

# The cuprate phase diagram: theory of the pseudogap metal, d-wave superconductivity, and charge order

Hong Kong University  
March 14, 2023  
Subir Sachdev

Maine Christos, Zhu-Xi Luo, Henry Shackleton, Ya-Hui Zhang,  
Mathias Scheurer, and S. S., arXiv:2302.07885  
Alexander Nikolaenko, Jonas v. Milczewski, Darshan G. Joshi,  
and S.S., arXiv:2211.10452

Talk online: [sachdev.physics.harvard.edu](https://sachdev.physics.harvard.edu)

PHYSICS

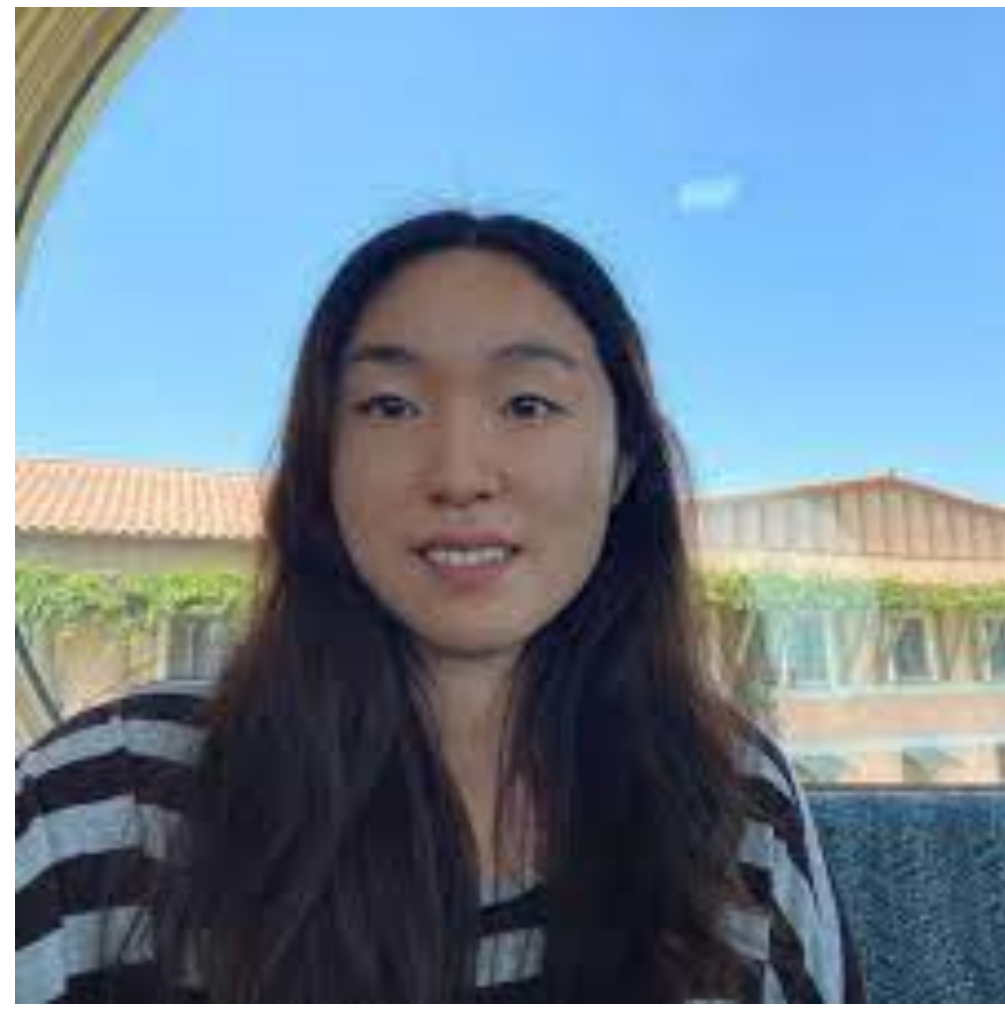


HARVARD





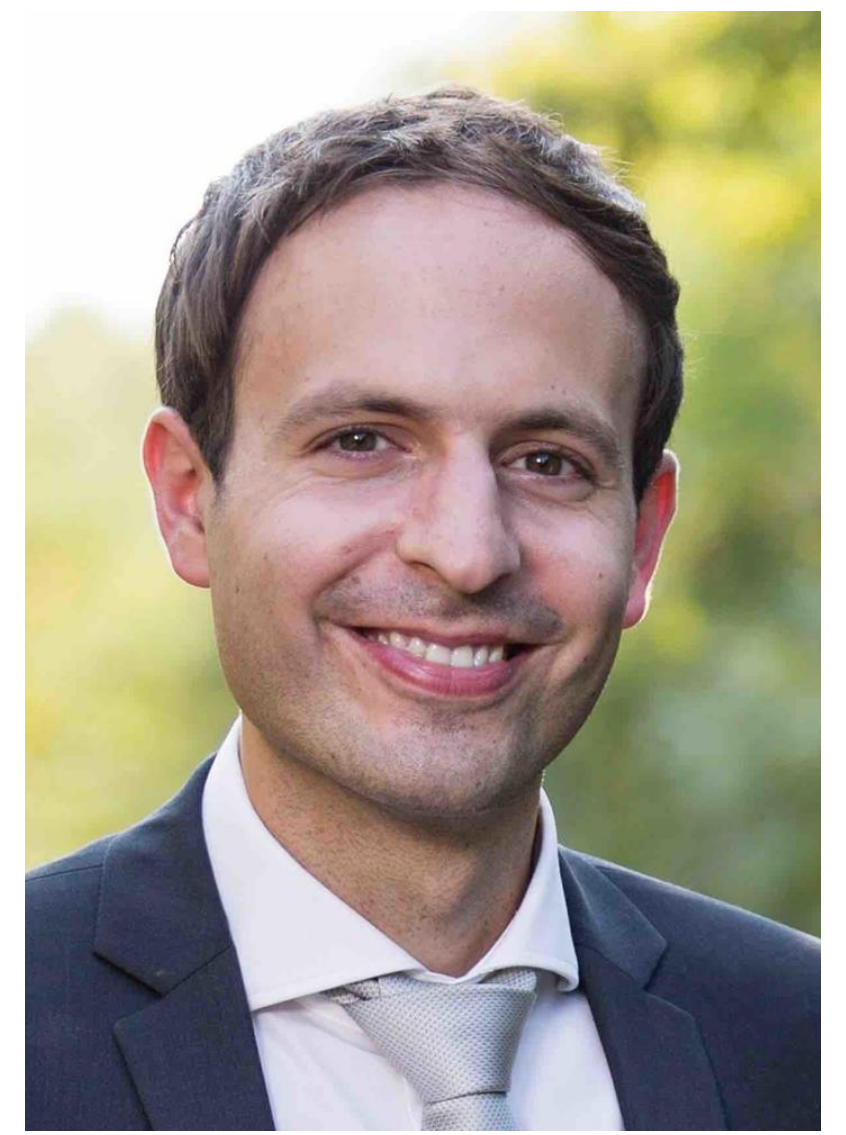
**Maine Christos**



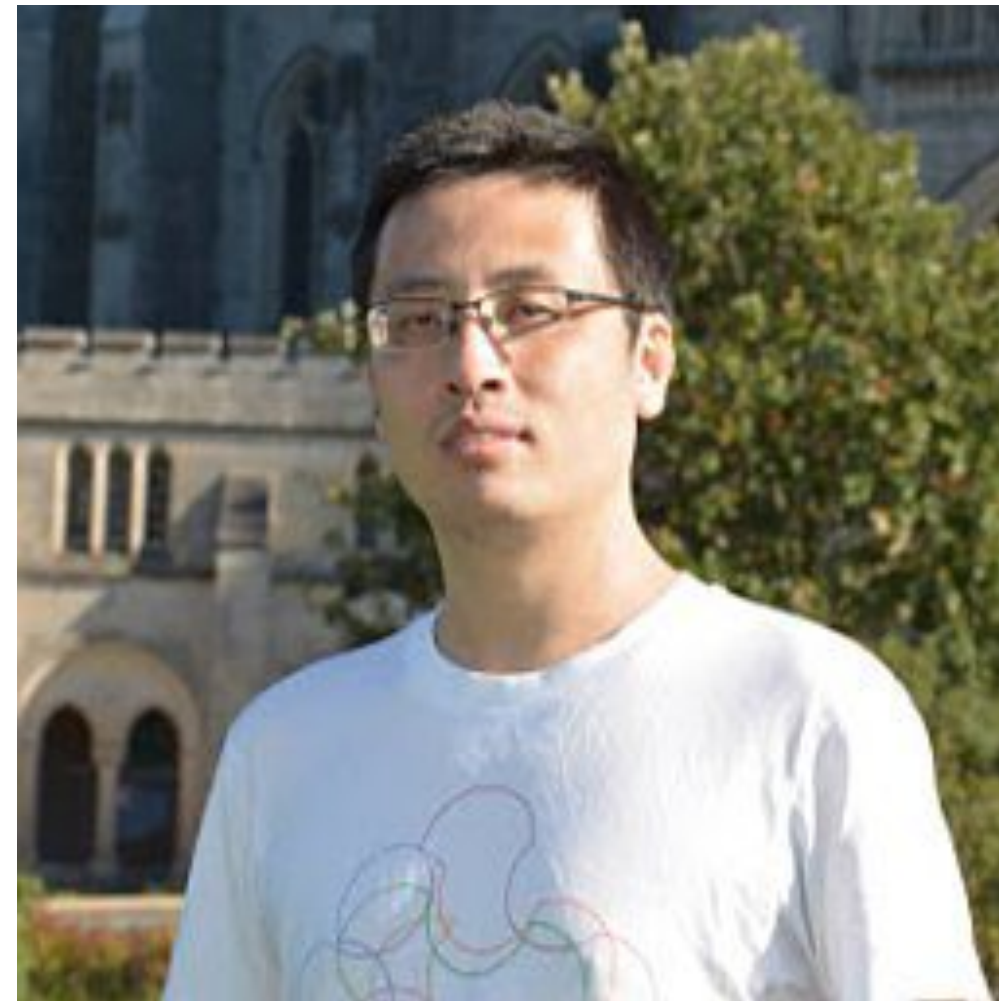
**Zhu-Xi Luo**



**Henry Shackleton**



**Mathias  
Scheurer**



**Ya-Hui  
Zhang**



**Alexander  
Nikolaenko**

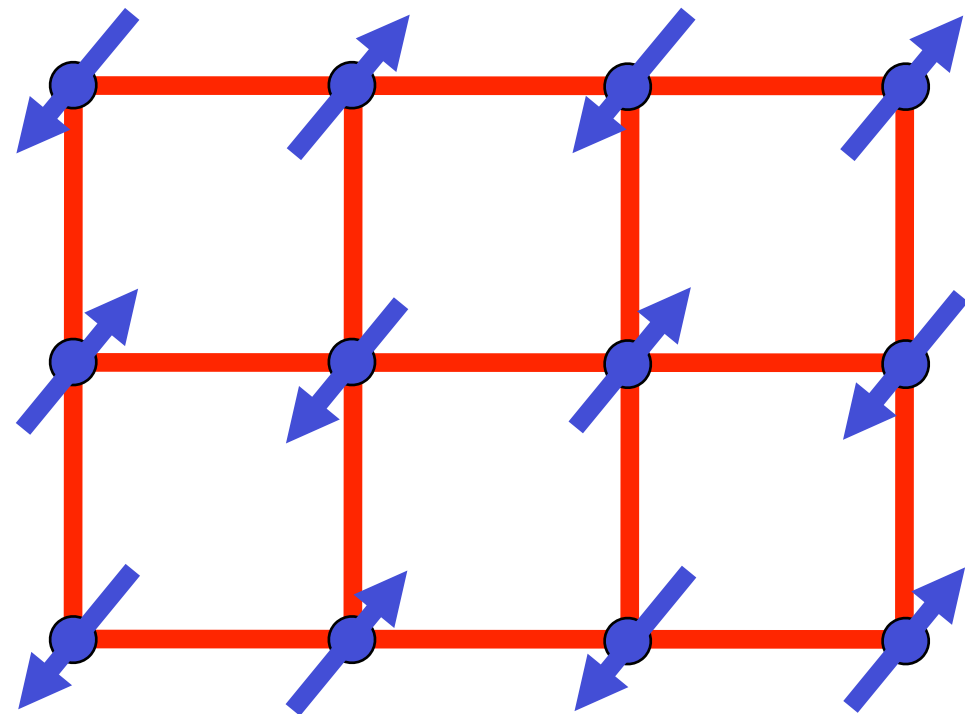


**Darshan Joshi**



**Jonas von Milczewski**

# Insulating $S=1/2$ antiferromagnet

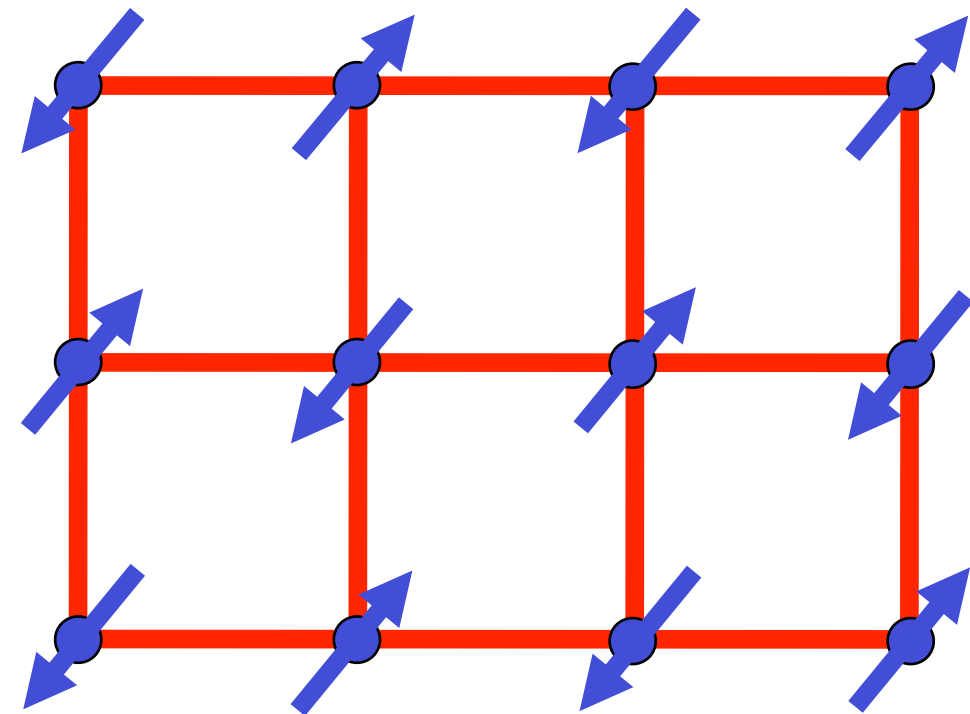


$$H = \sum_{i < j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

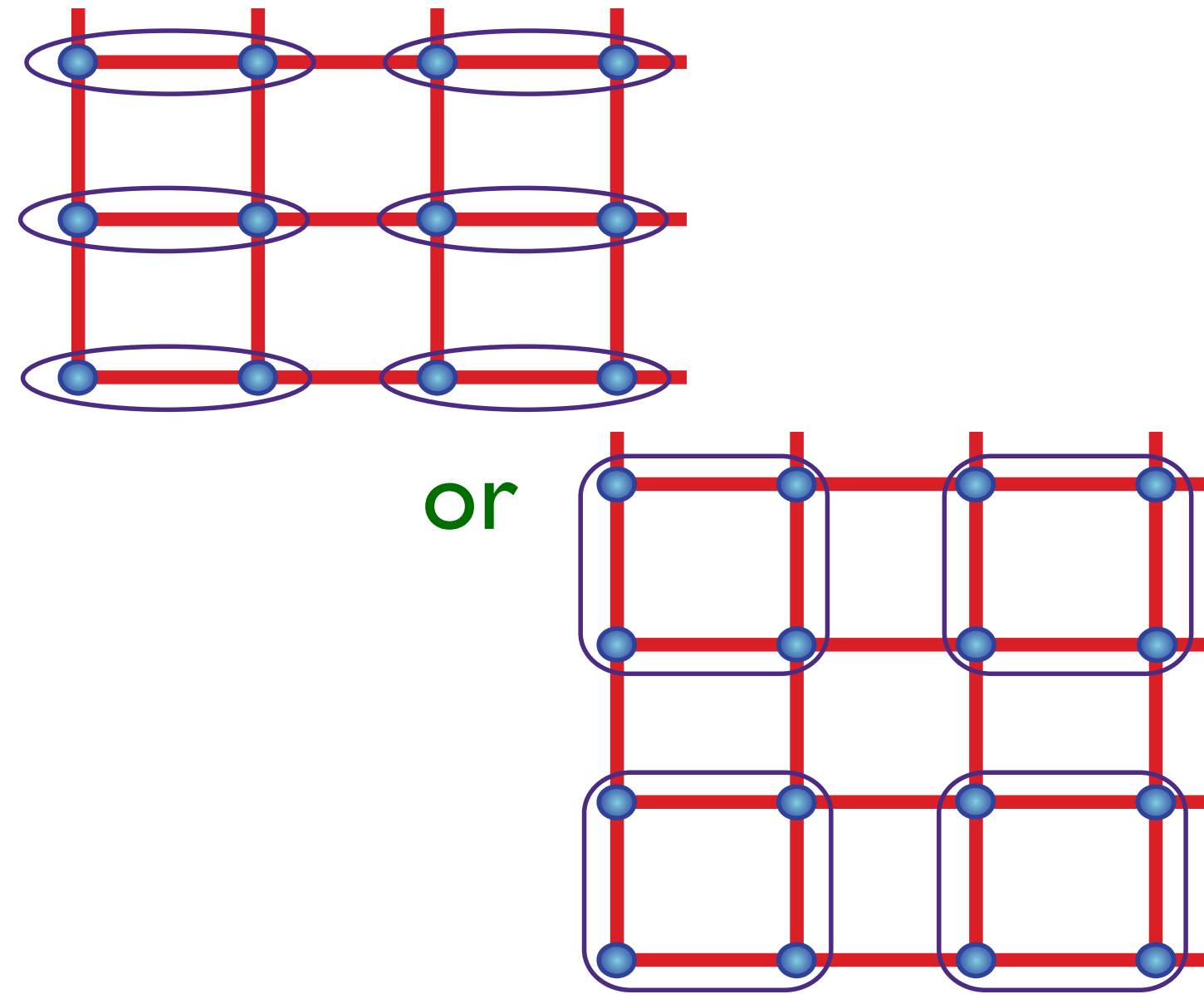
Schwinger bosons

$$\mathbf{S}_i = \frac{1}{2} b_{i\alpha}^\dagger \boldsymbol{\sigma}_{\alpha\beta} b_{i\beta}, \quad \sum_{\alpha=\uparrow,\downarrow} b_{i\alpha}^\dagger b_{i\alpha} = 1$$

# Insulating $S=1/2$ antiferromagnet



Higgs phase,  $\langle z_\alpha \rangle \neq 0$ :  
Néel order



Confining phase,  $\langle z_\alpha \rangle = 0$ :  
VBS order

$s$

$$H = \sum_{i < j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

Schwinger bosons

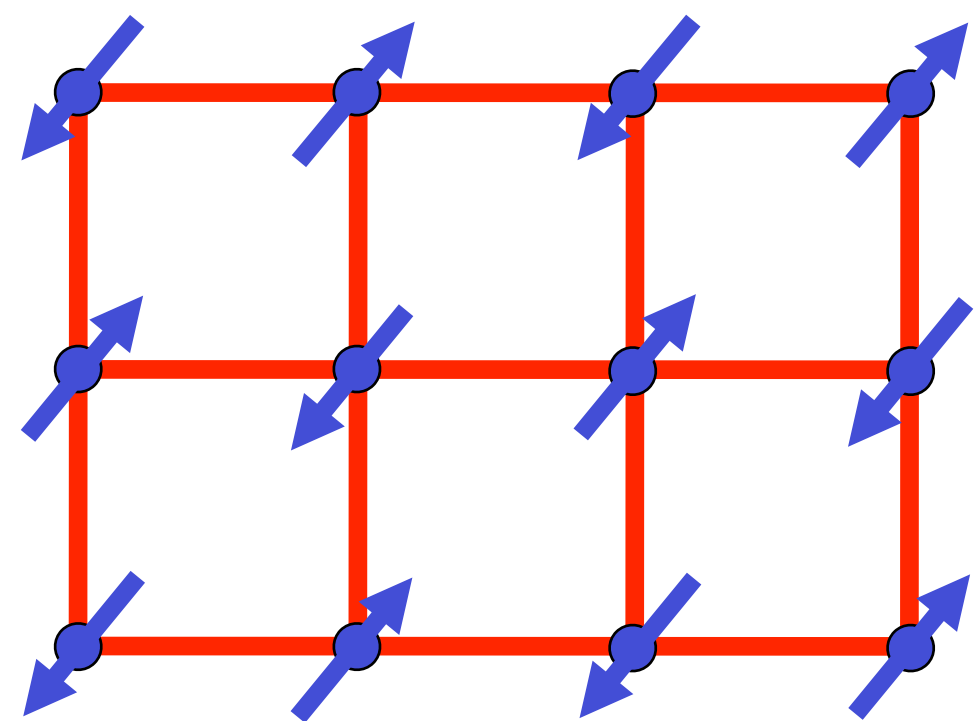
$$\mathbf{S}_i = \frac{1}{2} b_{i\alpha}^\dagger \boldsymbol{\sigma}_{\alpha\beta} b_{i\beta}, \quad \sum_{\alpha=\uparrow,\downarrow} b_{i\alpha}^\dagger b_{i\alpha} = 1$$

Low energy  $\mathbb{C}\mathbb{P}^1$  U(1) gauge theory

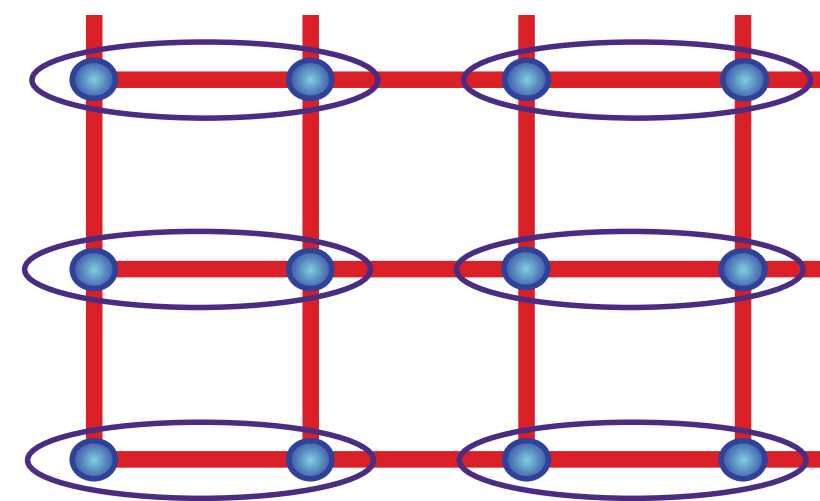
$$z_\alpha \sim b_{A\alpha} + \varepsilon_{\alpha\beta} b_{B\beta}$$

$$\mathcal{L} = |(\partial_\mu - ia_\mu)z_\alpha|^2 + s|z_\alpha|^2 + u|z_\alpha|^4 + \mathcal{L}_{\text{monopole}}$$

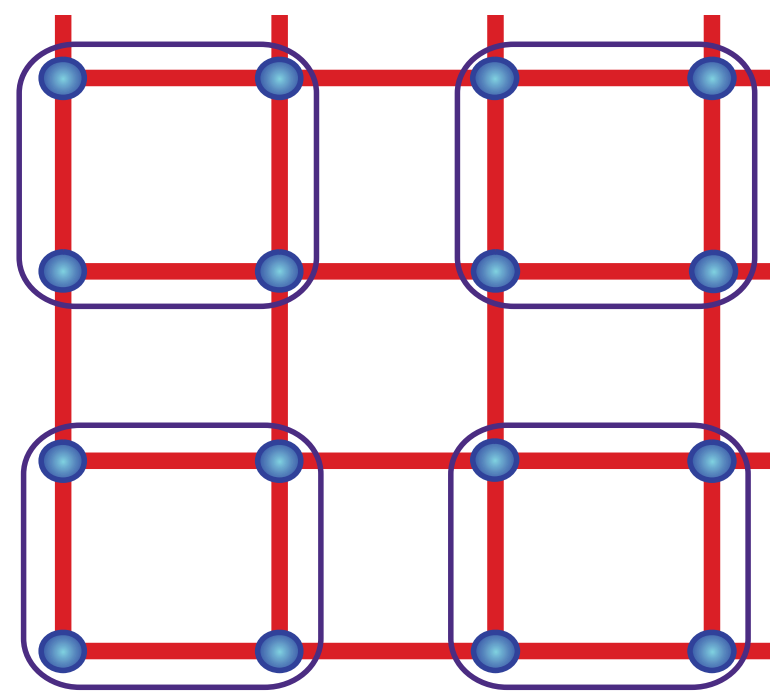
Insulating  $S=1/2$  antiferromagnet



Néel order



or



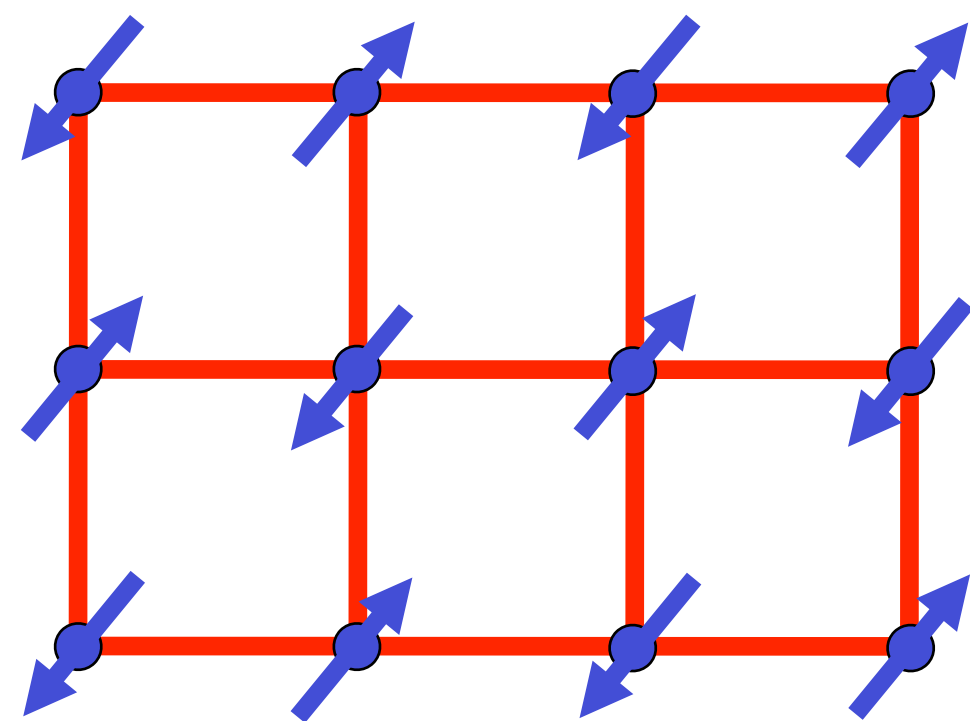
VBS order

$$H = \sum_{i < j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

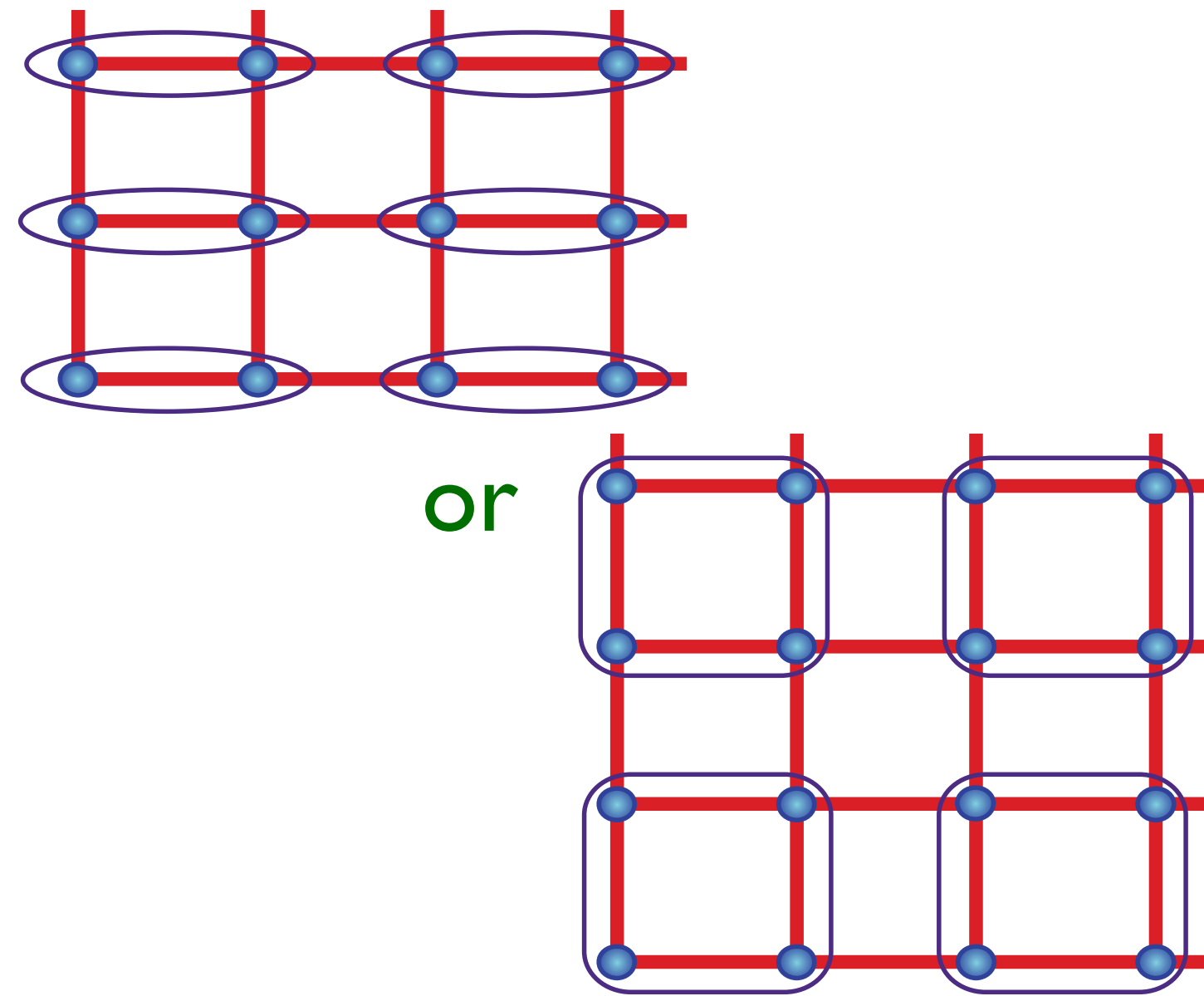
Schwinger fermions

$$\mathbf{S}_i = \frac{1}{2} f_{i\alpha}^\dagger \boldsymbol{\sigma}_{\alpha\beta} f_{i\beta}, \quad \sum_{\alpha=\uparrow,\downarrow} f_{i\alpha}^\dagger f_{i\alpha} = 1$$

# Insulating $S=1/2$ antiferromagnet



Confining phase:  
Néel order



Confining phase:  
VBS order

$$H = \sum_{i < j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

Schwinger fermions

$$\mathbf{S}_i = \frac{1}{2} f_{i\alpha}^\dagger \boldsymbol{\sigma}_{\alpha\beta} f_{i\beta}, \quad \sum_{\alpha=\uparrow,\downarrow} f_{i\alpha}^\dagger f_{i\alpha} = 1$$

$\pi$ -flux mean-field leads to a low energy theory of  $N_f = 2$  Dirac fermions  $\Psi_s$  coupled to an emergent SU(2) gauge field.

Dual to  $\mathbb{C}\mathbb{P}^1$  U(1) gauge theory!

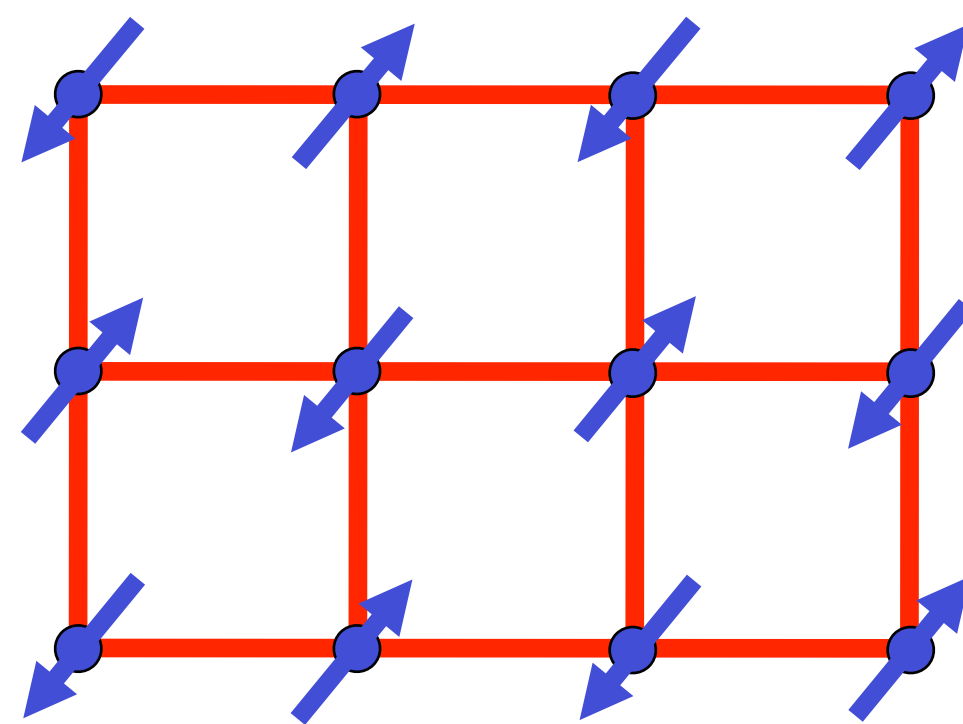
$$\mathcal{L} = i\bar{\Psi}_s \gamma_\mu D_\mu \Psi_s + \dots$$

Include charge fluctuations at half-filling: repulsive Hubbard model with pair-hopping

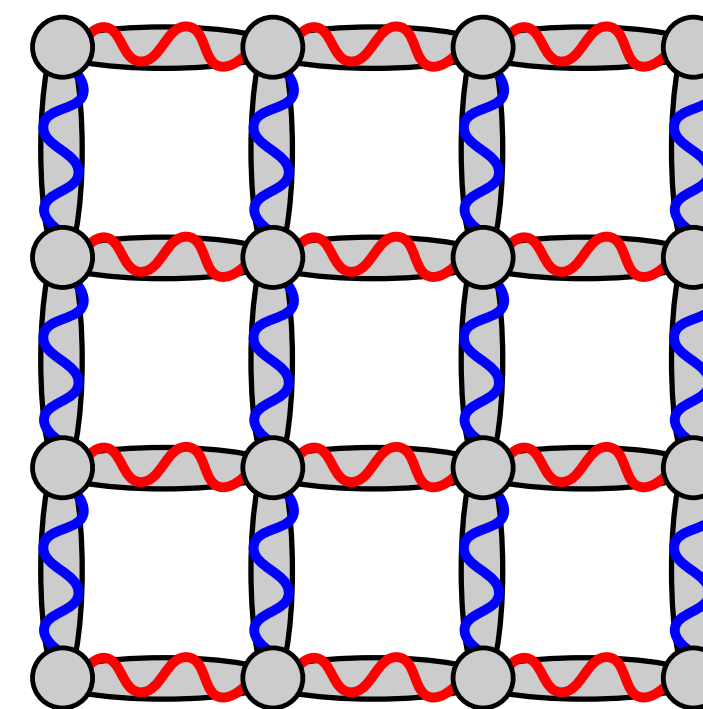
$$H = -\frac{t}{2} \sum_i K_i + U \sum_i (n_{i\uparrow} - \frac{1}{2})(n_{i\downarrow} - \frac{1}{2}) - W \sum_i K_i^2,$$

$$K_i = \sum_{\hat{e}=\pm\hat{x},\pm\hat{y}} \left( c_{i\alpha}^\dagger c_{i+\hat{e},\alpha} + c_{i+\hat{e},\alpha}^\dagger c_{i\alpha} \right)$$

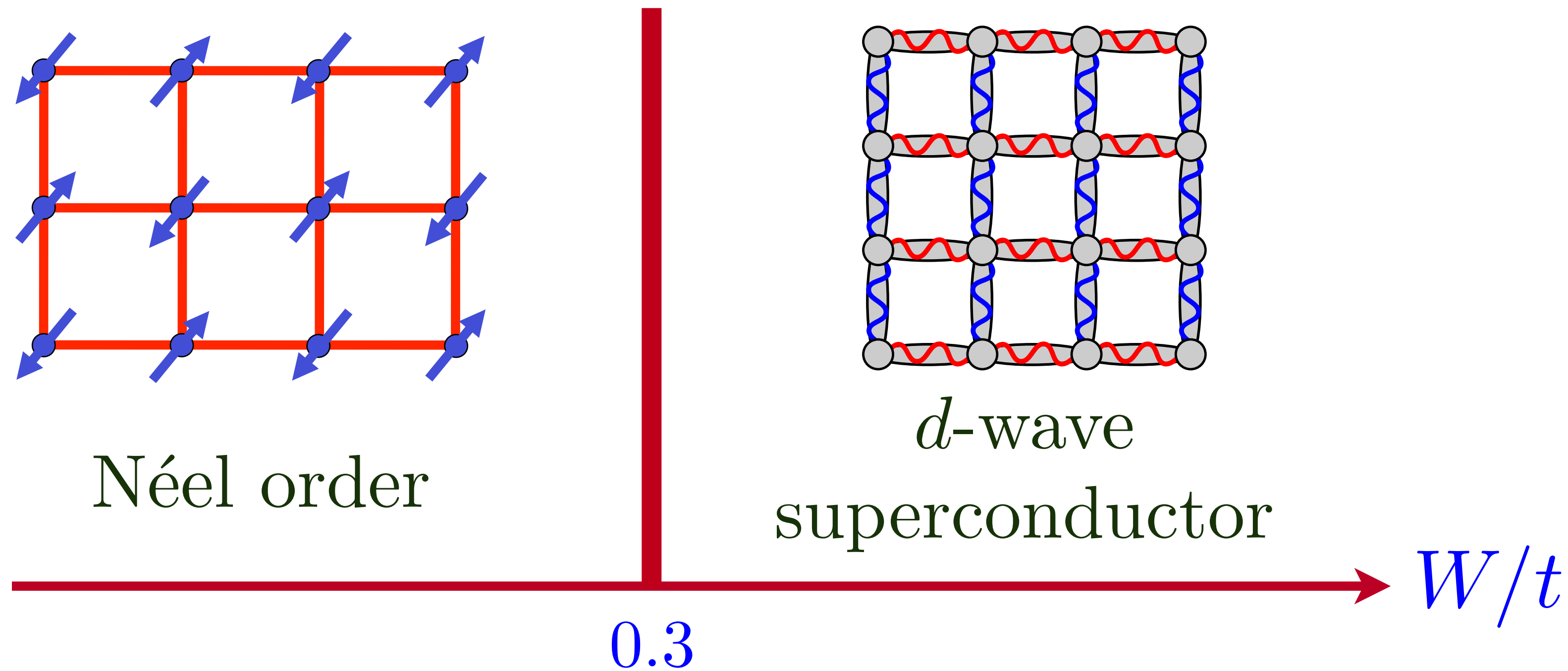
$$U/t = 4$$



Néel order



*d*-wave  
superconductor





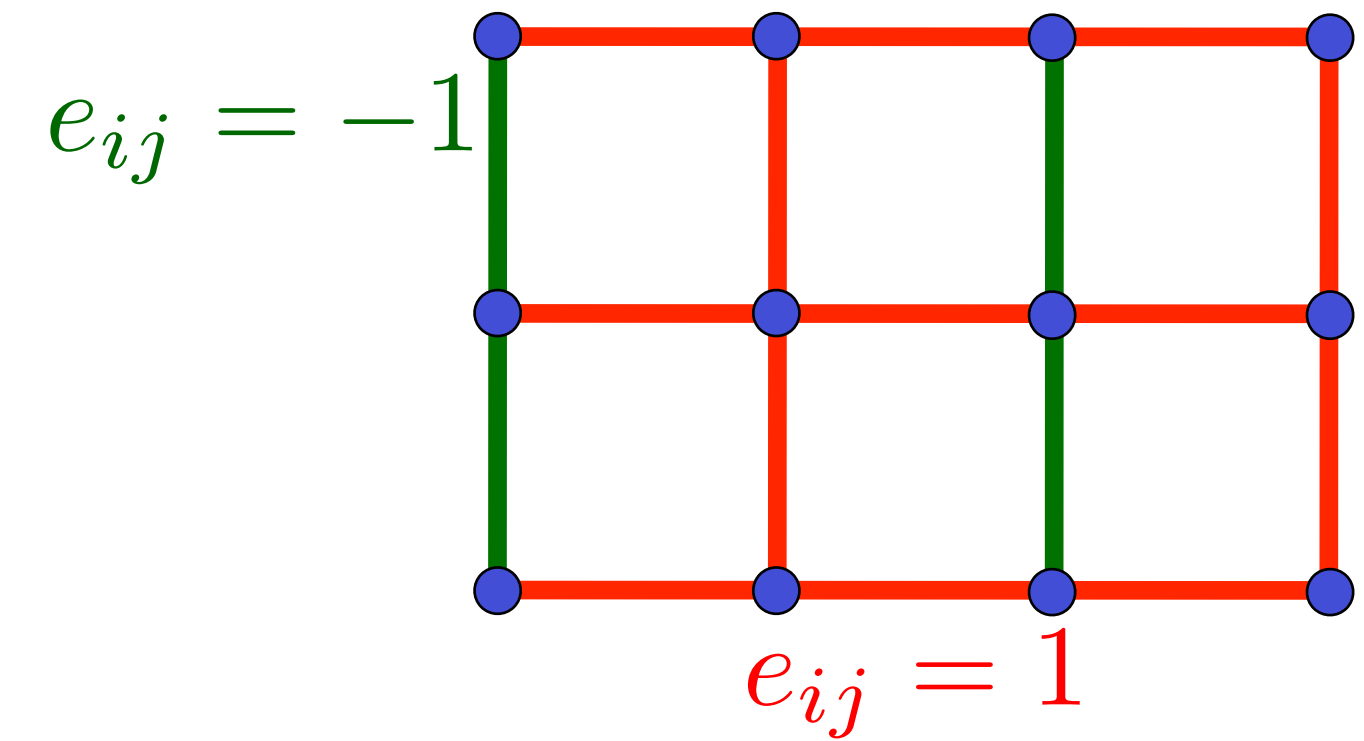


# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

- Begin with the  $\pi$ -flux spin liquid in the fermionic spinon description.

$$H_f = iJ \sum_{\langle ij \rangle} e_{ij} \left( f_{i\alpha}^\dagger f_{j\alpha} - f_{j\alpha}^\dagger f_{i\alpha} \right) = iJ \sum_{\langle ij \rangle} e_{ij} \left( \Psi_i^\dagger \Psi_j - \Psi_j^\dagger \Psi_i \right); \quad \Psi_i = \begin{pmatrix} f_{i\uparrow} \\ f_{i\downarrow} \end{pmatrix}$$

$H_f$  is invariant under distinct SU(2) rotations in spin and Nambu space.

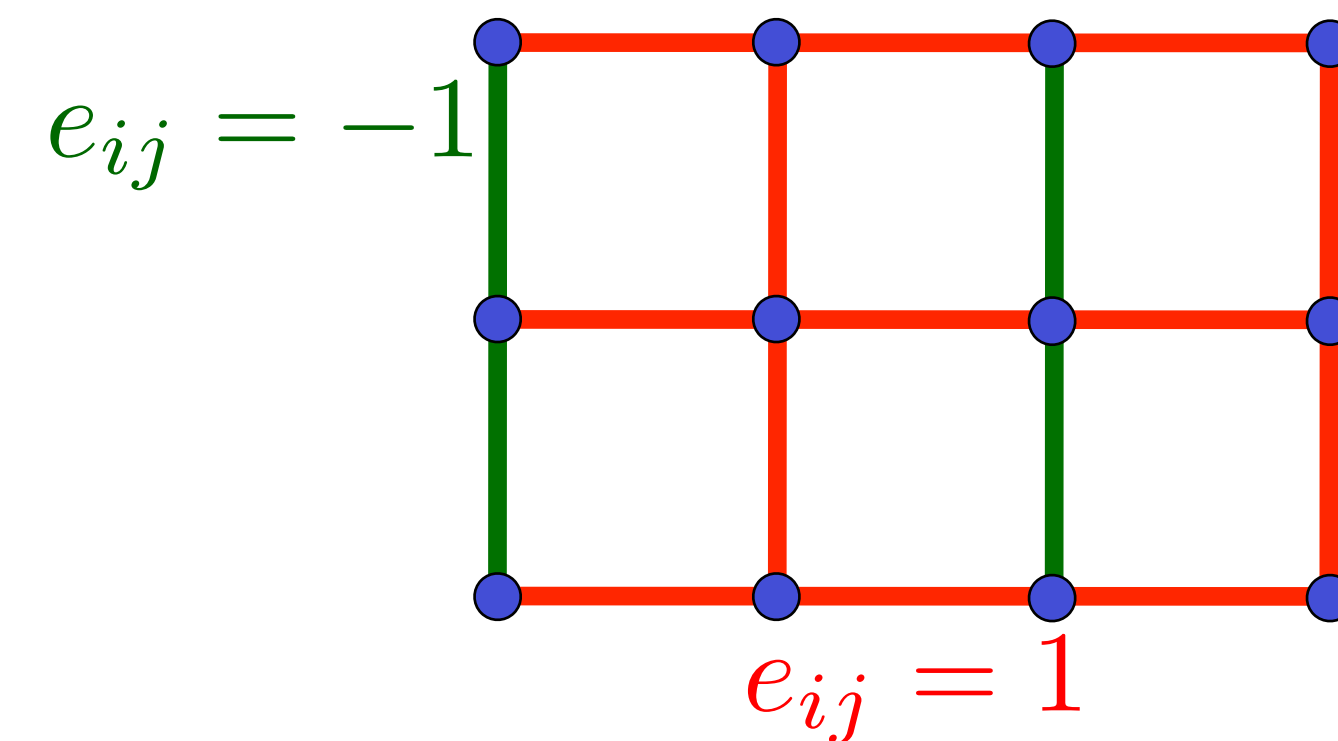


- We can fully confine the SU(2) gauge field by condensing a boson,  $B_i$ , which is a fundamental of gauge SU(2). To obtain superconductivity with charge  $2e$  pairs in the confining phase,  $B_i$  should also carry electromagnetic charge  $e$ .

# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

- Begin with the  $\pi$ -flux spin liquid in the fermionic spinon description.

$$H_f = iJ \sum_{\langle ij \rangle} e_{ij} \left( f_{i\alpha}^\dagger f_{j\alpha} - f_{j\alpha}^\dagger f_{i\alpha} \right) = iJ \sum_{\langle ij \rangle} e_{ij} \left( \Psi_i^\dagger \Psi_j - \Psi_j^\dagger \Psi_i \right); \quad \Psi_i = \begin{pmatrix} f_{i\uparrow} \\ f_{i\downarrow} \end{pmatrix}$$



$H_f$  is invariant under distinct SU(2) rotations in spin and Nambu space.

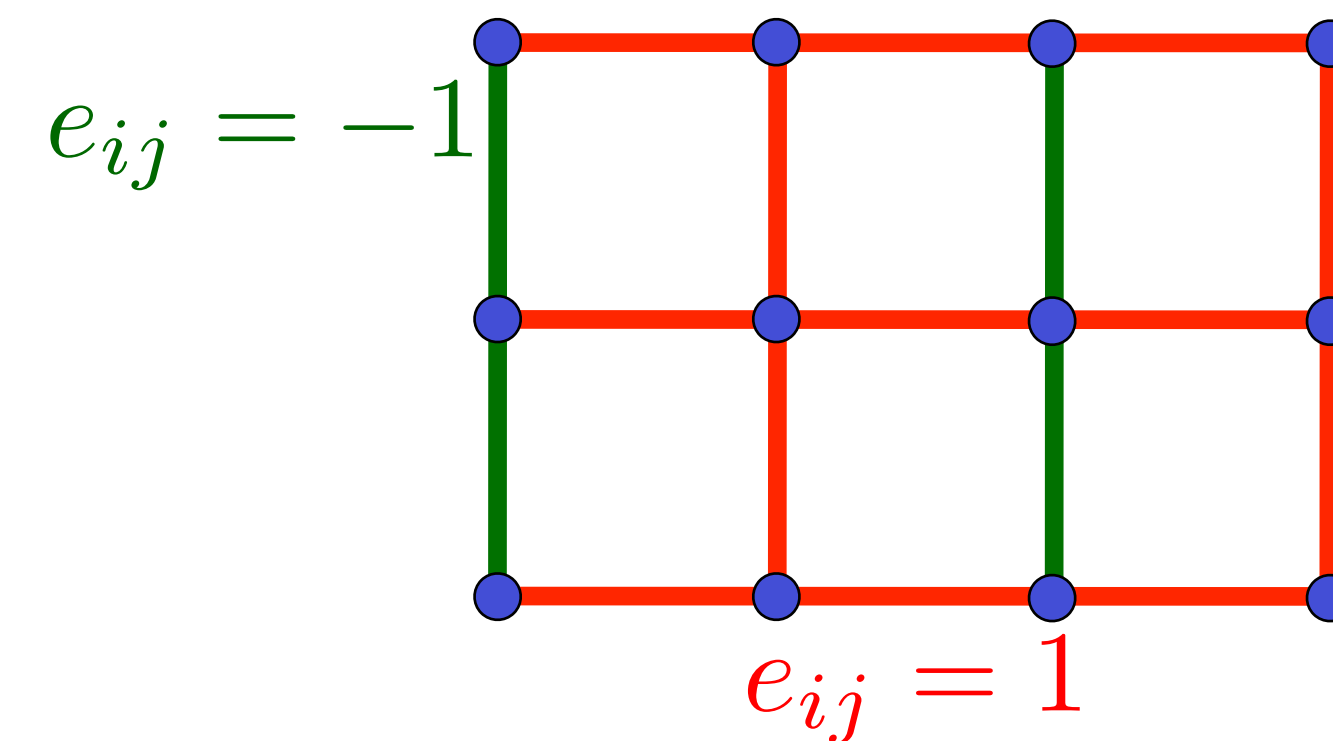
- We can fully confine the SU(2) gauge field by condensing a boson,  $B_i$ , which is a fundamental of gauge SU(2). To obtain superconductivity with charge  $2e$  pairs in the confining phase,  $B_i$  should also carry electromagnetic charge  $e$ .
- This uniquely identifies  $B_i$  as the ‘chargon’ of X.-G. Wen and P.A. Lee, PRL **76**, 503 (1996), related to the electrons  $c_{i\sigma}$  by

$$B_i = \begin{pmatrix} B_{1i} \\ B_{2i} \end{pmatrix} ; \quad \begin{pmatrix} c_{i\uparrow} \\ c_{i\downarrow}^\dagger \end{pmatrix} \sim \begin{pmatrix} B_{1i}^* & B_{2i}^* \\ -B_{2i} & B_{1i} \end{pmatrix} \begin{pmatrix} f_{i\uparrow} \\ f_{i\downarrow}^\dagger \end{pmatrix}$$

# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

- Begin with the  $\pi$ -flux spin liquid in the fermionic spinon description.

$$H_f = iJ \sum_{\langle ij \rangle} e_{ij} \left( f_{i\alpha}^\dagger f_{j\alpha} - f_{j\alpha}^\dagger f_{i\alpha} \right) = iJ \sum_{\langle ij \rangle} e_{ij} \left( \Psi_i^\dagger \Psi_j - \Psi_j^\dagger \Psi_i \right); \quad \Psi_i = \begin{pmatrix} f_{i\uparrow} \\ f_{i\downarrow} \end{pmatrix}$$



$H_f$  is invariant under distinct SU(2) rotations in spin and Nambu space.

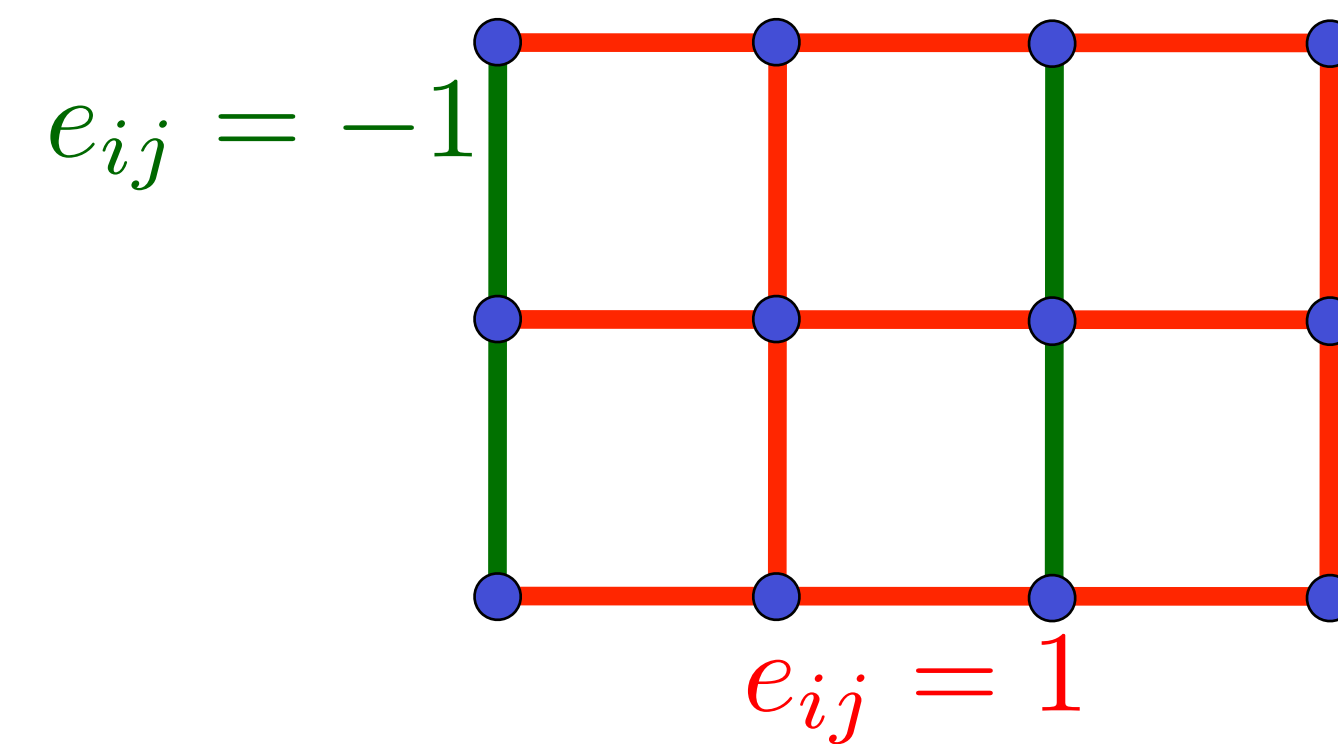
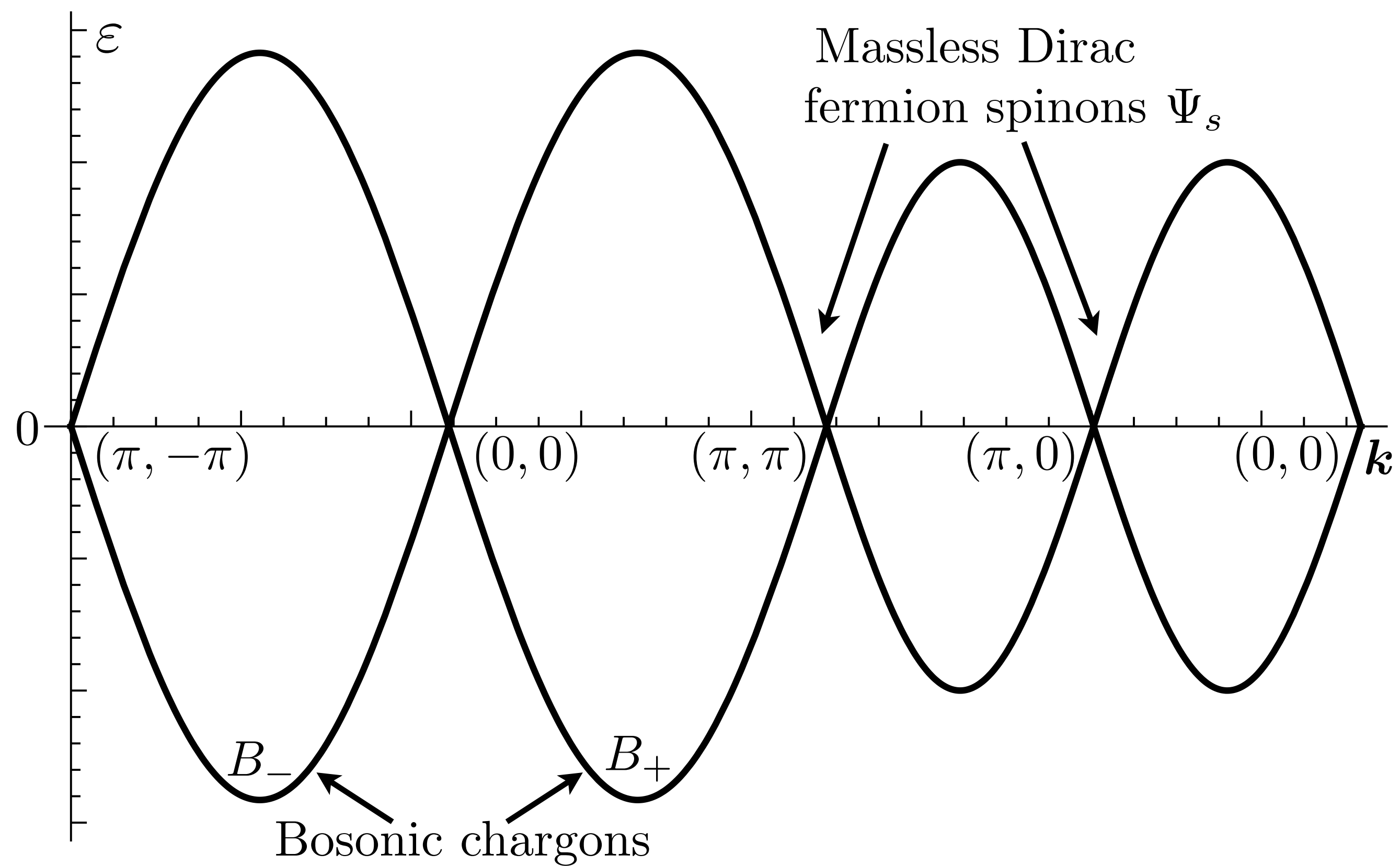
- We can fully confine the SU(2) gauge field by condensing a boson,  $B_i$ , which is a fundamental of gauge SU(2). To obtain superconductivity with charge  $2e$  pairs in the confining phase,  $B_i$  should also carry electromagnetic charge  $e$ .
- This uniquely identifies  $B_i$  as the ‘chargon’ of X.-G. Wen and P.A. Lee, PRL **76**, 503 (1996), related to the electrons  $c_{i\sigma}$  by

$$B_i = \begin{pmatrix} B_{1i} \\ B_{2i} \end{pmatrix}; \quad \begin{pmatrix} c_{i\uparrow} \\ c_{i\downarrow}^\dagger \end{pmatrix} \sim \begin{pmatrix} B_{1i}^* & B_{2i}^* \\ -B_{2i} & B_{1i} \end{pmatrix} \begin{pmatrix} f_{i\uparrow} \\ f_{i\downarrow}^\dagger \end{pmatrix}$$

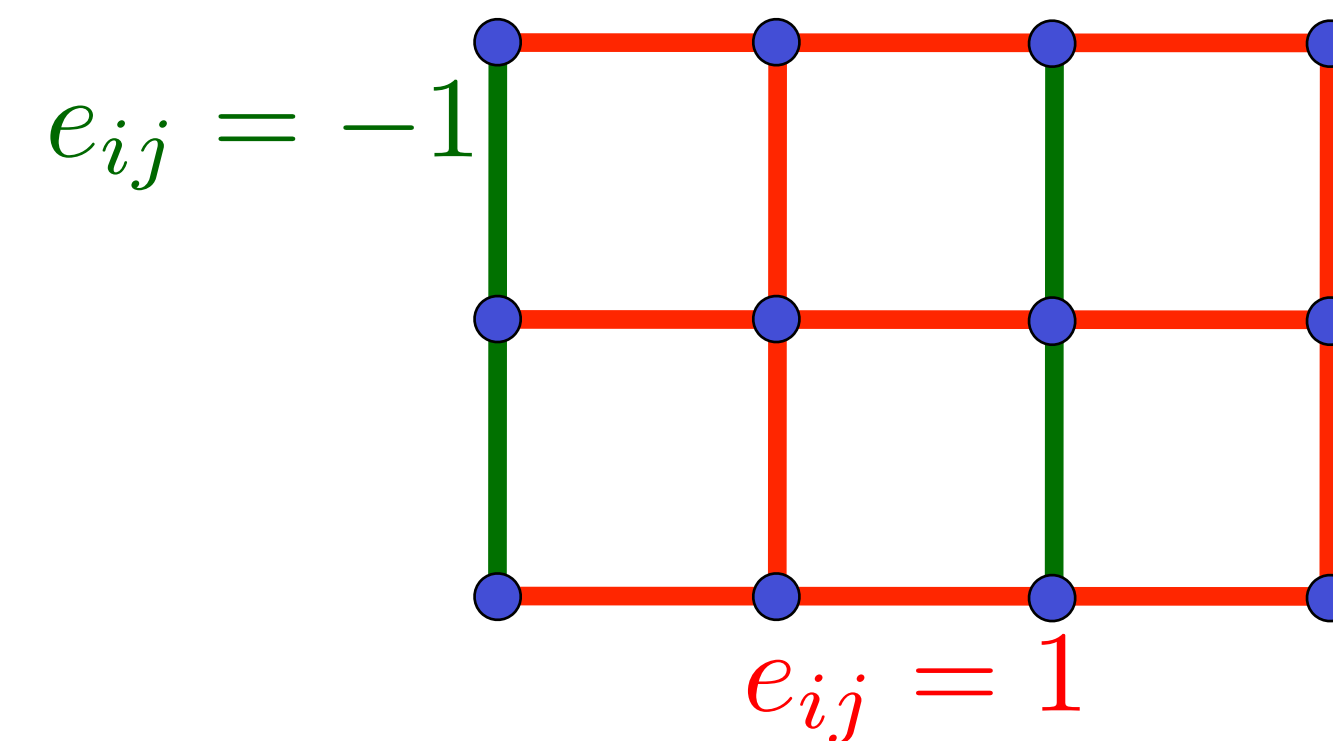
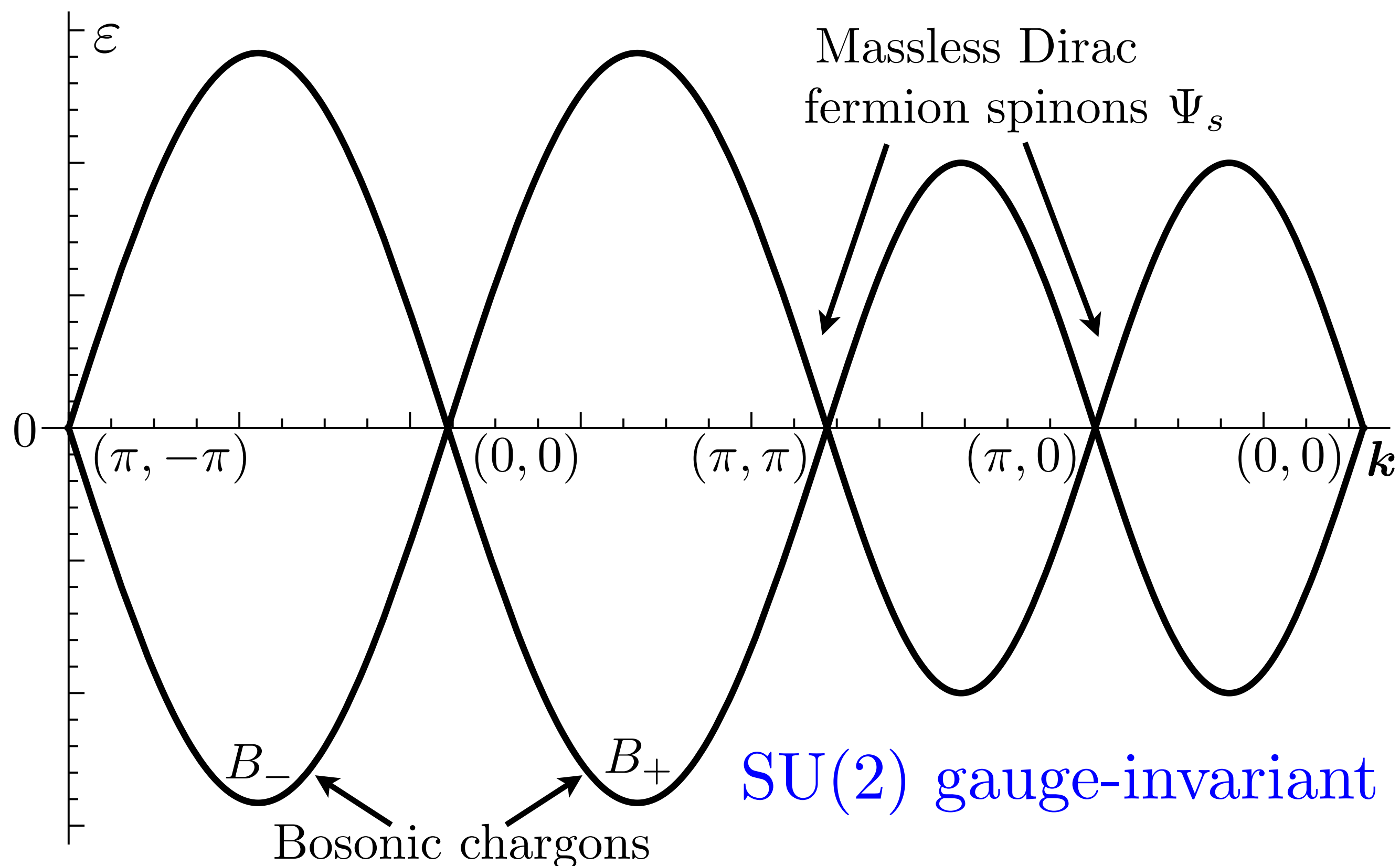
- Knowing the projective symmetry transformations of  $\Psi_i$ , we can deduce those of the  $B_i$ , and obtain the effective Hamiltonian for  $B_i$

$$H_B = r \sum_i B_i^\dagger B_i + iw \sum_{\langle ij \rangle} e_{ij} \left( B_i^\dagger B_j - B_j^\dagger B_i \right) + \dots$$

# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



SU(2) gauge-invariant order parameters of Higgs phases:

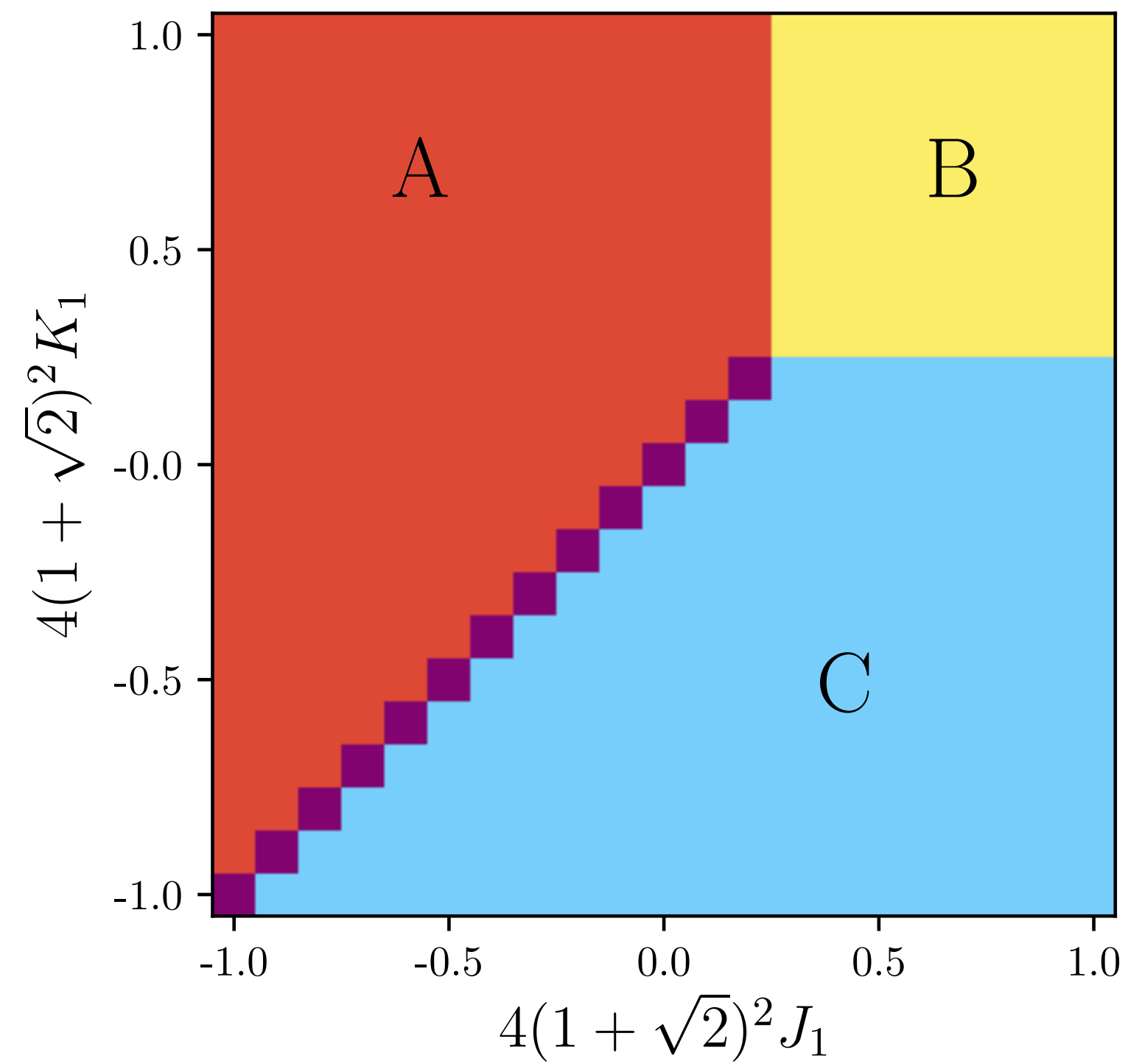
$$x\text{-CDW} : \rho_{(\pi,0)} = B_{a+}^* B_{a+} - B_{a-}^* B_{a-}$$

$$y\text{-CDW} : \rho_{(0,\pi)} = B_{a+}^* B_{a-} + B_{a-}^* B_{a+}$$

$$d\text{-density wave} : D = i (B_{a+}^* B_{a-} - B_{a-}^* B_{a+})$$

$$d\text{-wave superconductor} : \Delta = \varepsilon_{ab} B_{a+} B_{b-}$$

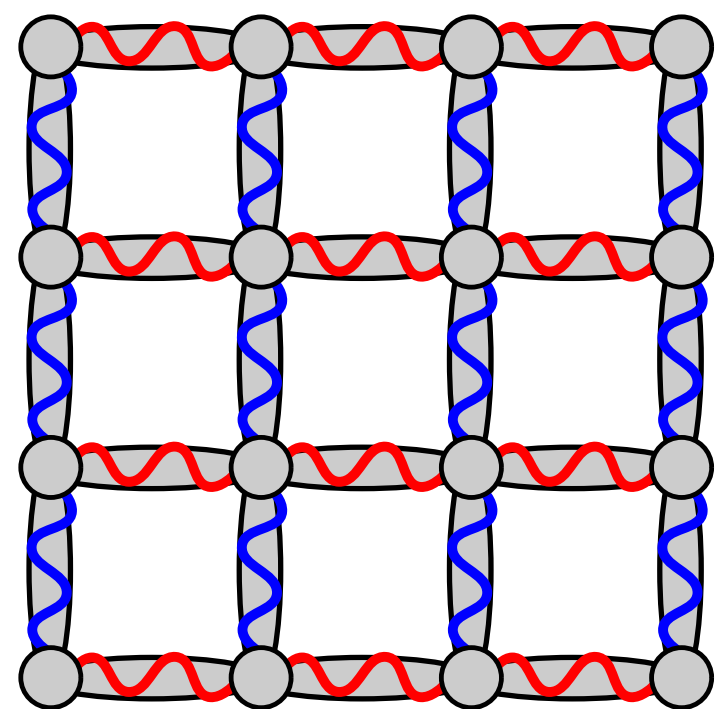
# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



$$\langle B \rangle \neq 0$$

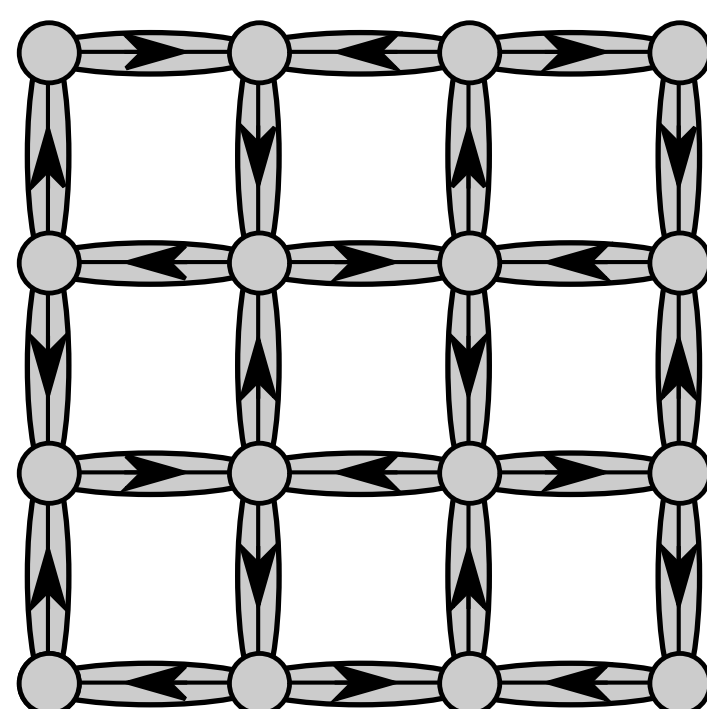
**Phase B**

*d*-wave SC



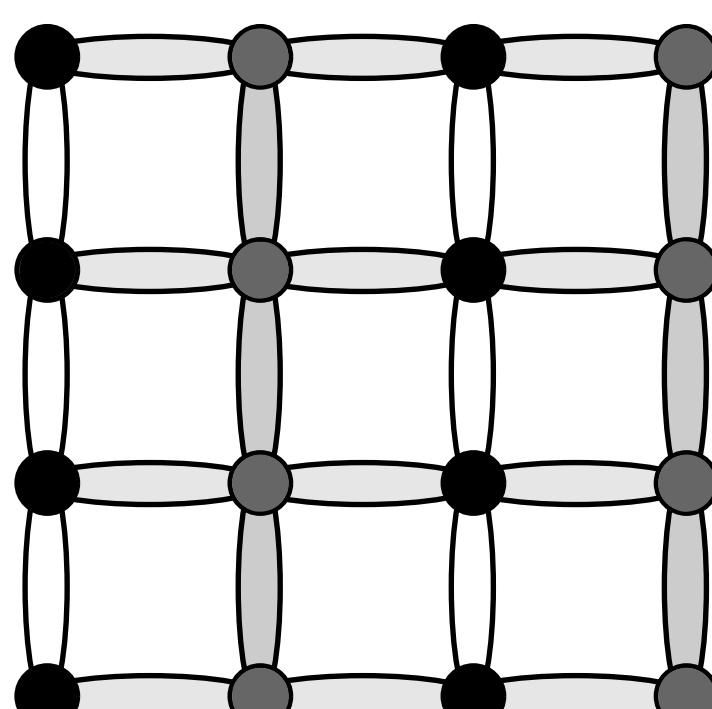
**Phase C**

*d*-density



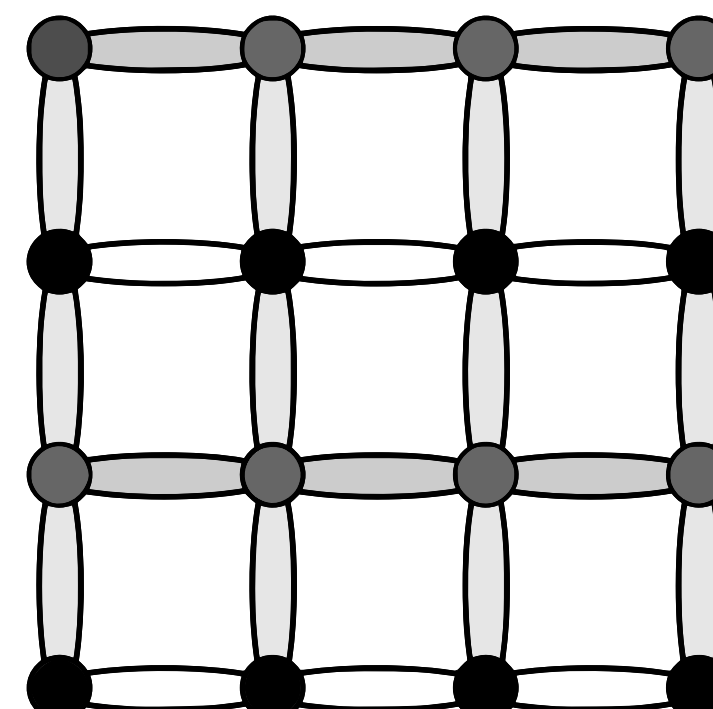
**Phase A**

$(\pi, 0)$  stripe

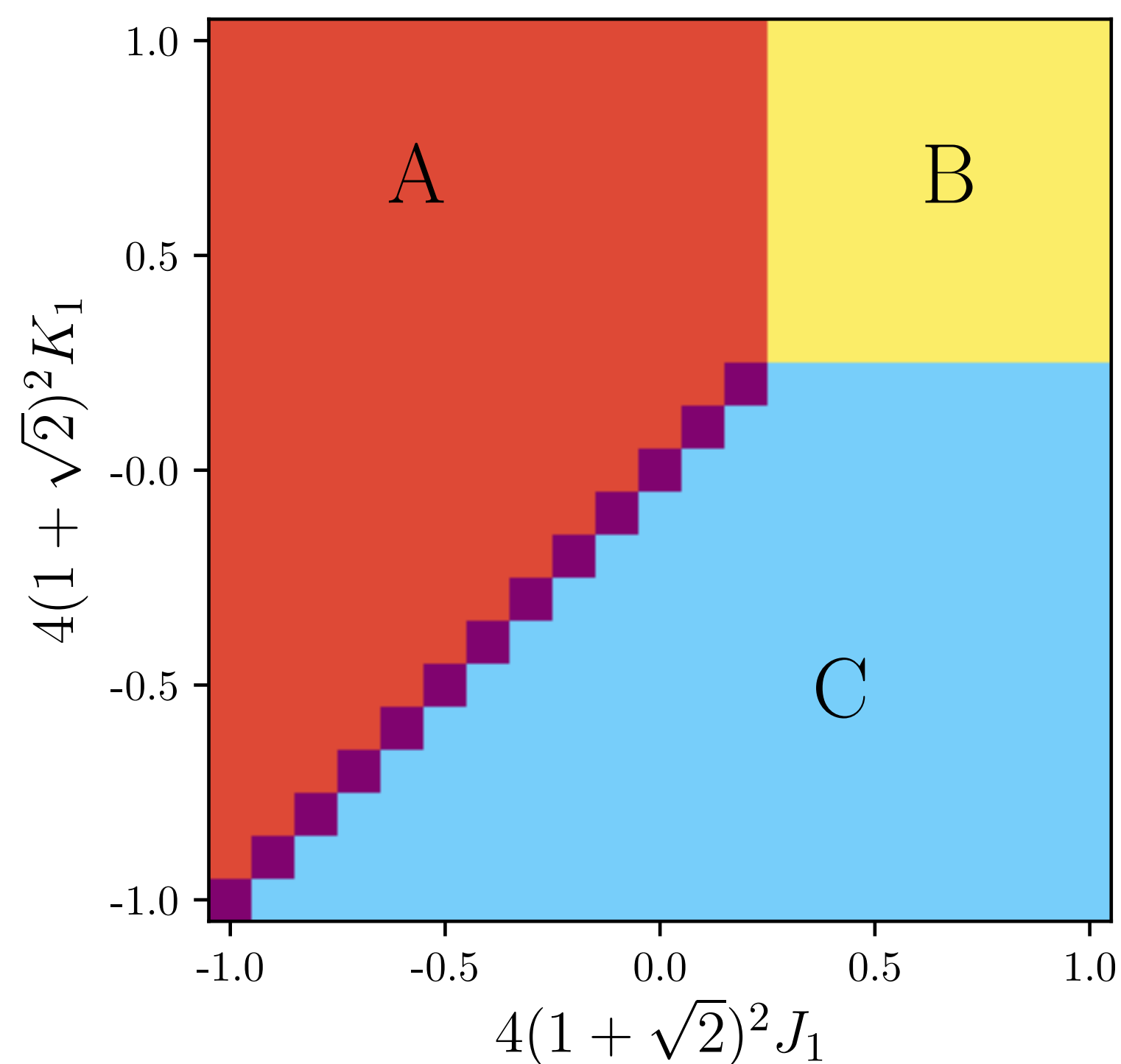


**Phase A**

$(0, \pi)$  stripe

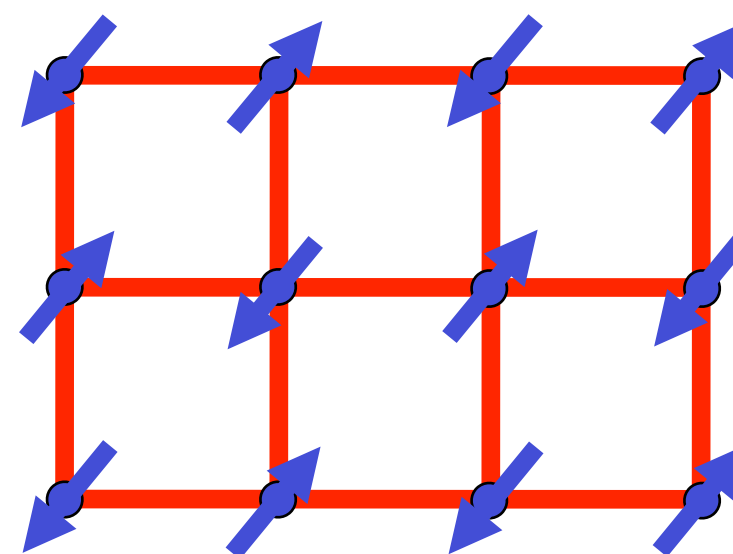


# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

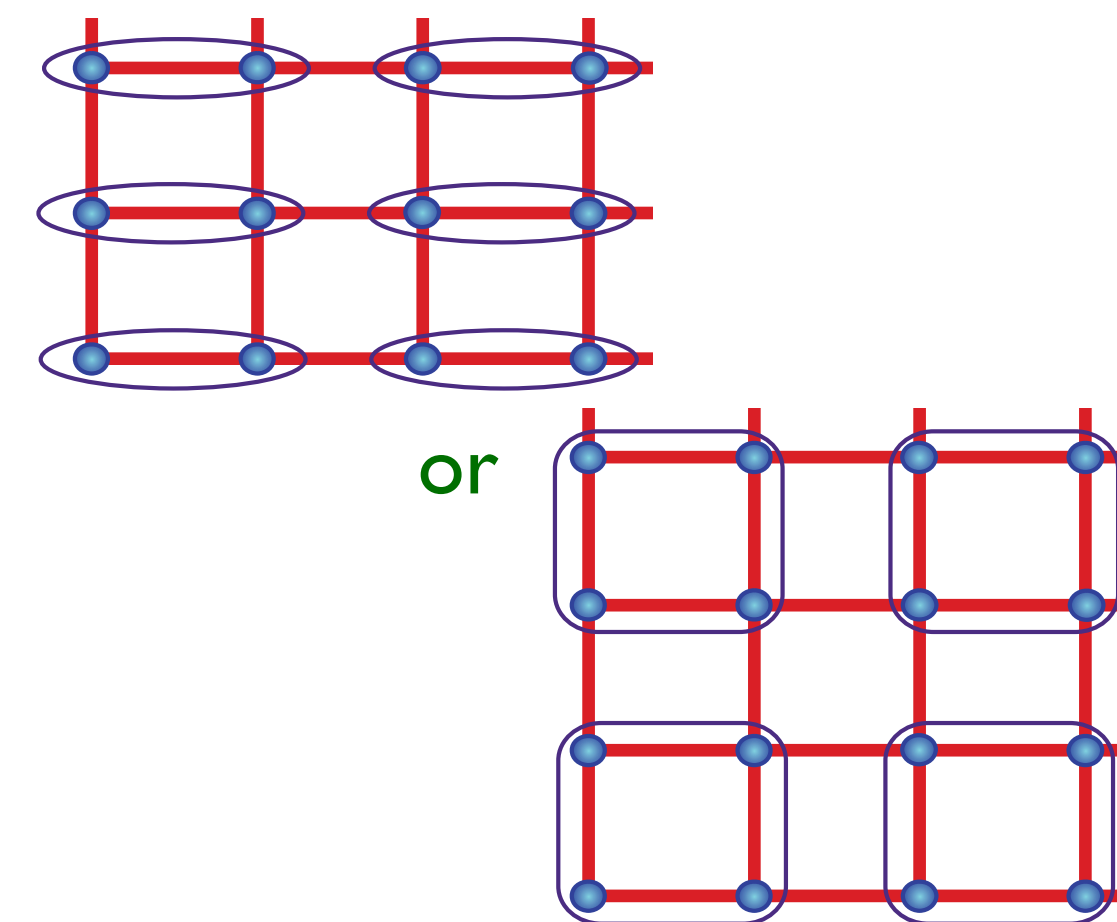


$$\langle B \rangle \neq 0$$

$$\langle B \rangle = 0$$



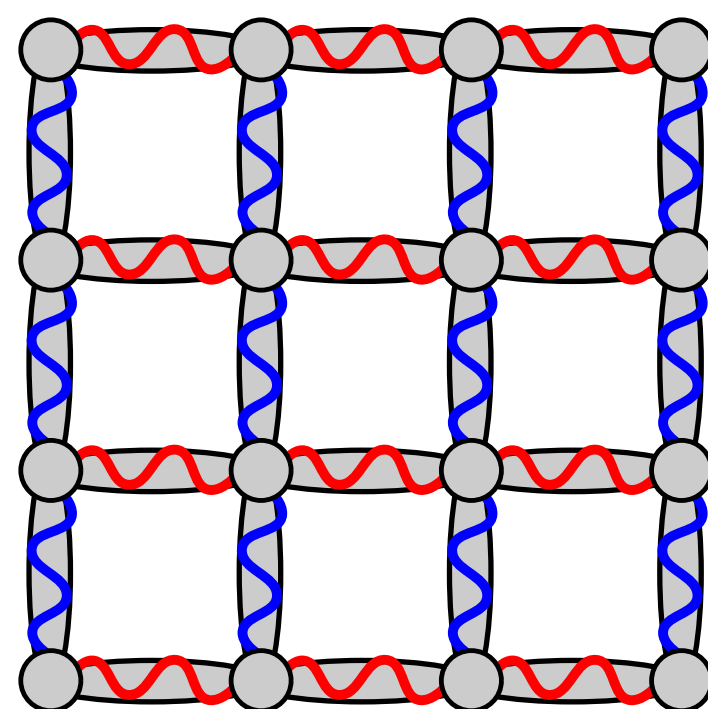
Confining phase:  
Néel order



Confining phase:  
VBS order

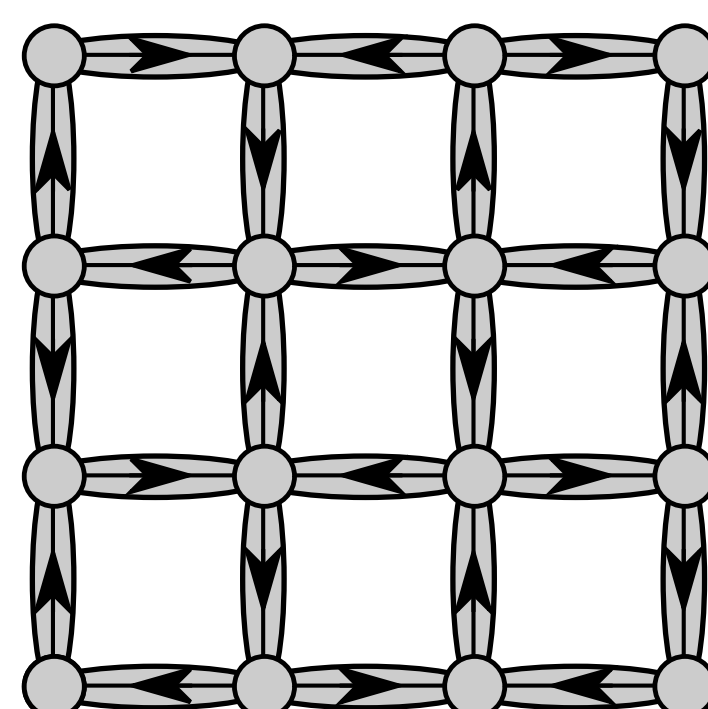
**Phase B**

*d*-wave SC



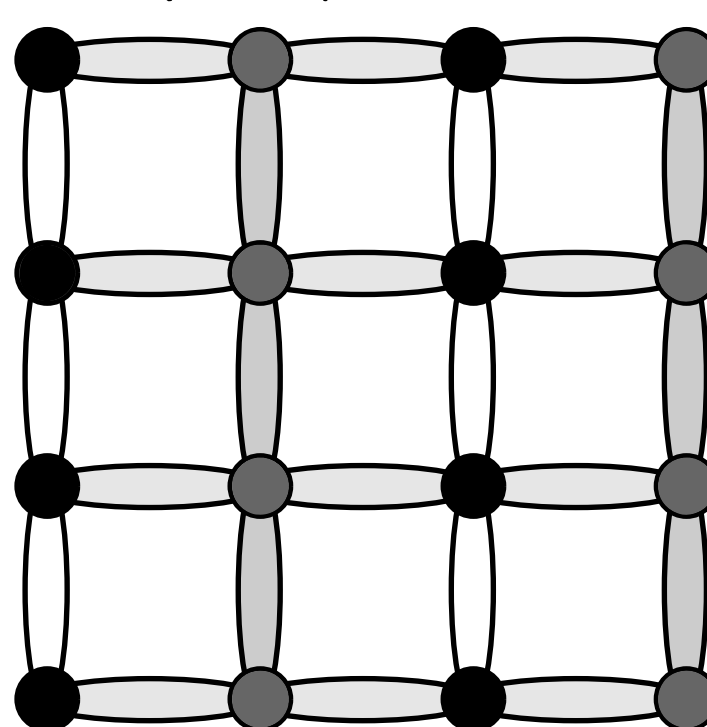
**Phase C**

*d*-density



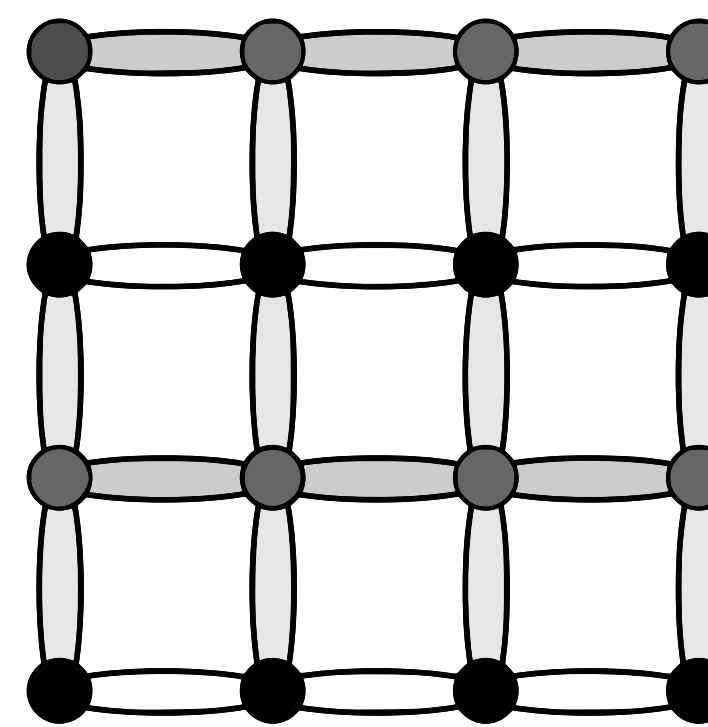
**Phase A**

$(\pi, 0)$  stripe

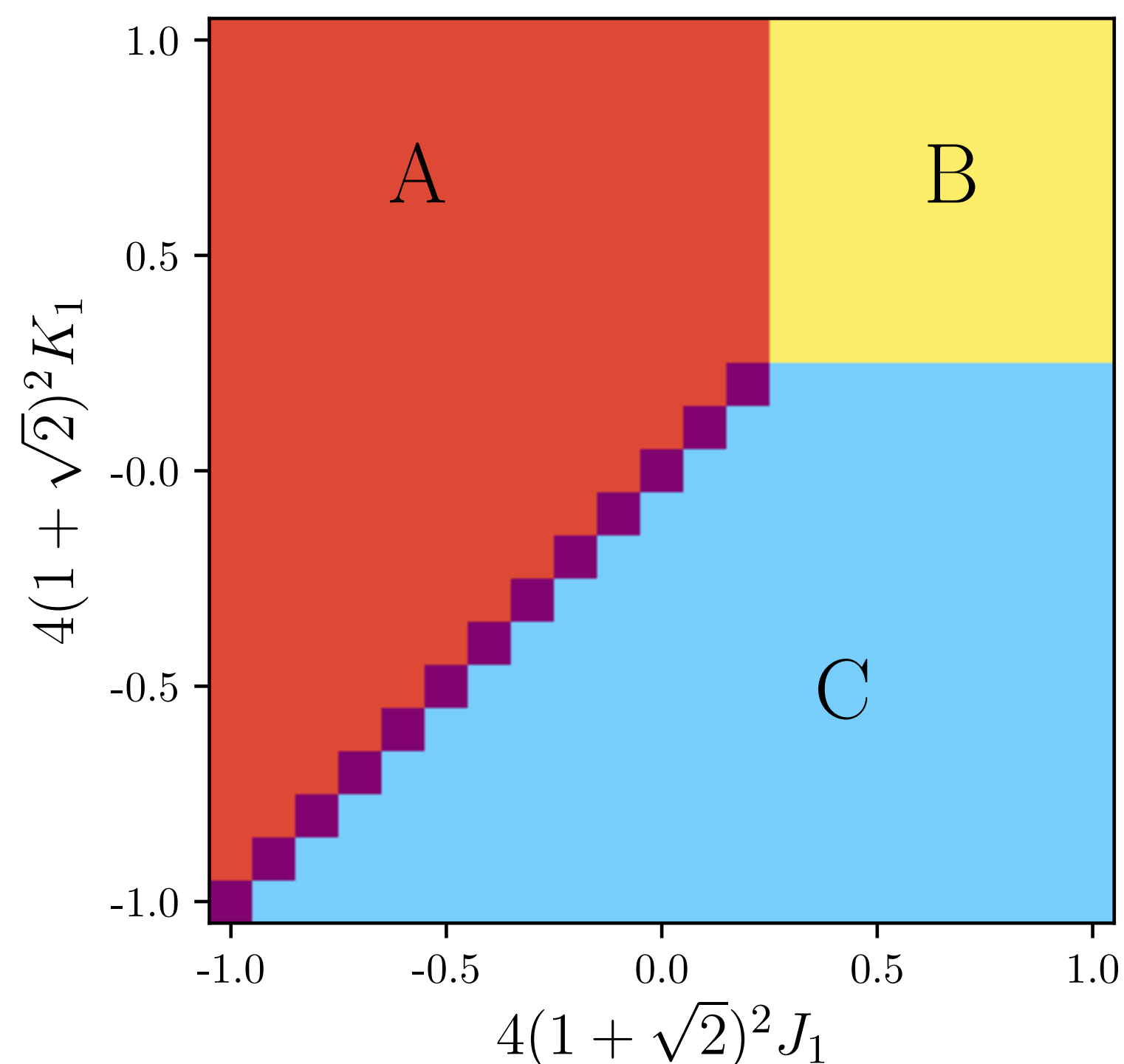


**Phase A**

$(0, \pi)$  stripe



# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



$$\langle B \rangle = 0$$

Confining phase.  
 $SO(5)_f$  broken.  
 Néel or  
 valence bond solid  
 order.

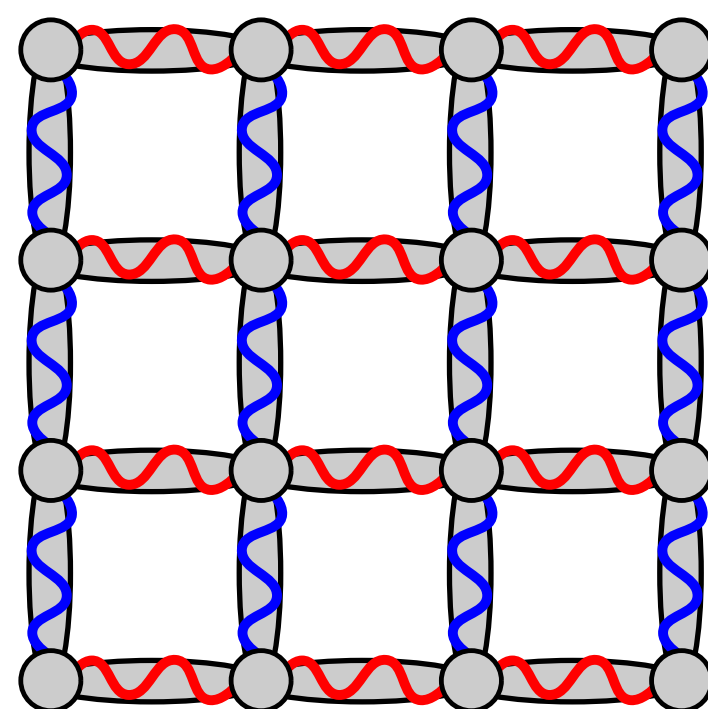
$$\langle B \rangle \neq 0$$

Higgs phase.  
 $SO(5)_b$  broken.  
*d*-wave superconductivity or  
 period-2 stripes or  
*d*-density wave order.

 $r$ 
 $r_c$ 

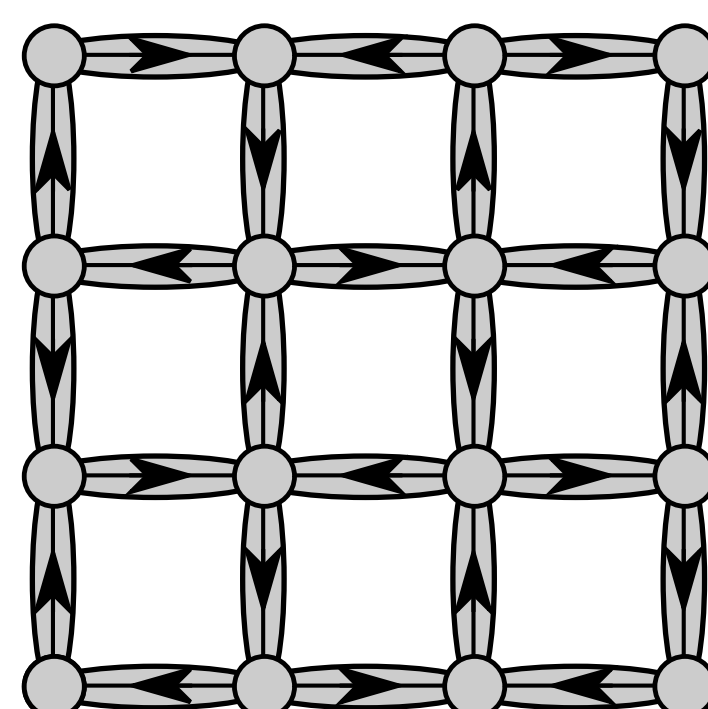
**Phase B**

*d*-wave SC



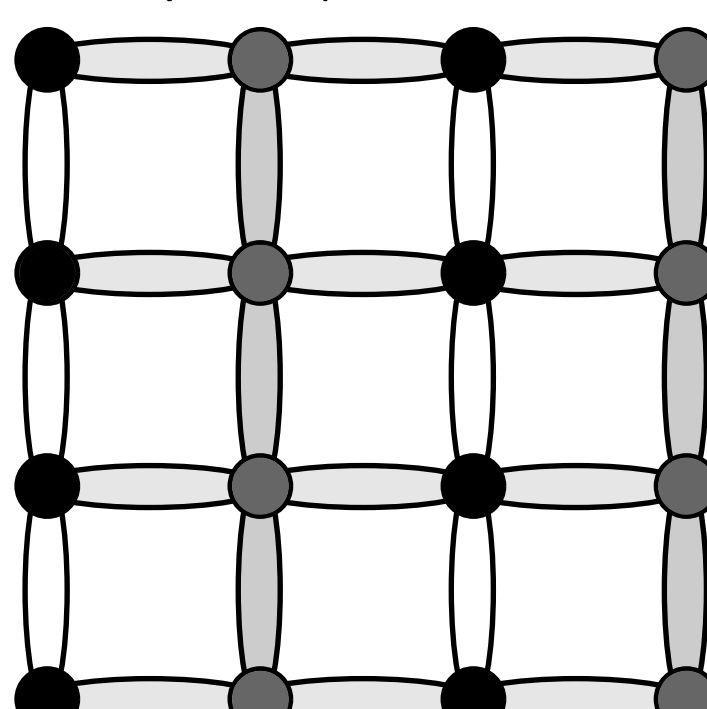
**Phase C**

*d*-density



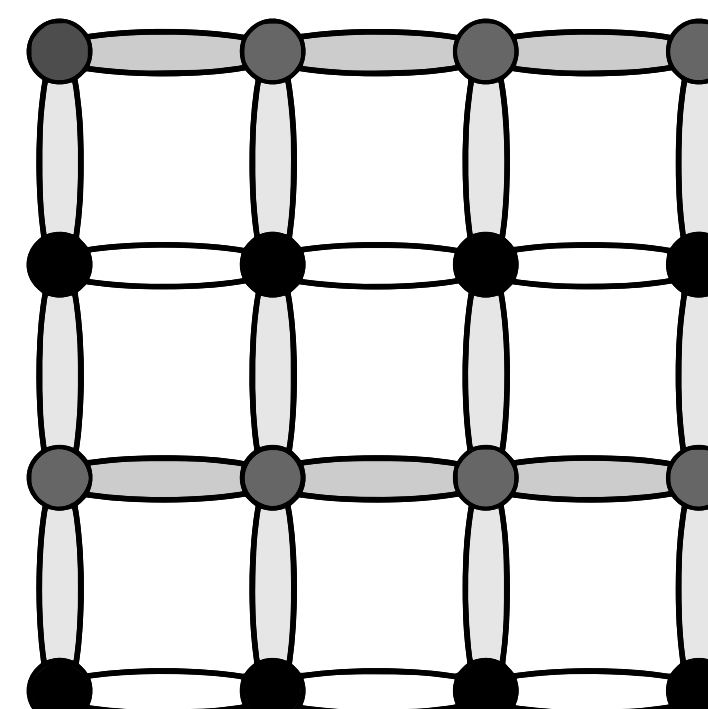
**Phase A**

$(\pi, 0)$  stripe



**Phase A**

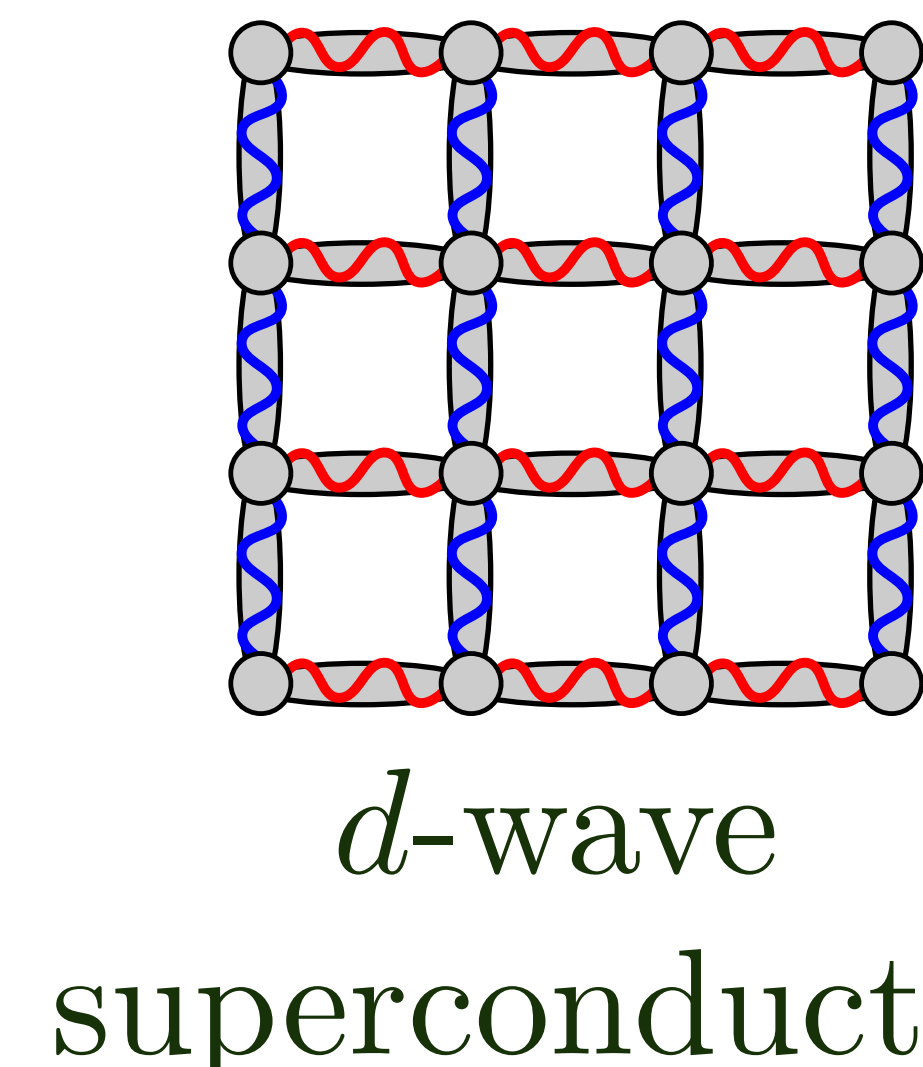
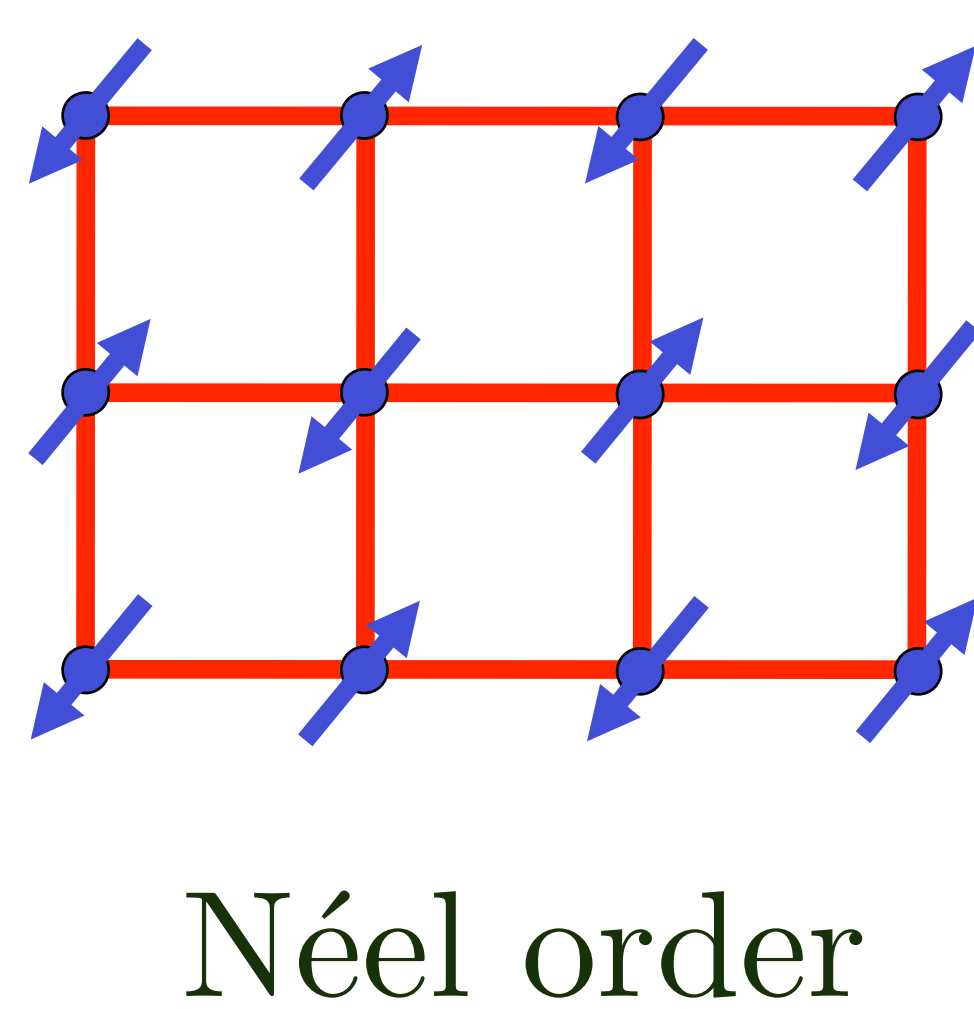
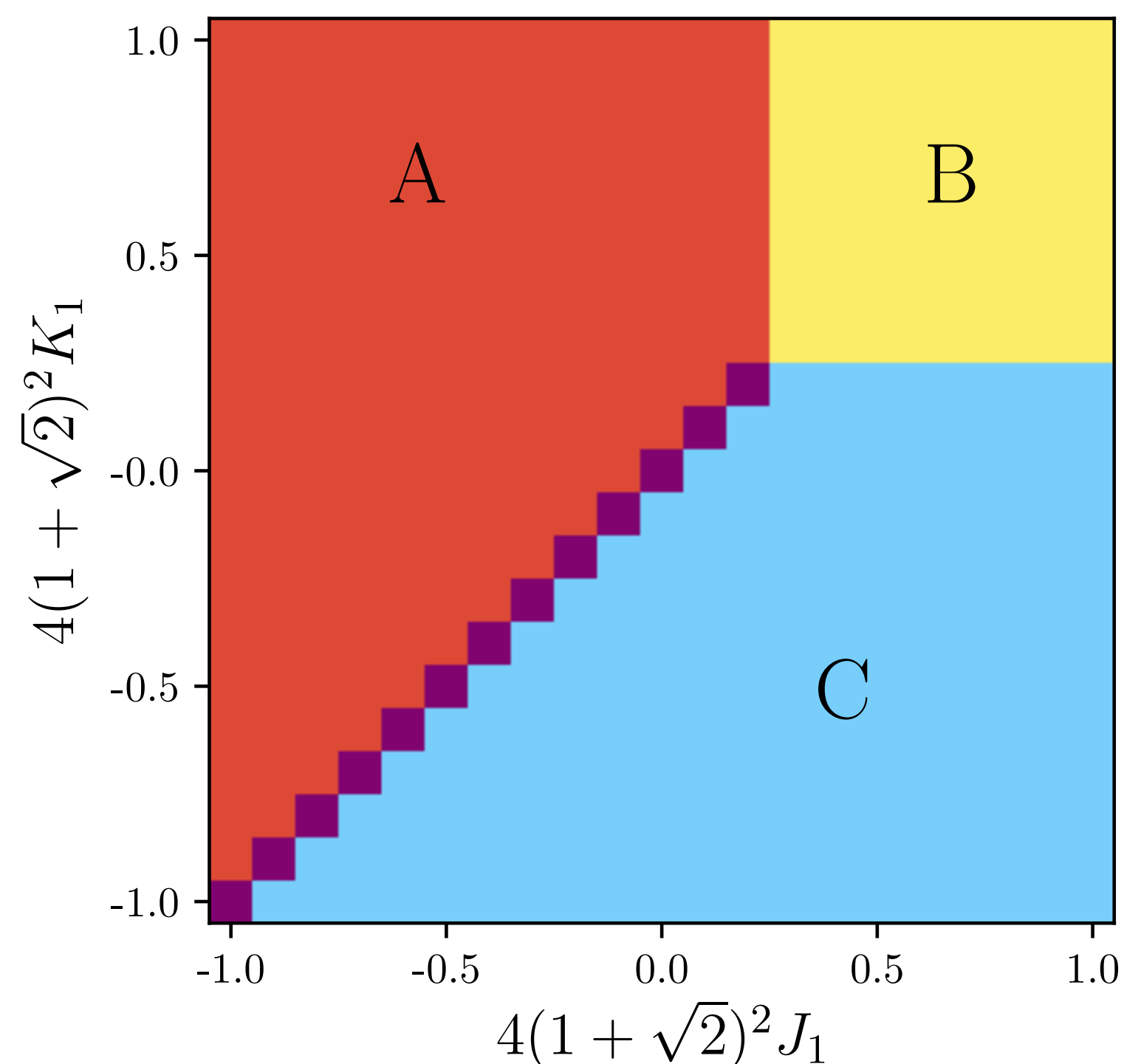
$(0, \pi)$  stripe



$$\mathcal{L} = i\bar{\Psi}_s \gamma_\mu D_\mu \Psi_s + |D_\mu B_s|^2 + r|B_s|^2 \dots$$

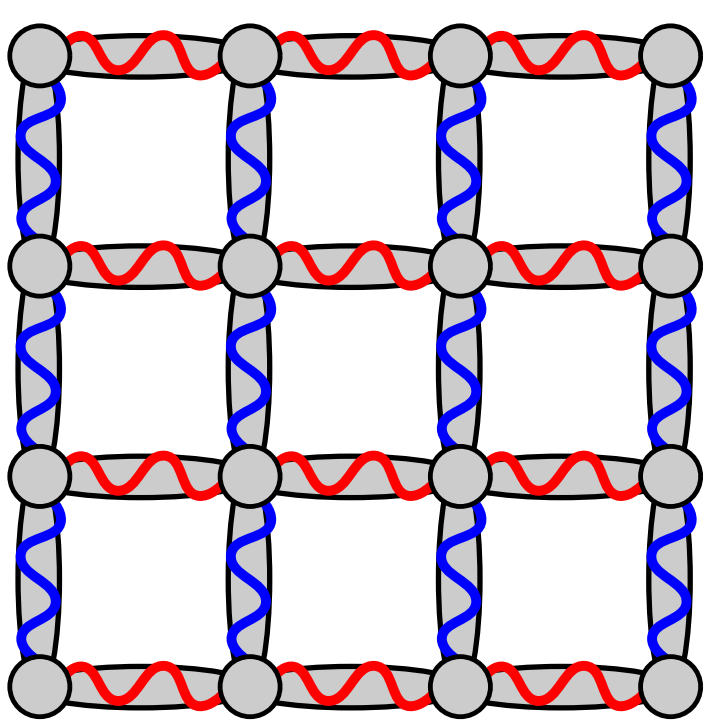
# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory

F. F. Assaad, M. Imada, and D. J. Scalapino, PRL 77, 4592 (1996)

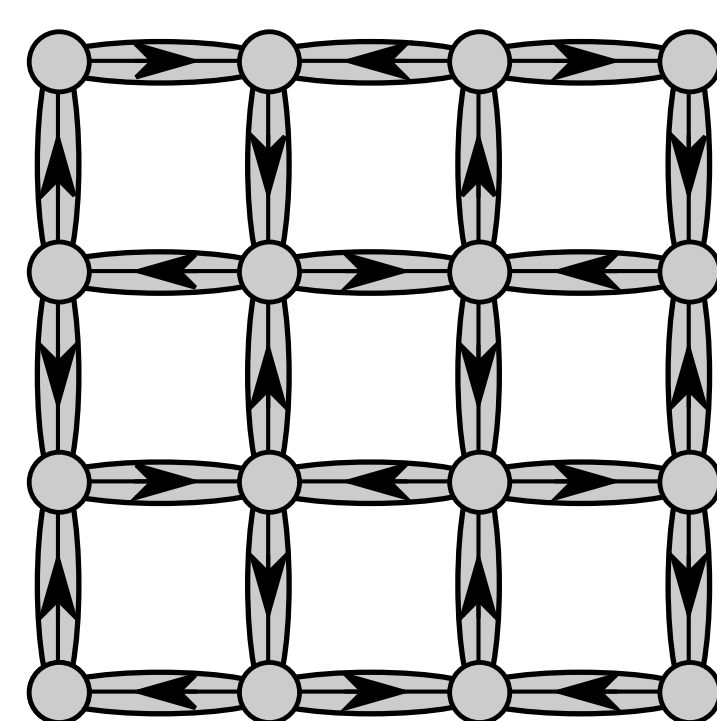


0.3  $\rightarrow W/t$

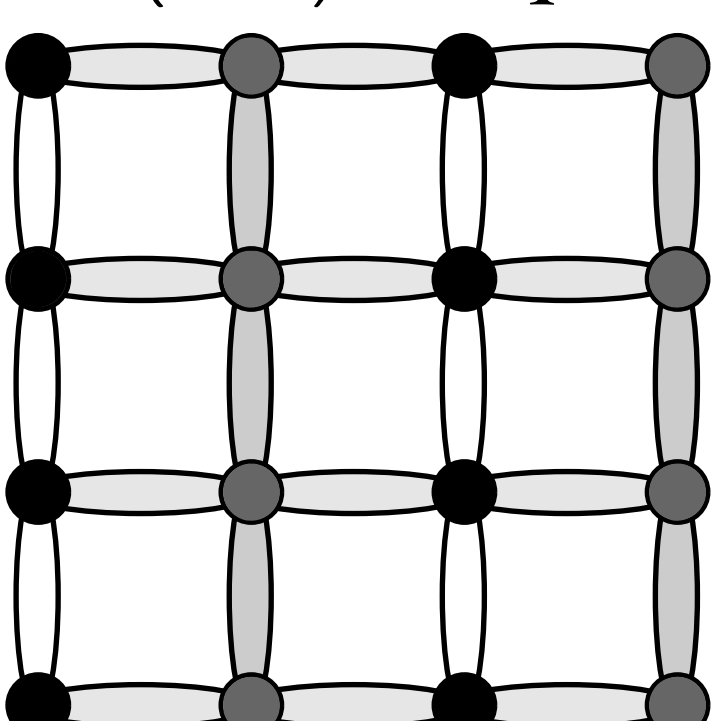
**Phase B**  
*d*-wave SC



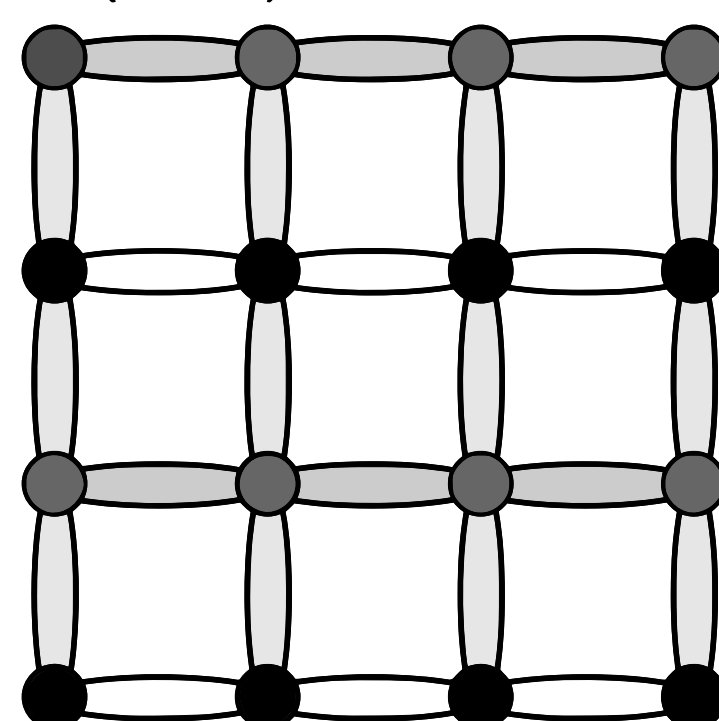
**Phase C**  
*d*-density



**Phase A**  
 $(\pi, 0)$  stripe

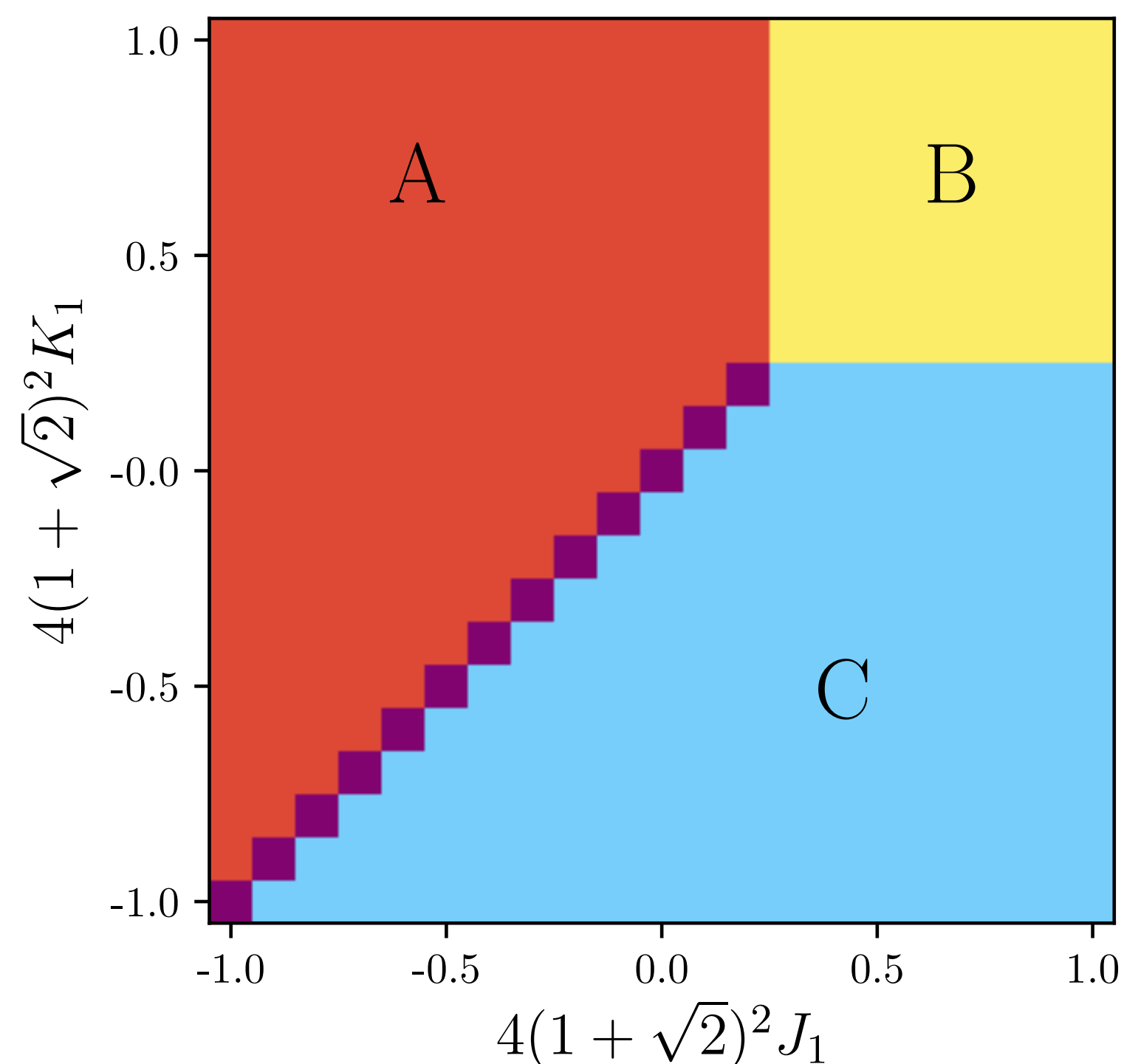


**Phase A**  
 $(0, \pi)$  stripe



$$\mathcal{L} = i\bar{\Psi}_s \gamma_\mu D_\mu \Psi_s + |D_\mu B_s|^2 + r|B_s|^2 \dots$$

# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



$$\langle B \rangle = 0$$

Confining phase.  
 $SO(5)_f$  broken.  
 Néel or  
 valence bond solid  
 order.

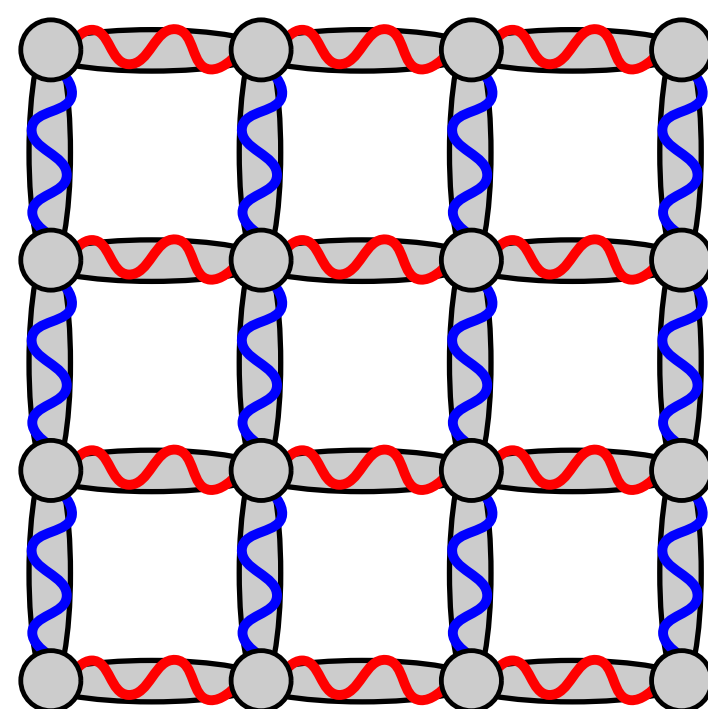
$$\langle B \rangle \neq 0$$

Higgs phase.  
 $SO(5)_b$  broken.  
*d*-wave superconductivity or  
 period-2 stripes or  
*d*-density wave order.

 $r$ 
 $r_c$ 

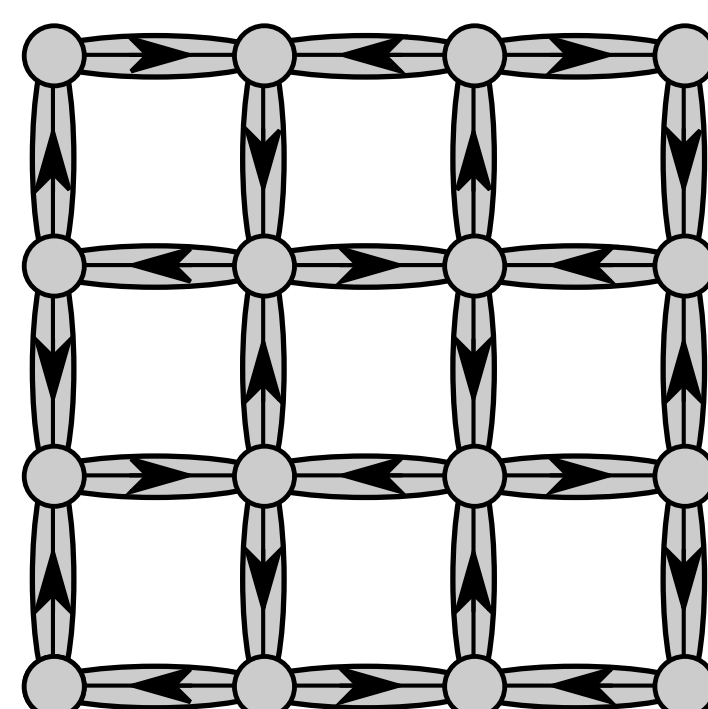
**Phase B**

*d*-wave SC



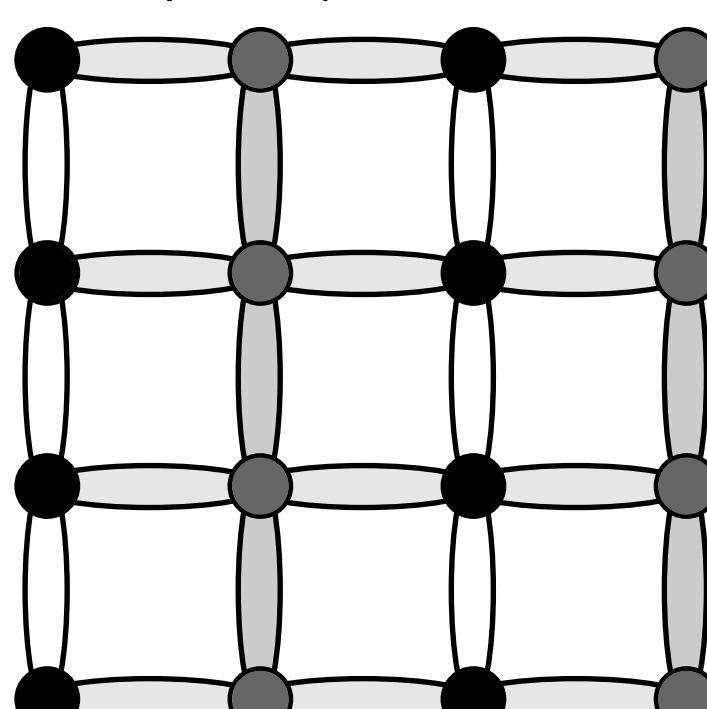
**Phase C**

*d*-density



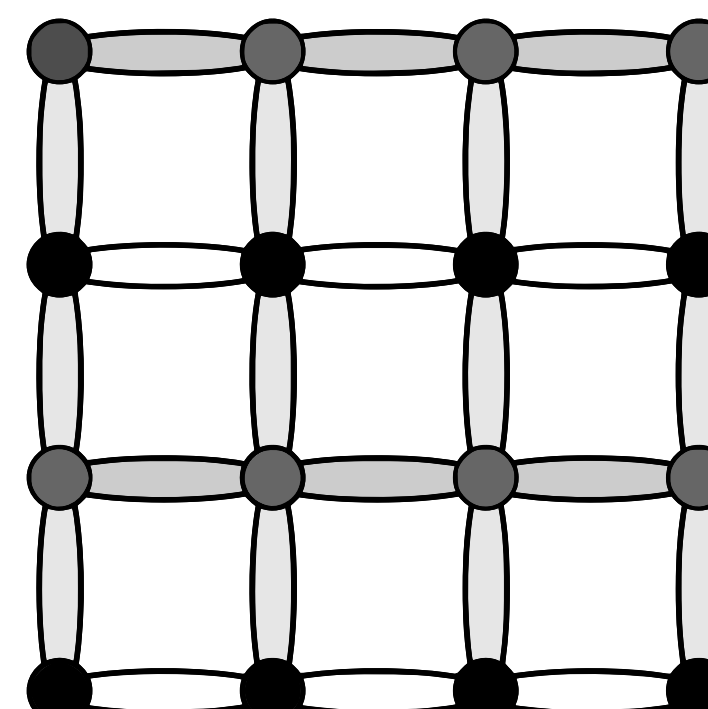
**Phase A**

$(\pi, 0)$  stripe



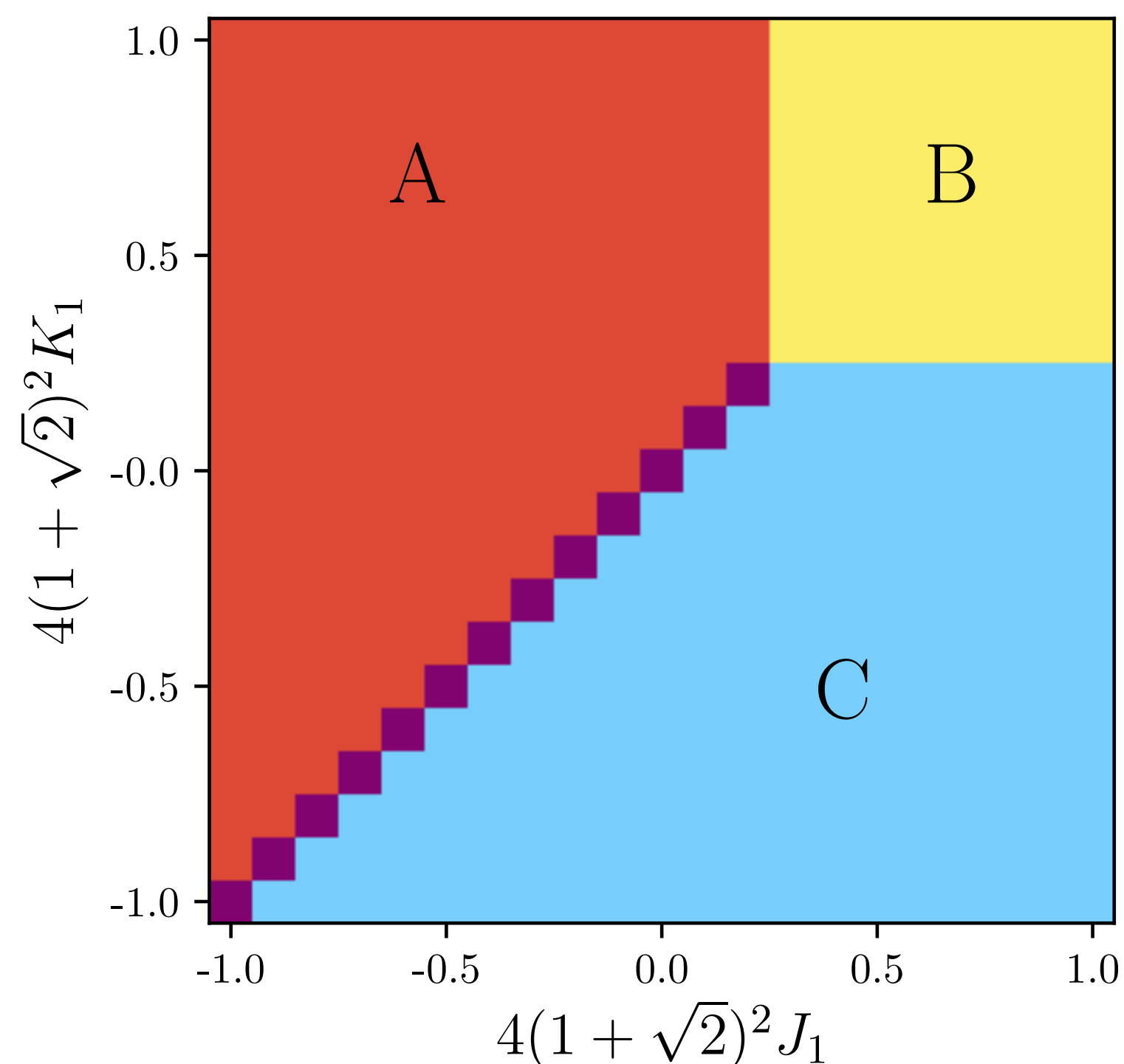
**Phase A**

$(0, \pi)$  stripe



$$\mathcal{L} = i\bar{\Psi}_s \gamma_\mu D_\mu \Psi_s + |D_\mu B_s|^2 + r|B_s|^2 \dots$$

# Include charge fluctuations at half-filling: confinement of SU(2) gauge theory



$$\langle B \rangle = 0$$

Confining phase.  
 $SO(5)_f$  broken.  
 Néel or  
 valence bond solid  
 order.

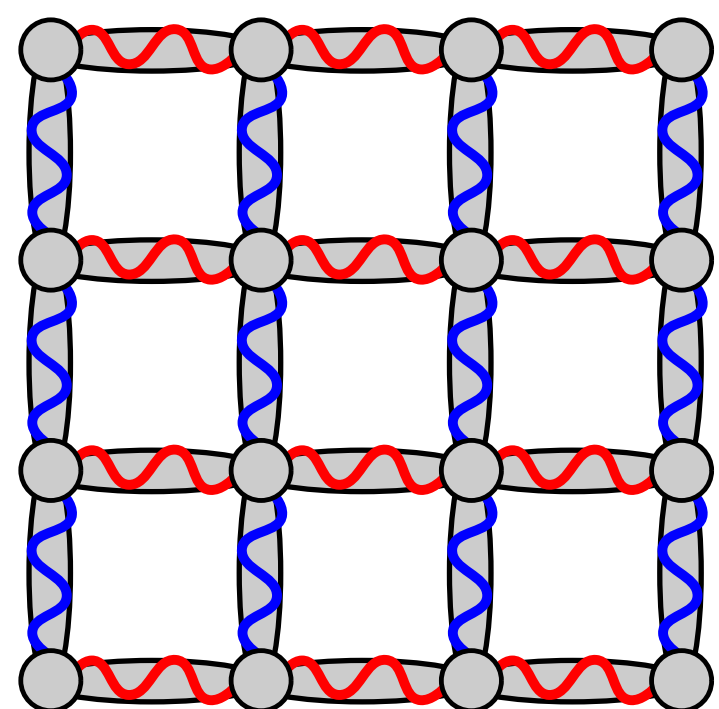
$$\langle B \rangle \neq 0$$

Higgs phase.  
 $SO(5)_b$  broken.  
*d*-wave superconductivity or  
 period-2 stripes or  
*d*-density wave order.



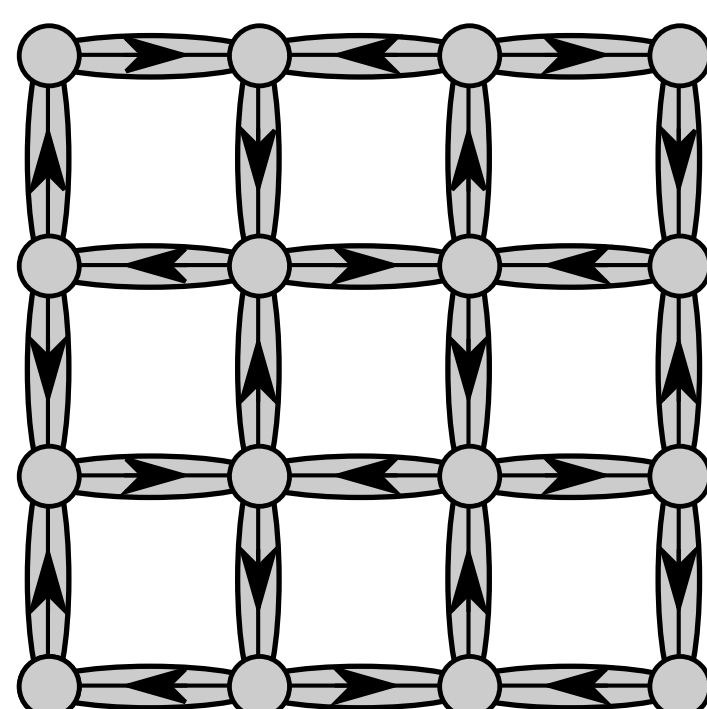
**Phase B**

*d*-wave SC



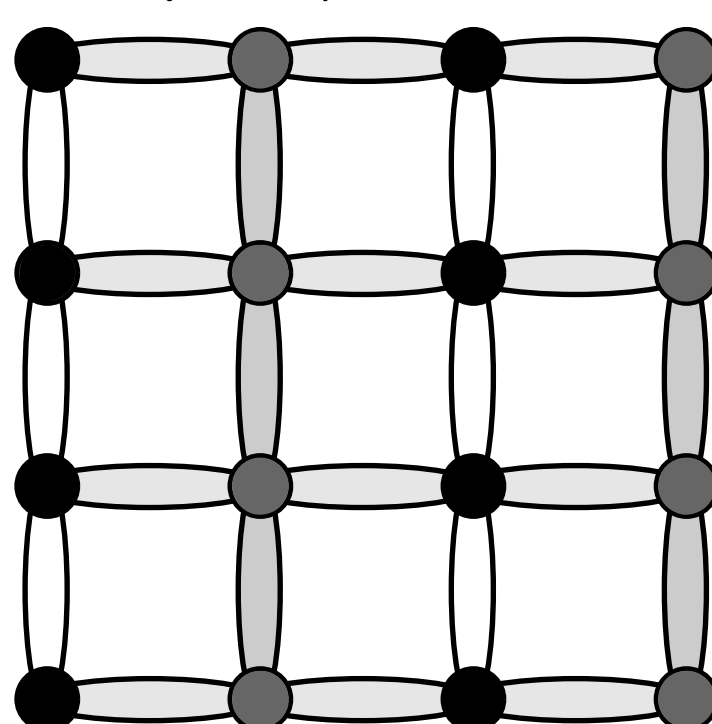
**Phase C**

*d*-density



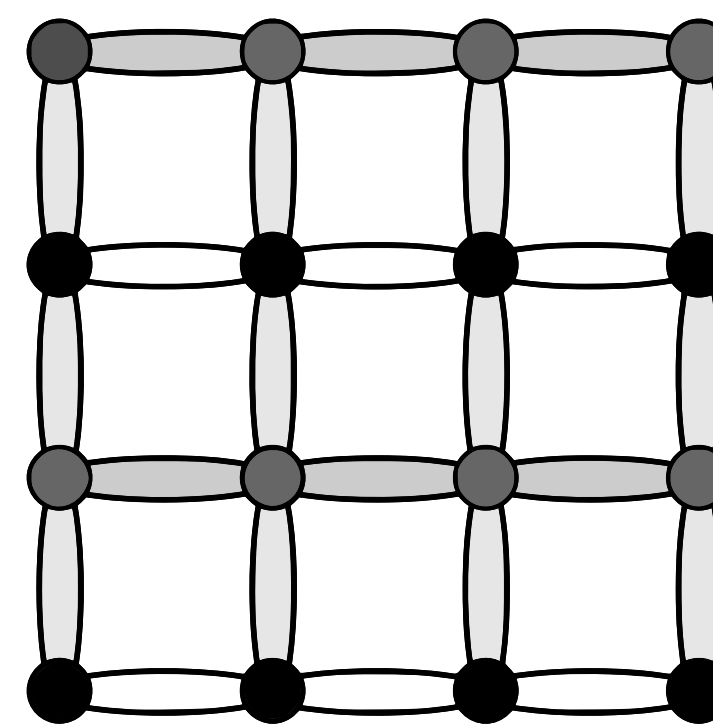
**Phase A**

$(\pi, 0)$  stripe



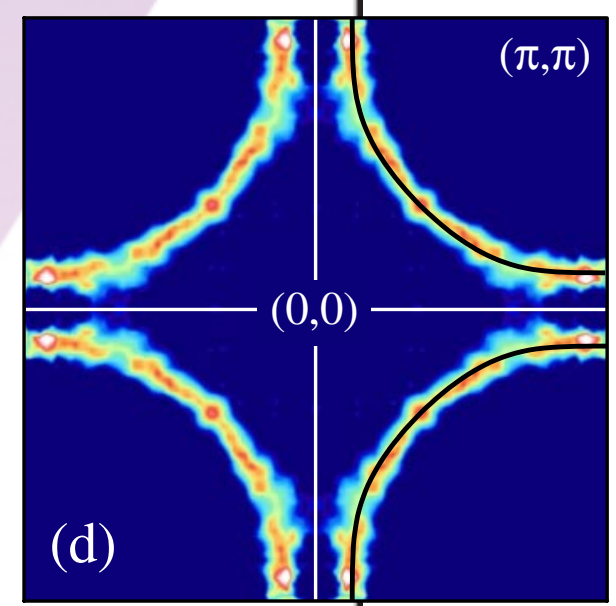
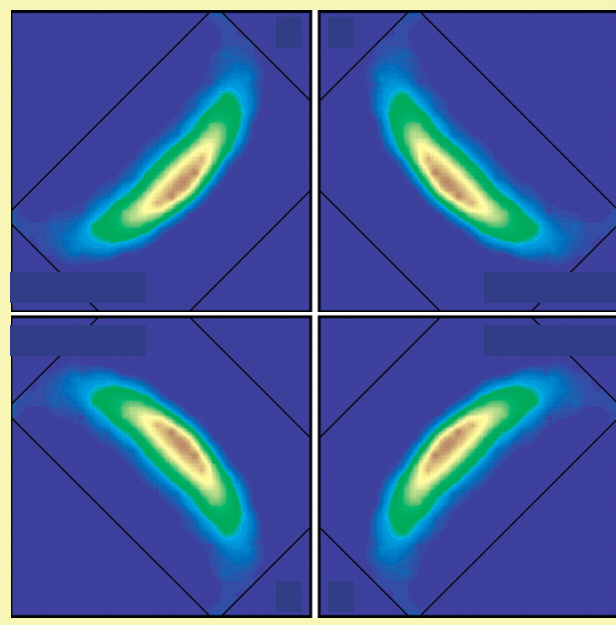
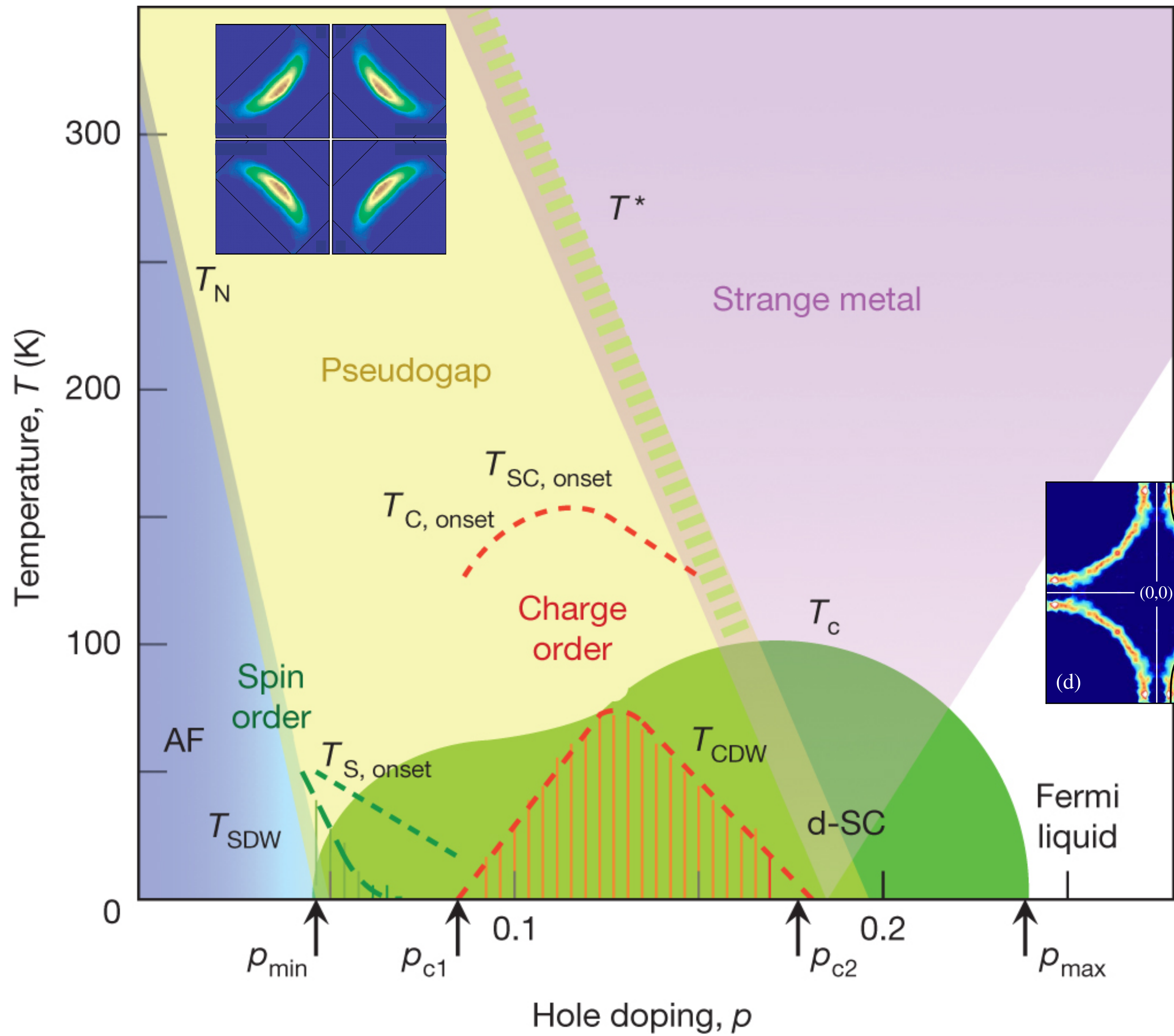
**Phase A**

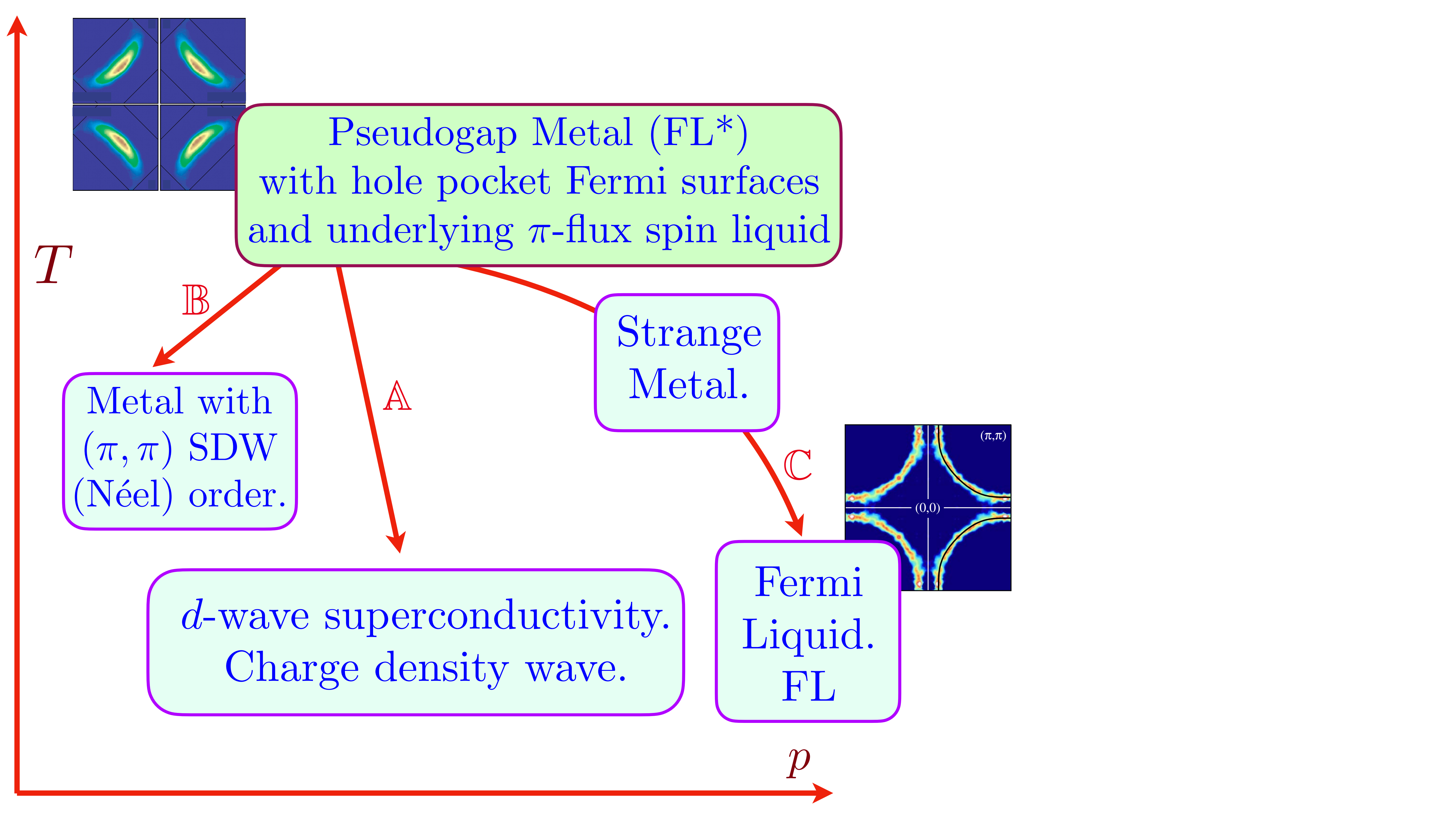
$(0, \pi)$  stripe



$$\mathcal{L} = i\bar{\Psi}_s \gamma_\mu D_\mu \Psi_s + |D_\mu B_s|^2 + r|B_s|^2 \dots$$

Possible CFT.  
 DQCP with  
 $SO(5)_f \times SO(5)_b$   
 symmetry.





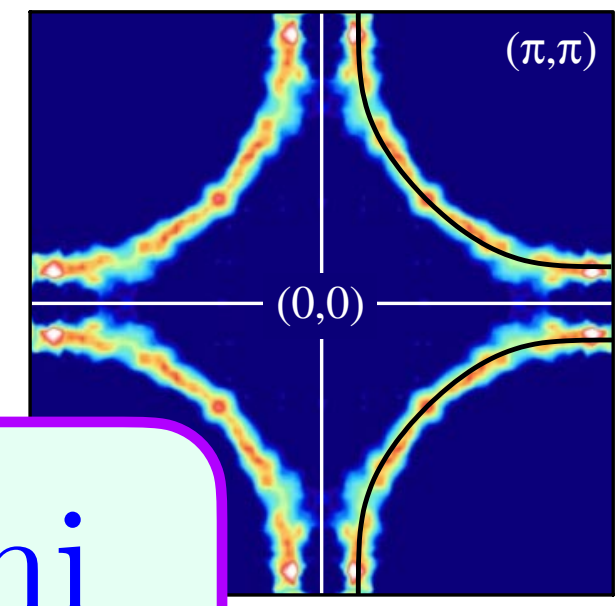
Pseudogap Metal (FL\*)  
with hole pocket Fermi surfaces  
and underlying  $\pi$ -flux spin liquid

Strange  
Metal.

Metal with  
 $(\pi, \pi)$  SDW  
(Néel) order.

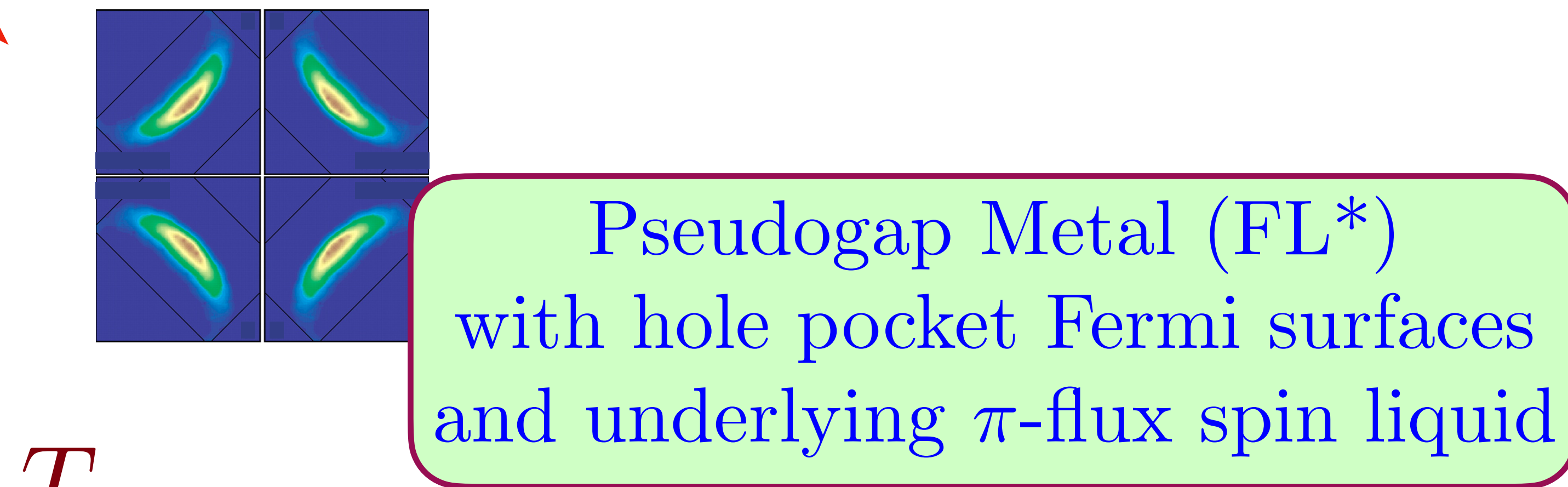
$d$ -wave superconductivity.  
Charge density wave.

Fermi  
Liquid.  
FL

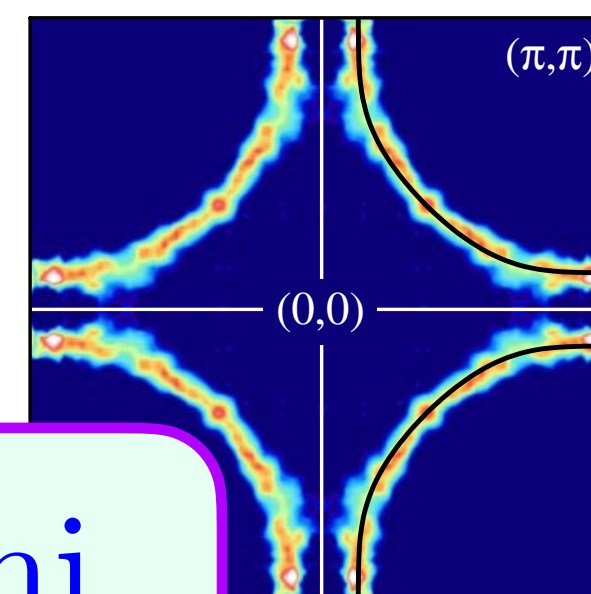


Arrow B

Condensation of  $z_\alpha$  in dual  $\mathbb{C}\mathbb{P}^1$   
U(1) gauge theory.



Strange  
Metal.



Fermi  
Liquid.  
FL

$d$ -wave superconductivity.  
Charge density wave.

$p$

Arrow A

Condensation of B in SU(2) gauge theory.

Longer-range couplings in  $H_B$  can lead to charge order with other periods

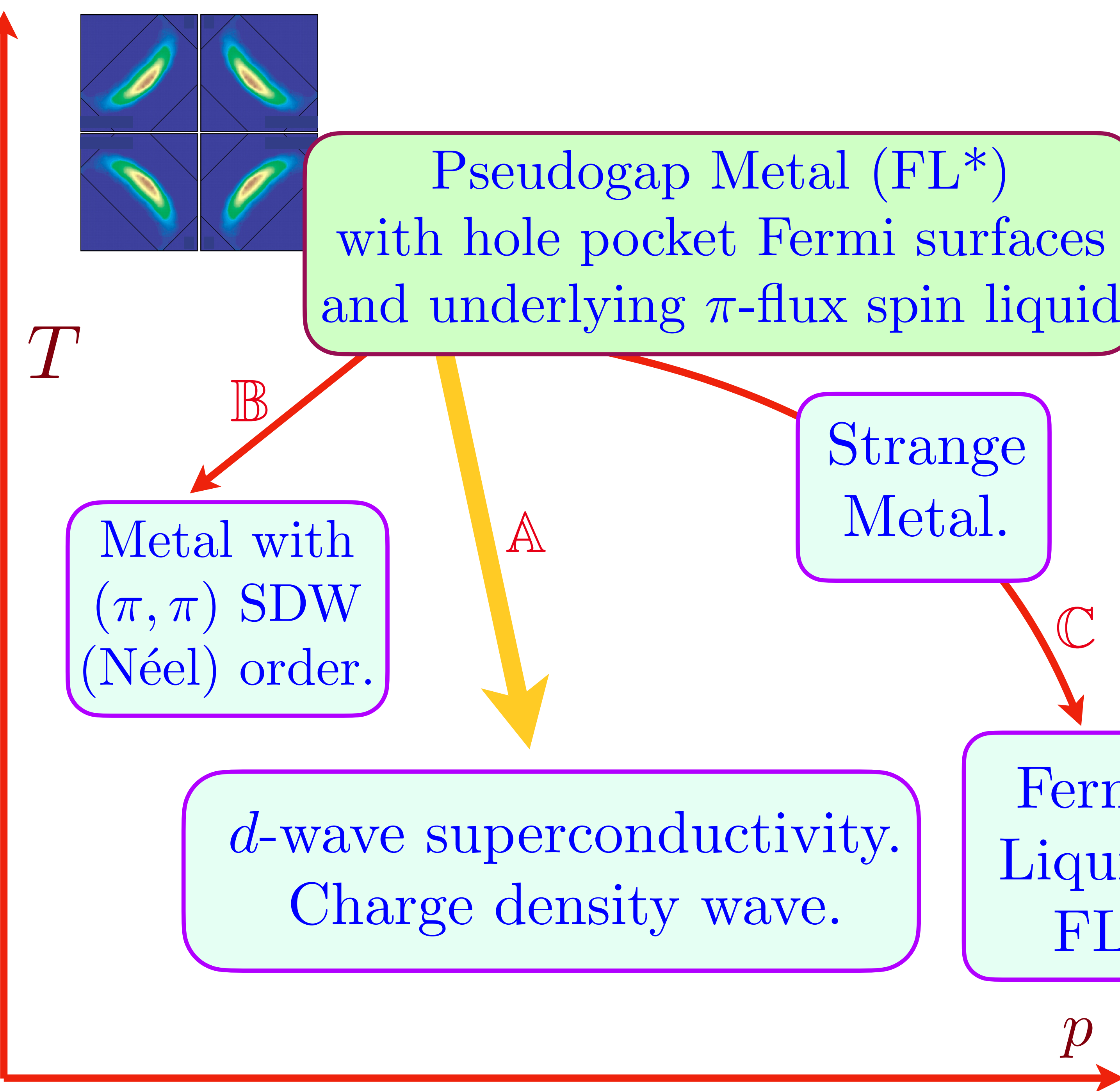
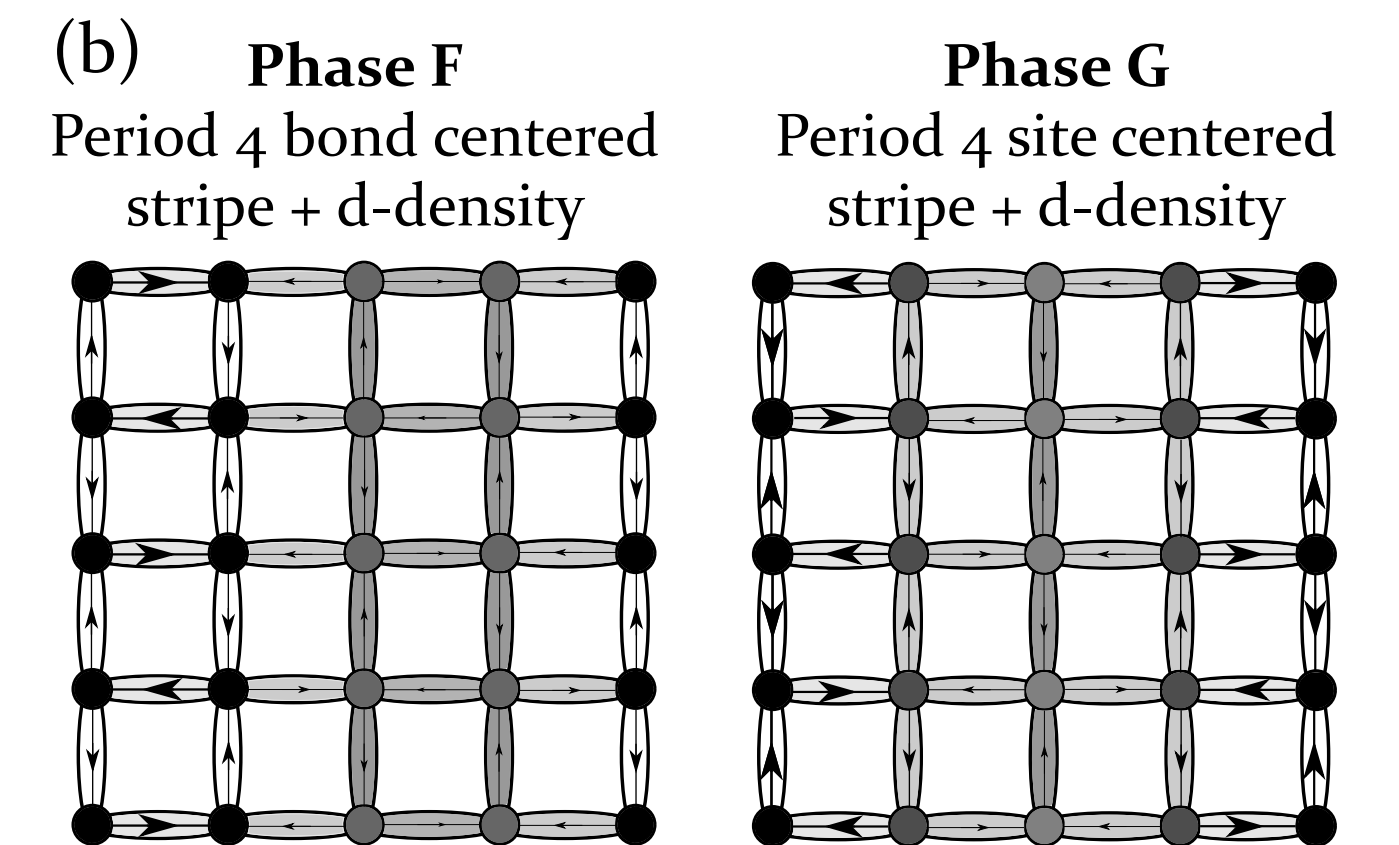
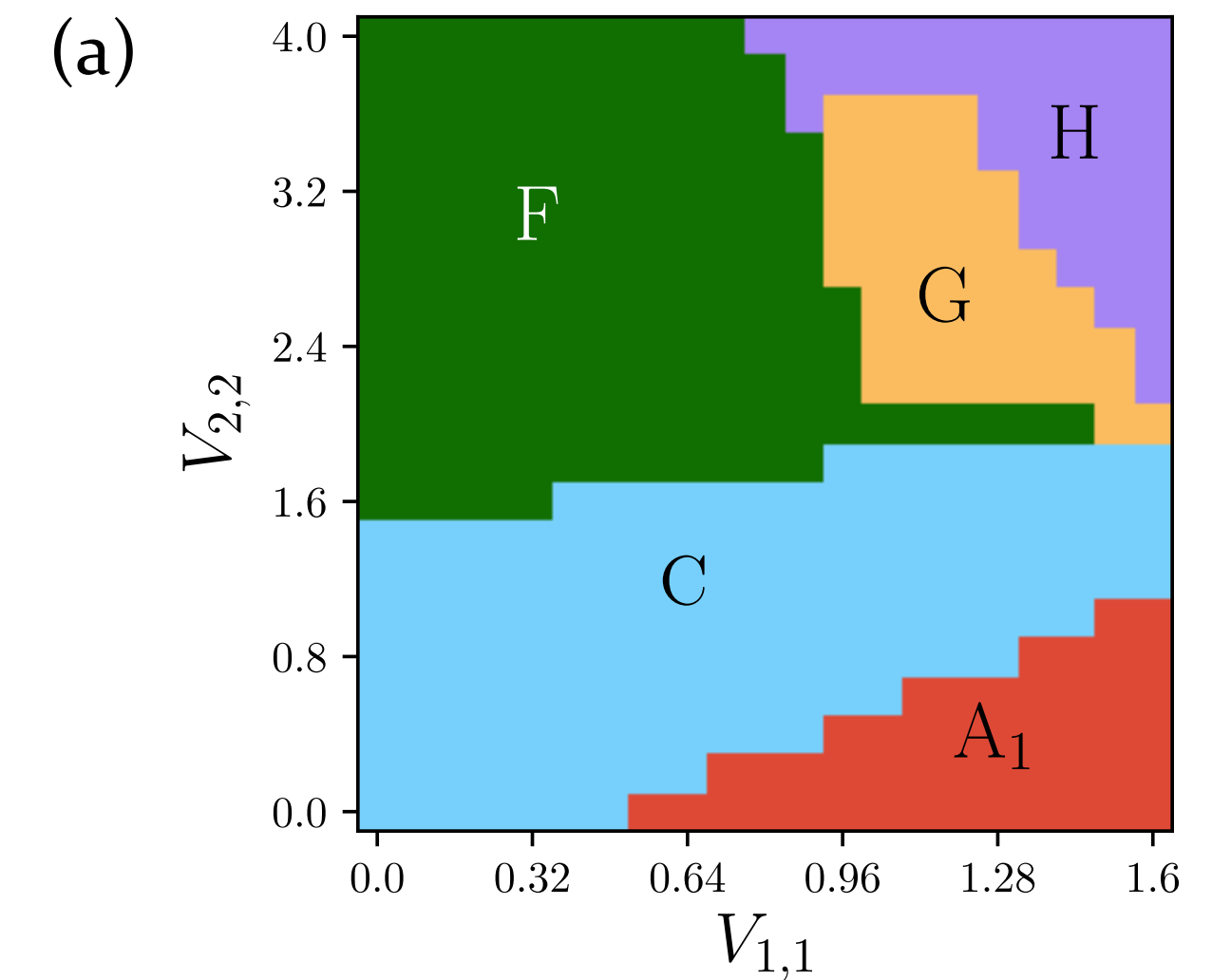
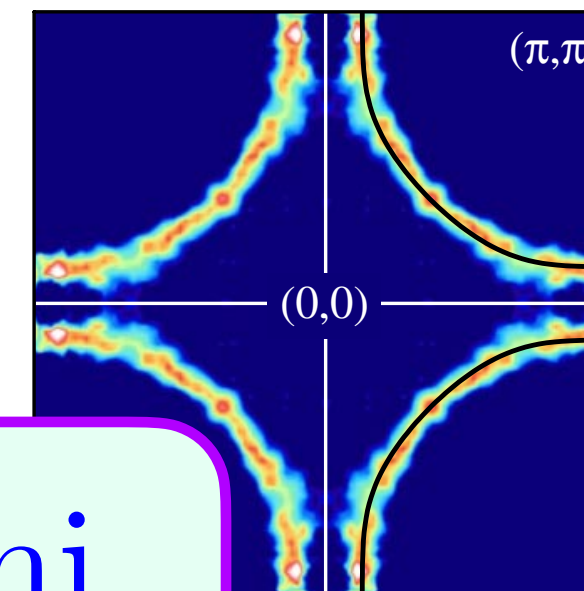
Pseudogap Metal (FL\*)  
with hole pocket Fermi surfaces  
and underlying  $\pi$ -flux spin liquid

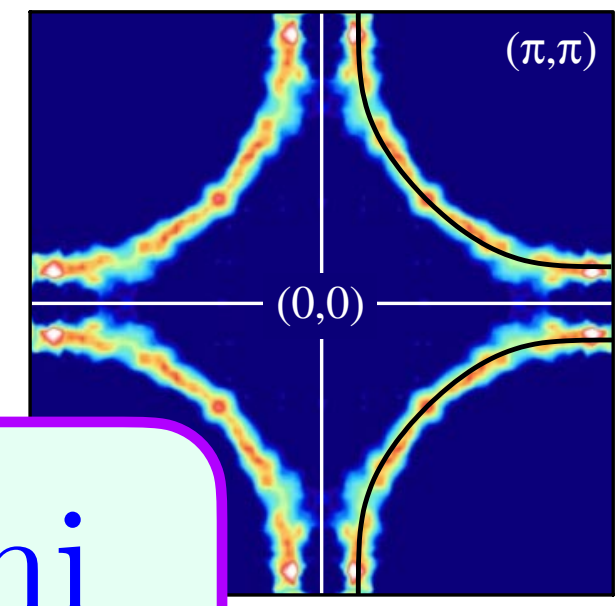
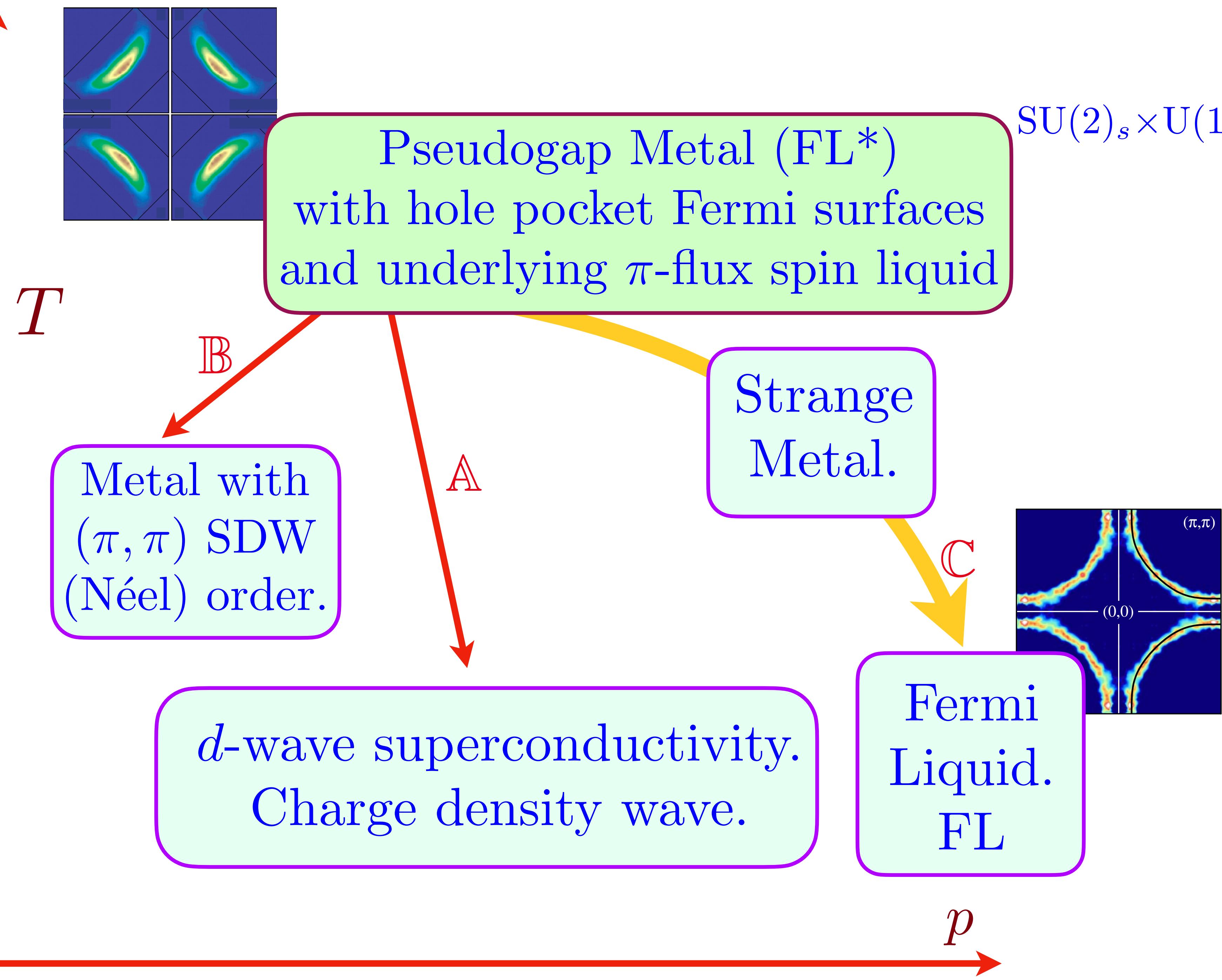
Strange  
Metal.

Metal with  
 $(\pi, \pi)$  SDW  
(Néel) order.

$d$ -wave superconductivity.  
Charge density wave.

Fermi  
Liquid.  
FL





Arrow A

Condensation of B in SU(2) gauge theory.

Longer-range couplings in  $H_B$  can lead to charge order with other periods

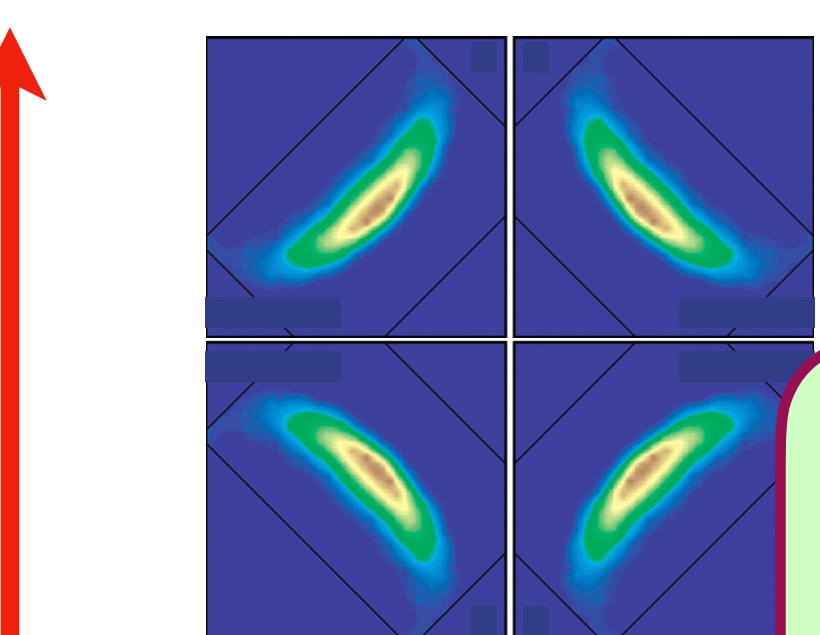
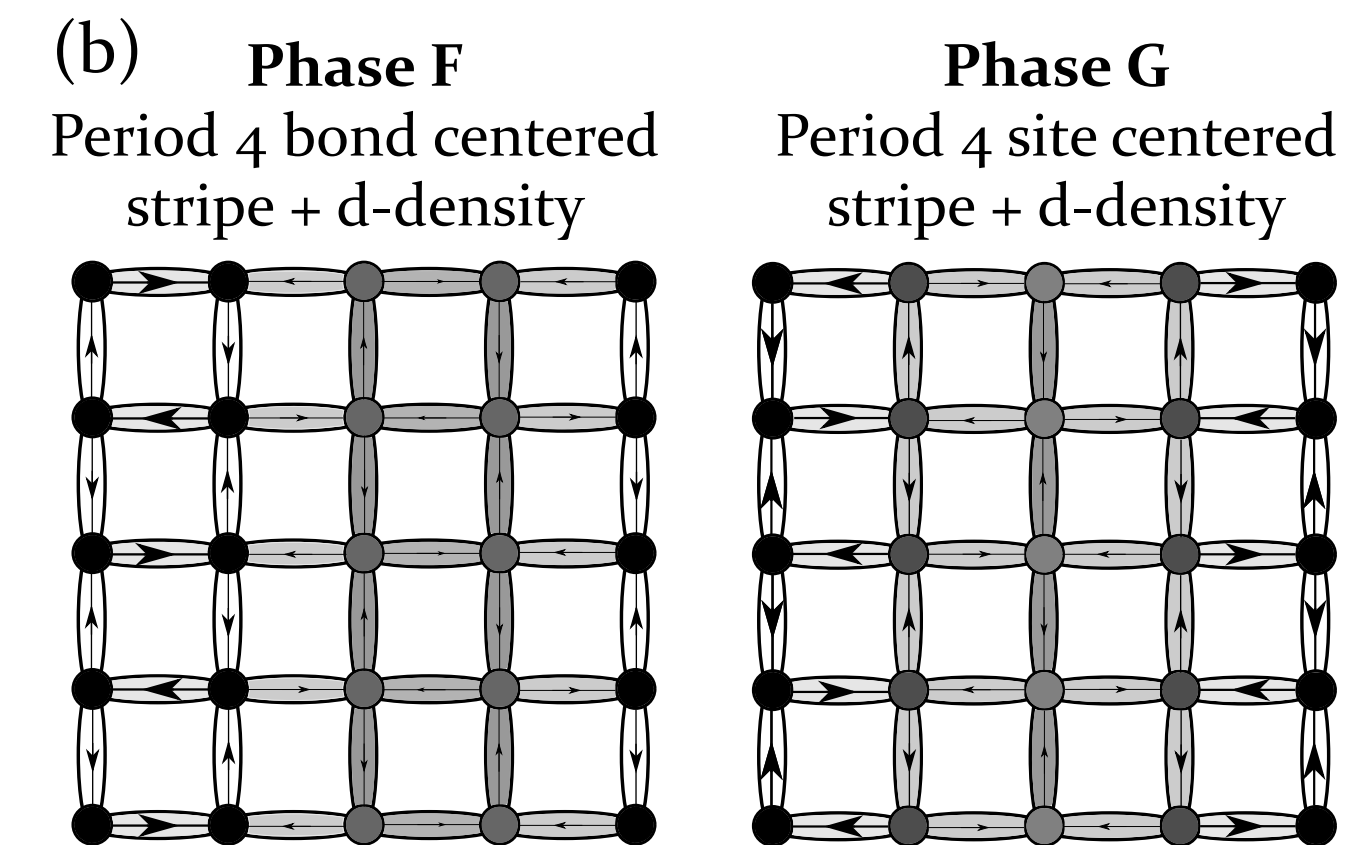
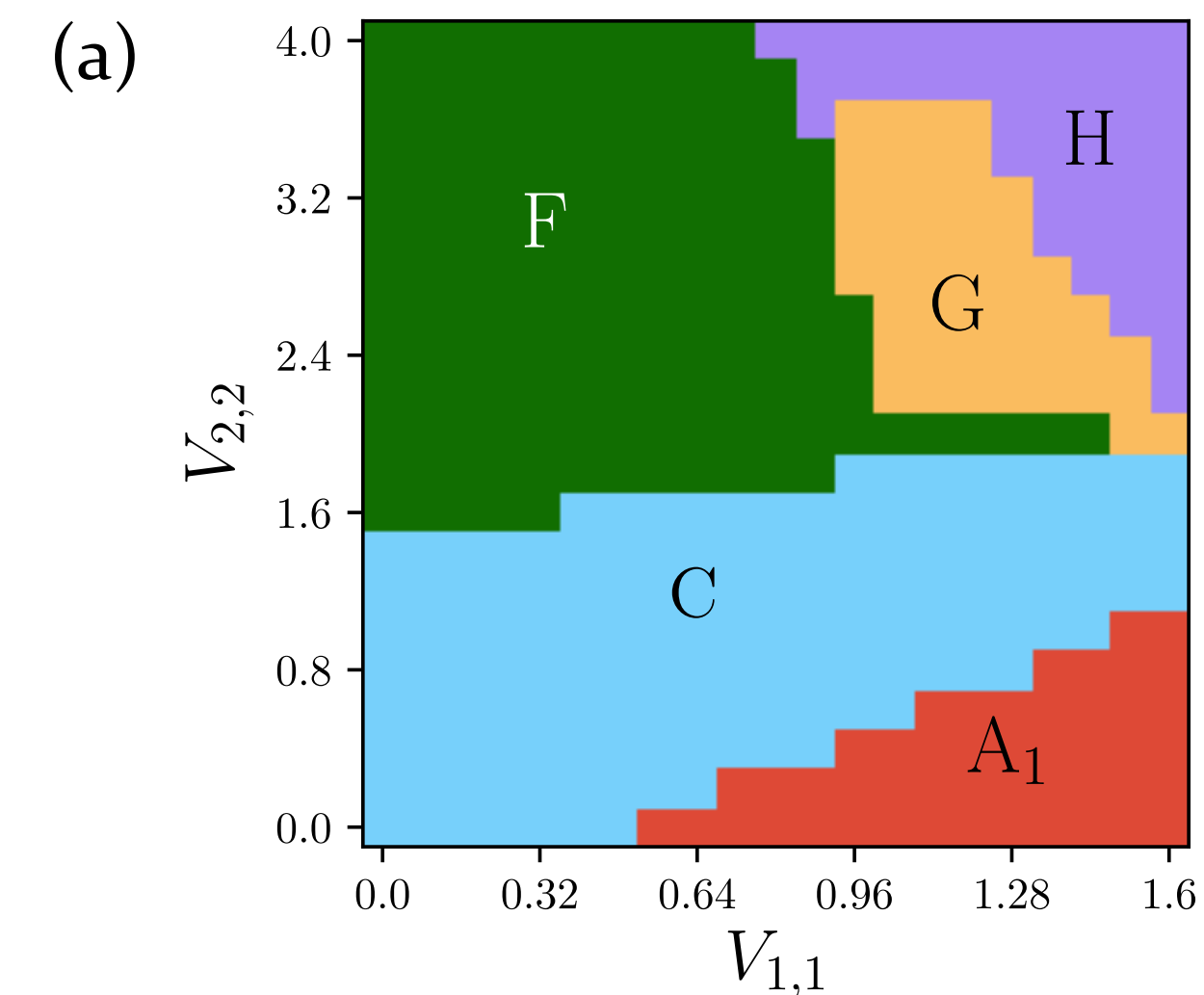
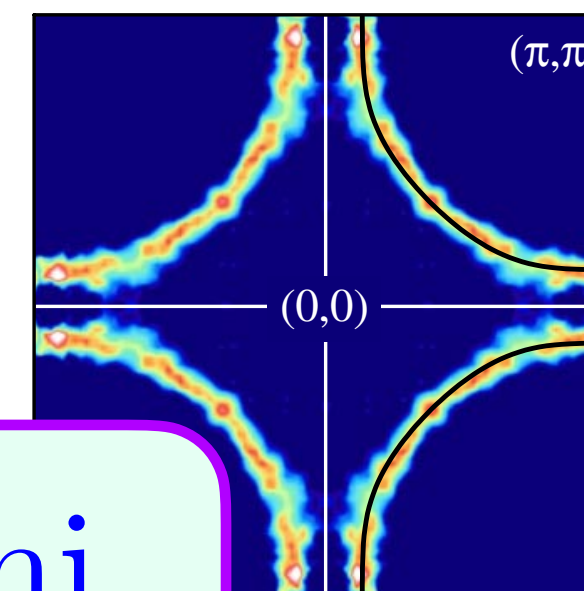
Pseudogap Metal (FL\*)  
with hole pocket Fermi surfaces  
and underlying  $\pi$ -flux spin liquid

Strange  
Metal.

Metal with  
 $(\pi, \pi)$  SDW  
(Néel) order.

$d$ -wave superconductivity.  
Charge density wave.

Fermi  
Liquid.  
FL



$T$

$p$

B

A

C