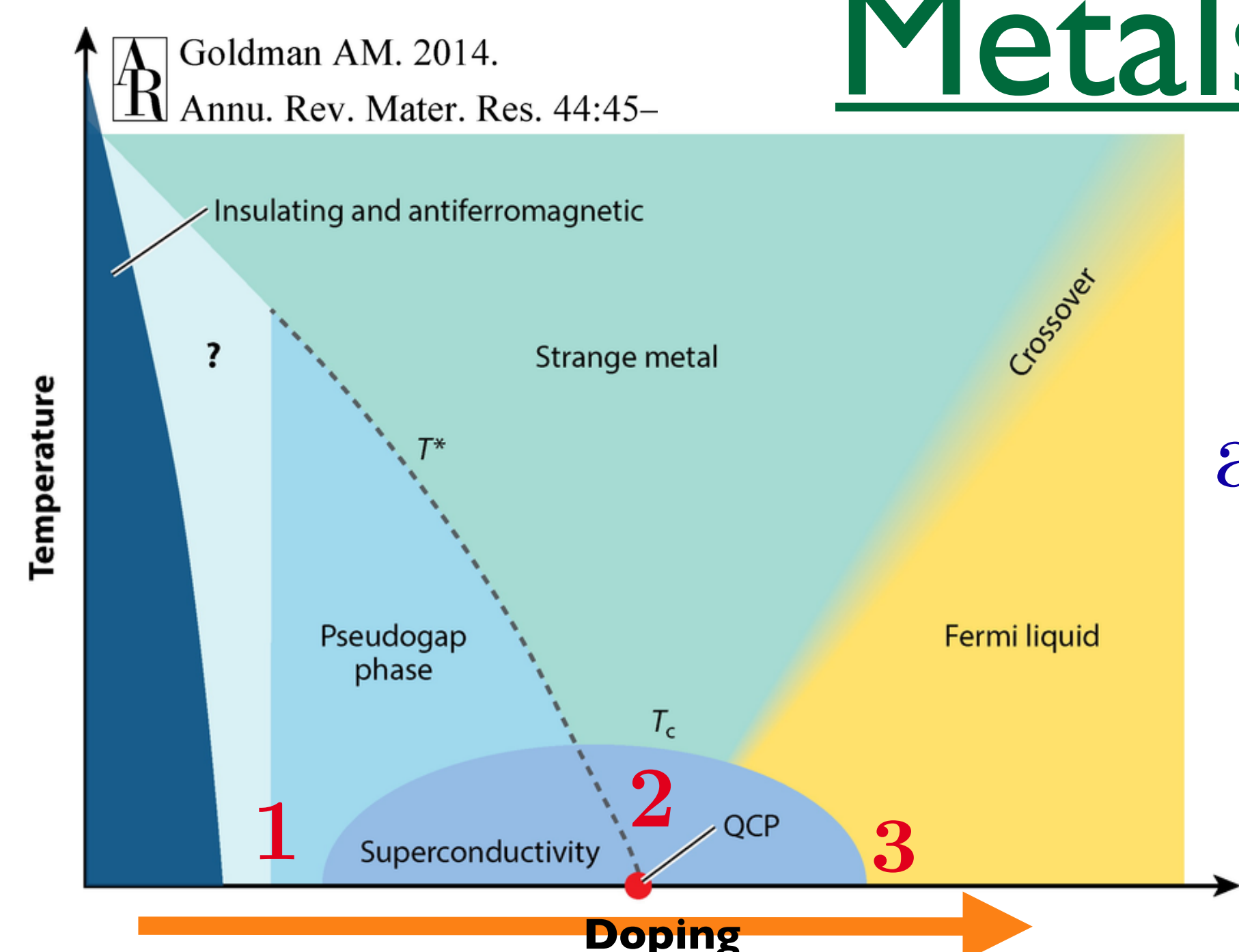
$$\mathcal{L} = \frac{1}{g} |(\partial_\mu - ia_\mu) z_\alpha|^2 + s_1 |P|^2 + s_2 |Q_a|^2$$



Small g with $\langle z_\alpha \rangle \neq 0$

Metals

SU(2) gauge theory with bosonic spinons R , adjoint Higgs field H^a , and fermionic chargons ψ .



$$c = R\psi$$

$$\sigma^a \Phi^a = R \sigma^b H^b R^\dagger$$

R is a 2×2 unitary matrix which transforms the spin components of the electron into a rotating reference frame. The Higgs field, H^a , measures the local magnetic order in the rotating reference frame.

Field	Symbol	Statistics	SU(2) _{gauge}	SU(2) _{spin}	U(1) _{e.m.charge}
Electron	c	fermion	1	2	-1
Magnetic moment	Φ^a	boson	1	3	0
Chargon	ψ	fermion	2	1	-1
Spinon	R	boson	2	2	0
Higgs	H^a	boson	3	1	0

SU(2) gauge theory coupled to $N_f = 2$ fundamental Dirac fermion spinons X , and adjoint Higgs fields Φ, Φ_1 , and Φ_2

$$\mathcal{L}_{\text{QCD}_3} = i \text{tr} (\bar{X} \gamma^\mu (\partial_\mu X + i X a_\mu)) + (D_\mu \Phi^a)^2 + s (\Phi^a)^2 + \lambda_2 \Phi^a \text{tr} (\sigma^a \bar{X} \mu^\nu X) + (D_\mu \Phi_1^a)^2 + \bar{s}_1 (\Phi_1^a)^2 + \lambda_3 \Phi_1^a \text{tr} (\sigma^a \bar{X} \mu^\nu \gamma^\tau \partial_\tau X) + (D_\mu \Phi_2^a)^2 + \bar{s}_2 (\Phi_2^a)^2 + \lambda_5 \Phi_2^a \text{tr} (\sigma^a \bar{X} \mu^\nu \gamma^\tau \partial_\tau X)$$

LGW-Hertz criticality of antiferromagnetism

S. Sachdev, M.A. Metlitski, Y. Qi, and C. Xu, Phys. Rev. B **80**, 155129 (2009)

