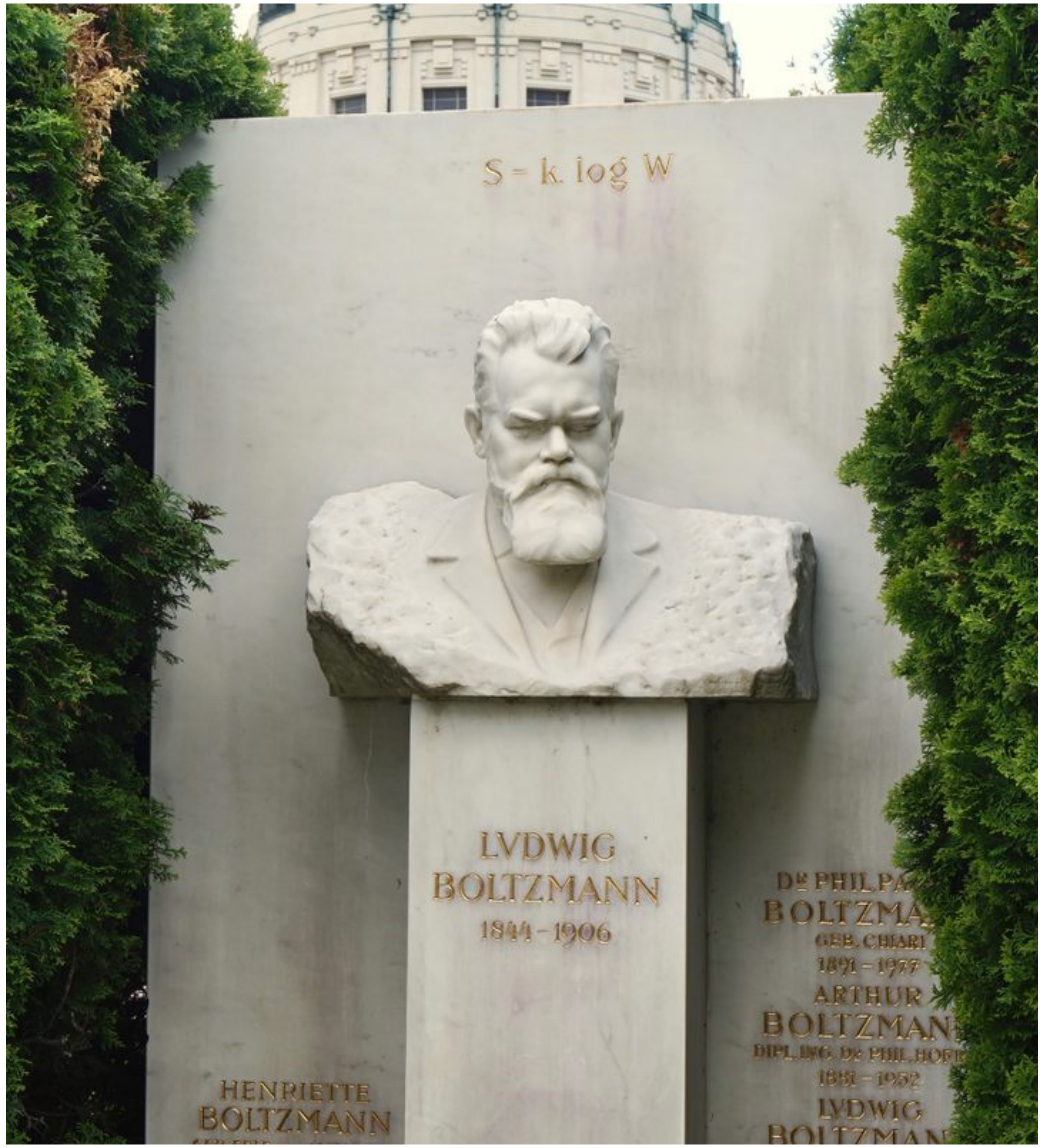
Statistical mechanics of strange metals and black holes

Kansas State University
April 25, 2022

Subir Sachdev

Talk online: sachdev.physics.harvard.edu

Foundations
by
Boltzmann



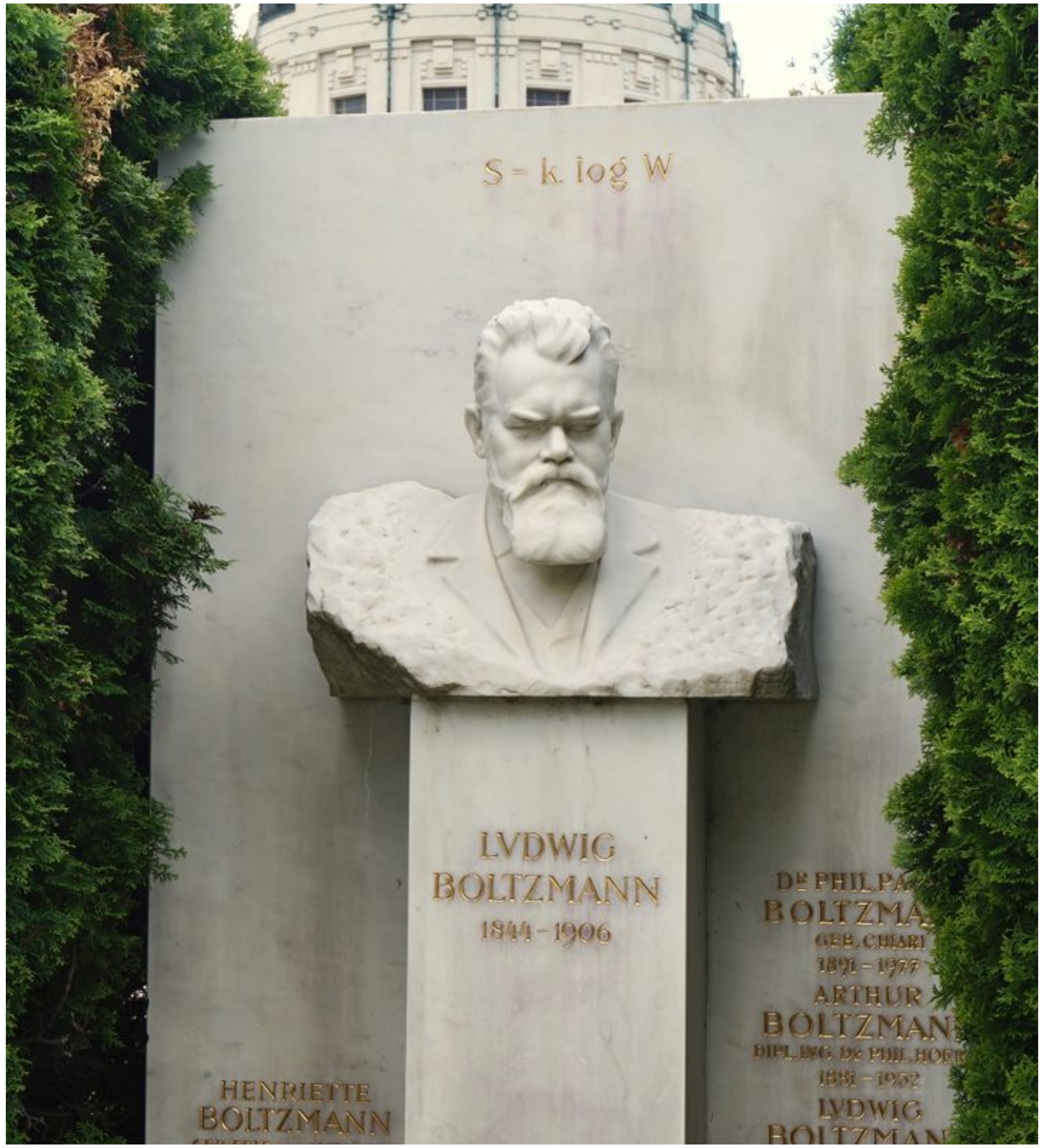
Ludwig Boltzmann
20 February 1844 - September 5, 1906
Vienna, Austria

Statistical interpretation of entropy (1870)

$$S = k_B \log W$$

Density of quantum states $D(E) = \exp(S(E)/k_B)$

Boltzmann equation (1872)
Dilute classical gas



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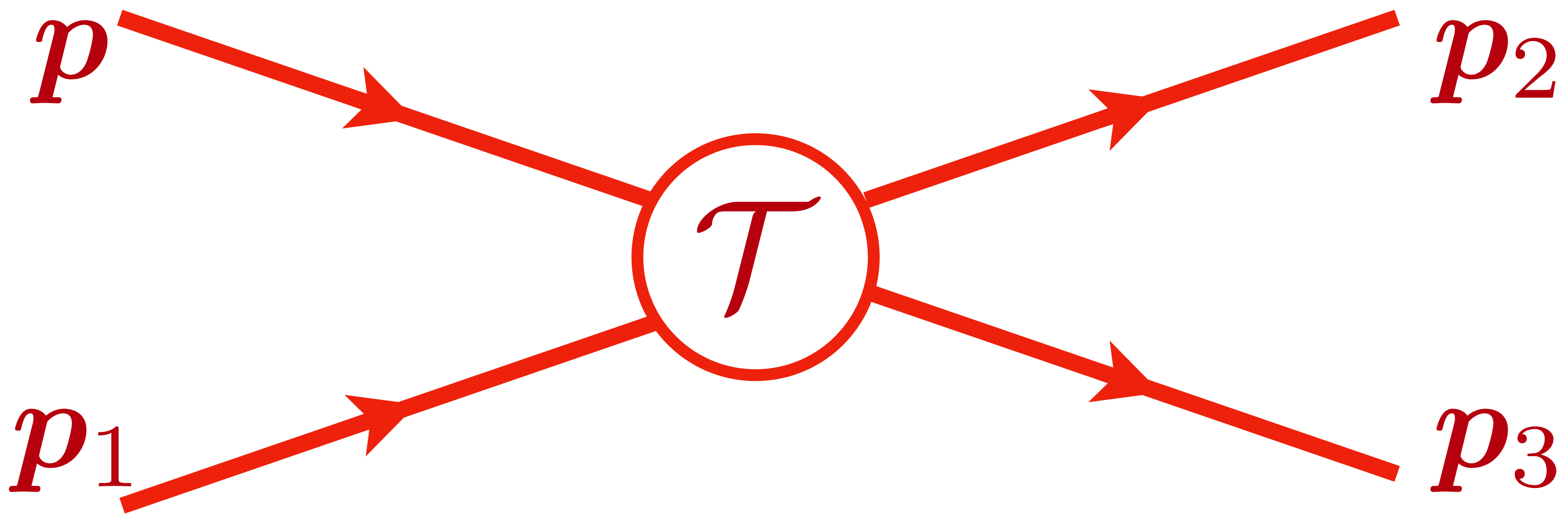
Molecular chaos: successive collisions are statistically independent

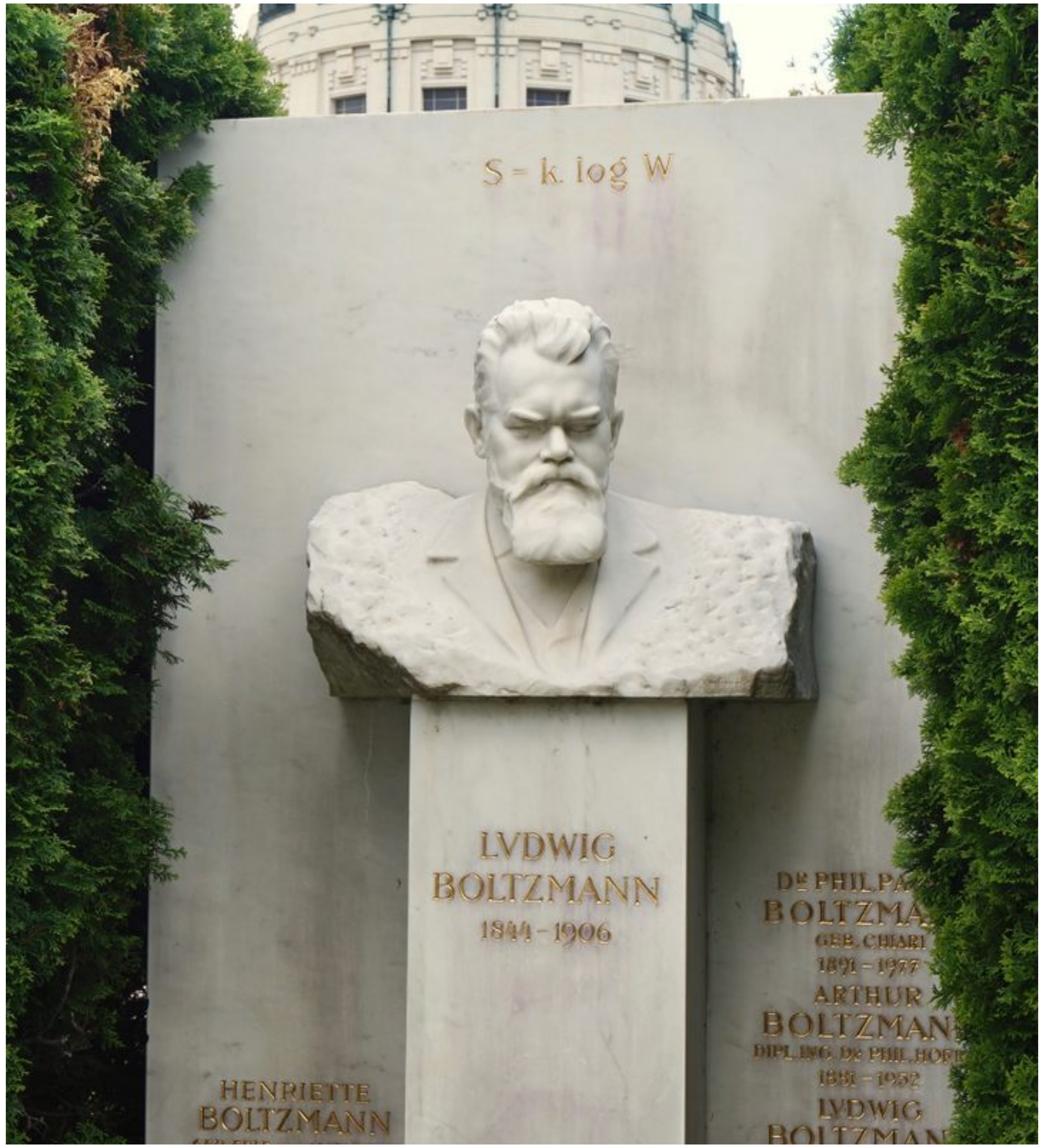
$$\frac{\partial f_{\boldsymbol{p}}}{\partial t} + \frac{\partial \varepsilon_{\boldsymbol{p}}}{\partial \boldsymbol{p}} \cdot \nabla_{\boldsymbol{r}} f_{\boldsymbol{p}} + \boldsymbol{F} \cdot \nabla_{\boldsymbol{p}} f_{\boldsymbol{p}} =$$

$$-2\pi \int_{\boldsymbol{p}_{1,2,3}} |\mathcal{T}|^2 \delta(\varepsilon_{\boldsymbol{p}} + \varepsilon_{\boldsymbol{p}_1} - \varepsilon_{\boldsymbol{p}_2} - \varepsilon_{\boldsymbol{p}_3}) \delta(\boldsymbol{p} + \boldsymbol{p}_1 - \boldsymbol{p}_2 - \boldsymbol{p}_3)$$

$$\times \left[f_{\boldsymbol{p}} f_{\boldsymbol{p}_1} - f_{\boldsymbol{p}_2} f_{\boldsymbol{p}_3} \right]$$



Quantum Boltzmann equation (Landau)
Dense gas of electrons



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Neglects quantum interference (entanglement)
between successive collisions

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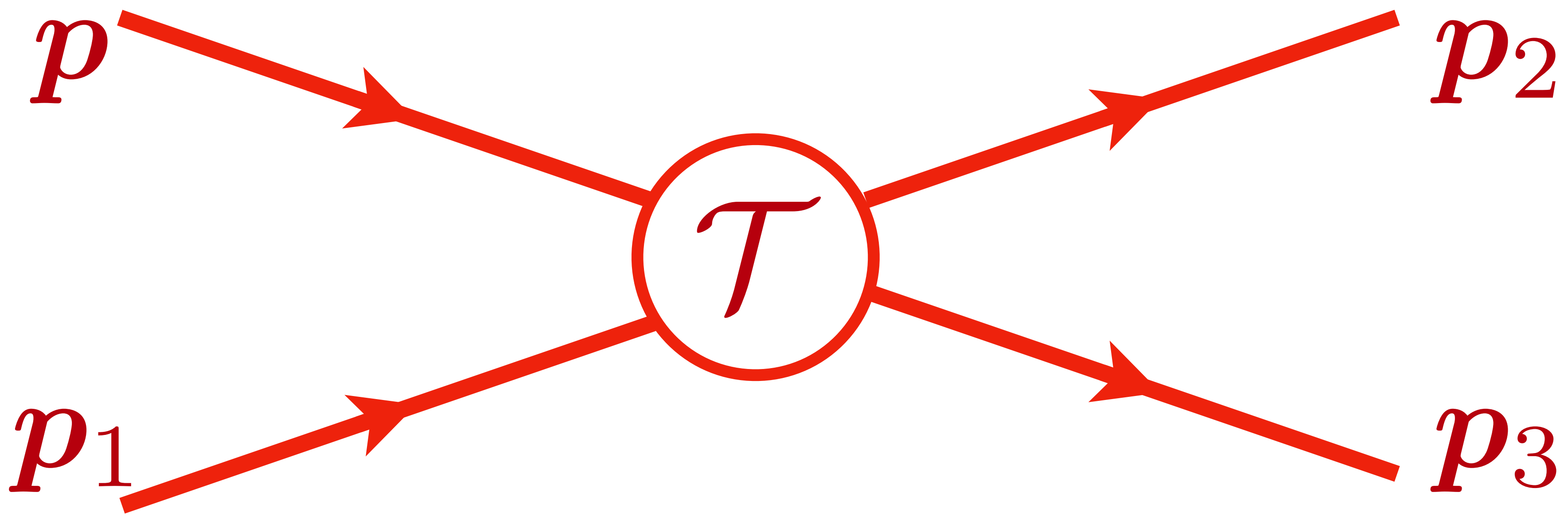
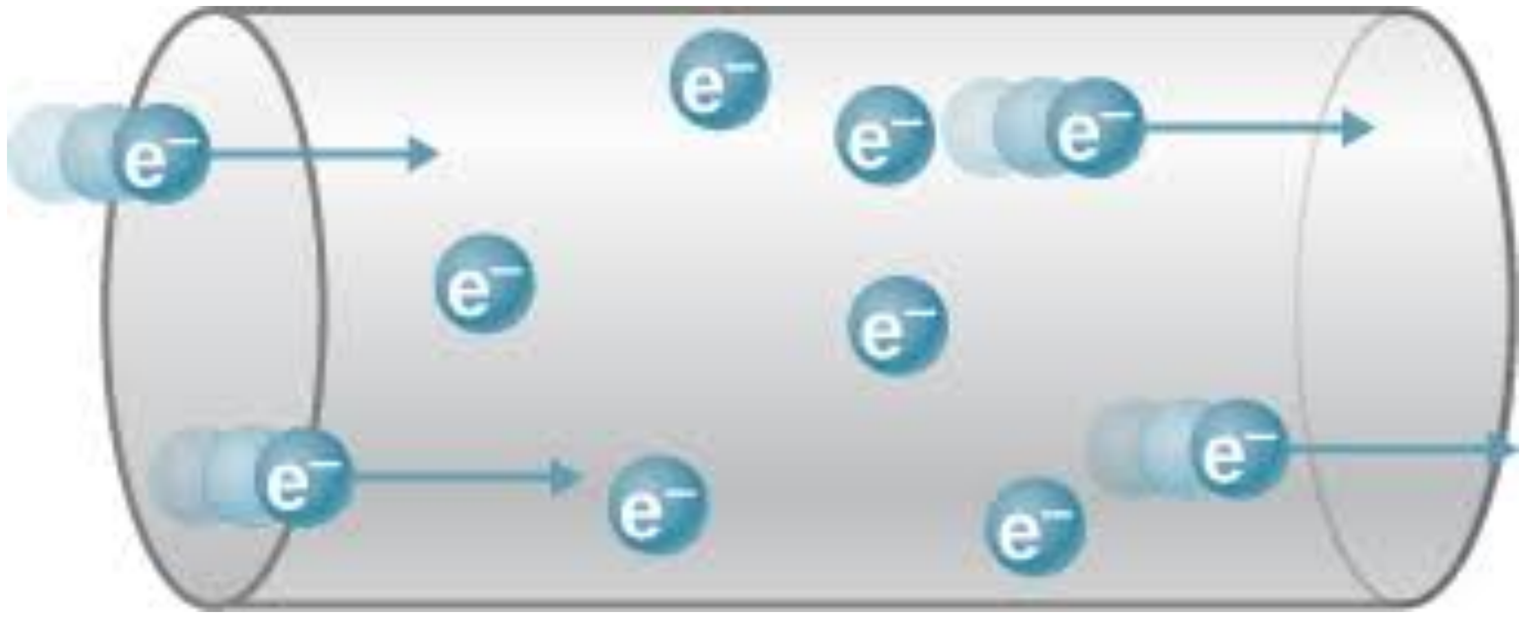
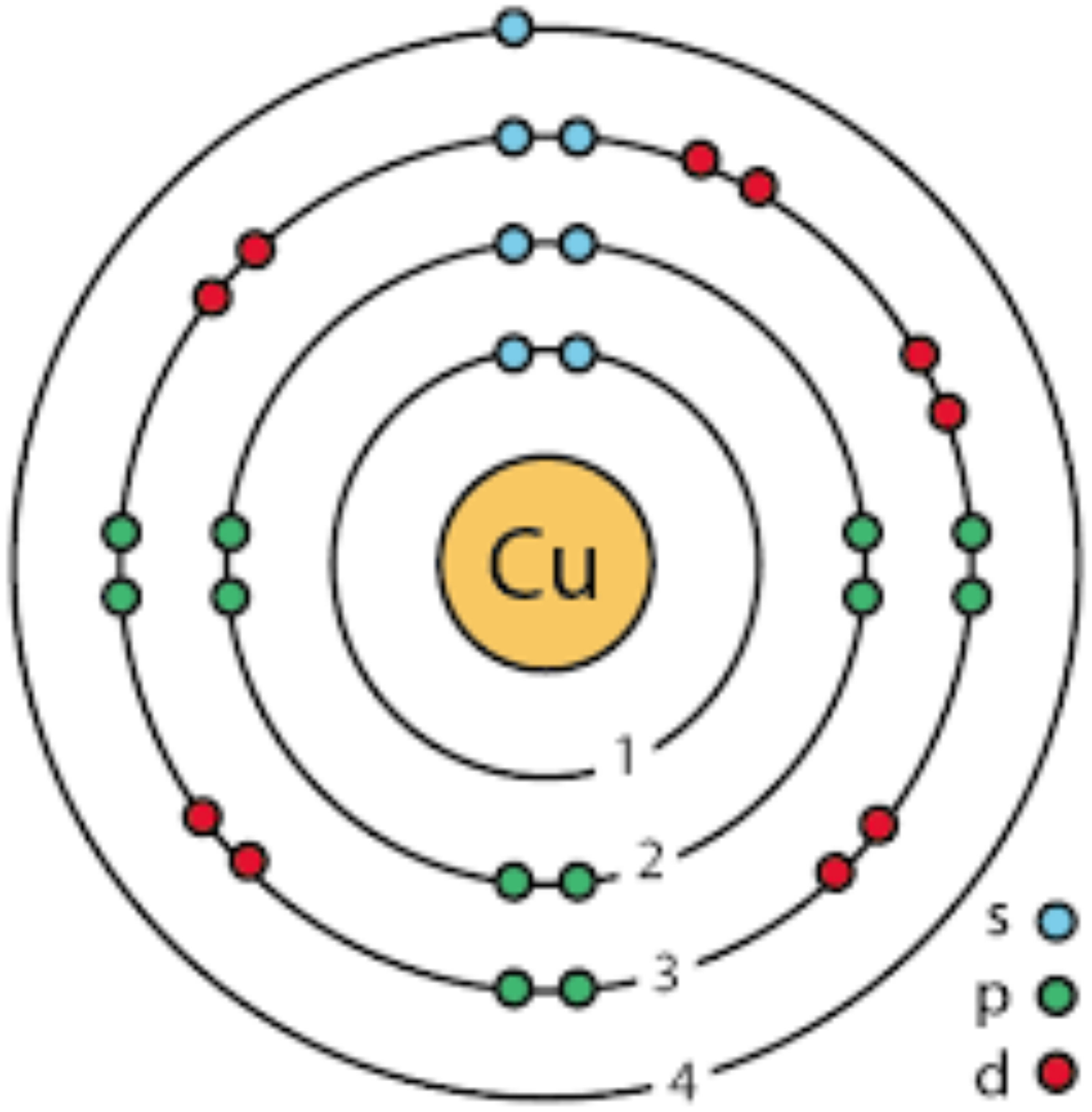
$$\times [f_{\boldsymbol{p}} f_{\boldsymbol{p}_1} (1 - f_{\boldsymbol{p}_2})(1 - f_{\boldsymbol{p}_3}) - f_{\boldsymbol{p}_2} f_{\boldsymbol{p}_3} (1 - f_{\boldsymbol{p}})(1 - f_{\boldsymbol{p}_1})]$$



Quantum theory of electrons: ordinary metals and strange metals

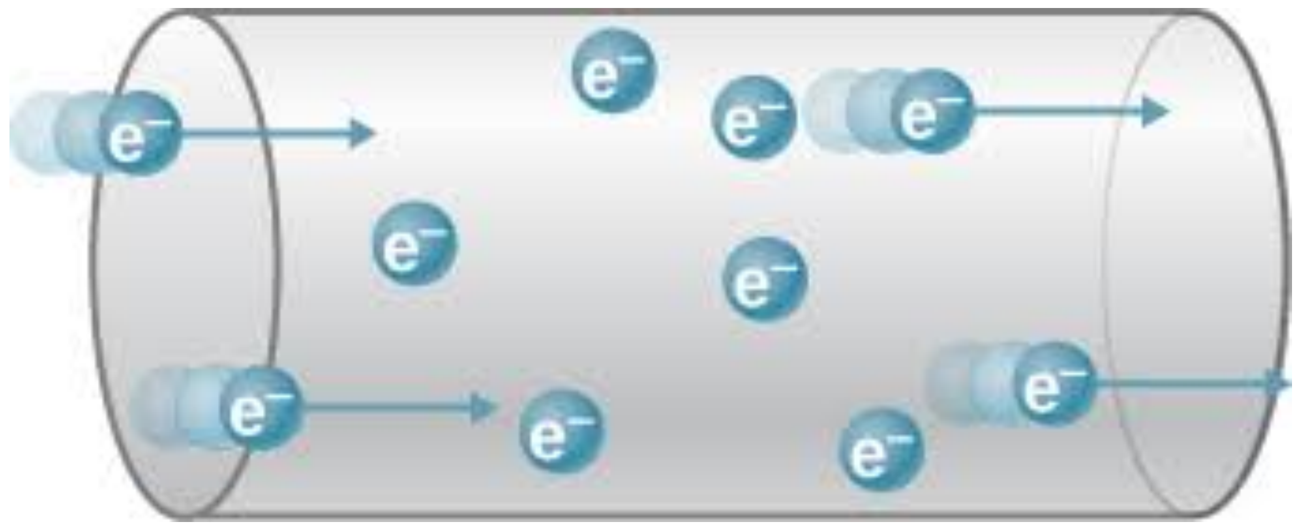
Copper





Each copper atom donates its outermost electron
These electrons move freely throughout the crystal and carry current

Current flow with electrons in Copper

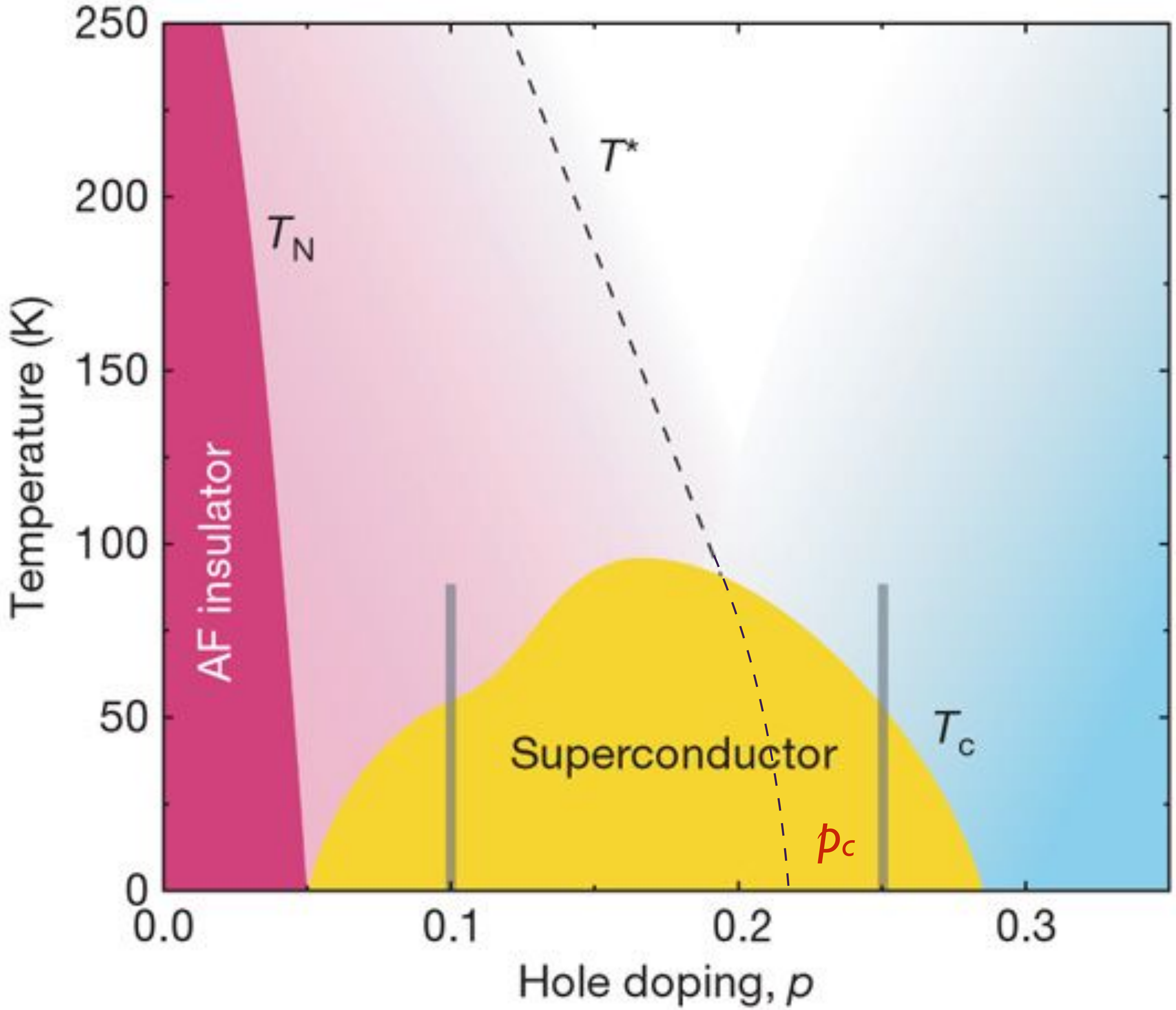


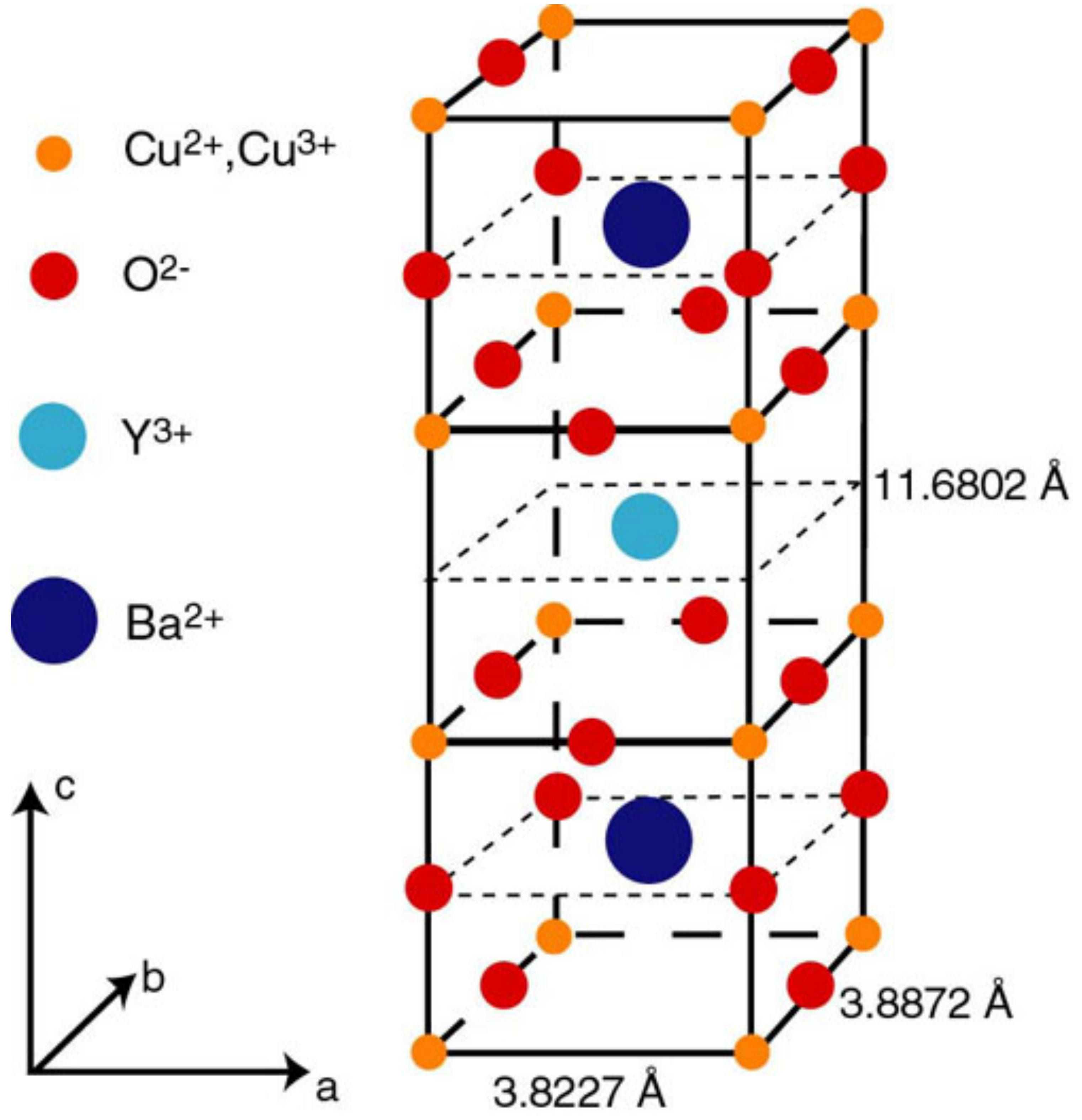
Flow of electrons described by Boltzmann equation $\Rightarrow$
typical scattering time $\tau \sim 1/T^2$, resistivity $\rho(T) = \rho(0) + AT^2$

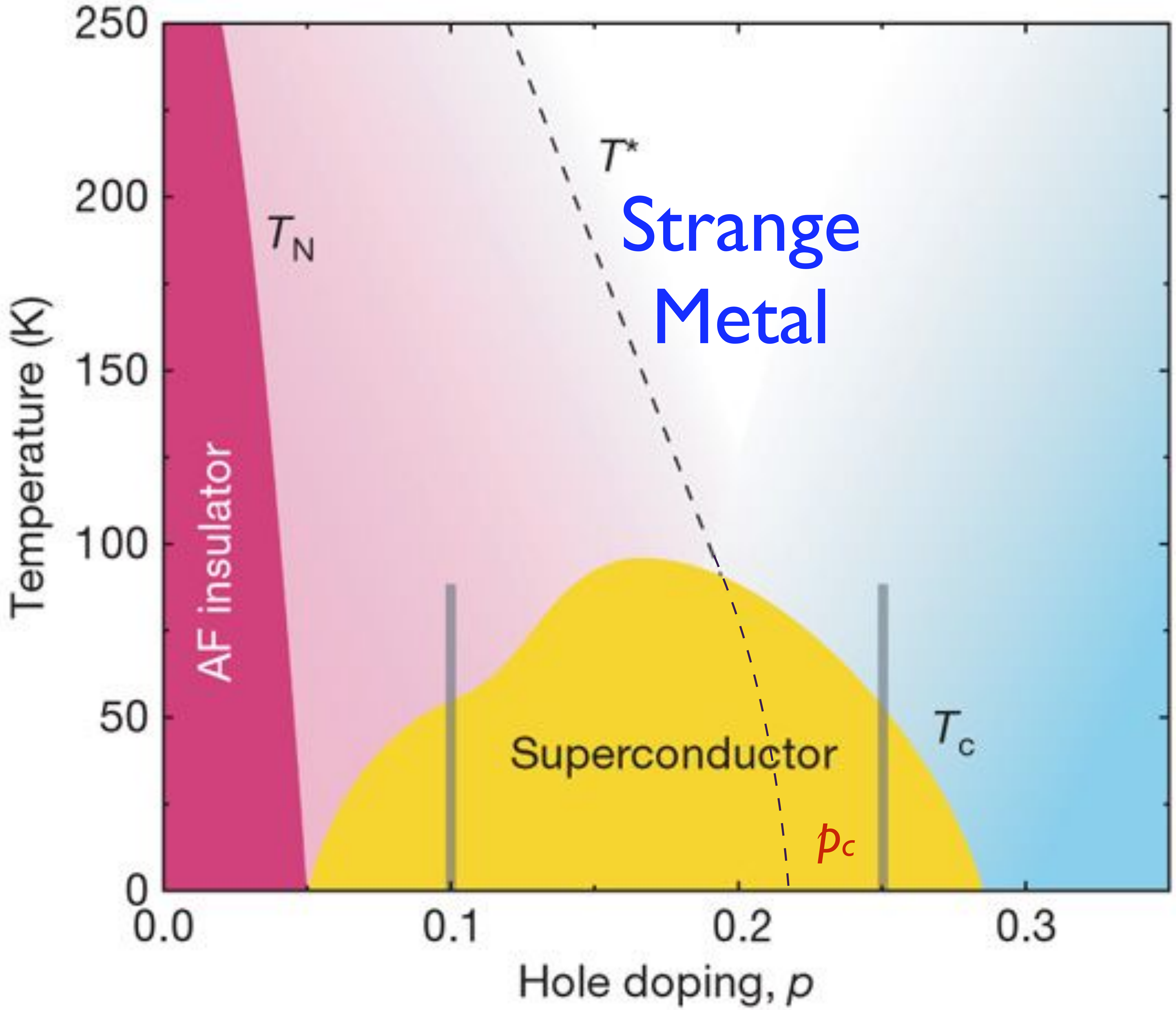
The time τ is much longer than a limiting 'Planckian time' $\dfrac{\hbar}{k_B T}$.

The long scattering time implies that individual electrons are well-defined.

The motion of electrons is 'ballistic' or 'integrable'
up to the long time τ, after which it is chaotic.


250
200
150
100
50
0
0.0
0.1
0.2
0.3
Temperature (K)
Hole doping, p
AF insulator
T_N
T*
Superconductor
T_c
p_c


Cu2+,Cu3+
O2-
Y3+
Ba2+
c
b
a
11.6802 Å
3.8872 Å
3.8227 Å


250
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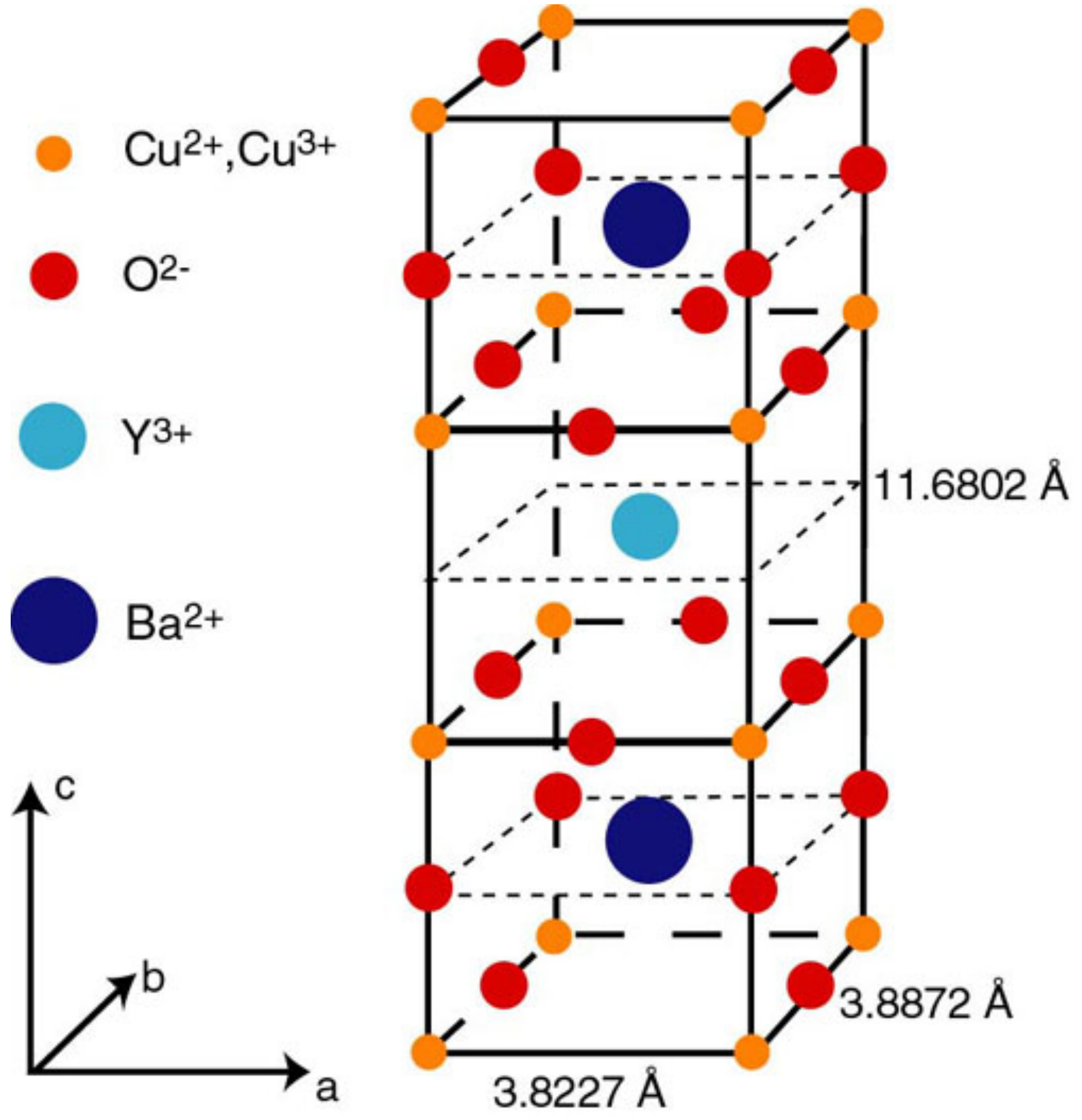
Temperature (K)
Hole doping, p
AF insulator
T_N
T*
Strange
Metal
Superconductor
T_c
p_c


Cu2+,Cu3+
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c
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11.6802 Å
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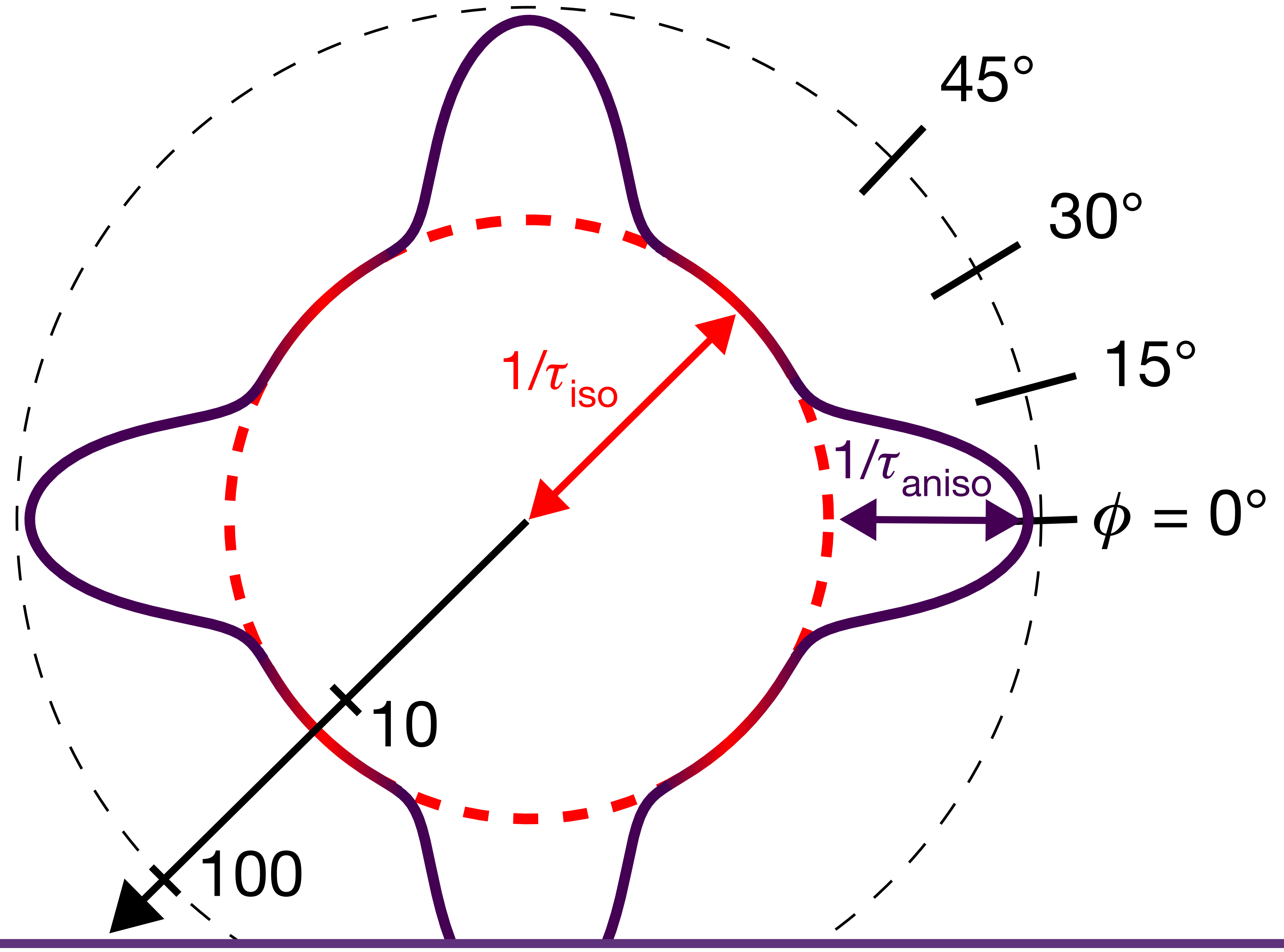
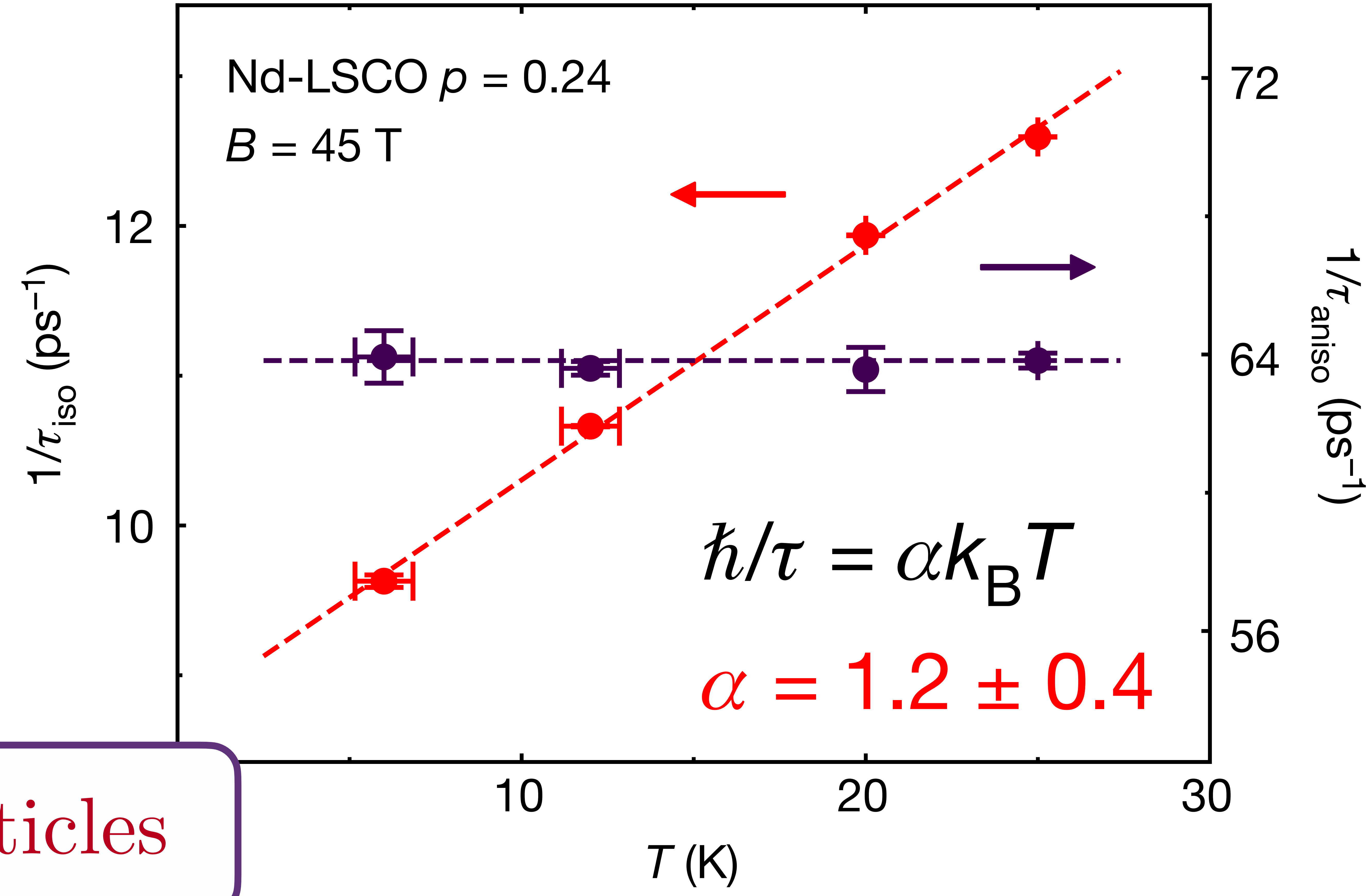
Linear-in temperature resistivity from an isotropic Planckian scattering rate

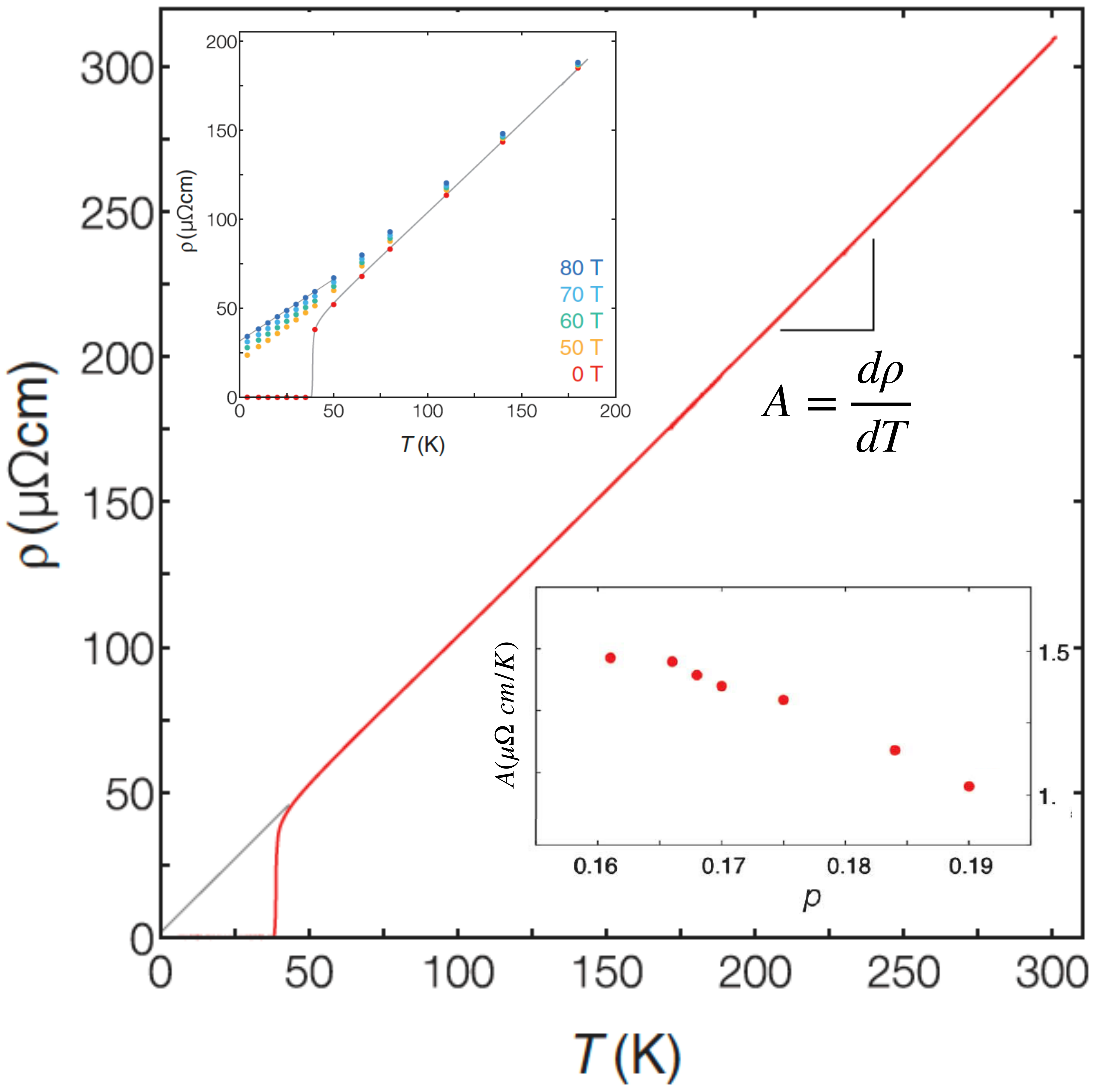
Nature **595**, 667-672 (2021)

G. Grissonnanche, Y. Fang, A. Legros, S. Verret, F. Laliberté, C. Collignon, J. Zhou, D. Graf, P. Goddard, L. Taillefer, B. J. Ramshaw

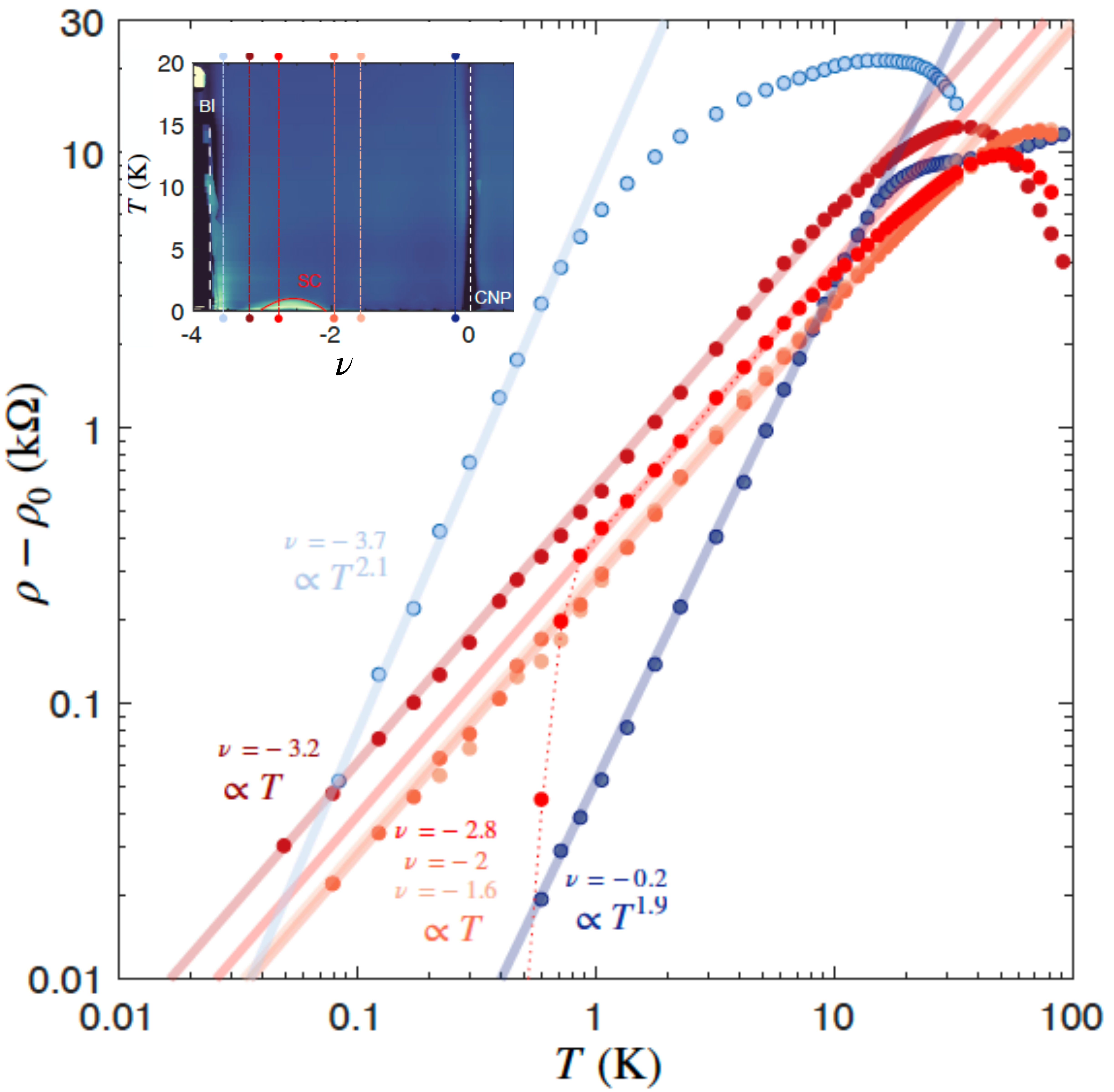




Current flow without quasiparticles



LSCO: Giraldo-Gallo et al. 2018



MATBG: Jaoui et al. 2021

Questions

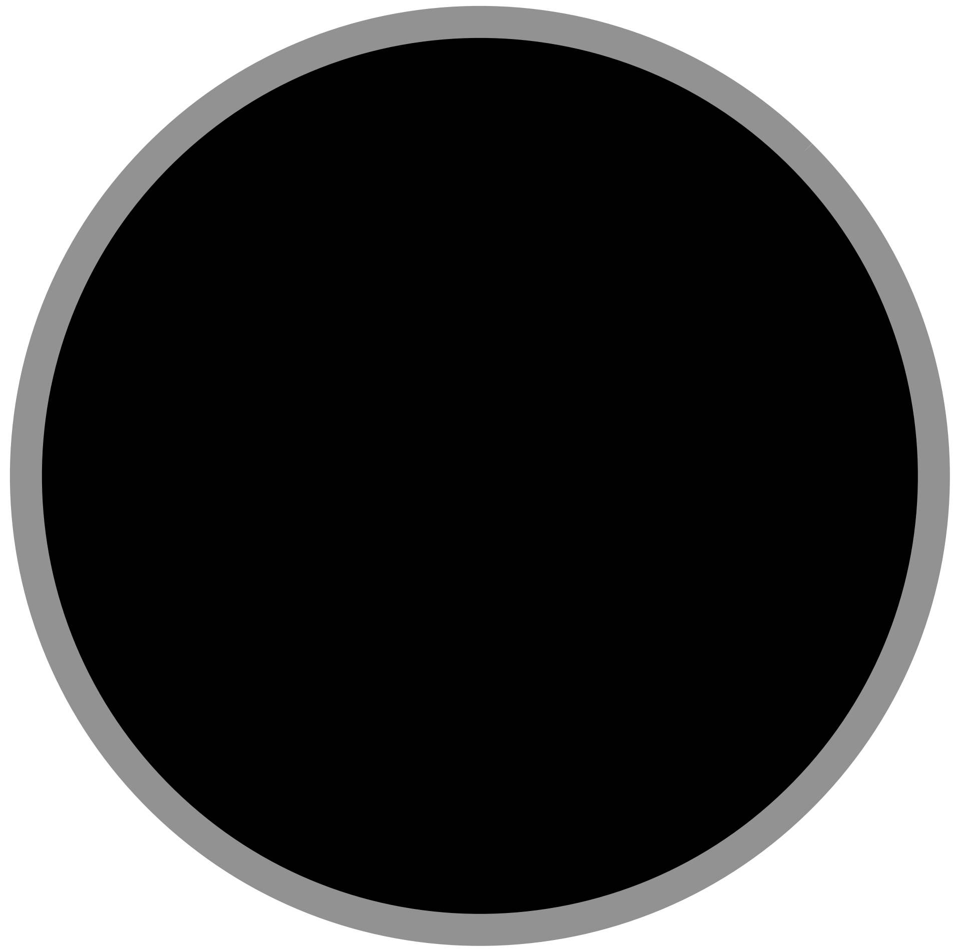
- Needed: A theory for current flow in a ‘strange metal’ with an entangled soup of electrons.
- Needed: theory for collision time in resistivity $\sim \hbar/(k_B T)$.
- Needed: theory for the appearance of superconductivity in such a ‘strange metal’.

Quantum black holes and holography

Black Holes

Objects so dense that light is gravitationally bound to them.

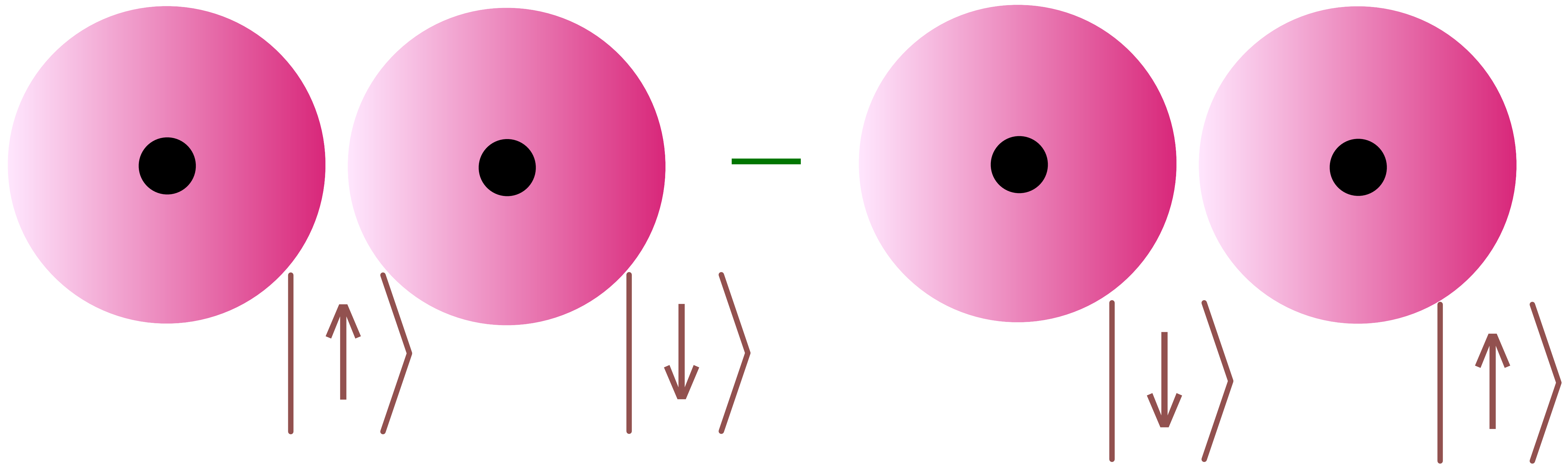
Horizon radius $R = \frac{2GM}{c^2}$



G Newton's constant, c velocity of light, M mass of black hole

For $M =$ earth's mass, $R \approx 9\, mm$!

Quantum Entanglement across a black hole horizon



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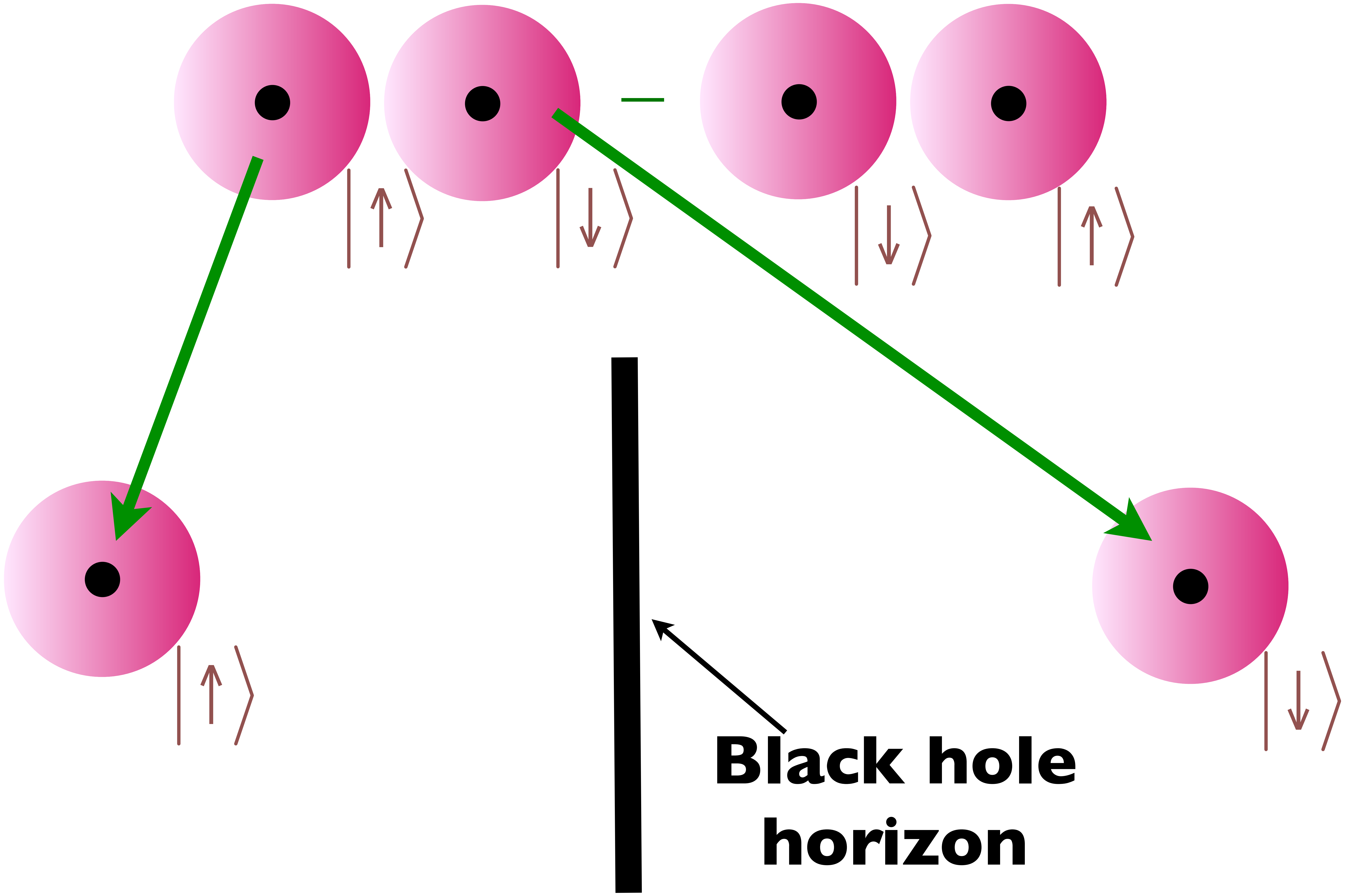
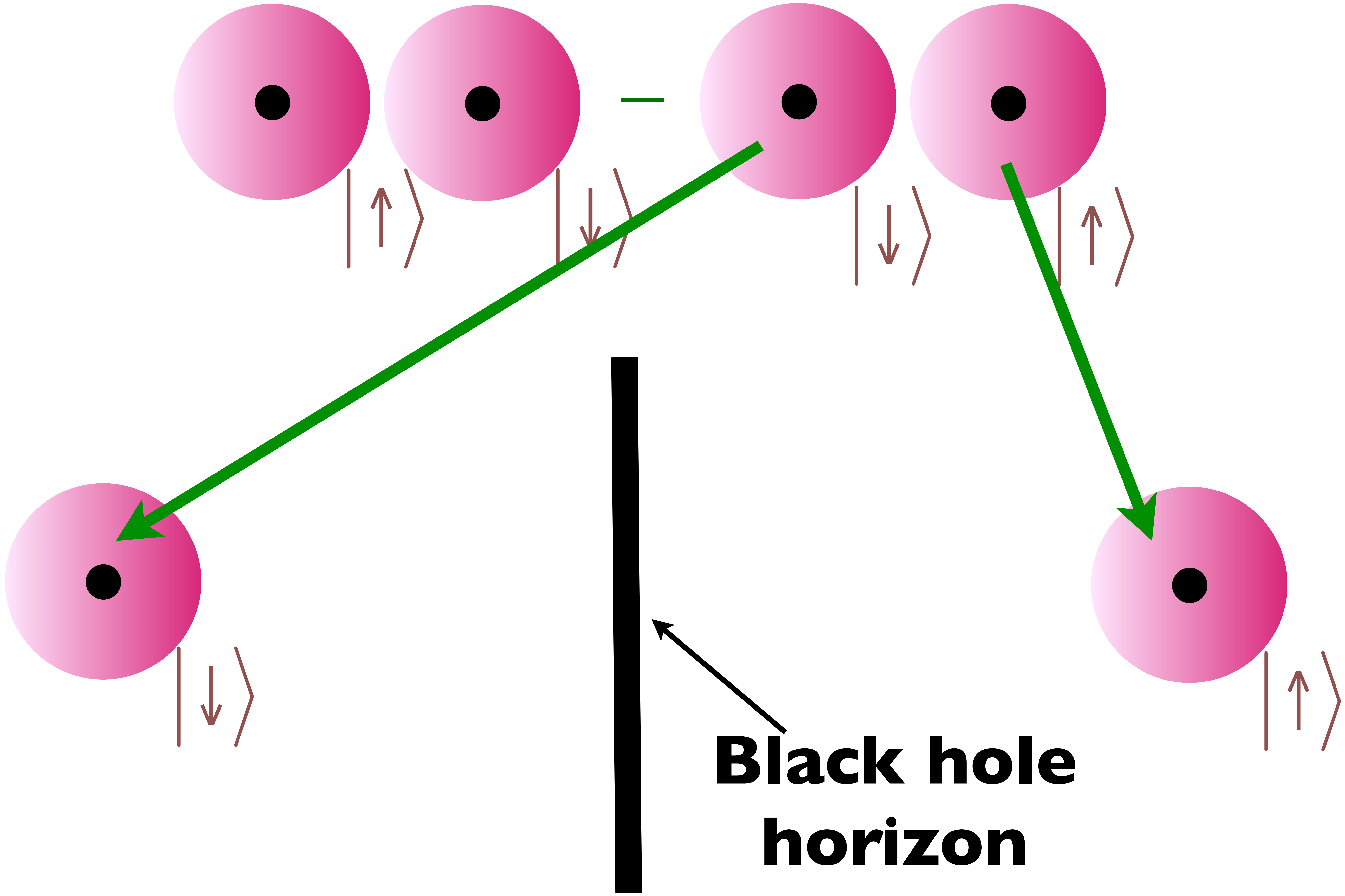


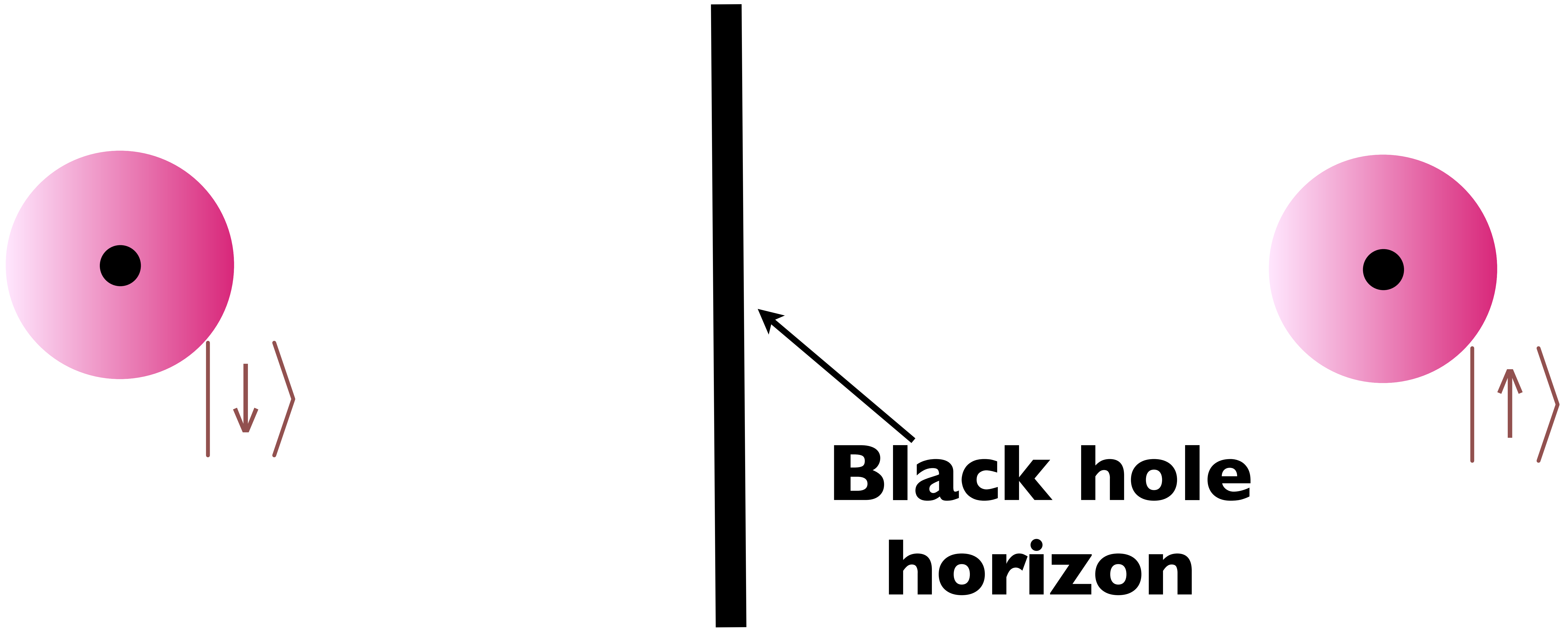
Quantum Entanglement across a black hole horizon



Quantum Entanglement across a black hole horizon

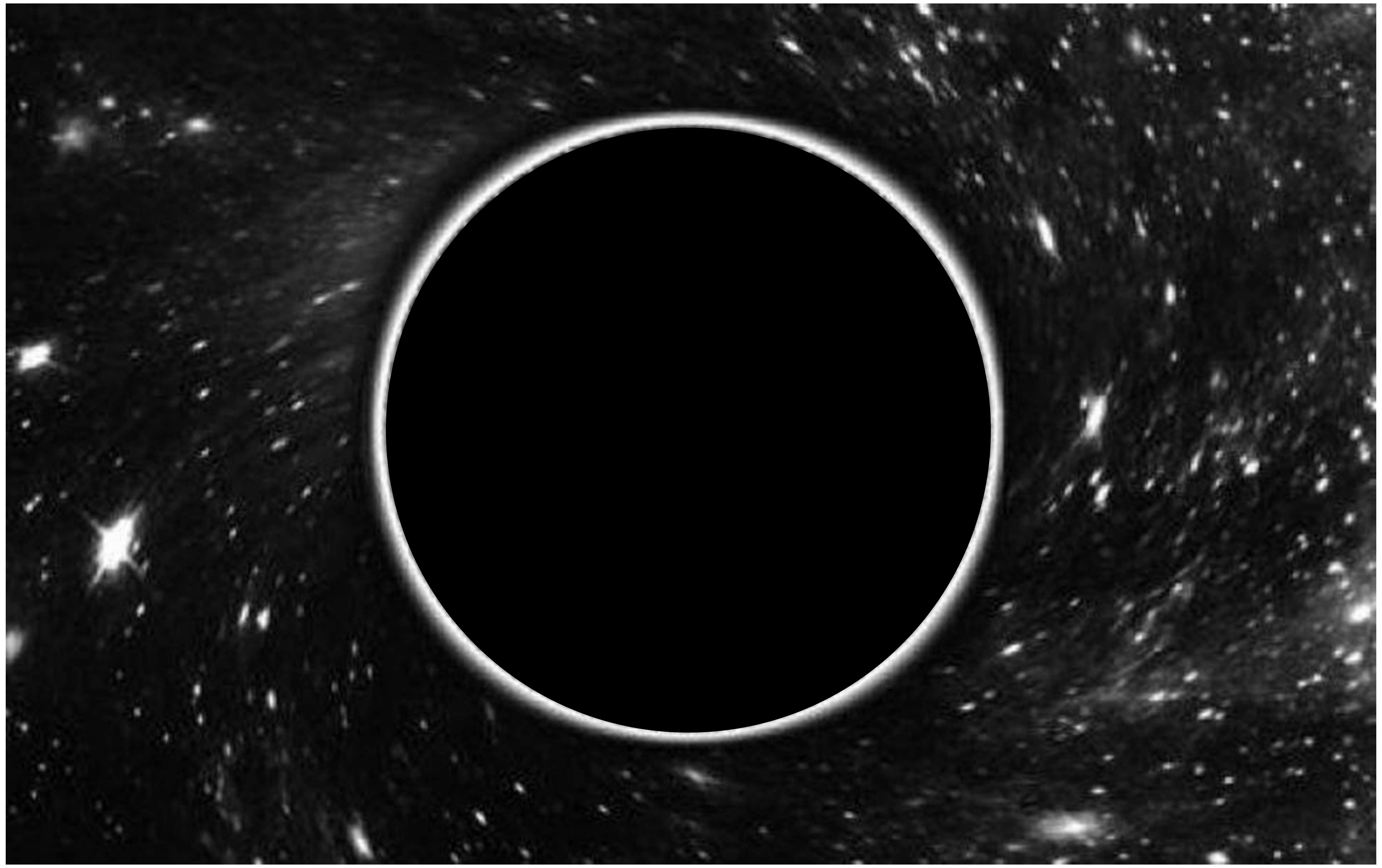
Hawking (1975) used other arguments to show that black hole horizons have a temperature

(The entanglement reasoning: to an outside observer, the state of the electron inside the black hole cannot be known, and so the outside electron is in a random state.)



Quantum Black holes

- Black holes have an entropy and a temperature, $T_H = \hbar c^3/(8\pi G M k_B)$.
- The entropy is proportional to their surface area. $S = A k_B c^3/(4G\hbar)$.



T.B. BAKKER / DR. J.P.VAN DER SCHAAR

J. D. Bekenstein, PRD **7**, 2333 (1973)
S.W. Hawking, Nature **248**, 30 (1974)

Quantum Black holes

Bohr-Sommerfeld semiclassical quantum theory of a black hole in d spatial dimensions

$$\mathcal{Z} = \int \mathcal{D}g_{\mu\nu}\mathcal{D}a_\mu \exp\left(-\frac{1}{\hbar}\int d^d x \int_0^{\hbar/(k_B T)} d\tau \sqrt{g}\,\mathcal{L}_d[g_{\mu\nu}, a_\mu]\right)$$

$g_{\mu\nu} \Rightarrow$ spacetime metric, $\quad g = \det(g_{\mu\nu})$

$a_\mu \Rightarrow$ Electromagnetic gauge field

$\mathcal{L}_d \Rightarrow$ *Classical* Einstein-Maxwell action

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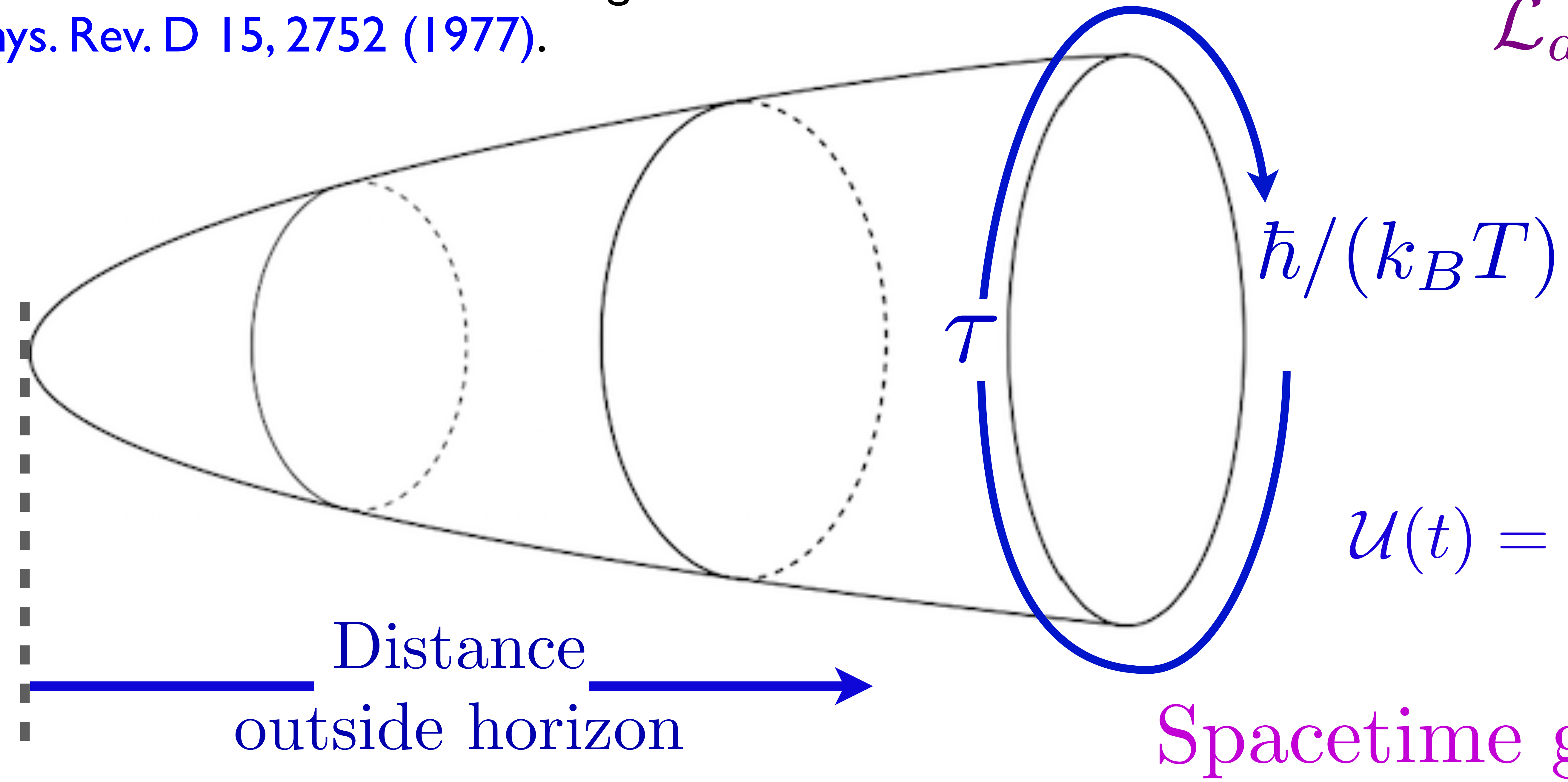
G.W. Gibbons and S.W. Hawking, Phys. Rev. D 15, 2752 (1977).



Evaluate path integral at black hole saddle point

$$\mathcal{U}(t) = \exp\left(-i\mathcal{H}t/\hbar\right) \quad \Leftrightarrow \quad \mathcal{Z} = \mathrm{Tr}\exp\left(-\mathcal{H}/(k_BT)\right)$$

Spacetime geometry of a black hole in imaginary time τ

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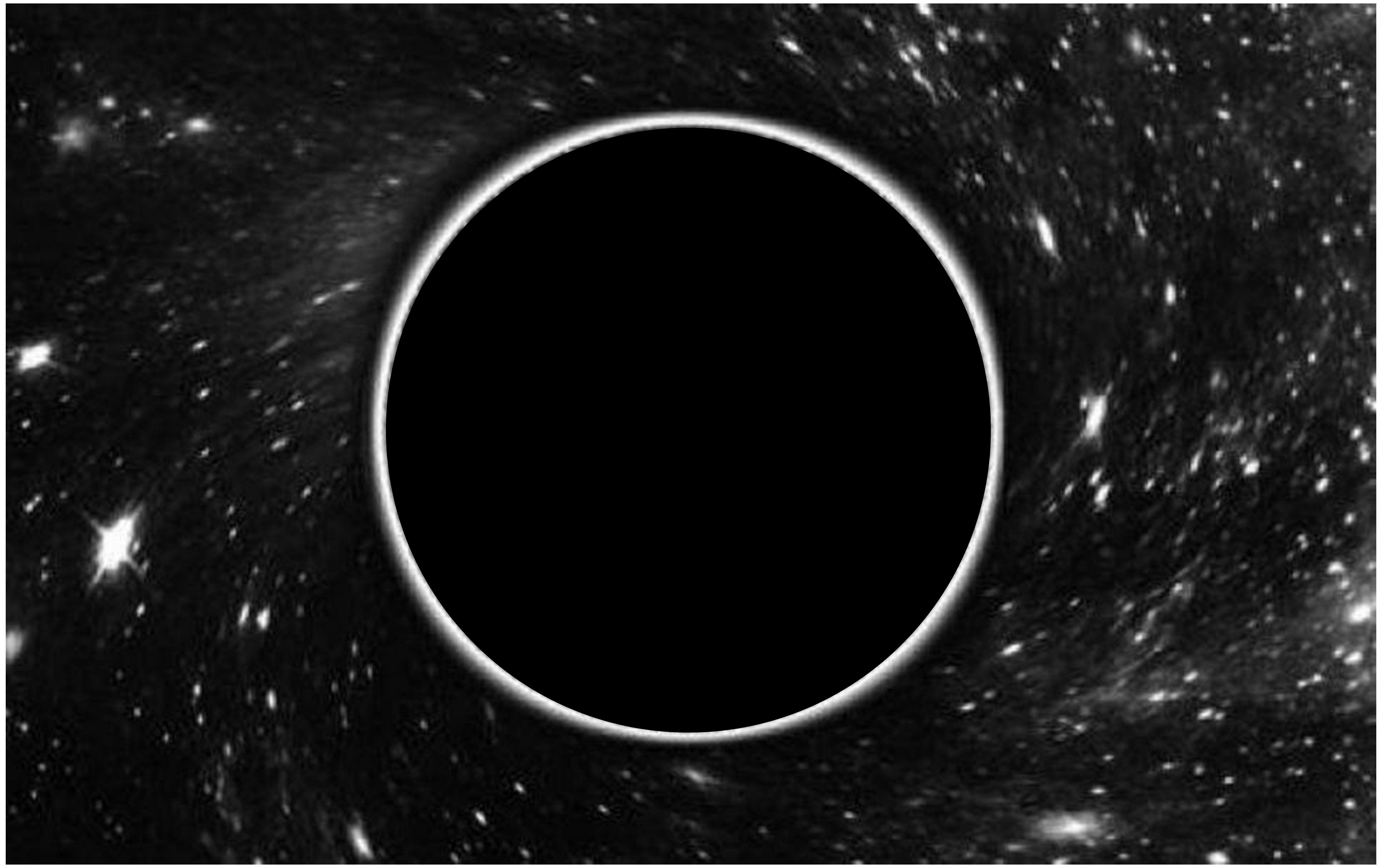

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<u>Remarkable features:</u>

- Entropy is finite.
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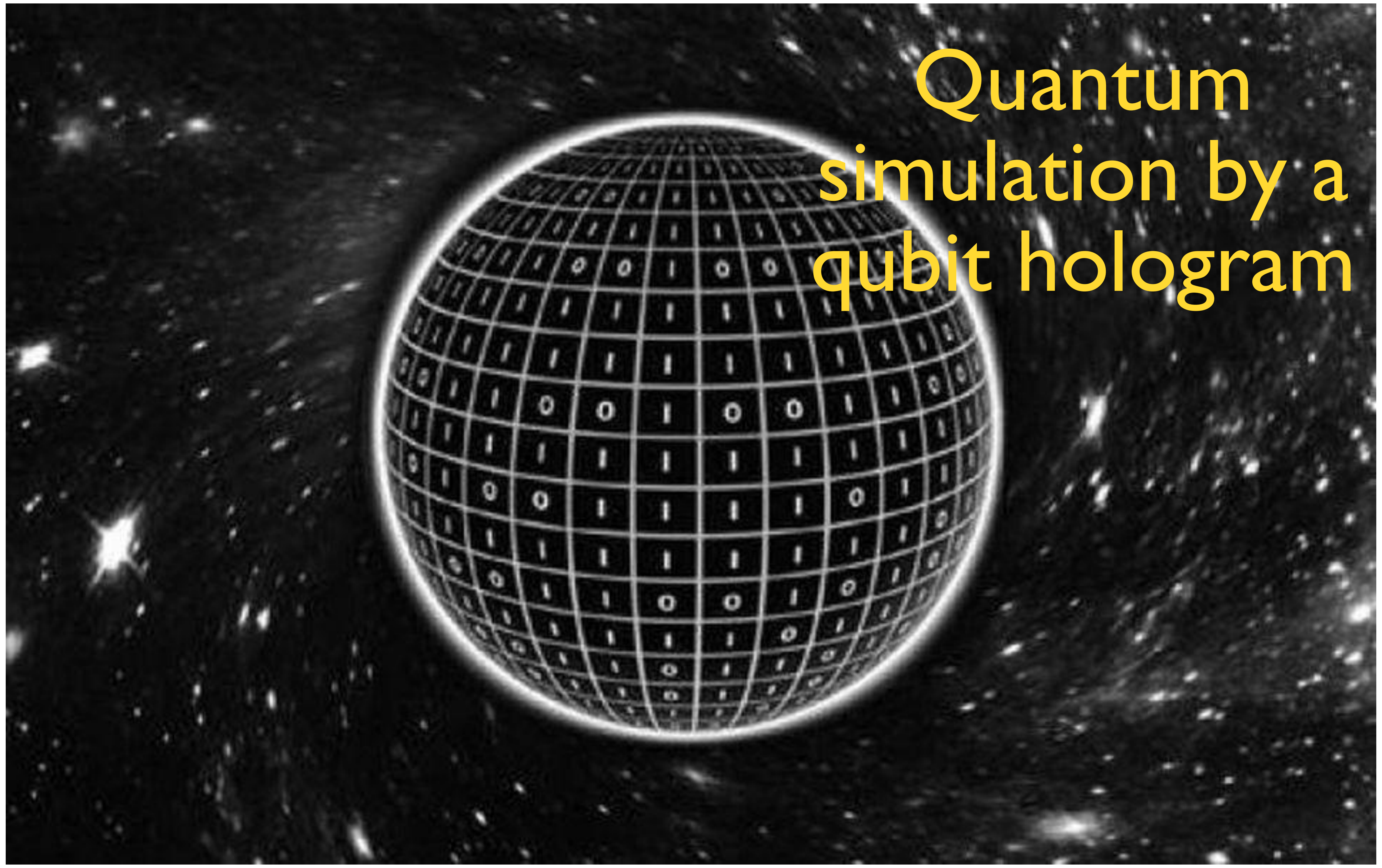

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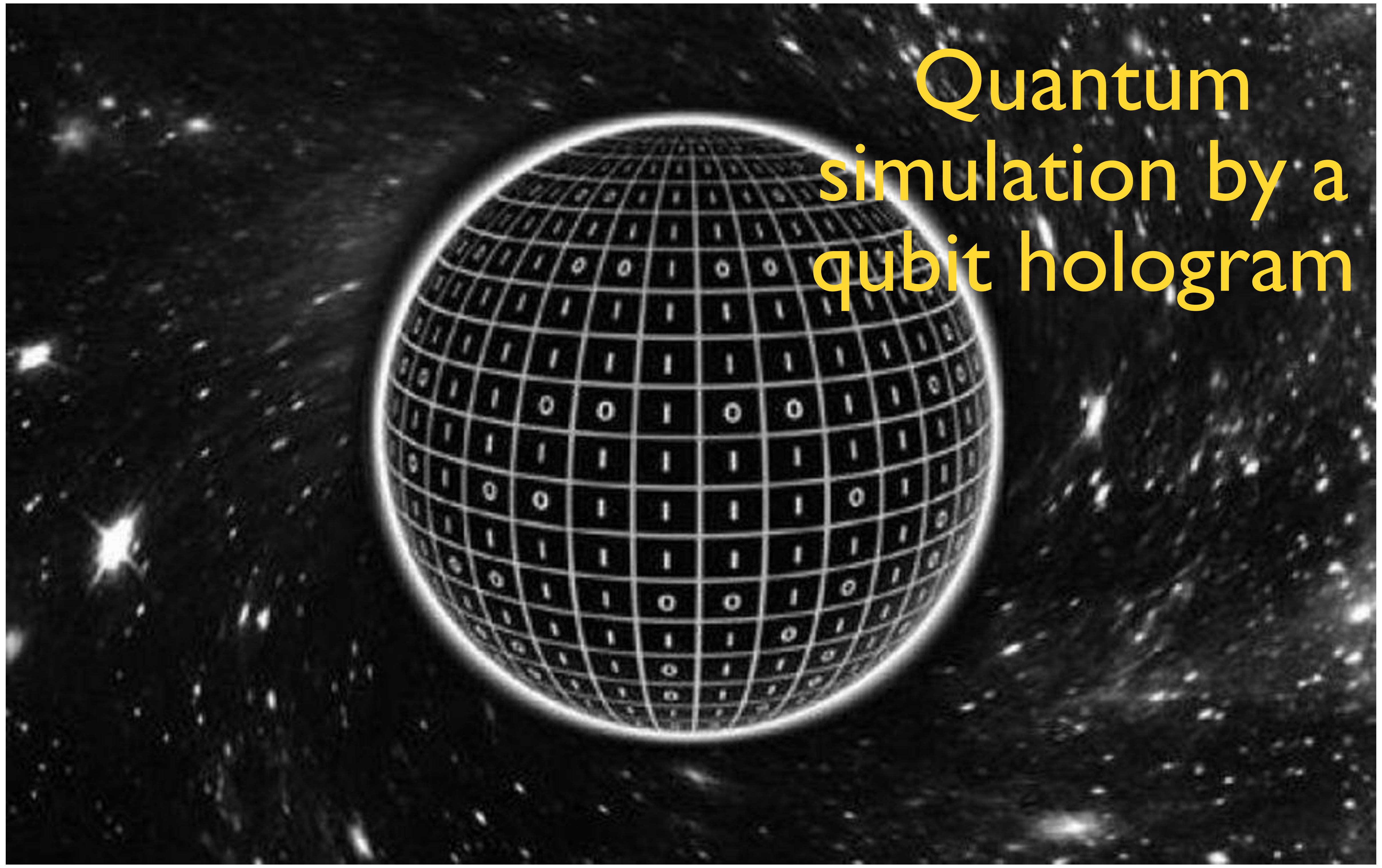
- They relax to thermal equilibrium in a time $\sim 8\pi G M/c^3$



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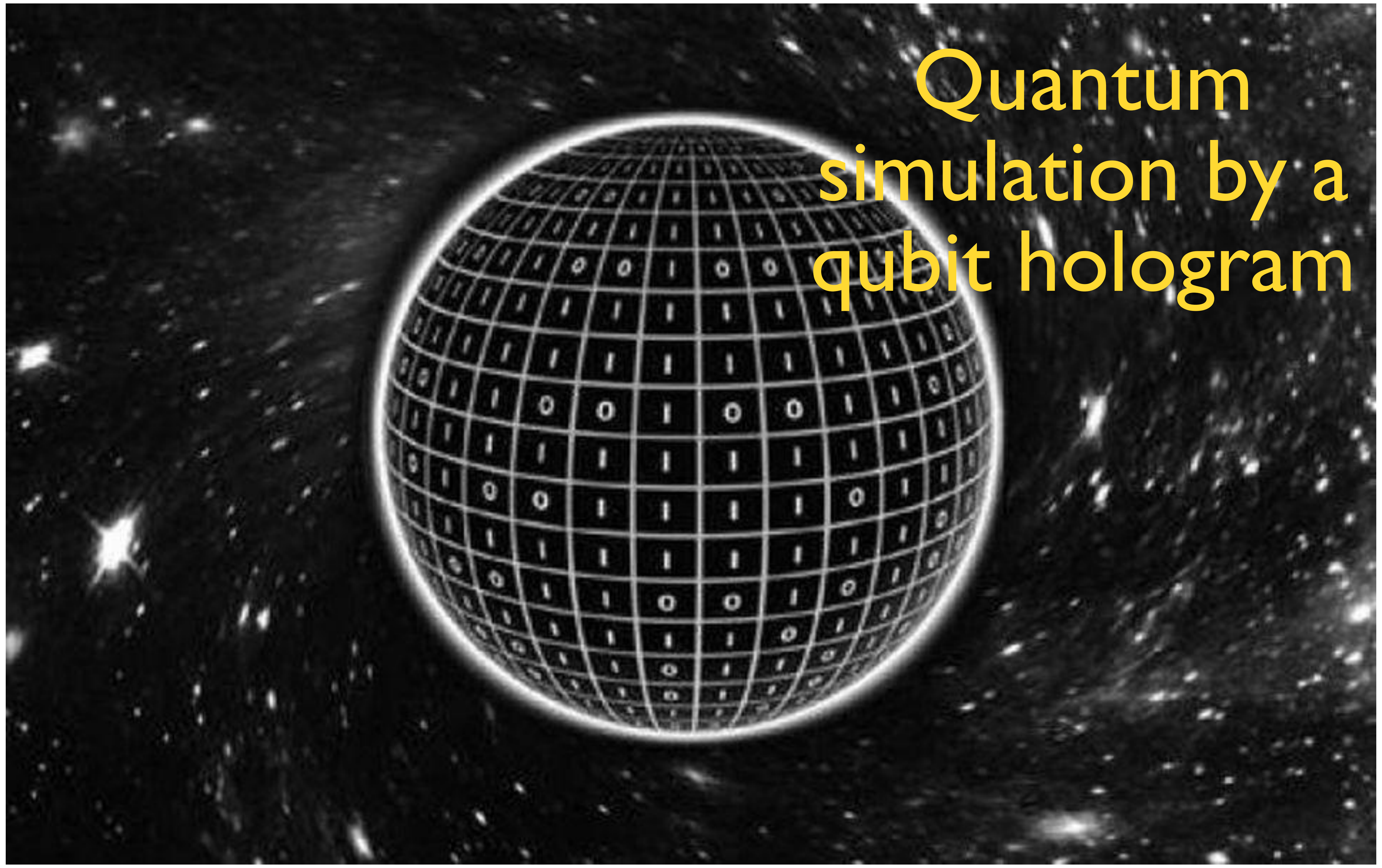
- The entropy is proportional to their surface area. $S = A k_B c^3/(4G\hbar)$.
- They relax to thermal equilibrium in a time $\sim 8\pi GM/c^3 = \hbar/(k_B T_H)$ which is Planckian!



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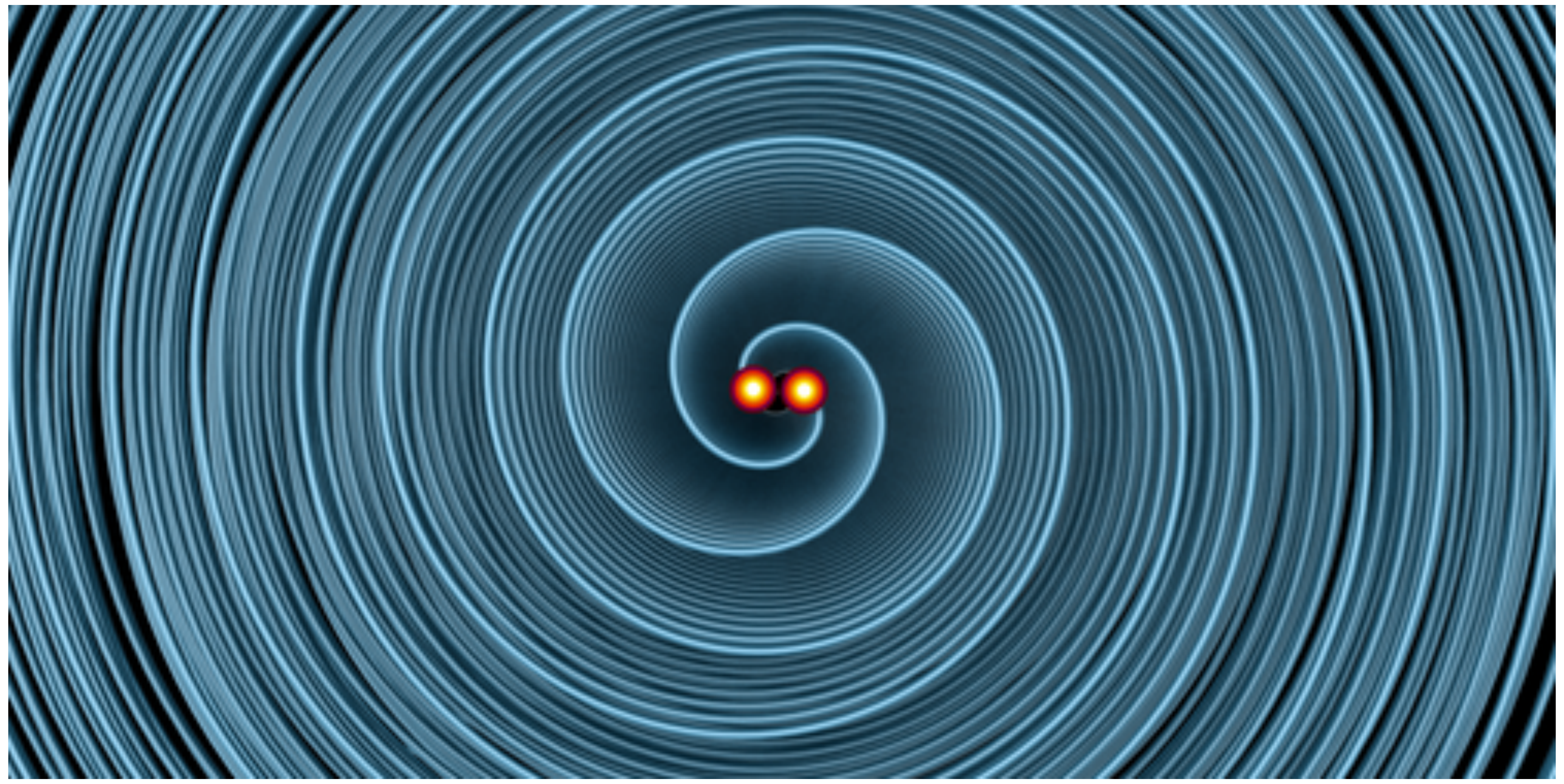
Black Holes Obey Information-Emission Limits

April 22, 2021 • *Physics* 14, s47 –Christopher Crockett

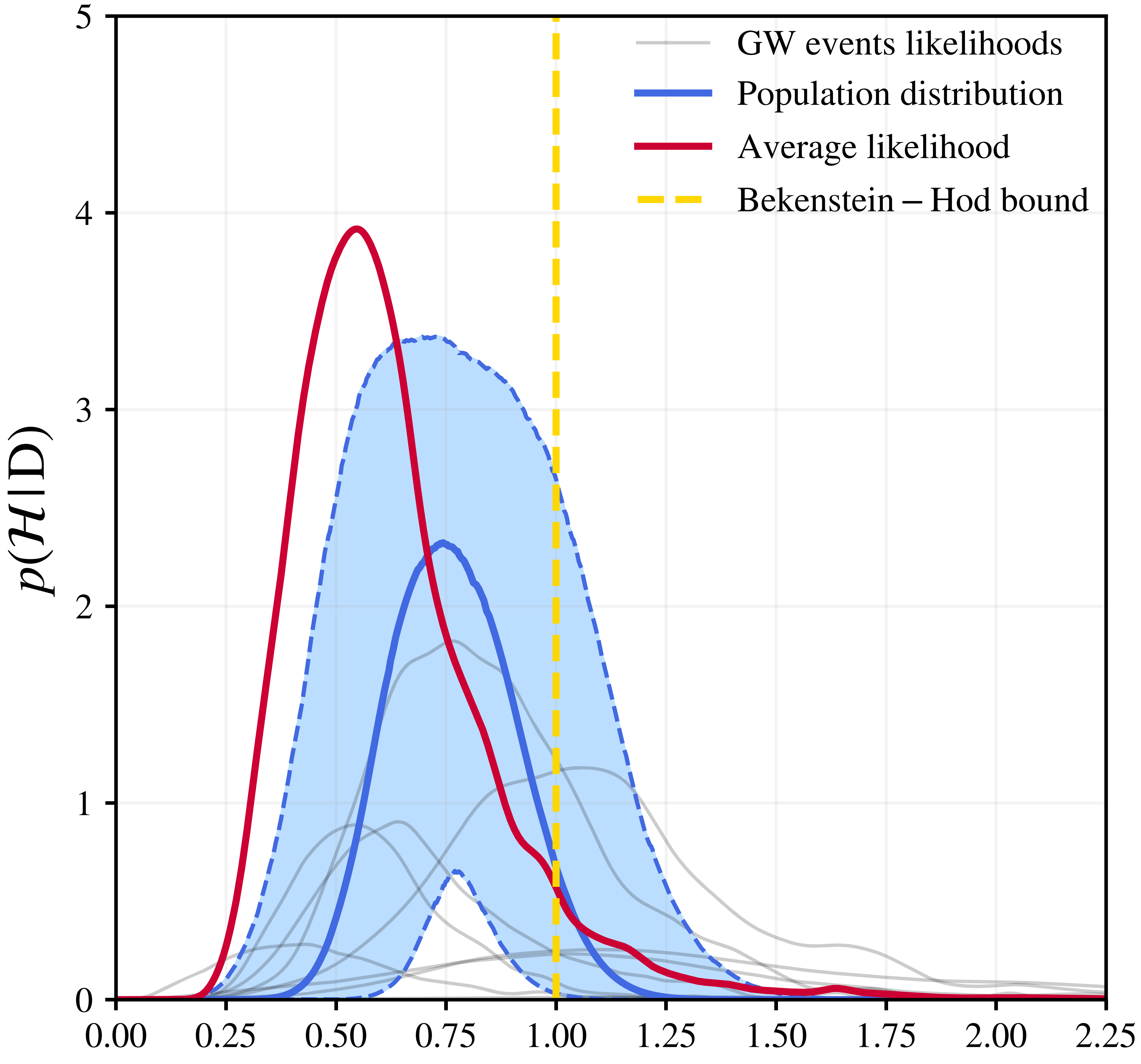
G. Carullo, D. Laghi, J. Veitch, W. Del Pozzo, Phys. Rev. Lett. **126**, 161102 (2021)

An analysis of the gravitational waves emitted from black hole mergers confirms that black holes are the fastest known information dissipaters.



Gravity wave observations of 8 different black holes show a relaxation time

$$\tau \sim \frac{\hbar}{k_B T}$$



$$\mathcal{H} = \frac{1}{\pi}\frac{\hbar/\tau}{k_B T}$$

Questions

- Can Hawking's Bohr-Sommerfield theory of black hole entropy be interpreted in terms of a $D(E)$ of a Schrödinger-Heisenberg quantum system?

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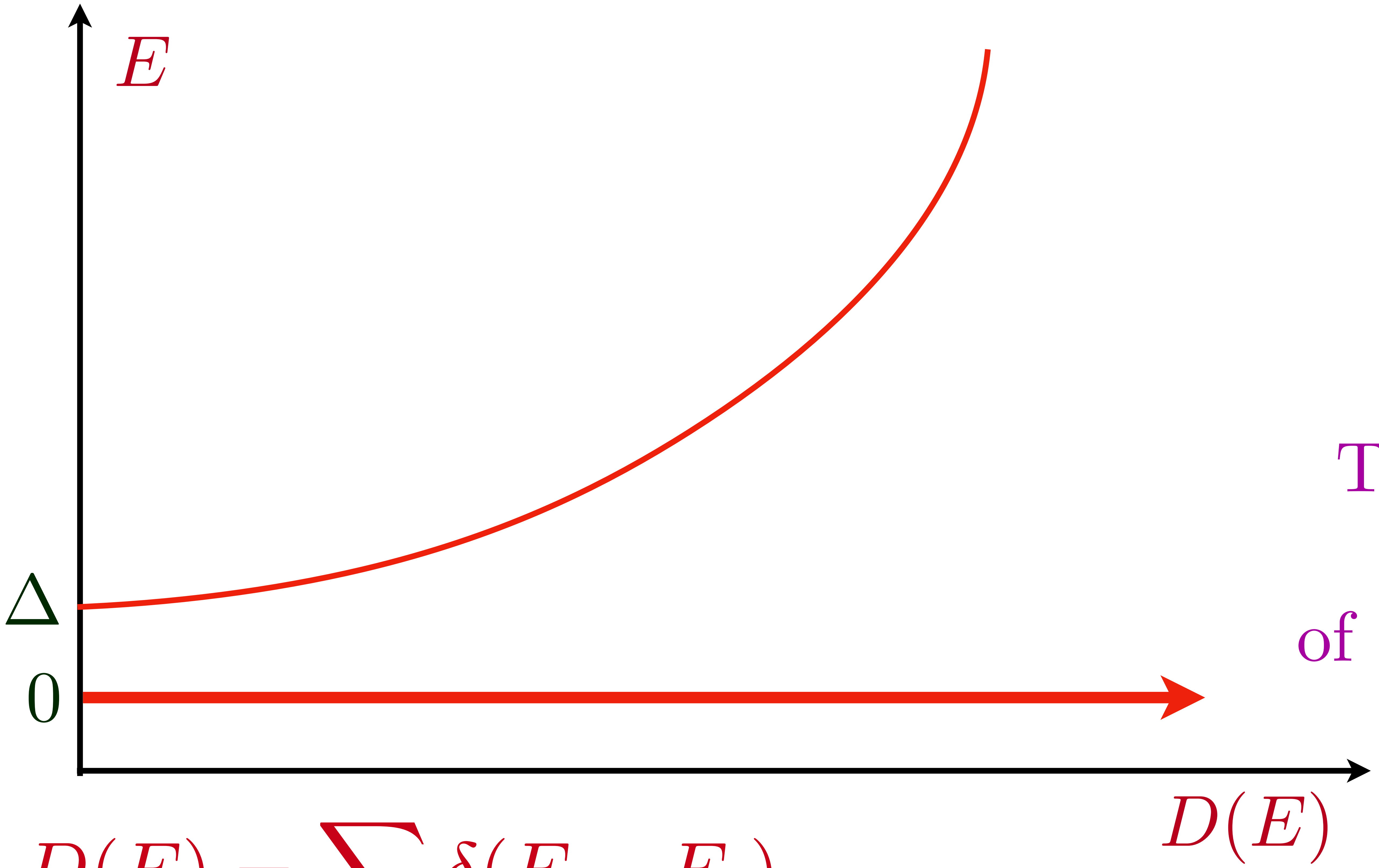


$$D(E) = \sum_i \delta(E - E_i)$$
$$= \exp\left(S/k_B\right)\delta(E) + \ldots$$

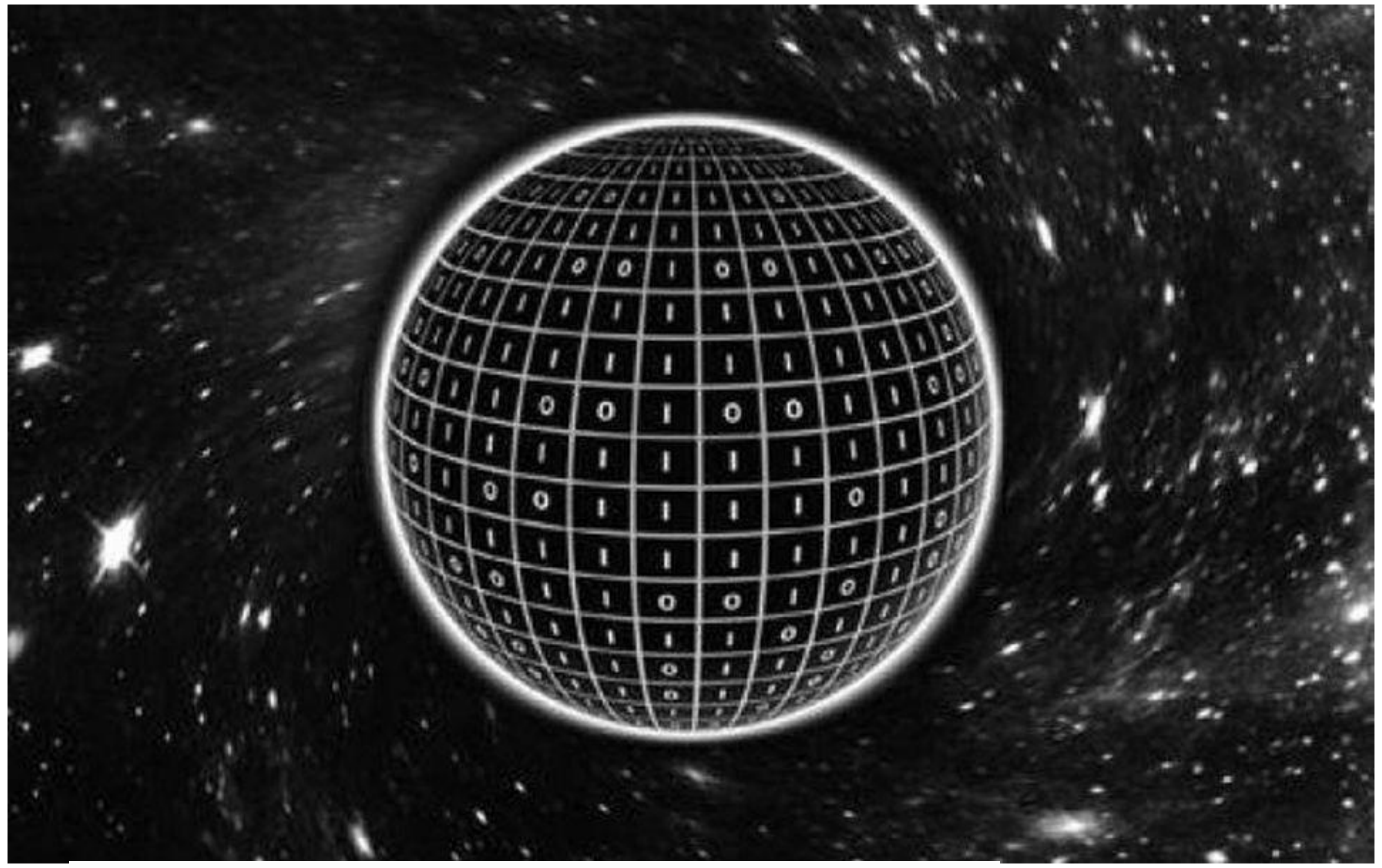
Black holes in string theory realize the entropy as an exact, exponentially large, degeneracy of the ground state.
This is not the generic behavior of the semiclassical path integral of Einstein gravity, as we will see …

Questions

- Can Hawking's Bohr-Sommerfield theory of black hole entropy be interpreted in terms of a $D(E)$ of a Schrödinger-Heisenberg quantum system?

- Can black hole entropy be understood 'holographically', as that of a unitary quantum system in one lower spatial dimension with a finite number of states?

G. 'tHooft (1993); L. Susskind (1993); Maldacena (1998)



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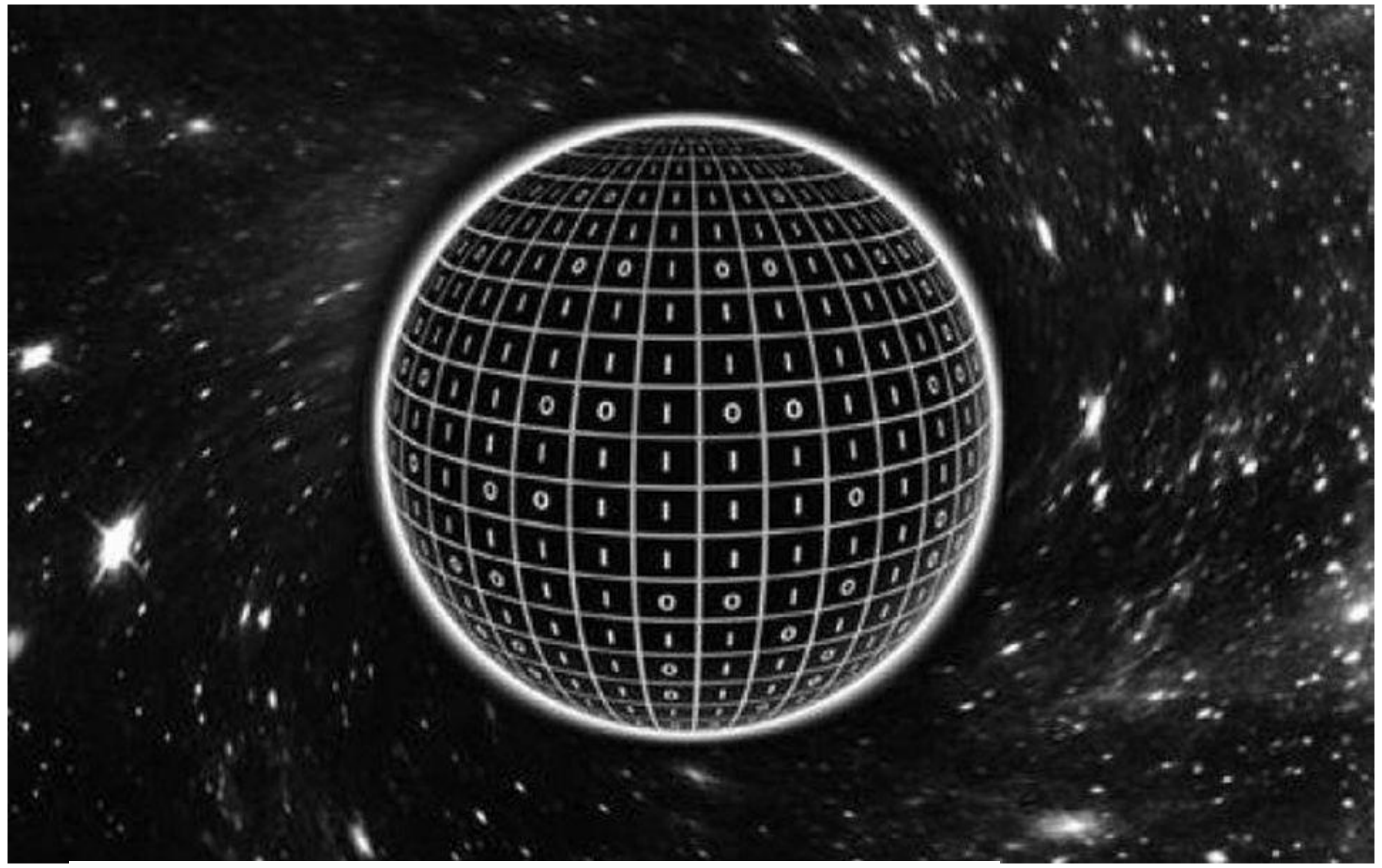
- The unitary quantum system cannot have particle-like excitations if it is to reproduce the rapid Planckian dynamics at the rate $k_B T/\hbar$.



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- Can black hole entropy be understood 'holographically', as that of a unitary quantum system in one lower spatial dimension with a finite number of states?

- The unitary quantum system cannot have particle-like excitations if it is to reproduce the rapid Planckian dynamics at the rate $k_B T/\hbar$.

- Can we compute the evolution of the entropy as the black hole evaporates? Is it that of an evaporating unitary quantum system?

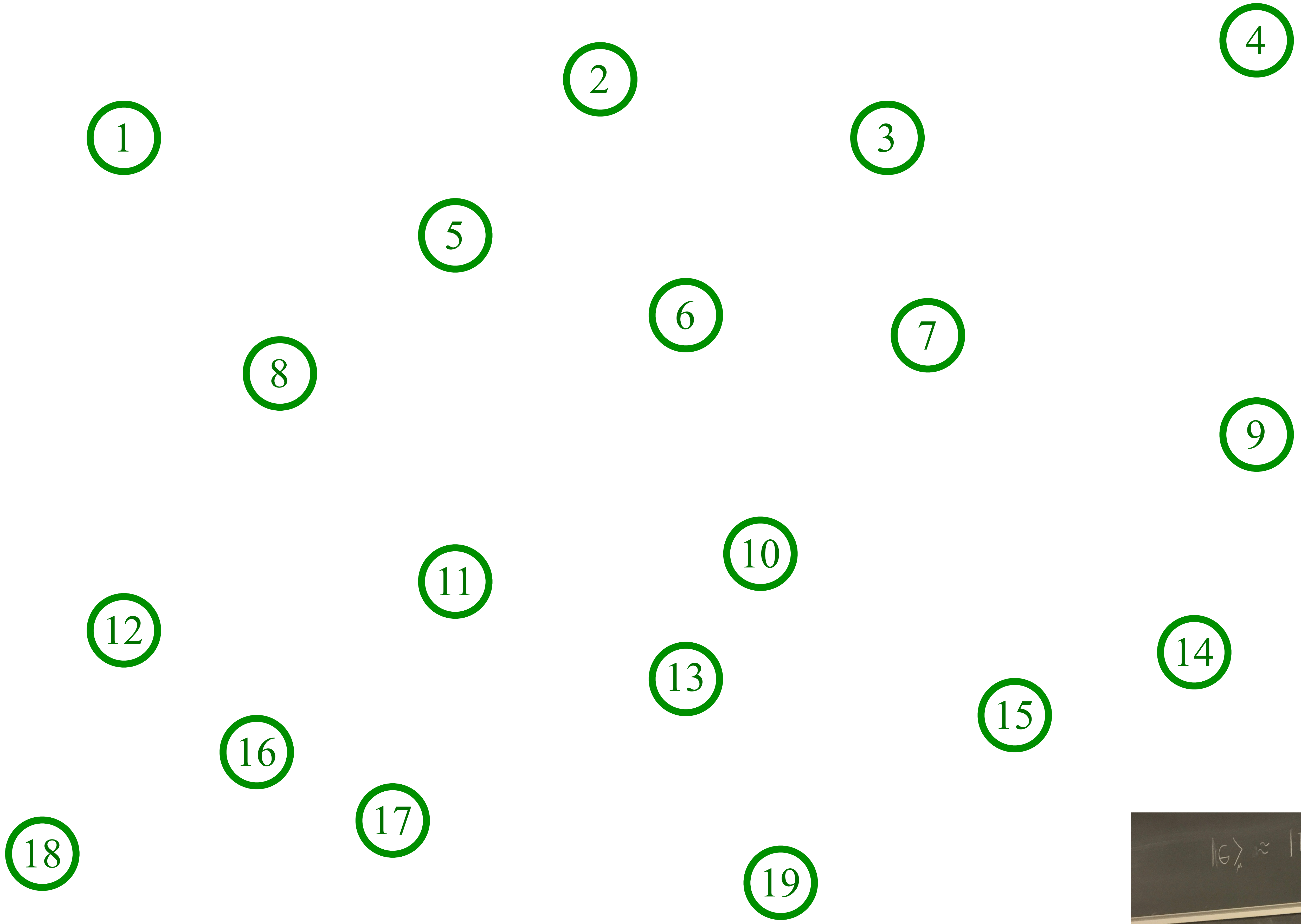
Sachdev-Ye-Kitaev Model

A solvable model of multi-particle entanglement which accounts for quantum interference between successive collisions:

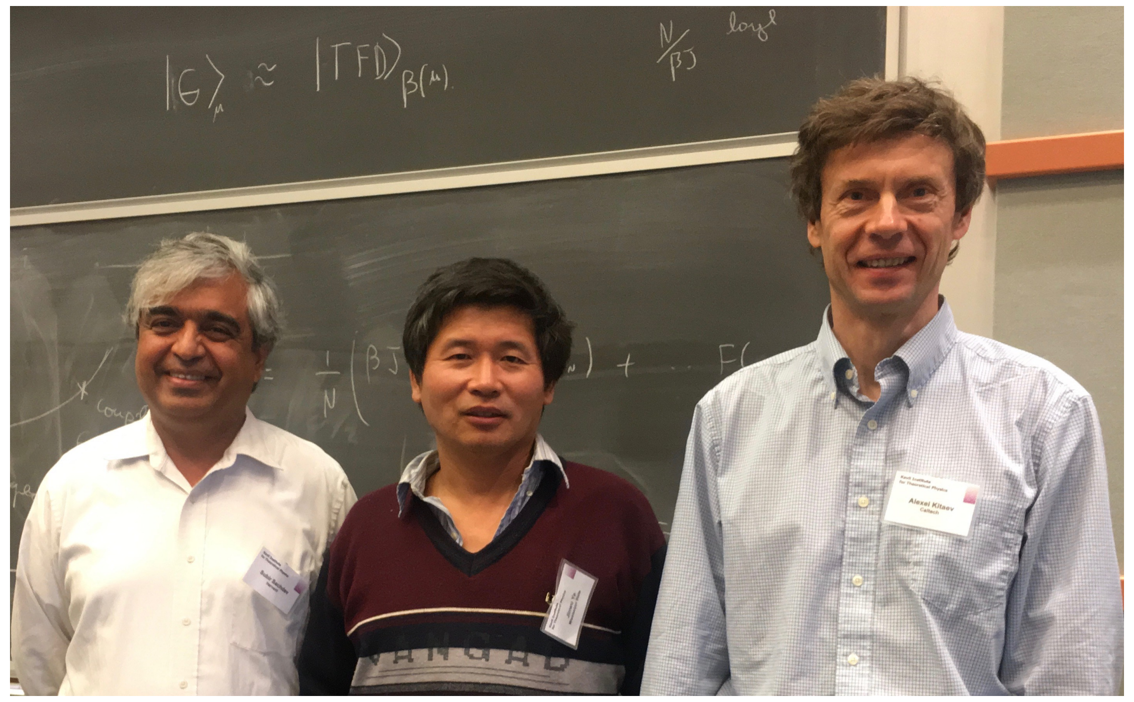
leading to a metal with no particle-like excitations

The SYK model

Sachdev, Ye (1993); Kitaev (2015)



Pick a set of random positions



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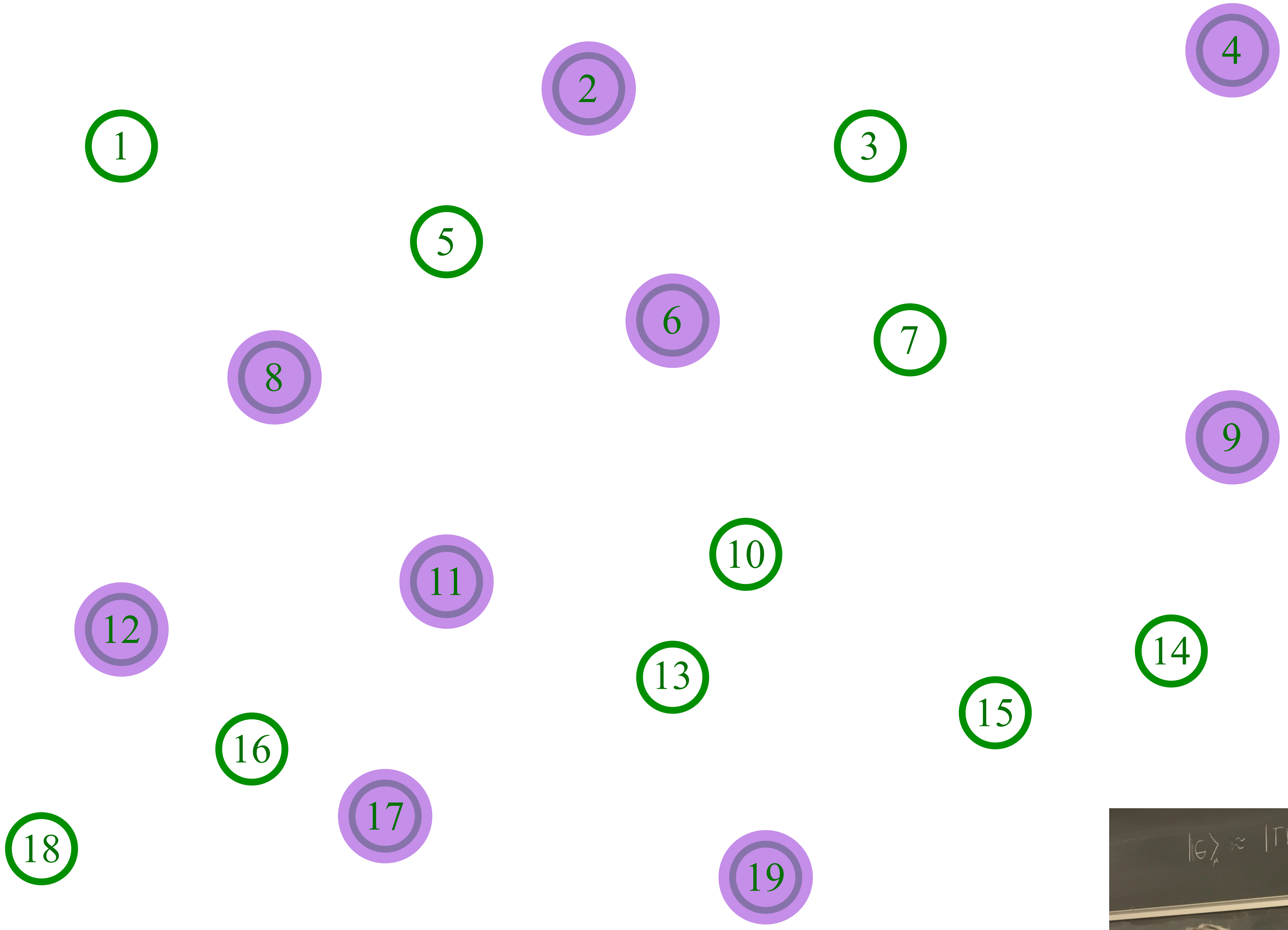
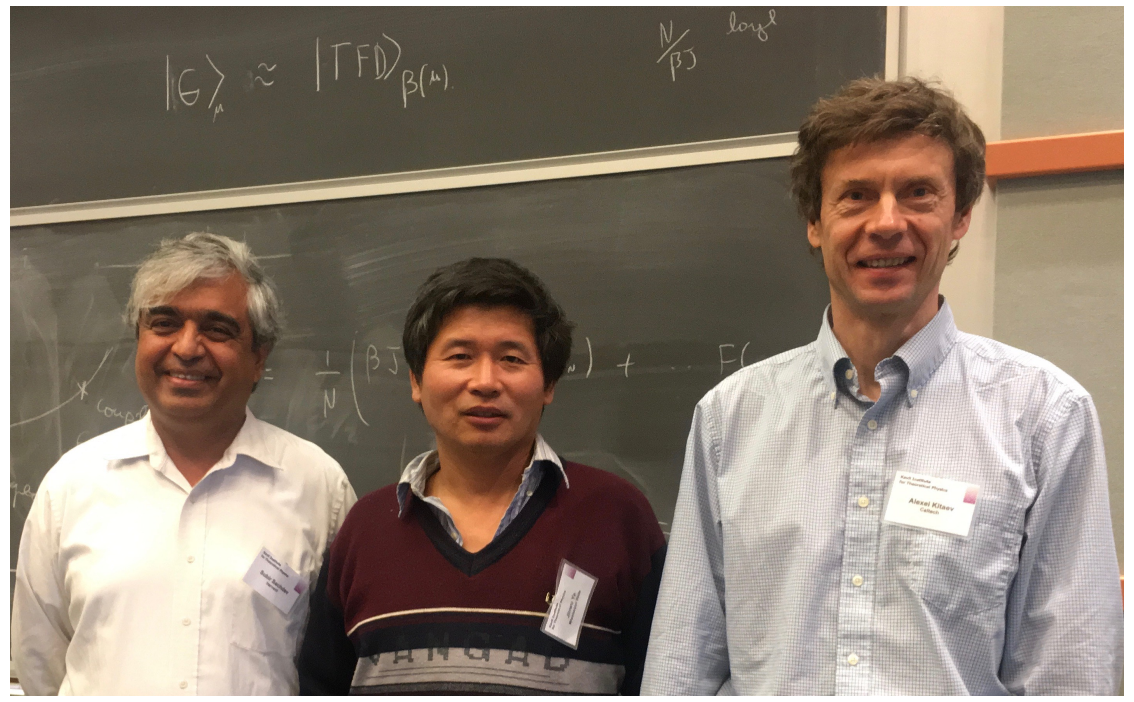


Place electrons randomly on some sites



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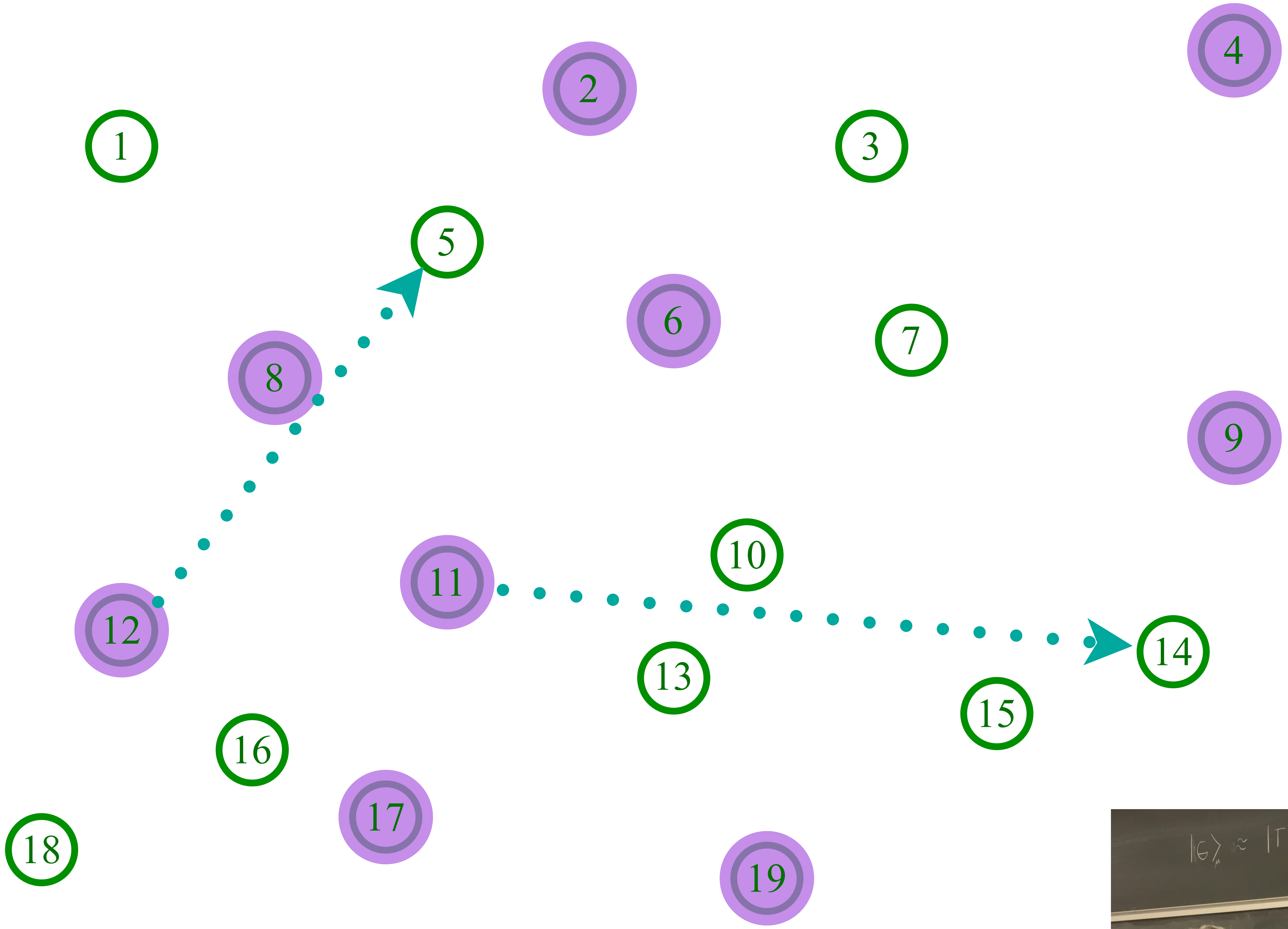
$$U_{11,12;5,14}$$



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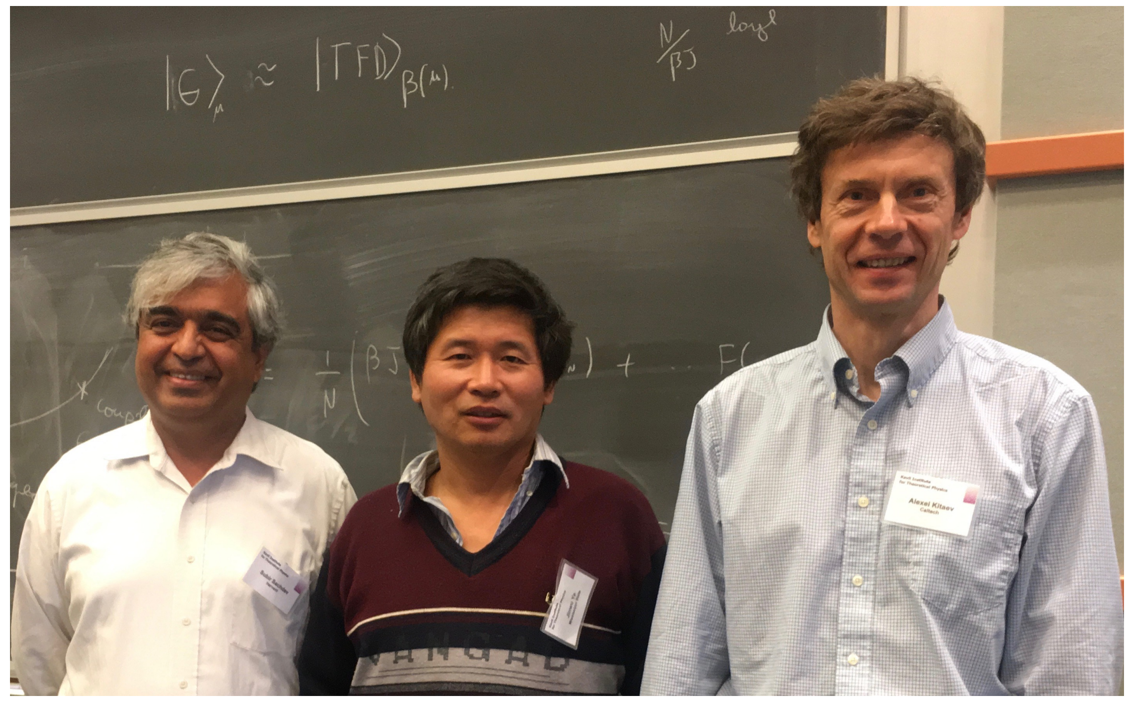

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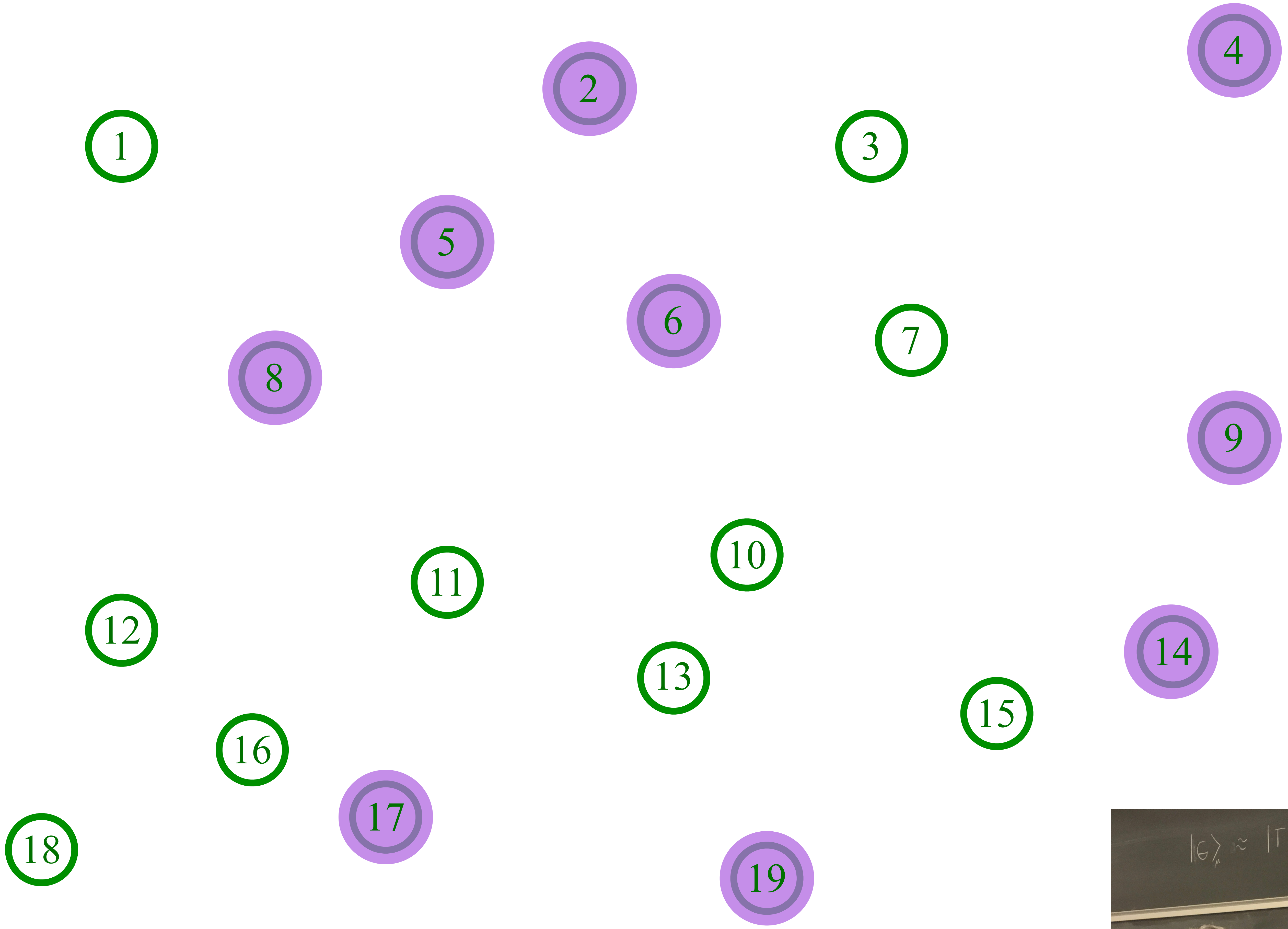
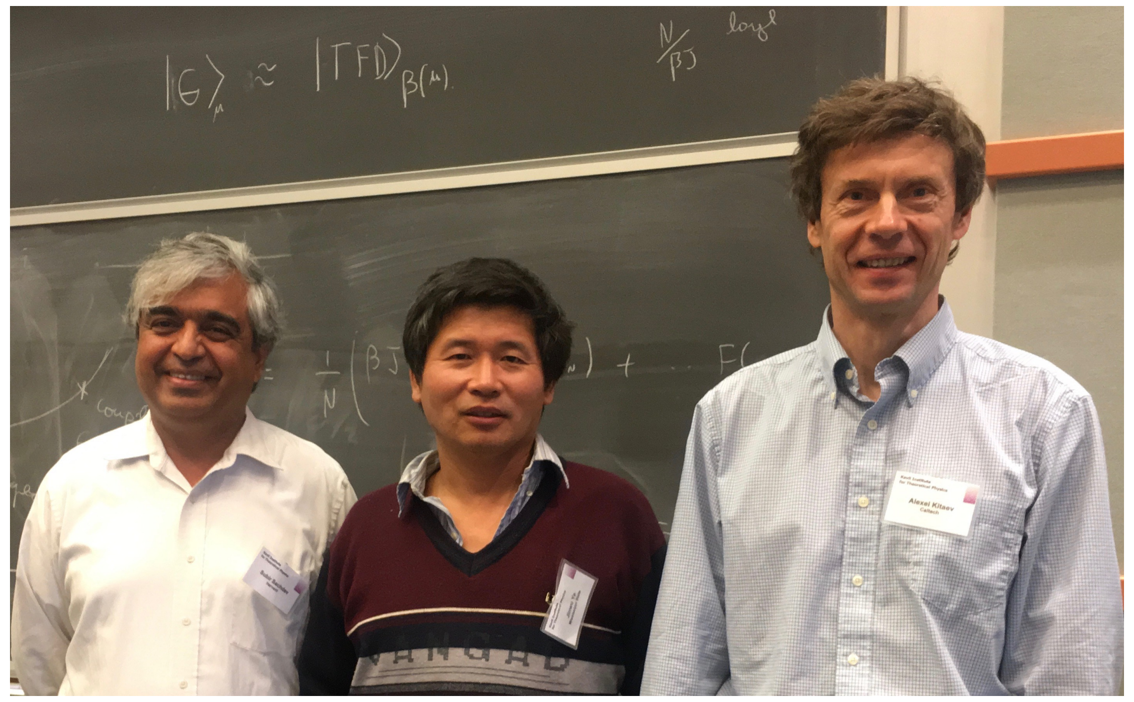


Entangle electrons pairwise randomly



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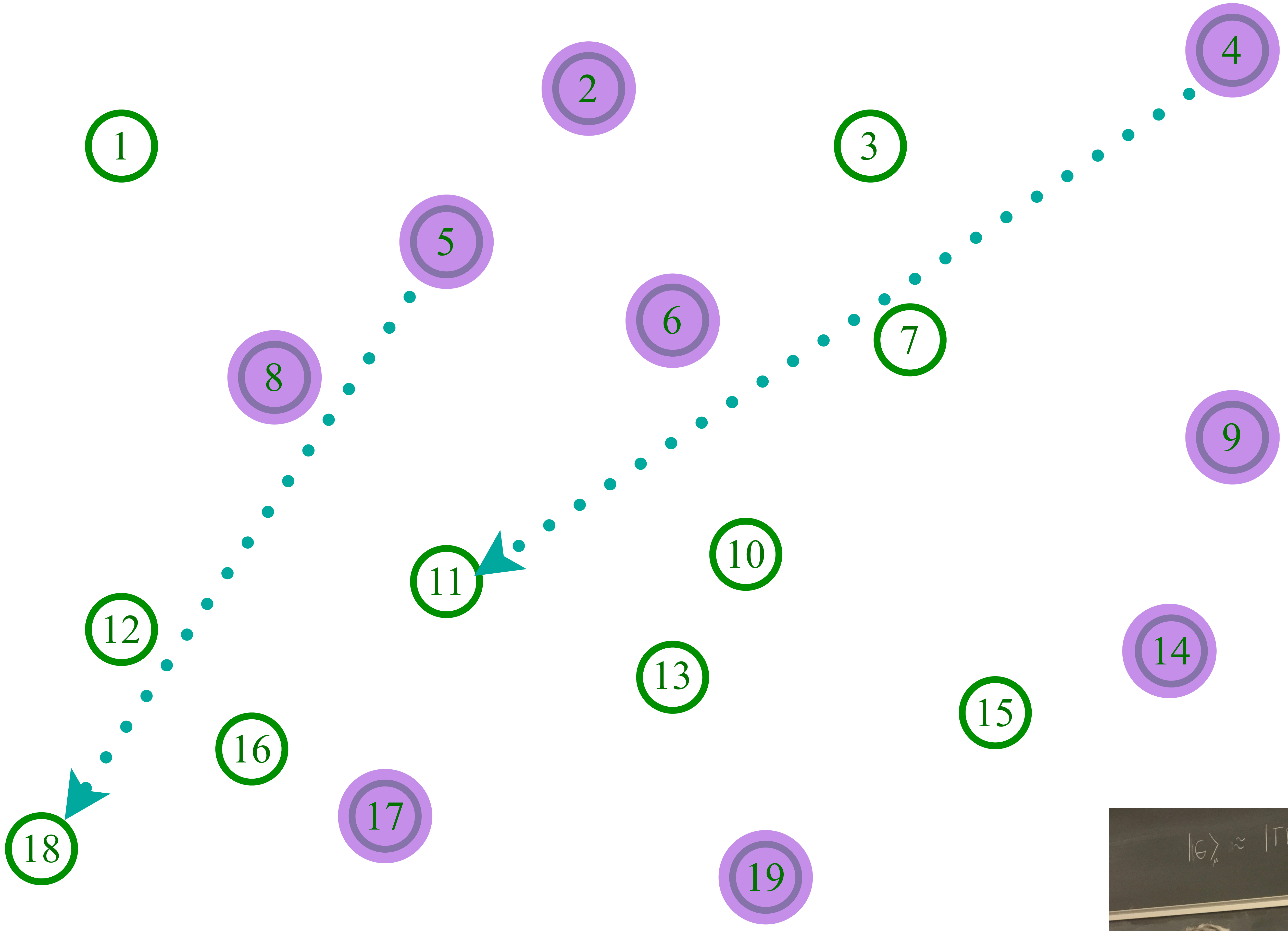
$U_{4,5;11,18}$



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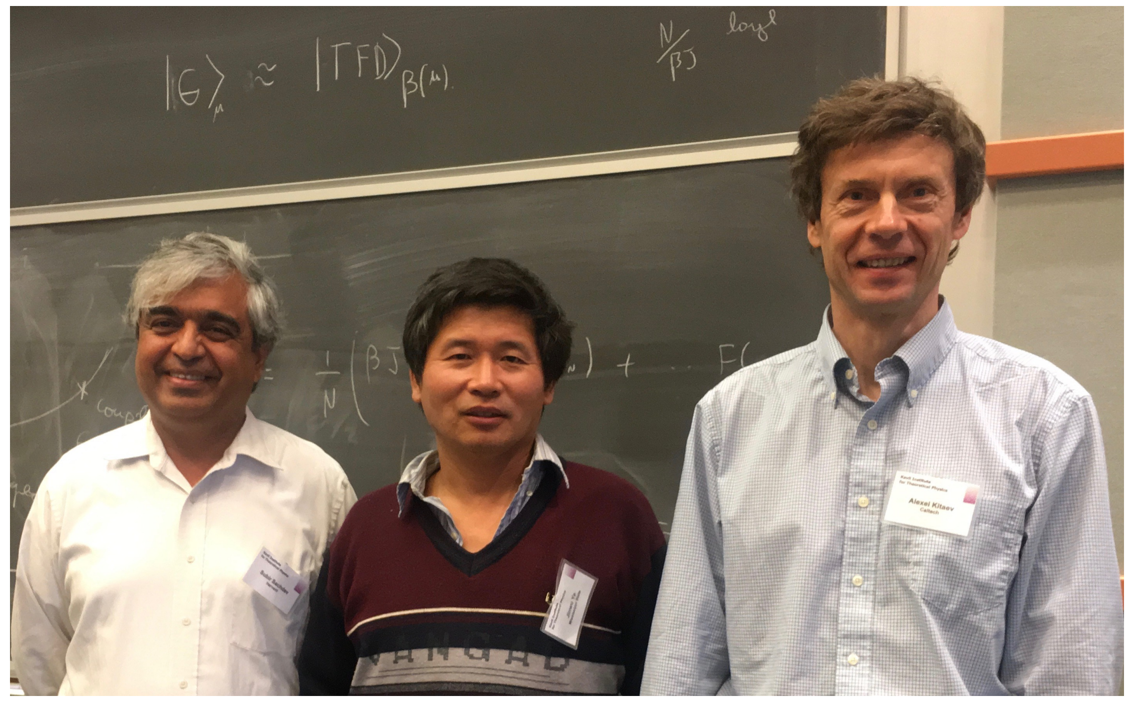

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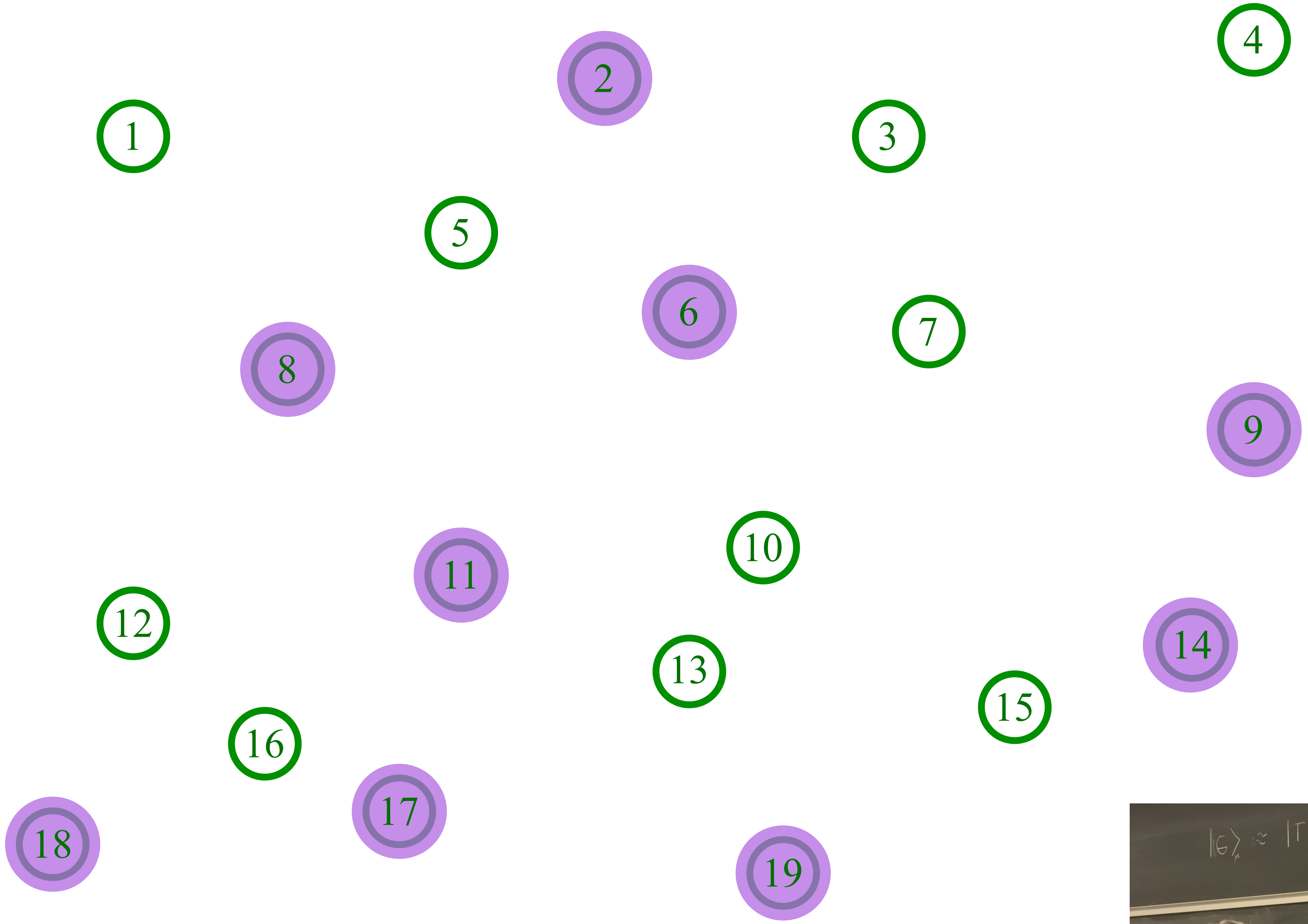
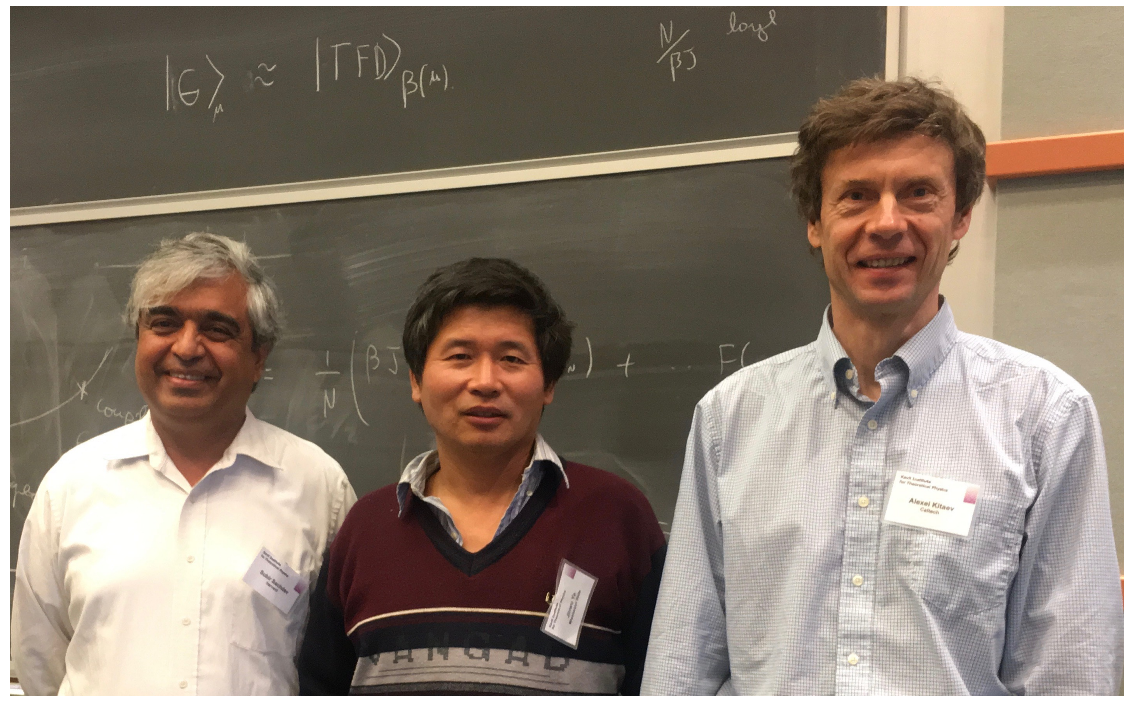
$U_{4,5;11,18}$



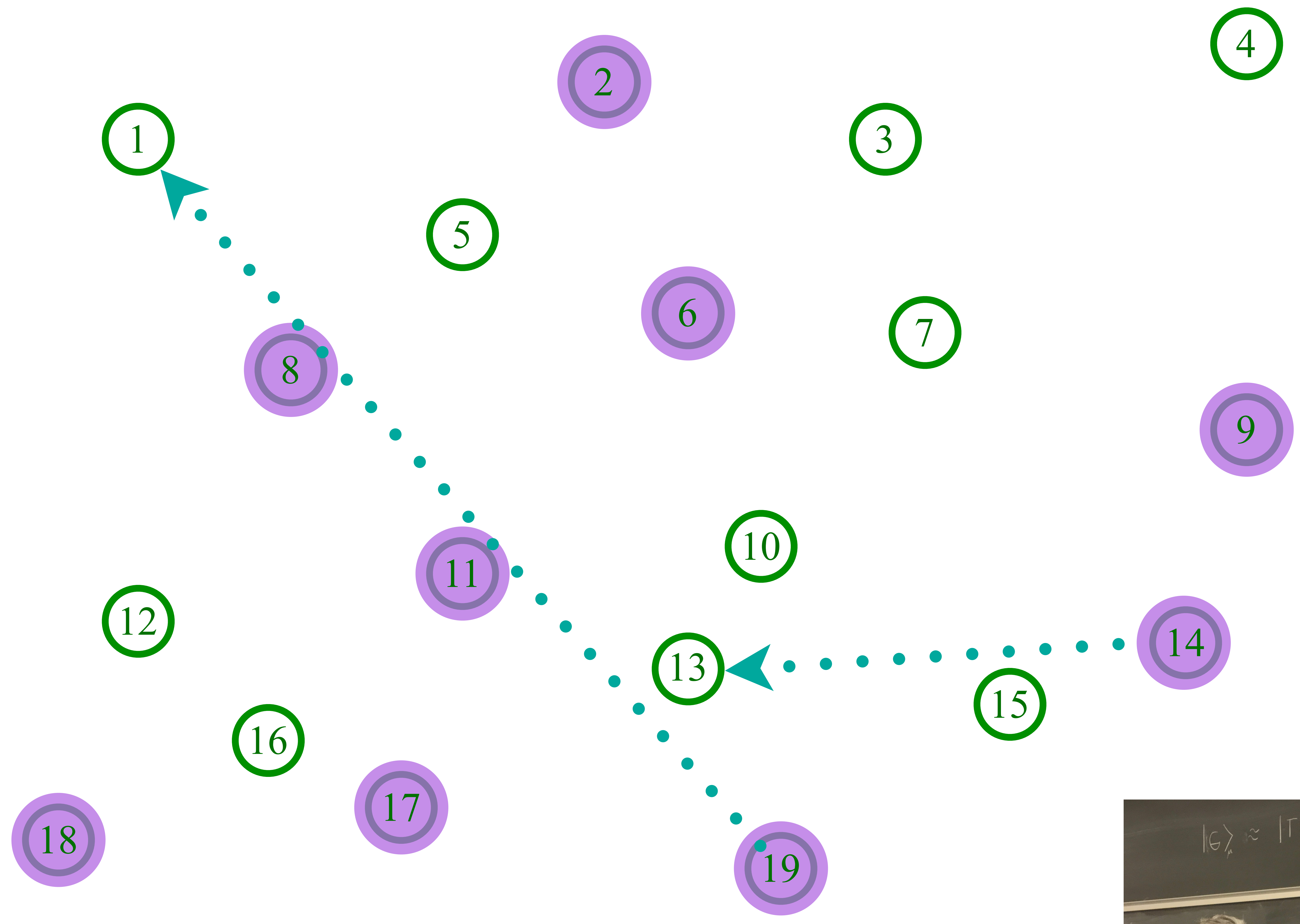
Entangle electrons pairwise randomly



The SYK model

Sachdev, Ye (1993); Kitaev (2015)

$U_{14,19;1,13}$



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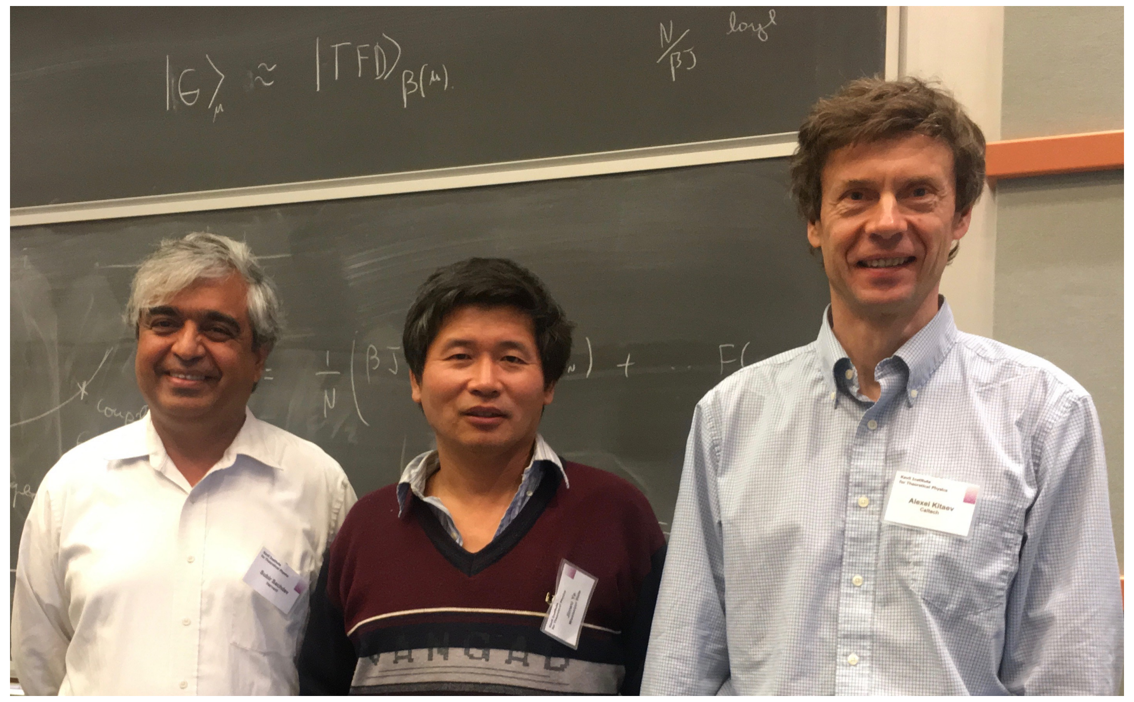

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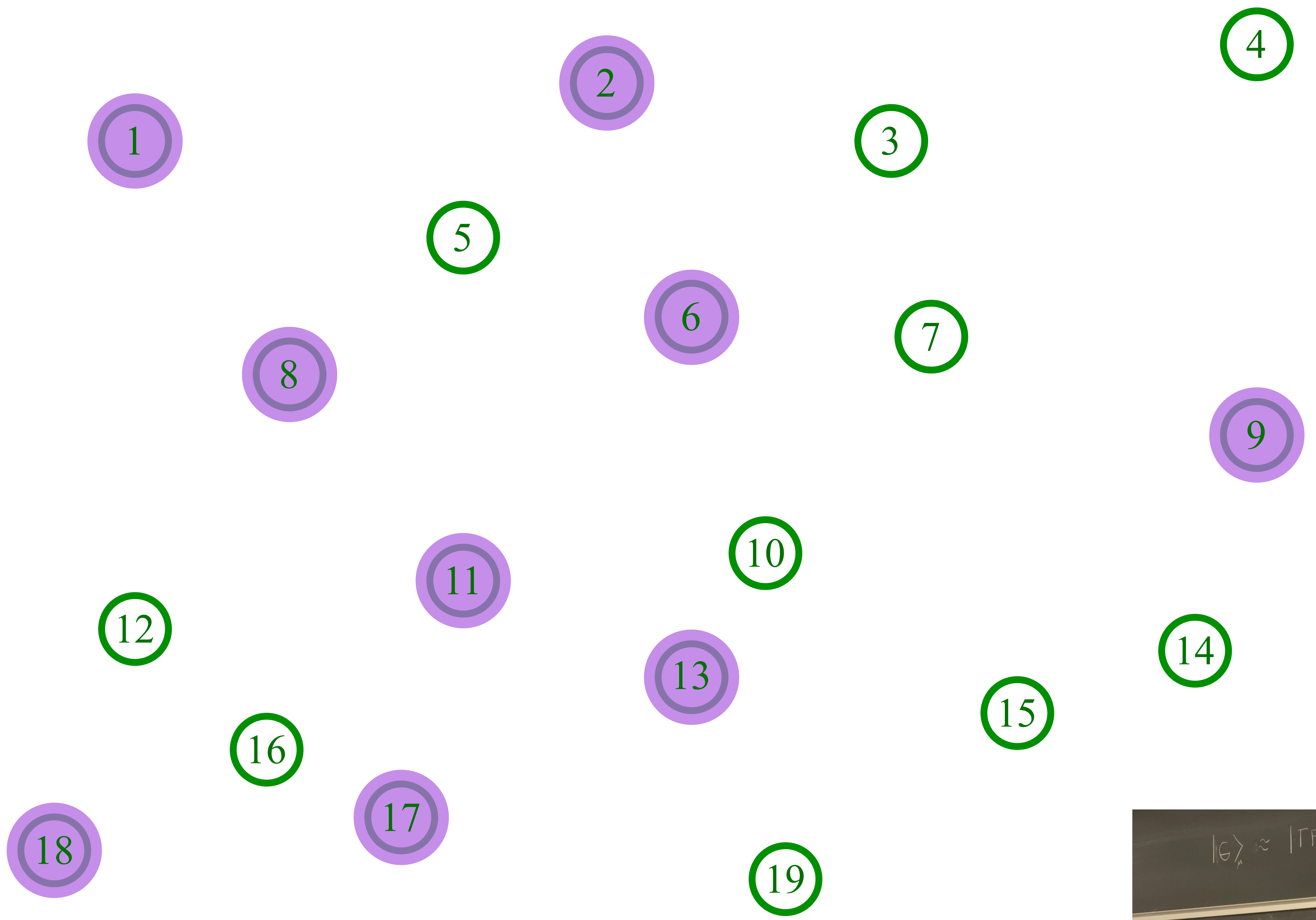

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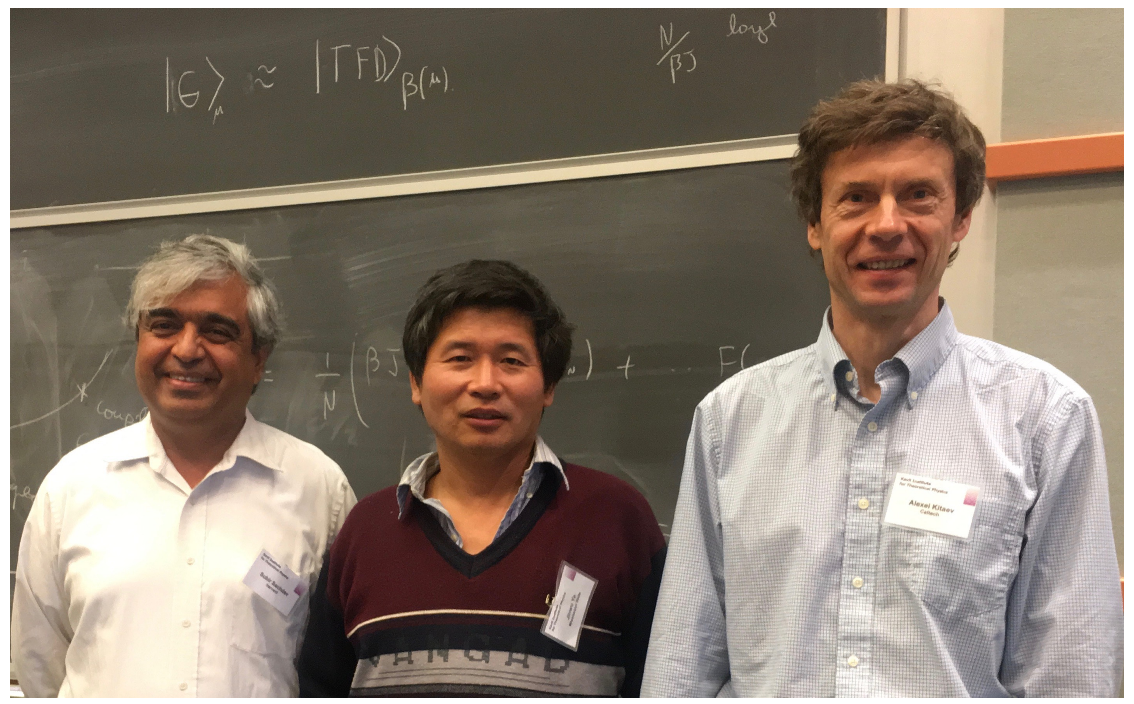

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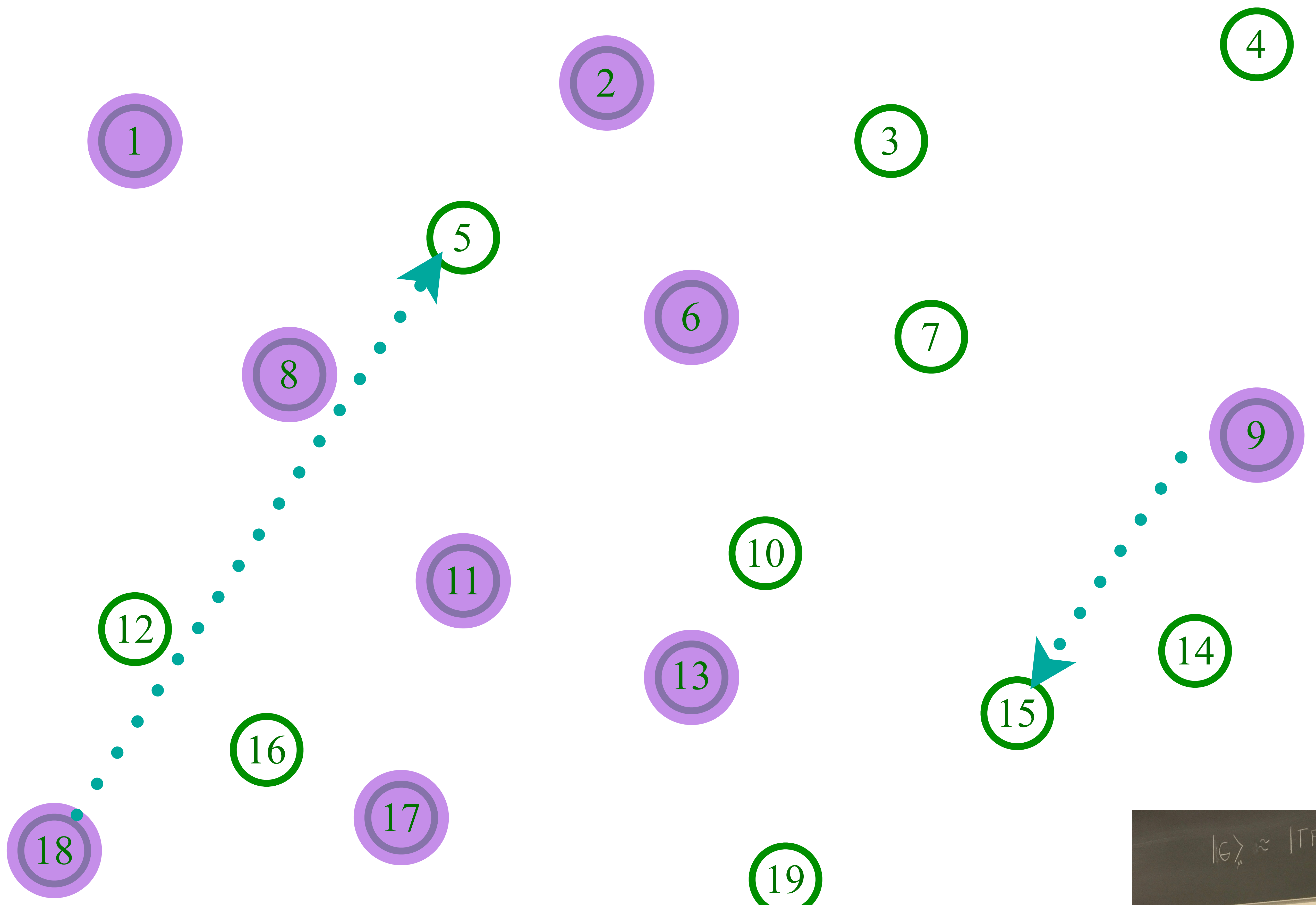

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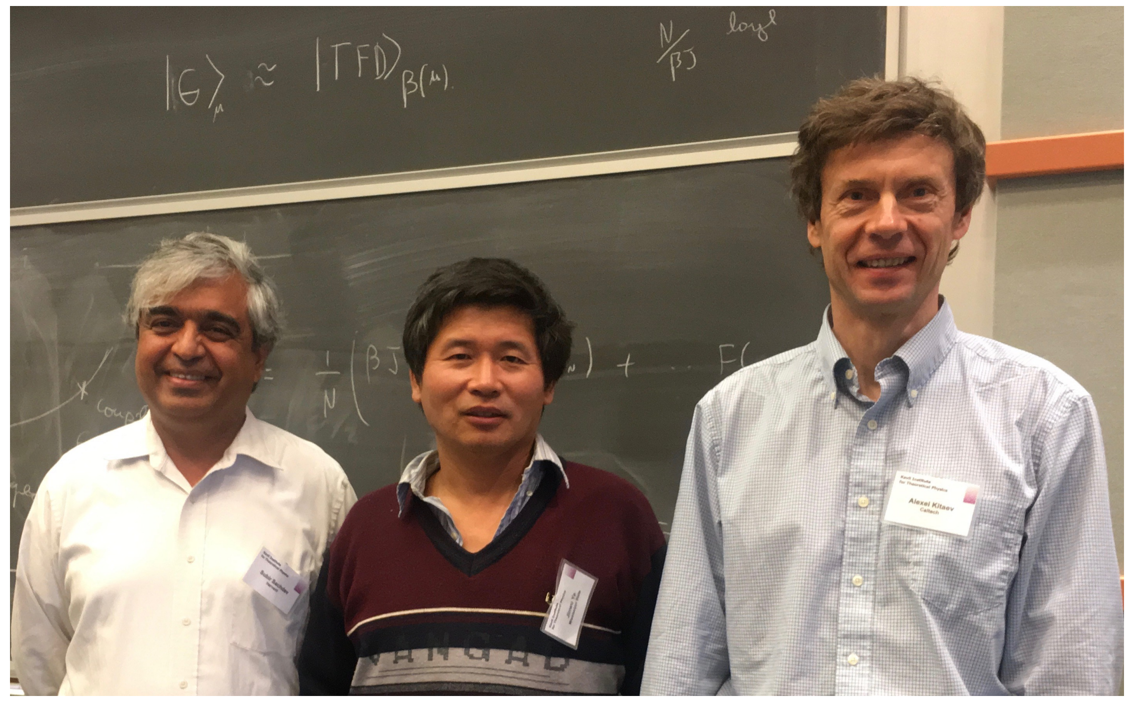

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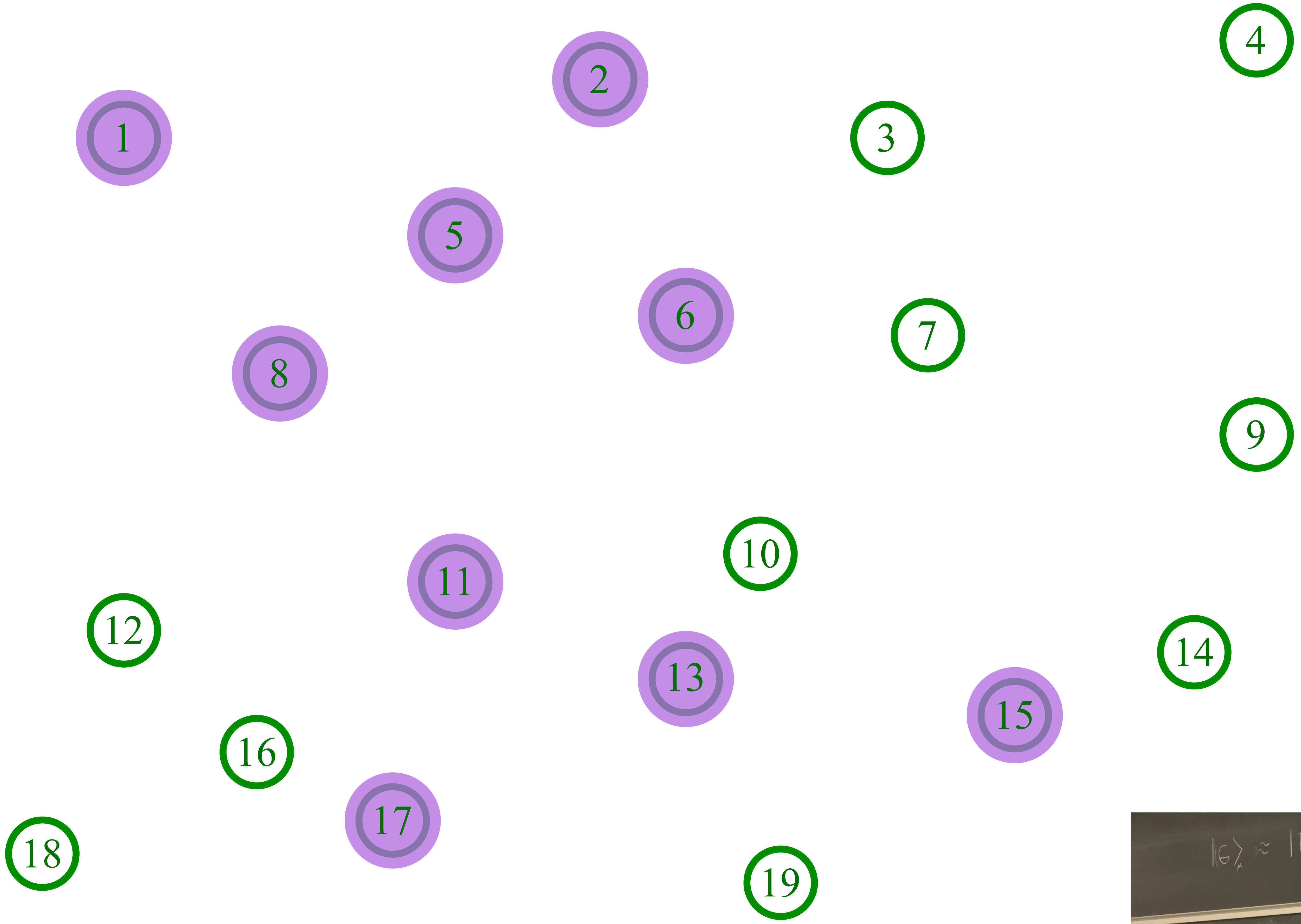

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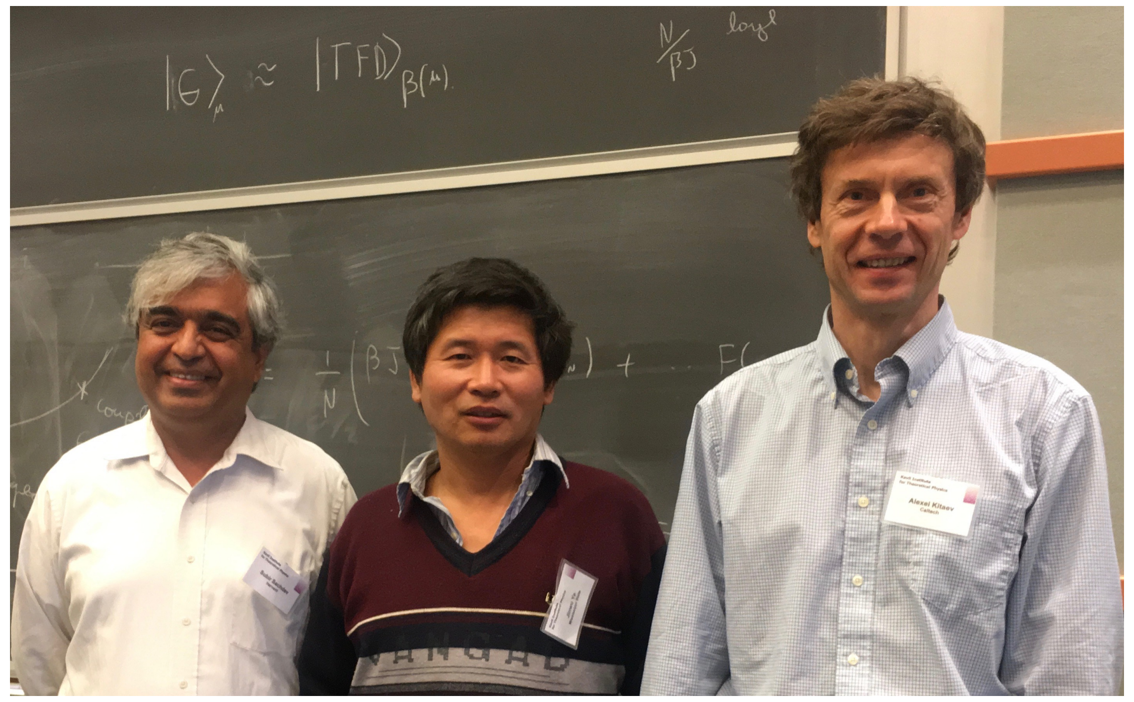

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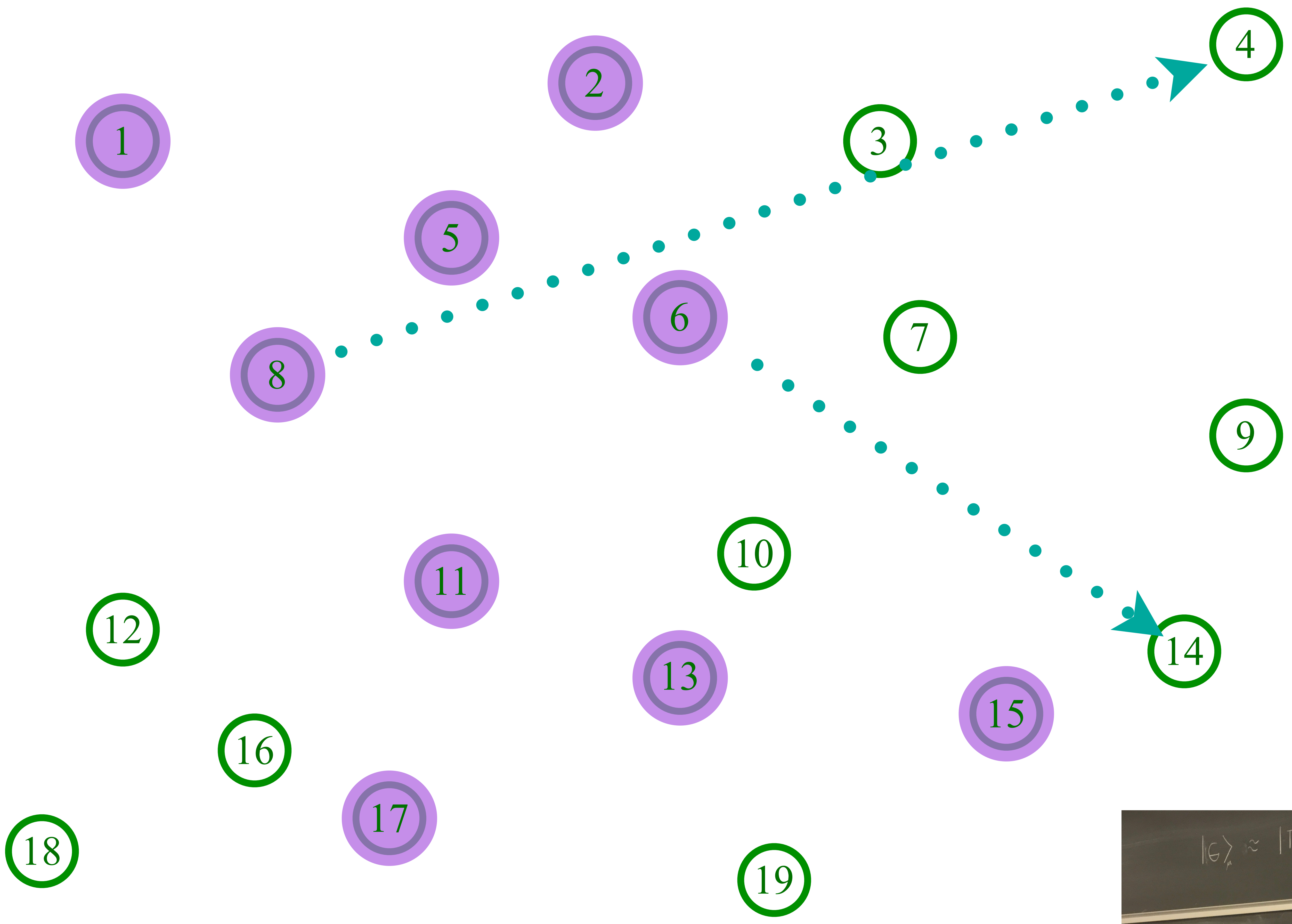
$$U_{6,8;4,14}$$



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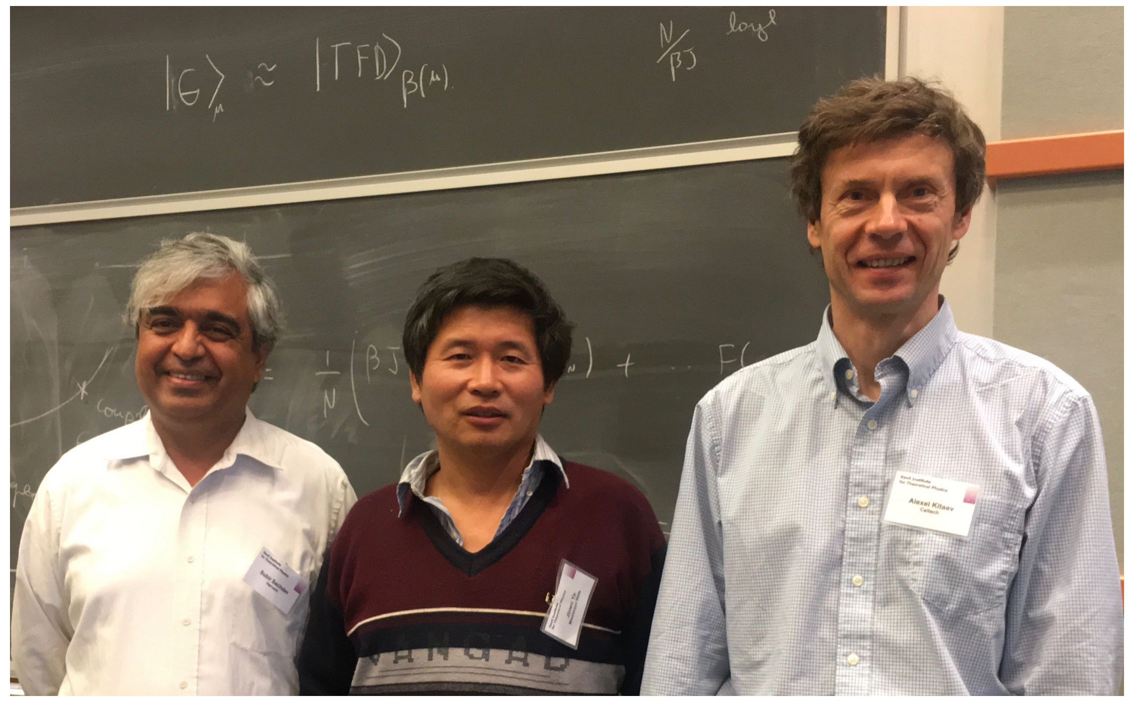

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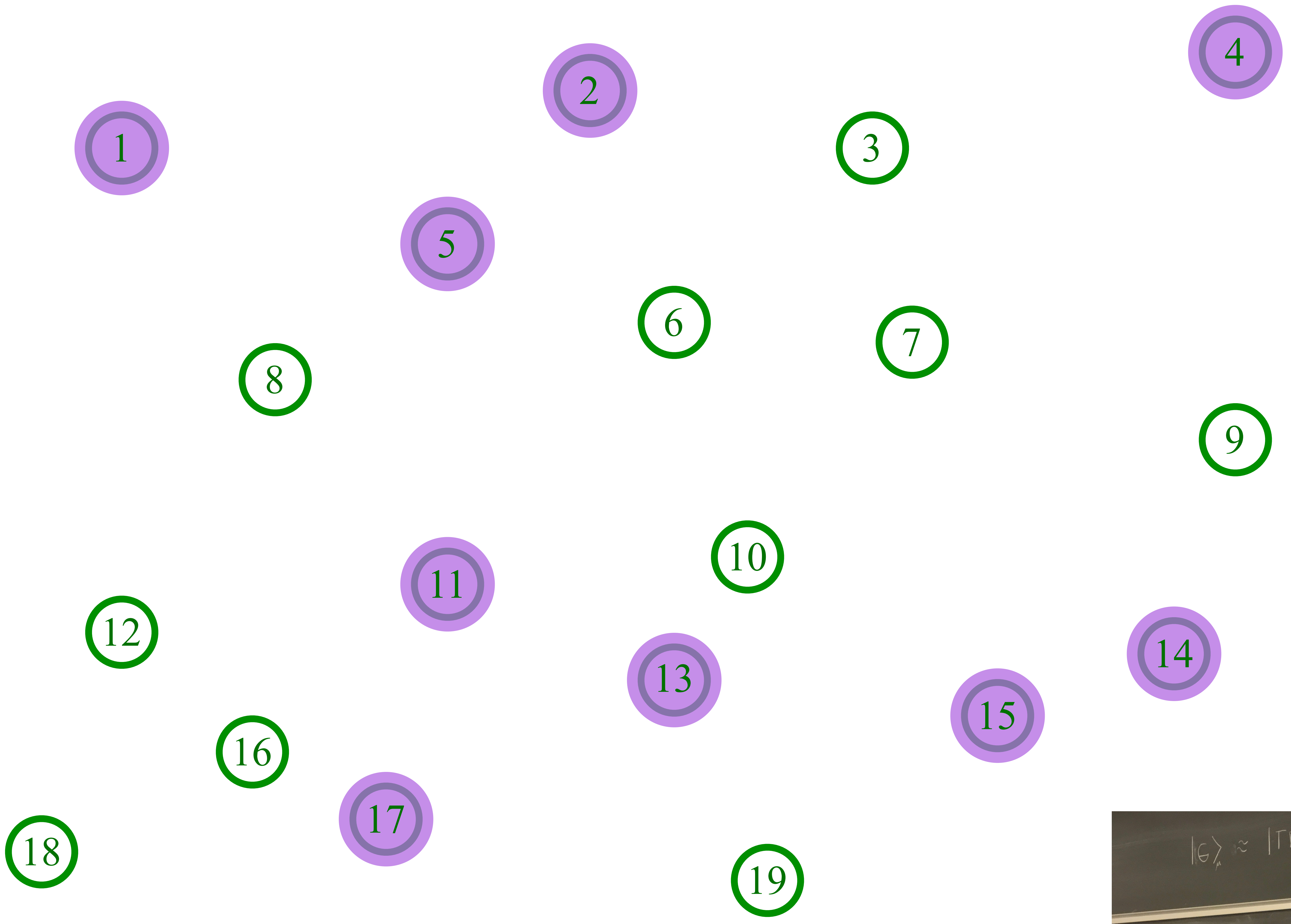

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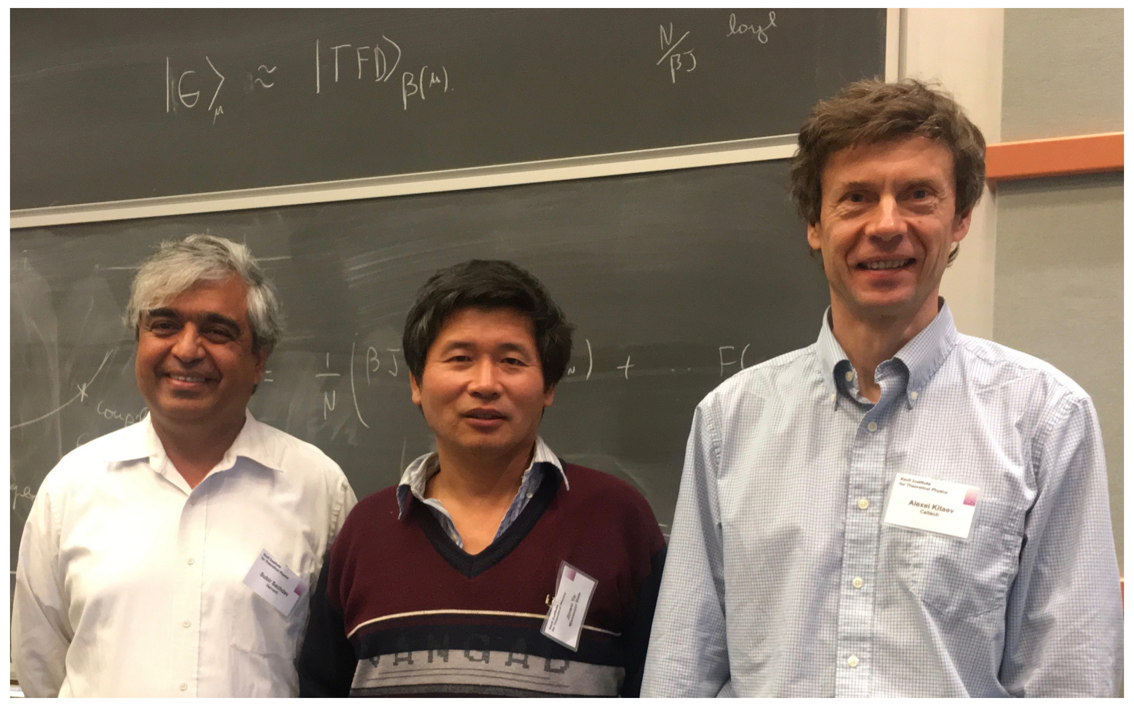


The Sachdev-Ye-Kitaev (SYK) model

(See also: the "2-Body Random Ensemble" in nuclear physics; did not obtain the large *N* limit;
T.A. Brody, J. Flores, J.B. French, P.A. Mello, A. Pandey, and S.S.M. Wong, Rev. Mod. Phys. **53**, 385 (1981))

$$H = \frac{1}{(2N)^{3/2}} \sum_{\alpha,\beta,\gamma,\delta=1}^{N} U_{\alpha\beta;\gamma\delta}\, c_\alpha^\dagger c_\beta^\dagger c_\gamma c_\delta - \mu \sum_\alpha c_\alpha^\dagger c_\alpha$$

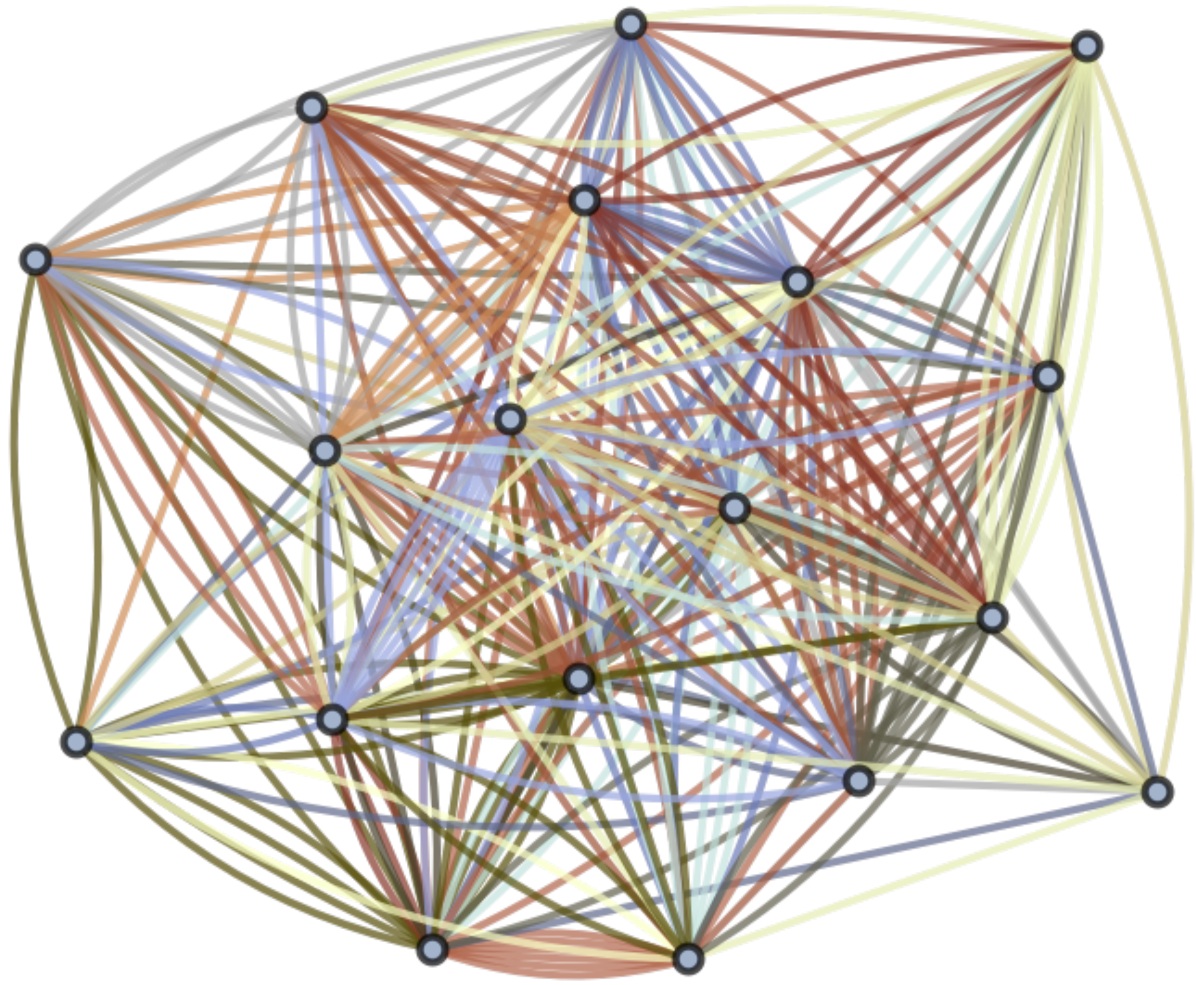
$$c_\alpha c_\beta + c_\beta c_\alpha = 0 \quad , \quad c_\alpha c_\beta^\dagger + c_\beta^\dagger c_\alpha = \delta_{\alpha\beta}$$

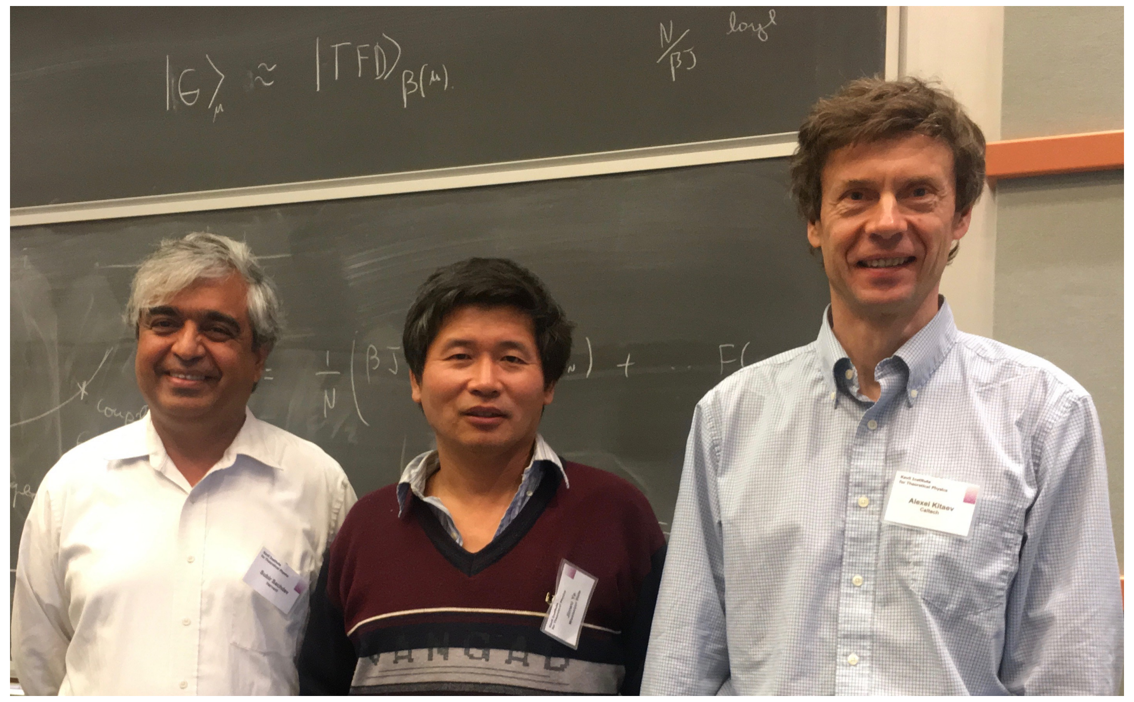
$$\mathcal{Q} = \frac{1}{N} \sum_\alpha c_\alpha^\dagger c_\alpha$$

$U_{\alpha\beta;\gamma\delta}$ are independent random variables with $\overline{U_{\alpha\beta;\gamma\delta}} = 0$ and $\overline{|U_{\alpha\beta;\gamma\delta}|^2} = U^2$
$N \to \infty$ yields critical strange metal.

S. Sachdev and J. Ye, PRL **70**, 3339 (1993)

A. Kitaev, unpublished; S. Sachdev, PRX **5**, 041025 (2015)





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- Planckian time dynamics with a relaxation time $\sim \hbar/(k_B T)$ when $k_B T \ll U$.

D. Chowdhury, A. Georges, O. Parcollet, S. Sachdev, arXiv: 2109.05037, review article

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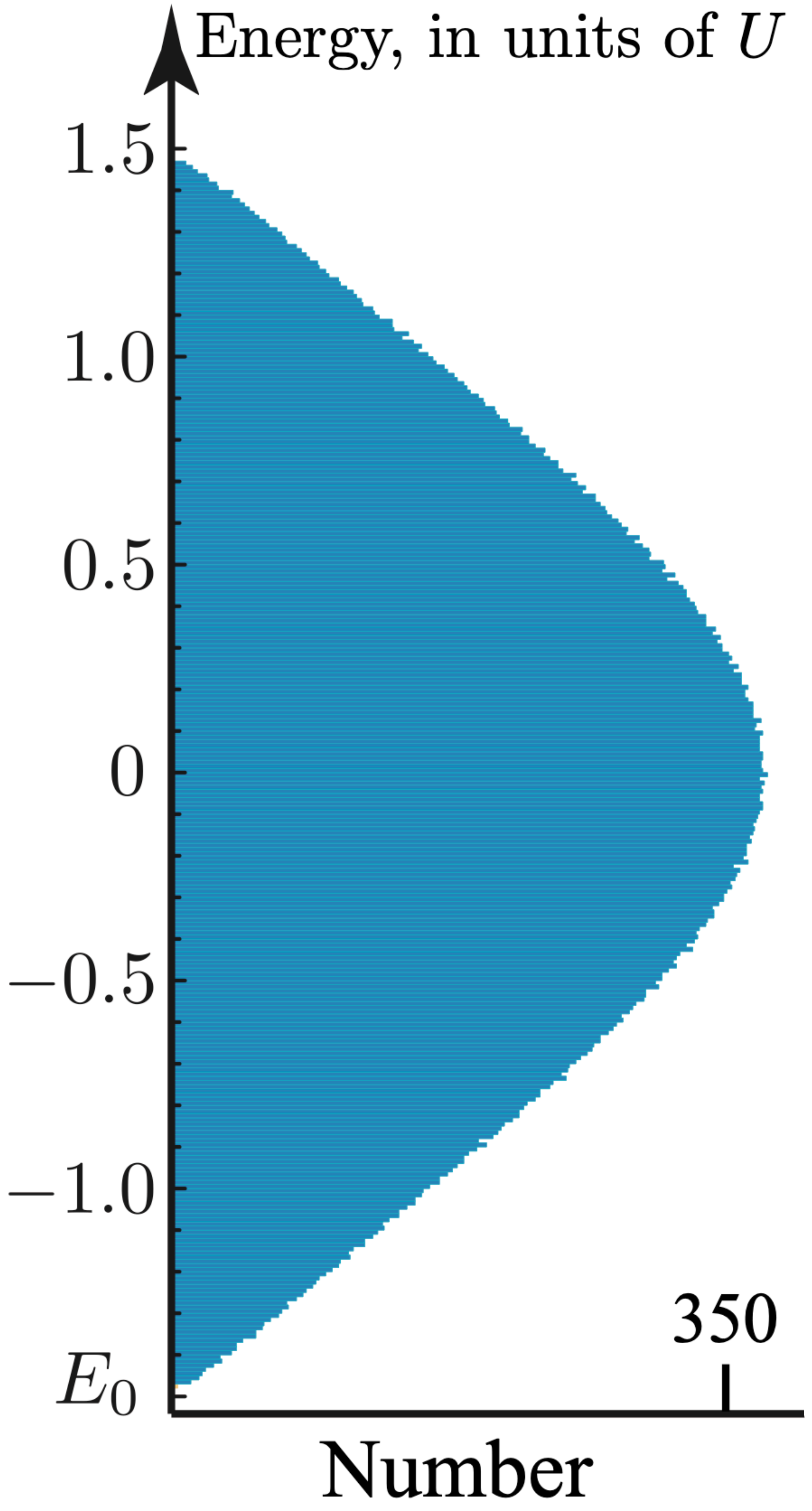
The SYK model

$$D(E) = \sum_i \delta(E - E_i); \quad E_0 + E_i \Rightarrow \text{Many body eigenvalue}$$



Many-body density of states

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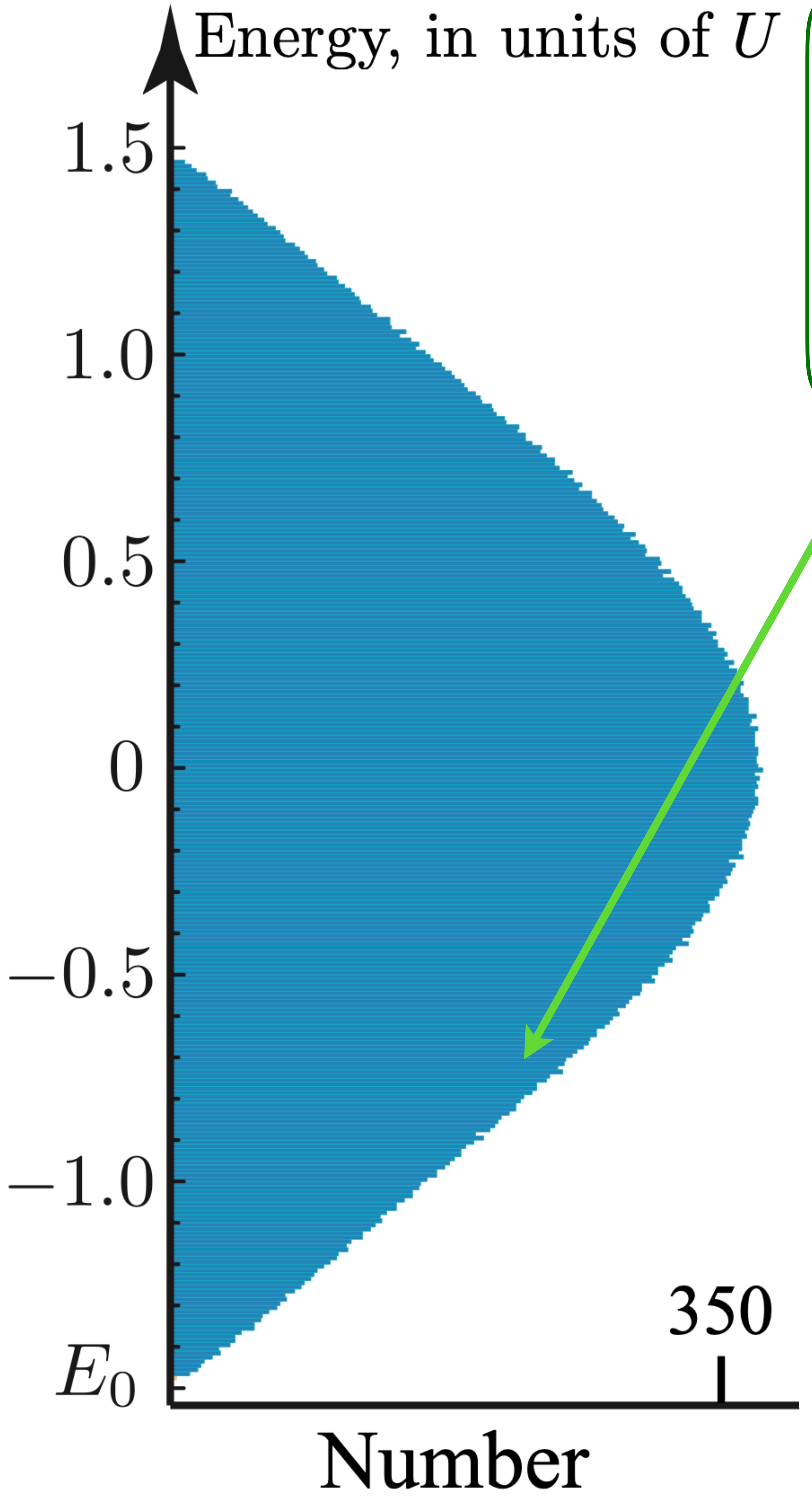

$$D(E) \sim e^{S(E)}$$
$$= e^{N s_0 + \sqrt{2N\gamma E}}$$
$$S(T \to 0) = N(s_0 + \gamma T)$$

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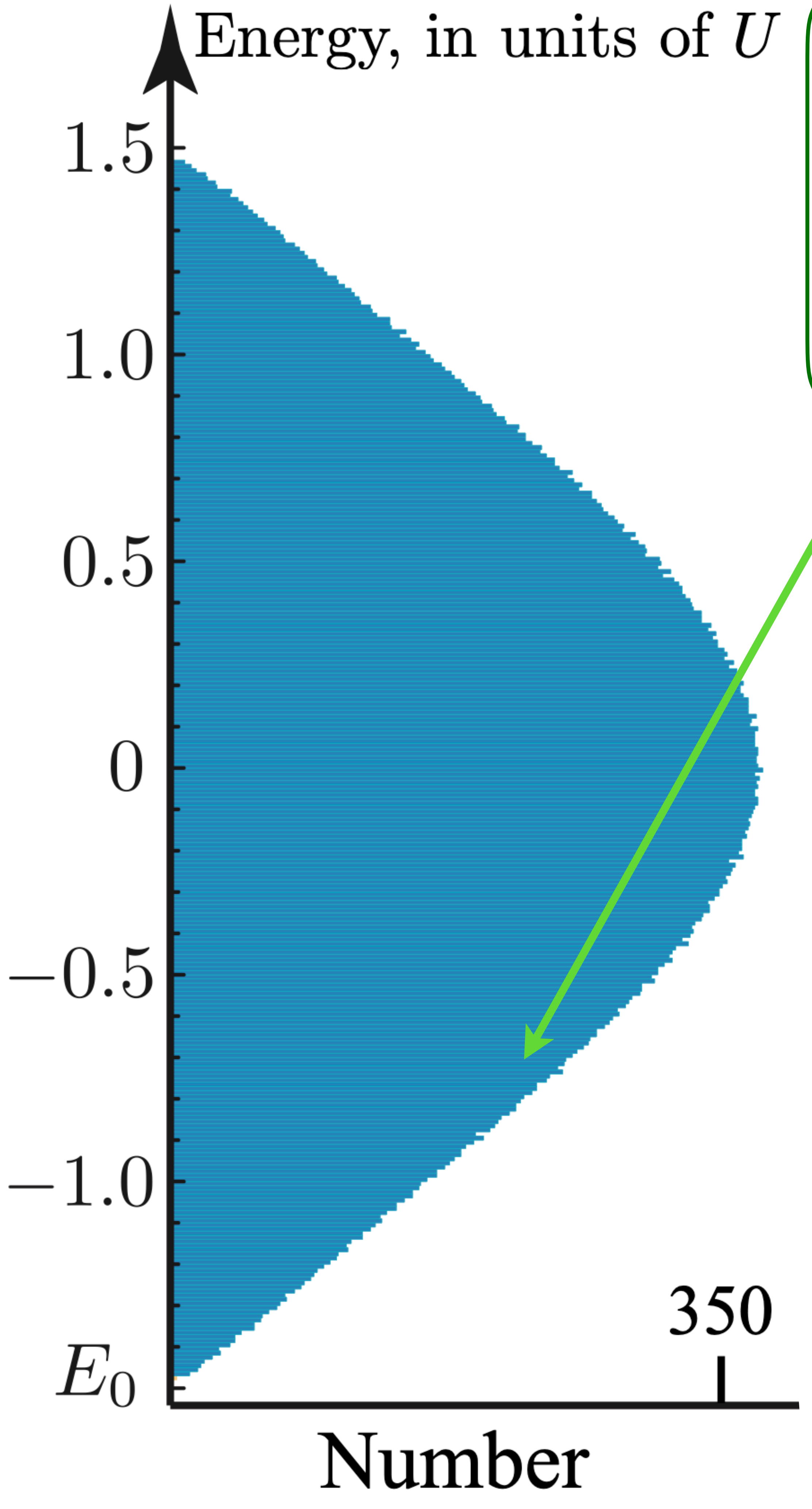

$$D(E) \sim e^{S(E)}$$
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$$s_0 = 0.464848\ldots$$

A. Georges, O. Parcollet, and S. Sachdev, PRB **63**, 134406 (2001)

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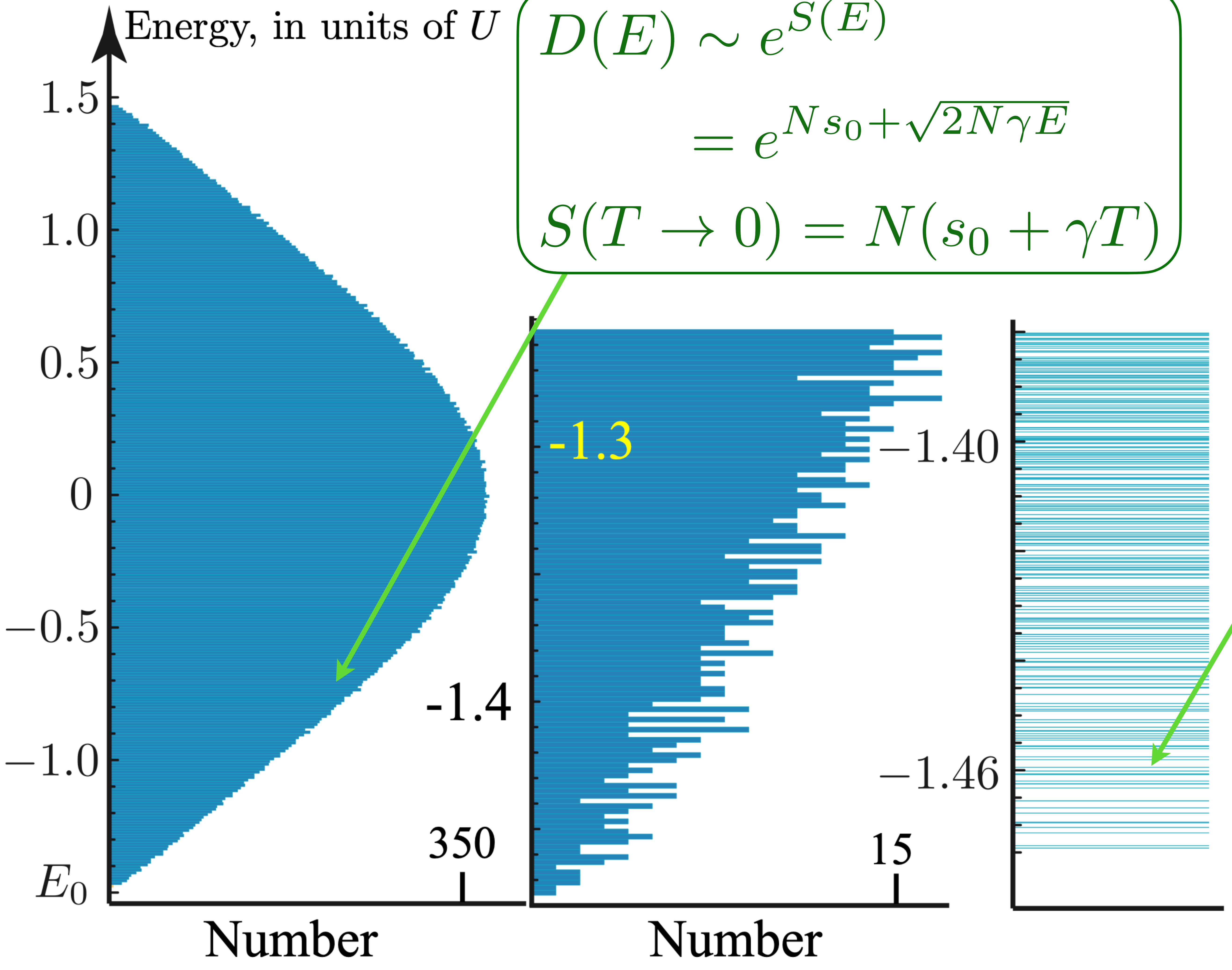

$$D(E) \sim e^{S(E)} = e^{Ns_0 + \sqrt{2N\gamma E}}$$

$$S(T \to 0) = N(s_0 + \gamma T)$$

$$D(E) \sim 2\, e^{Ns_0} \sqrt{2N\gamma E}$$

No particle-like decomposition: wavefunctions change chaotically from one state to the next.

$$s_0 = 0.464848\ldots$$

A. Georges, O. Parcollet, and S. Sachdev, PRB **63**, 134406 (2001)

Maldacena, Stanford (2016); Bagrets *et al.* (2017); Cotler *et al.* (2017); Kitaev, Suh (2017); Stanford, Witten (2017)

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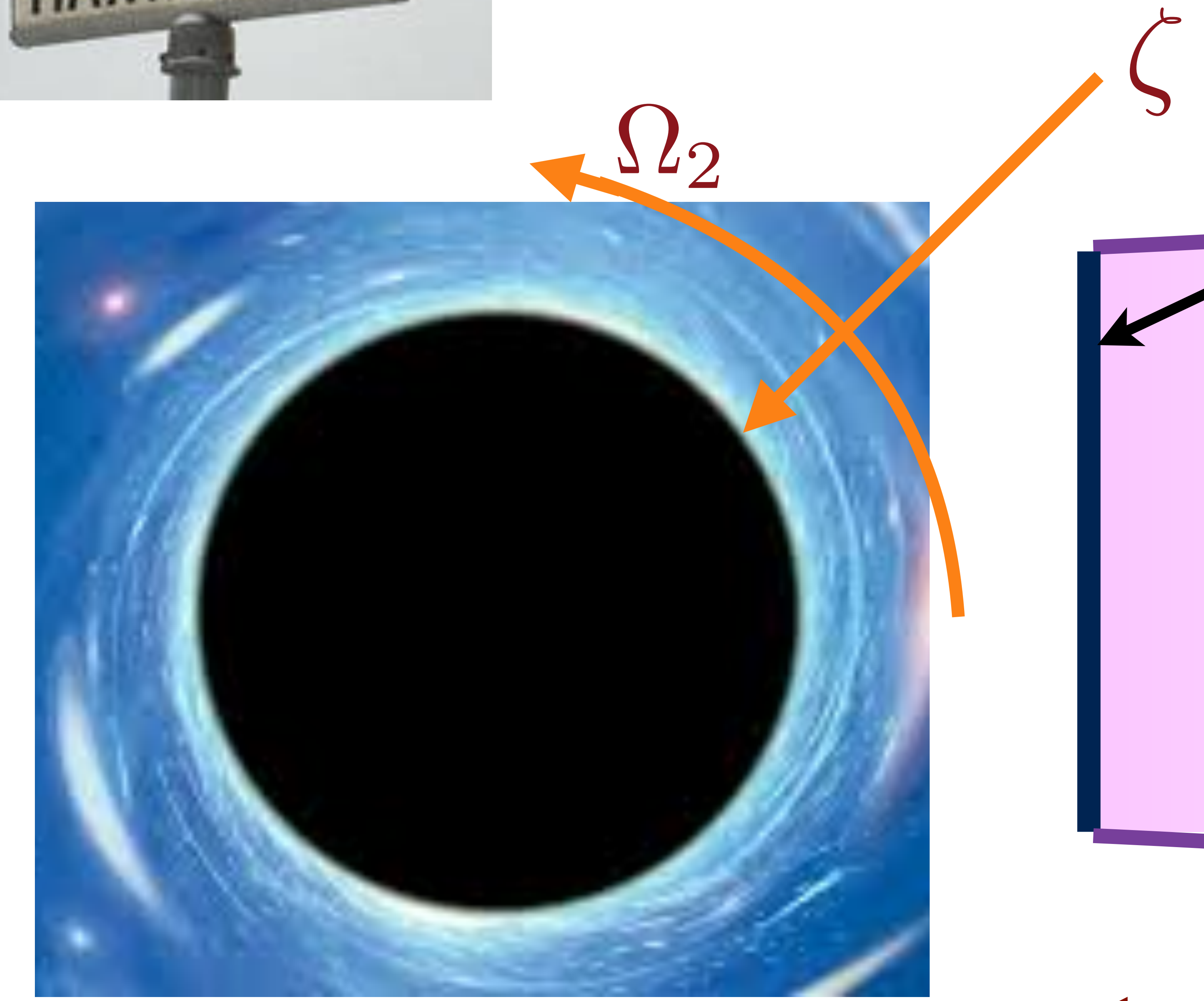
- The $D(E)$ is determined by a time-reparameterization $\tau \to f(\tau)$ mode (similar to the graviton being fluctuations of the spacetime metric), and a phase mode $\phi(\tau)$:

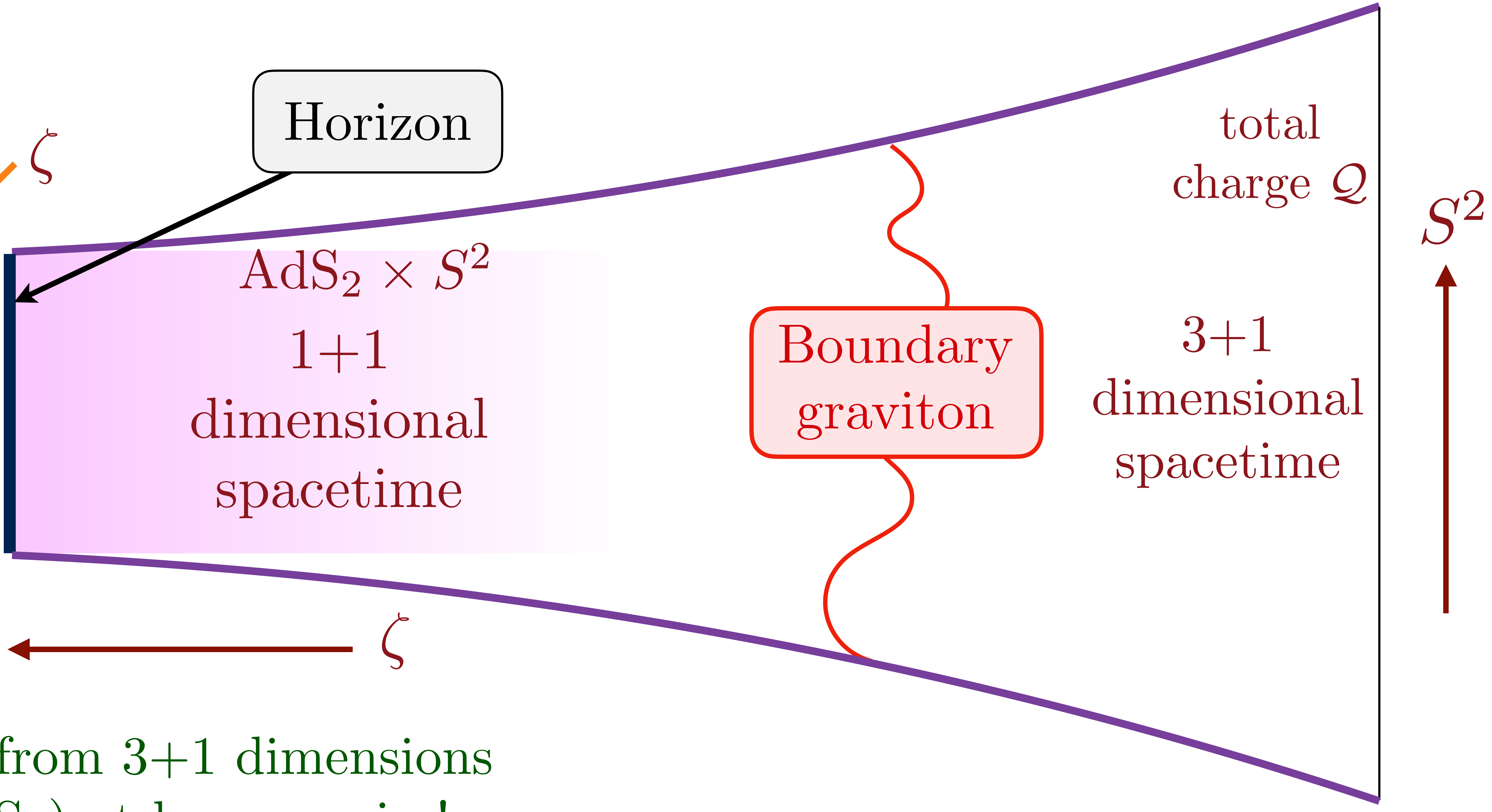
$$\mathcal{Z}_{SYK} = e^{N s_0} \int \mathcal{D}f \mathcal{D}\phi \exp\left(-\frac{1}{\hbar} \int_0^{\hbar/(k_B T)} d\tau \, \mathcal{L}_{SYK}[f, \phi] \right)$$

D. Chowdhury, A. Georges, O. Parcollet, S. Sachdev, arXiv: 2109.05037, review article

From the SYK model to charged black holes

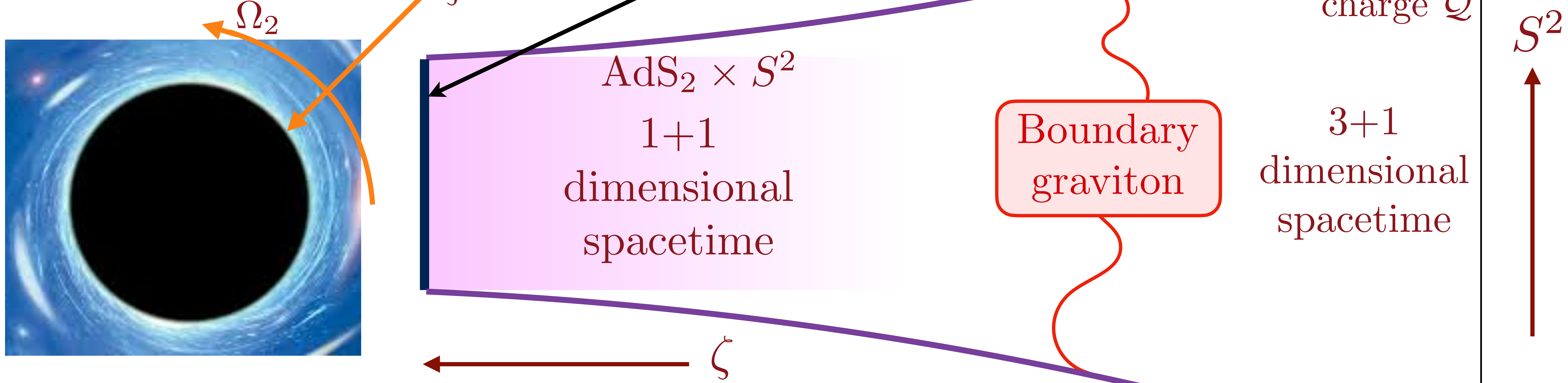
Reissner-Nordstrom black hole of Einstein-Maxwell theory





Dimensional reduction from 3+1 dimensions to 1+1 dimensions (AdS_2) at low energies!

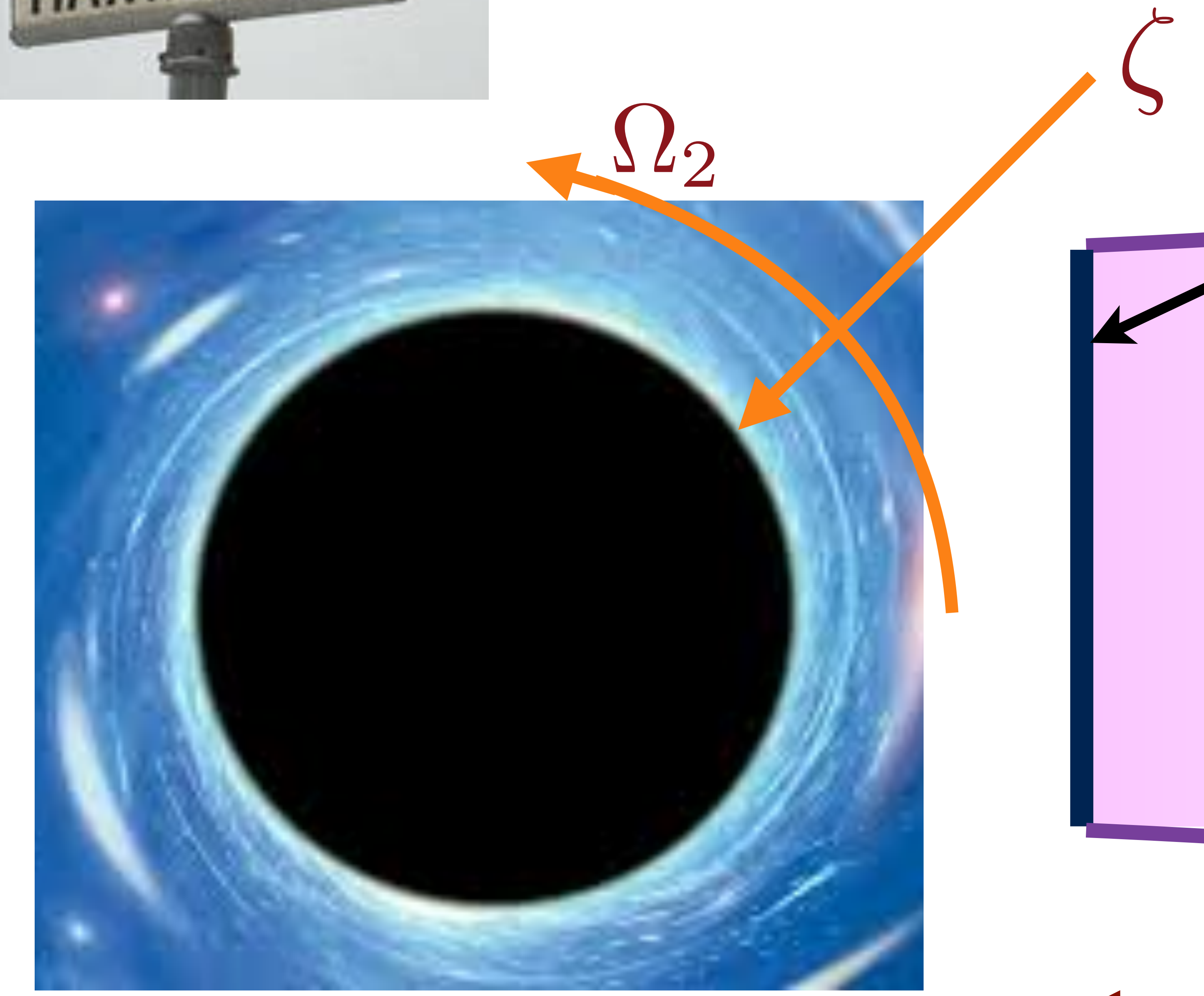
Reissner-Nordstrom black hole of Einstein-Maxwell theory

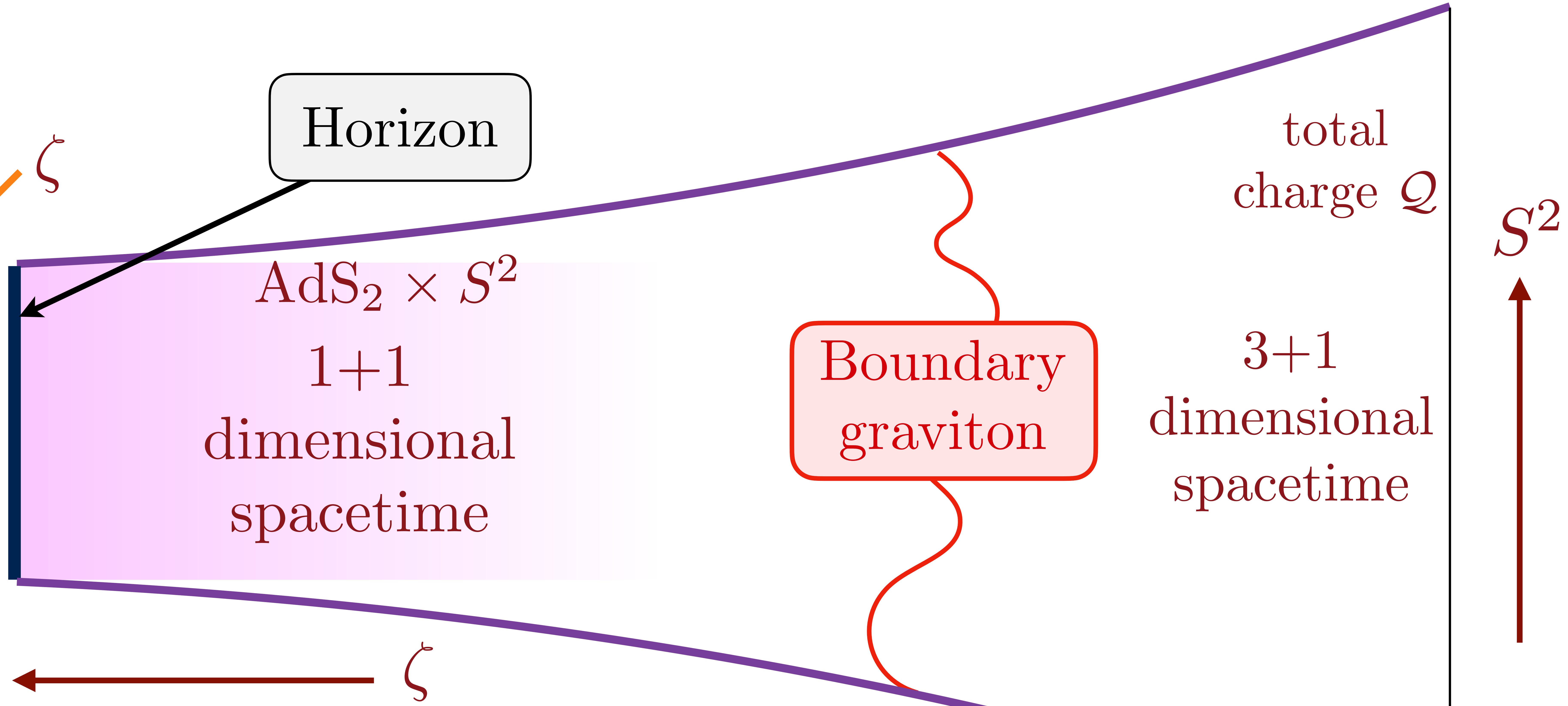


Dimensional reduction from 3+1 dimensions
to 1+1 dimensions (AdS_2) at low energies!

Is there a mapping to a quantum system with Planckian dynamics in 0+1 dimensions?

Reissner-Nordstrom black hole of Einstein-Maxwell theory





Yes!

The low energy theory of a charged black hole is *exactly* the low energy theory of time reparameterizations of the SYK model.

Quantum theory of charged black holes

The near-horizon 1+1 dimensional theory of a charged black hole

$$\mathcal{Z}_{\mathcal{Q}} = \int \mathcal{D}g_{\mu\nu}\mathcal{D}a_\mu \exp\left(-\frac{1}{\hbar}\int d\zeta \int_0^{\hbar/(k_B T)} d\tau \sqrt{g}\, \mathcal{L}_1[g_{\mu\nu}, a_\mu]\right)$$

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$$\Updownarrow$$

$$\mathcal{Z}_{SYK} = e^{N s_0} \int \mathcal{D}f \mathcal{D}\phi \exp\left(-\frac{1}{\hbar}\int_0^{\hbar/(k_B T)} d\tau\, \mathcal{L}_{SYK}[f, \phi]\right)$$

after relating the boundary component of $g_{\mu\nu}$ to f,
and the boundary value of a_τ to ϕ.

Sachdev (2010); Kitaev (2015); Sachdev (2015); Bagrets, Altland, Kamenev (2016); Maldacena, Stanford, Yang (2016); Moitra, Trivedi, Vishal (2018) ; Gaikwad, Joshi, Mandal, Wadia (2018); Iliesiu, Turaci (2020)

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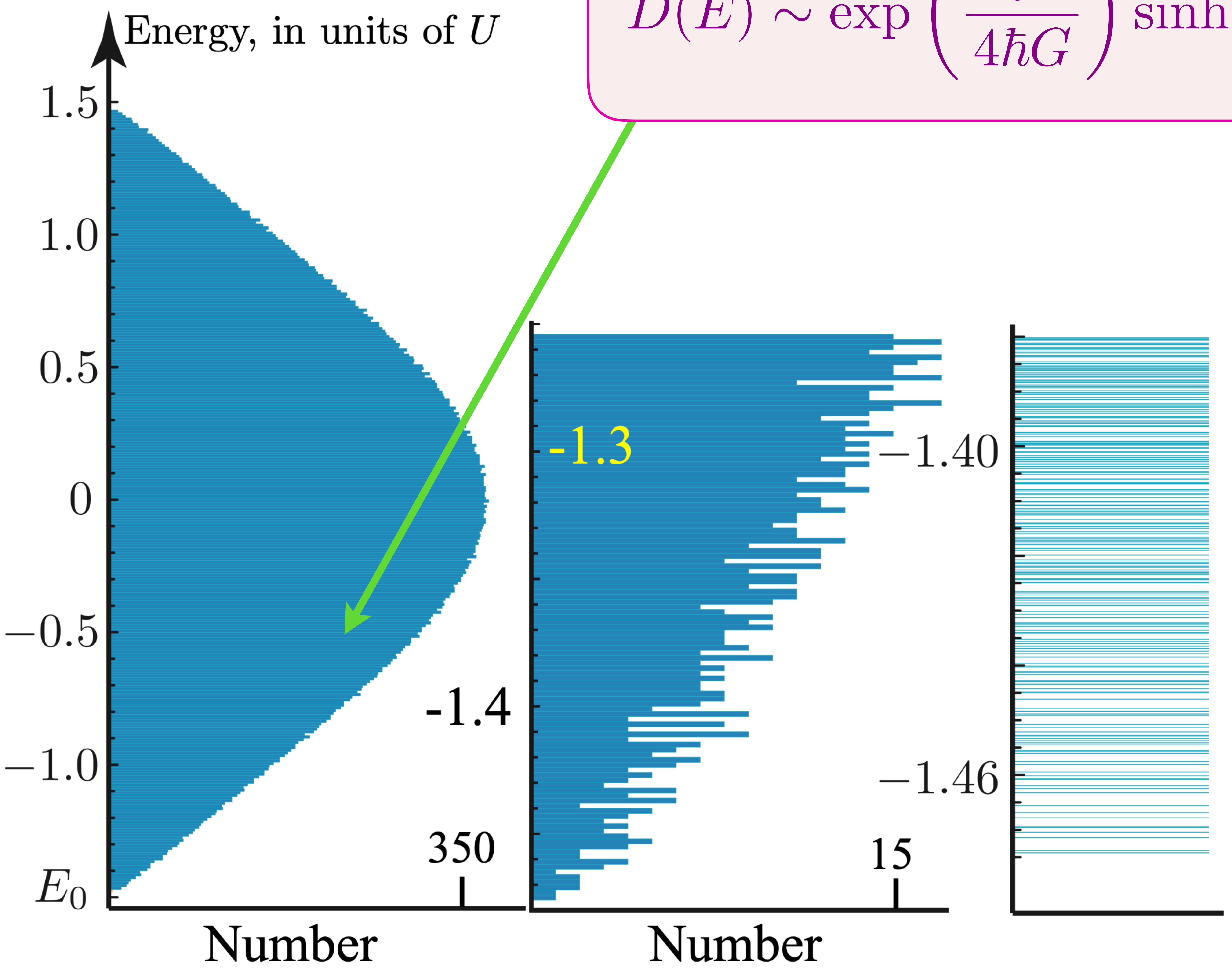
after relating the boundary component of $g_{\mu\nu}$ to f,
and the boundary value of a_τ to ϕ.

$$\frac{S(T)}{k_B} = \frac{1}{\hbar G}\left(\frac{A_0 c^3}{4} + \frac{\sqrt{\pi} A_0^{3/2} c^2}{2}\frac{k_B T}{\hbar}\right) - \frac{3}{2}\ln\left(\frac{\sqrt{c^5/\hbar G}}{k_B T/\hbar}\right) + \ldots .$$

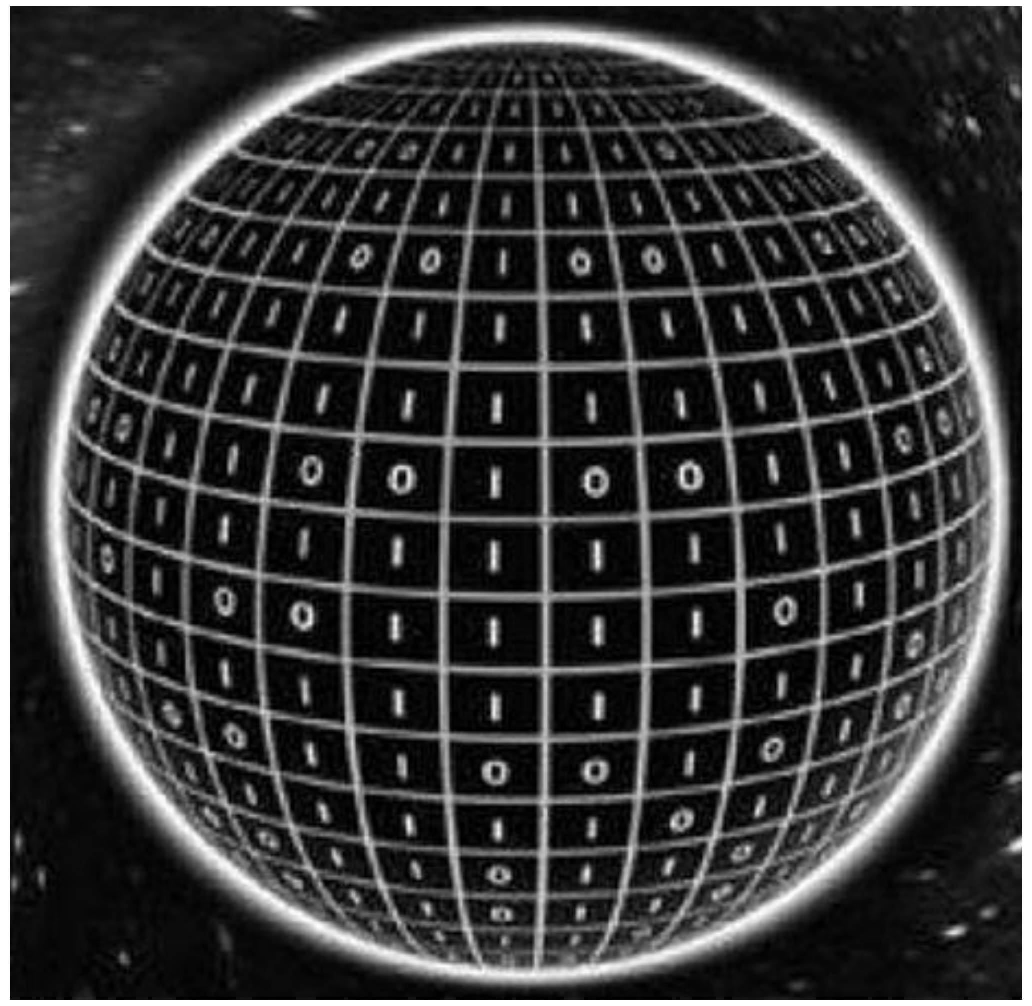
where A_0 is the area of the black hole horizon at $T = 0$.

Quantum theory of charged black holes

$$D(E) \sim \exp\left(\frac{A_0 c^3}{4\hbar G}\right) \sinh\left(\left[\frac{\sqrt{\pi} A_0^{3/2} c^2}{\hbar^2 G} E\right]^{1/2}\right)$$

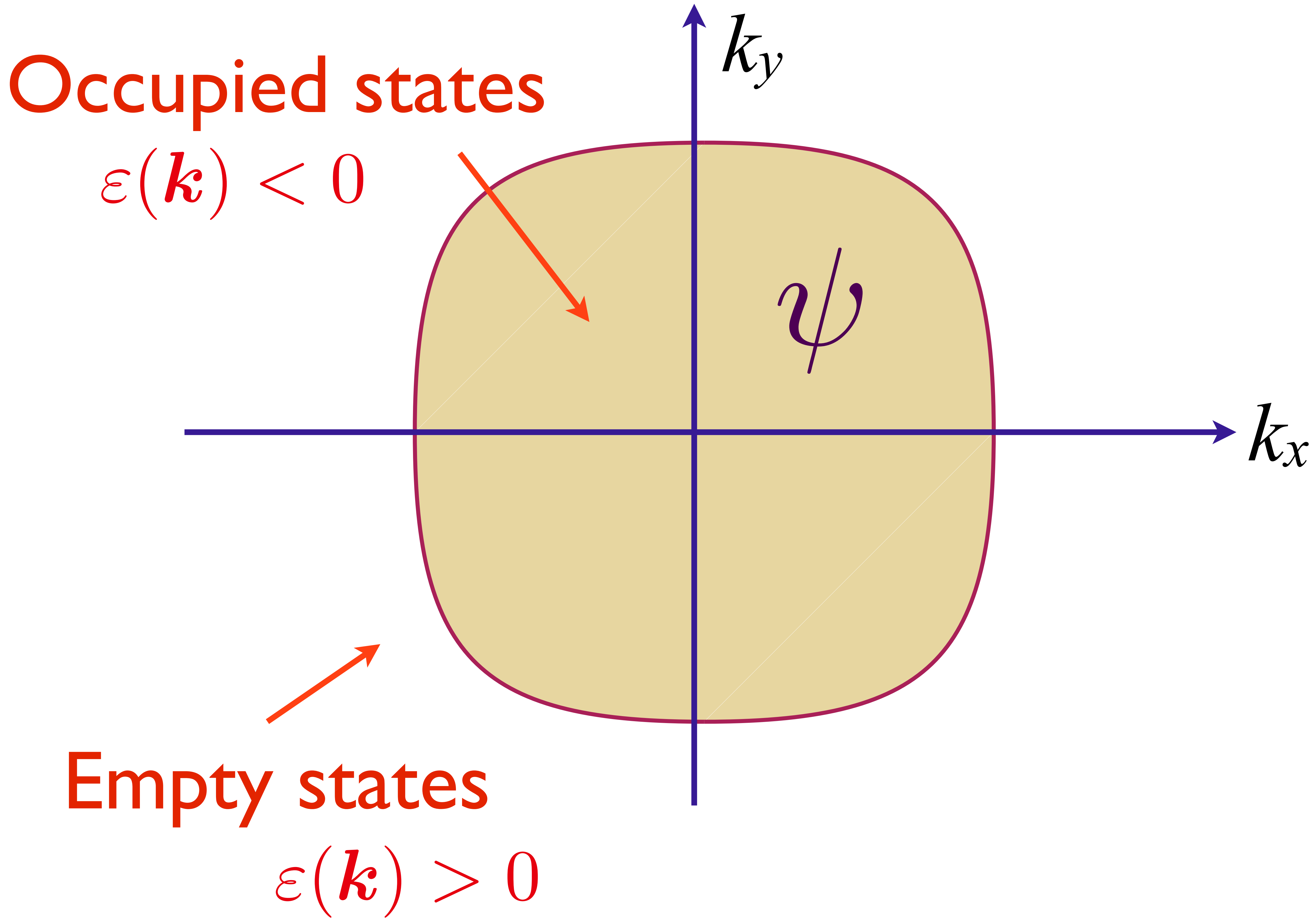


Same lower energy coarse-grained density of states in a model of interacting (fermionic) qubits with a discrete spectrum!



From the SYK model to linear-T resistivity

Fermi surface coupled to a critical boson



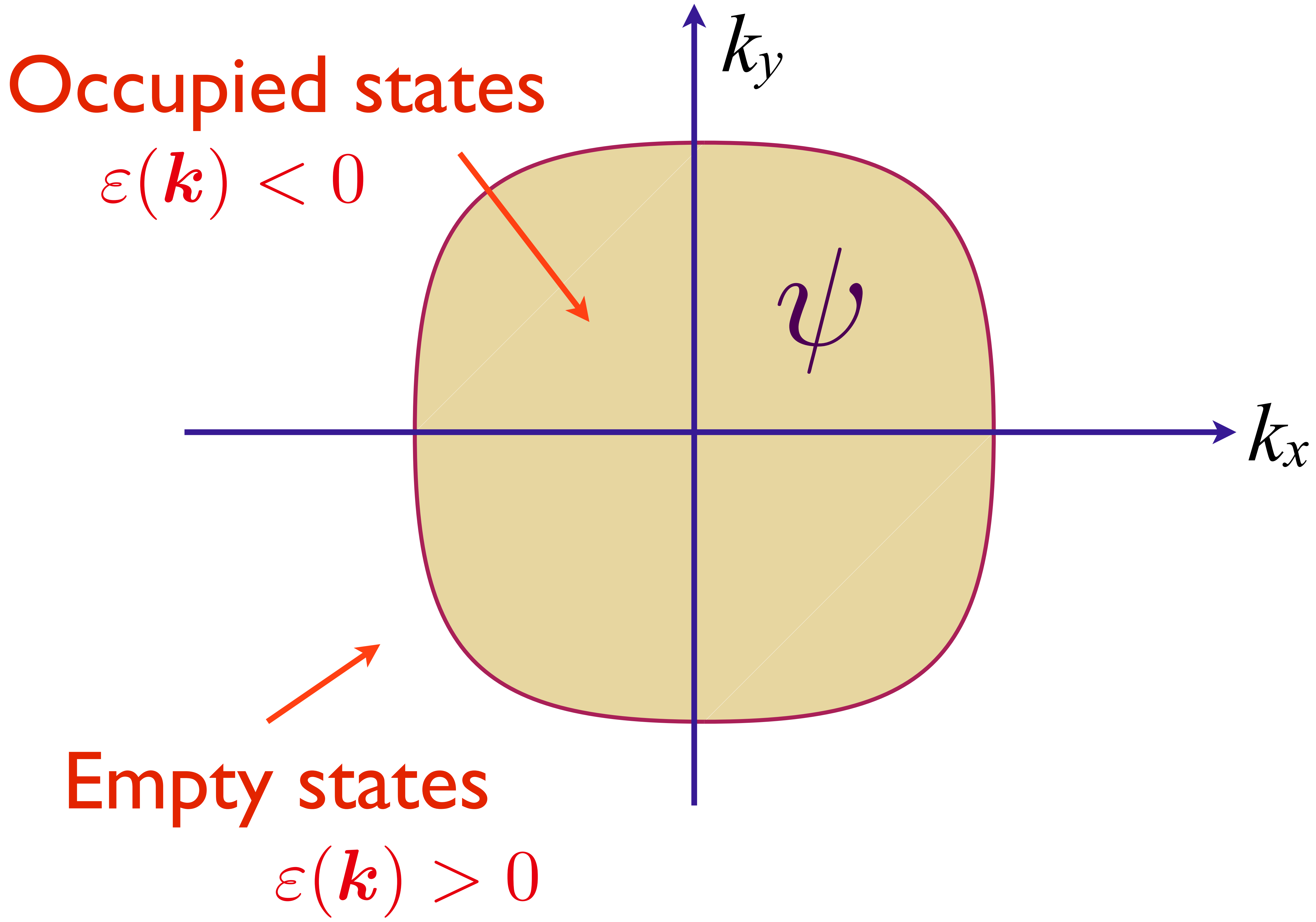
+

a critical boson

ϕ

- Nematic order
- Ferromagnetic order
- Transverse component of abelian or non-abelian gauge field

Fermi surface coupled to a critical boson



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"Yukawa" coupling: $g \int d^2r d\tau\, \psi^\dagger(r,\tau)\psi(r,\tau)\phi(r,\tau)$

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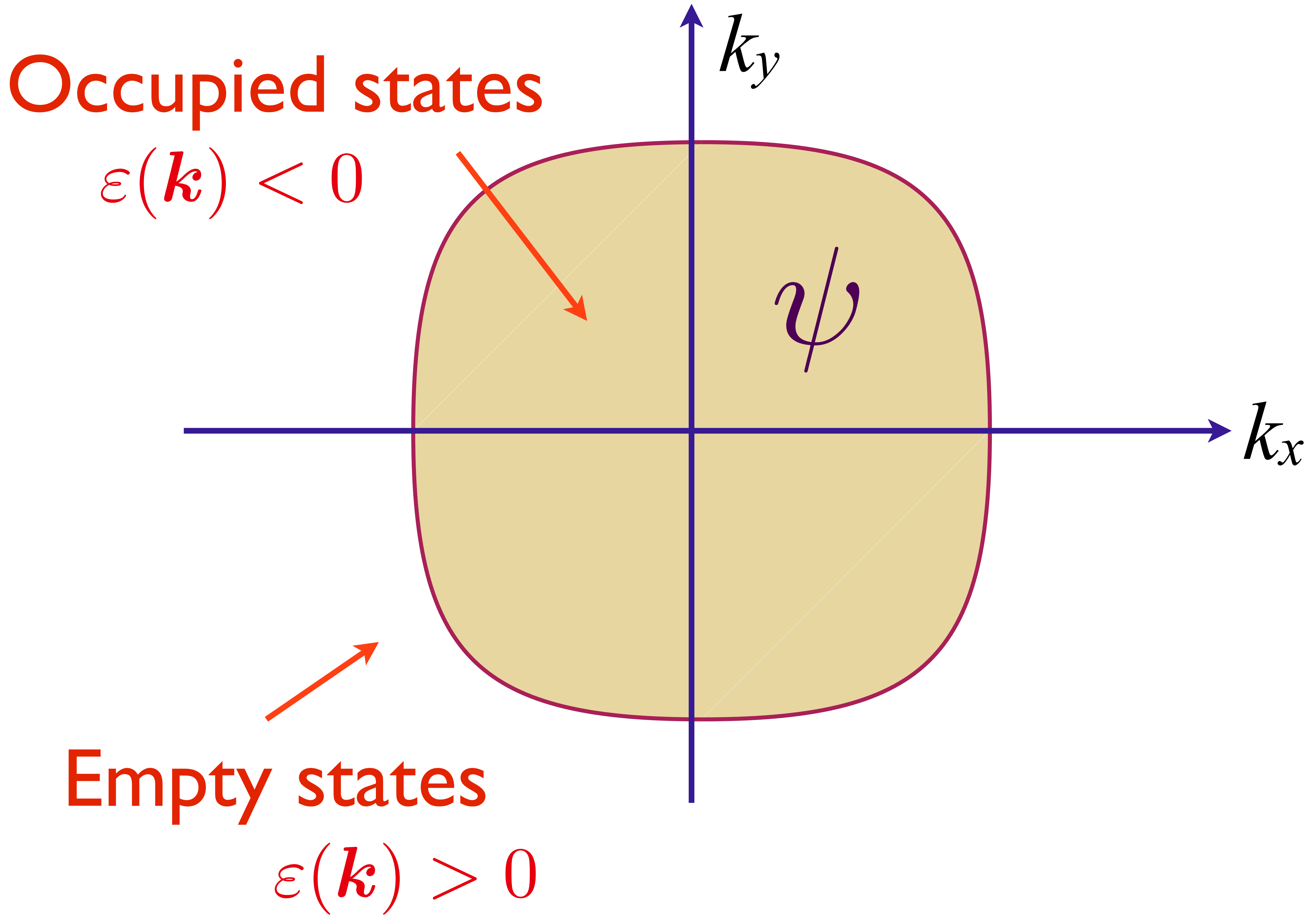

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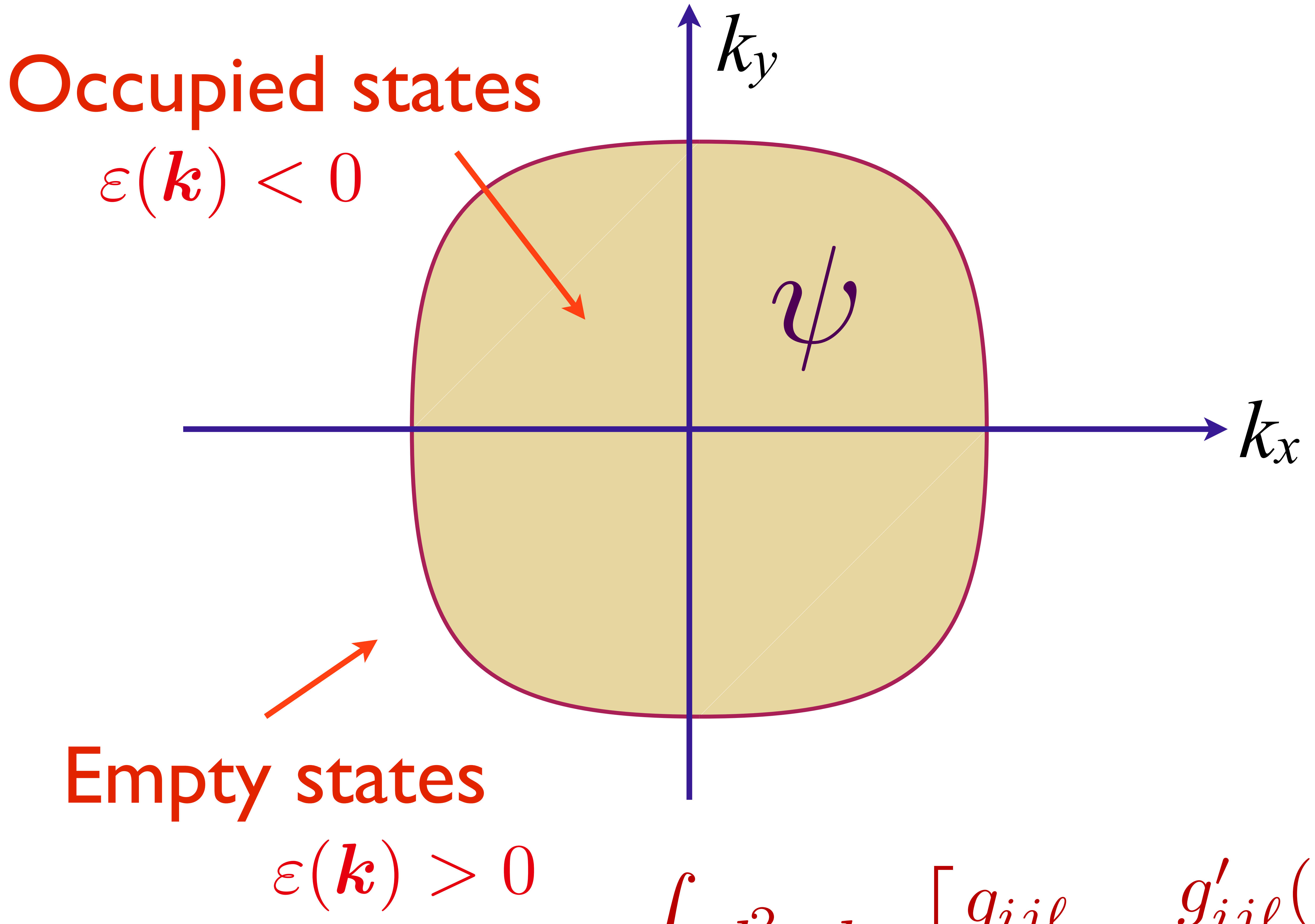
"Yukawa" coupling: $\displaystyle \frac{g_{ij\ell}}{N} \int d^2r d\tau \, \psi_i^\dagger(r,\tau)\psi_j(r,\tau)\phi_l(r,\tau)$

Random couplings in flavor space leads to large N theory of a strange metal, with zero resistivity

Fermi surface coupled to a critical boson



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$$\int d^2r d\tau \left[\frac{g_{ij\ell}}{N} + \frac{g'_{ij\ell}(r)}{N}\right] \psi_i^\dagger(r,\tau)\psi_j(r,\tau)\phi_l(r,\tau)$$

Random couplings in flavor *and* position space leads to large N theory of a strange metal, with linear-T resistivity

Aavishkar Patel, Haoyu Guo, Ilya Esterlis, S.S. arXiv: 2203.04990

Summary

- SYK: a solvable model without particle-like excitations, exhibiting thermalization and many-body chaos in a time of order $\hbar/(k_B T)$, independent of microscopic energy scales.

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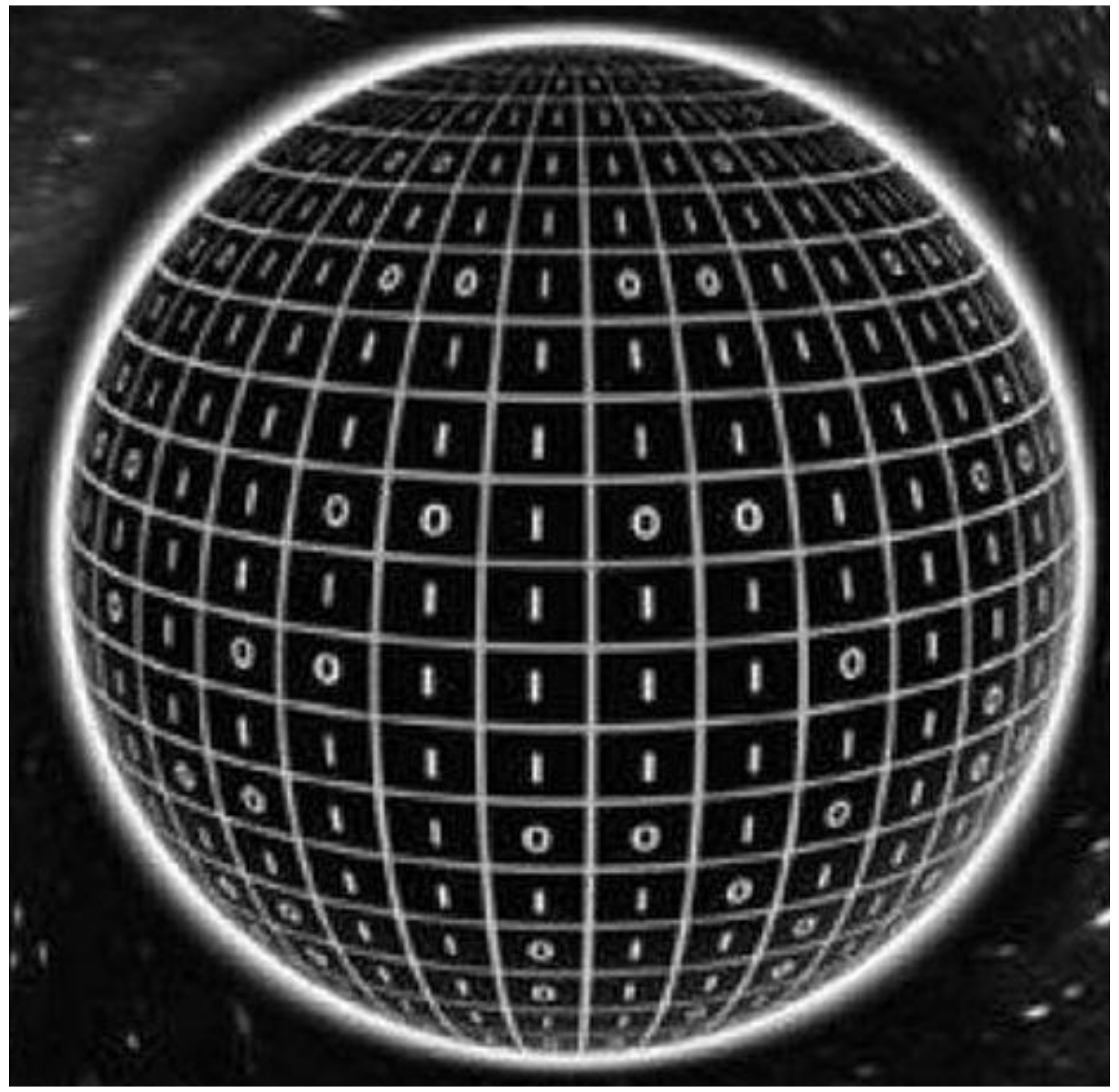
- Low energy theory of time reparameterizations is the theory of the boundary graviton in 1+1 dimensional quantum gravity on AdS_2.

Summary

- The semiclassical entropy of Einstein gravity is reproduced by a unitary quantum system with a discrete spectrum. Further work along these lines has led to progress on the Page curve describing the time evolution of the entropy of an evaporating black hole.



Summary

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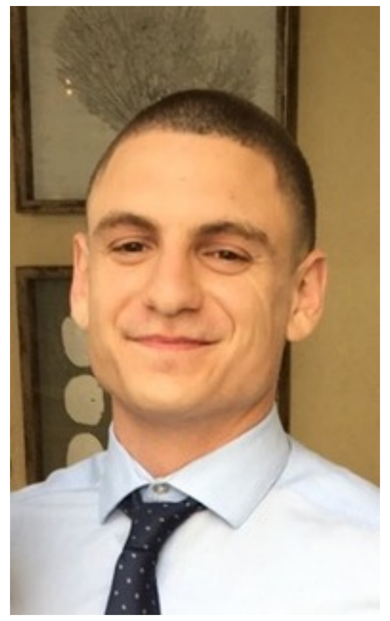
- Linear-T resistivity arises from spatially random interactions in a two-dimensional quantum-critical metal.





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