

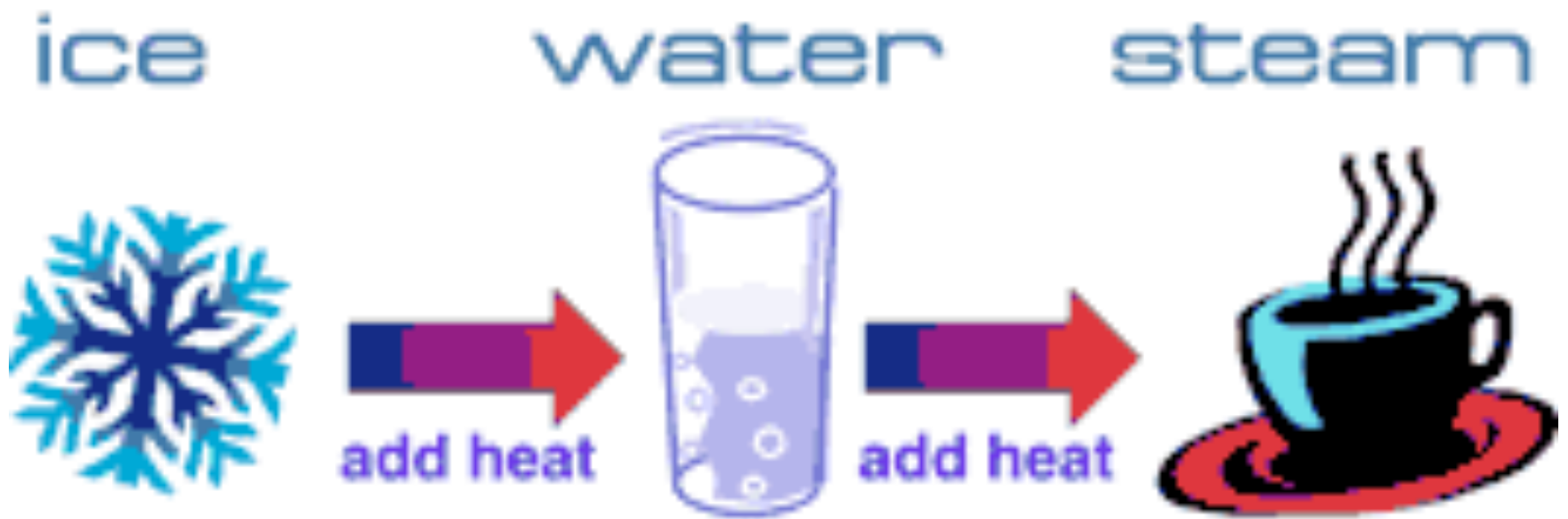
Quantum Criticality and Black Holes

Subir Sachdev

Talk online at <http://sachdev.physics.harvard.edu>



What is a “phase transition” ?



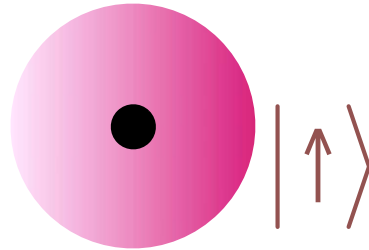
A change in the collective properties of a macroscopic number of atoms

What is a “*quantum phase transition*” ?

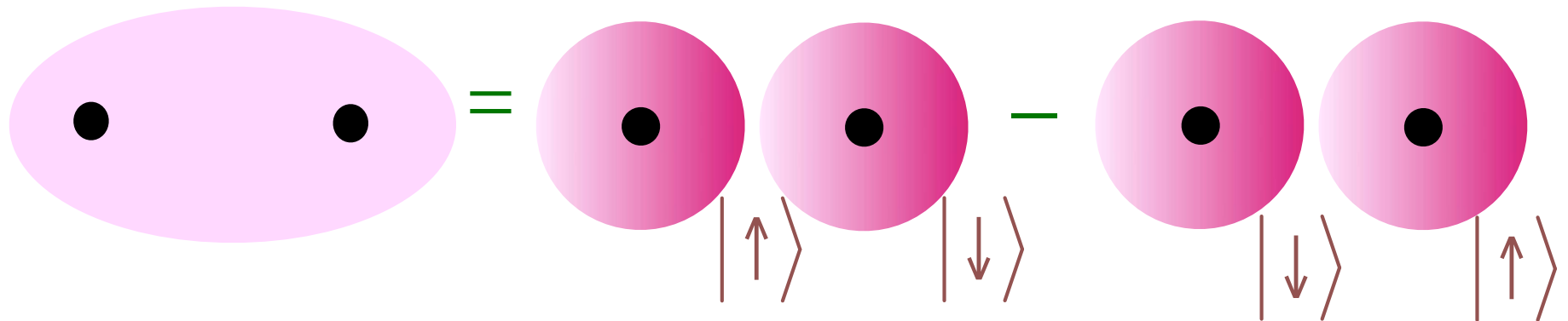
Change in the nature of
entanglement in a macroscopic
quantum system.

Entanglement

Hydrogen atom:



Hydrogen molecule:



$$= \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Superposition of two electron states leads to non-local correlations between spins

Outline

1. Superfluid-insulator quantum transitions

Experiments on ultracold atoms

2. Theory of quantum-critical transport

Collisionless- t_0 -hydrodynamic crossover of conformal field theories

3. Entanglement of valence bonds

Deconfined criticality in antiferromagnets

4. Nernst effect in the cuprate superconductors

Quantum criticality and dyonic black holes

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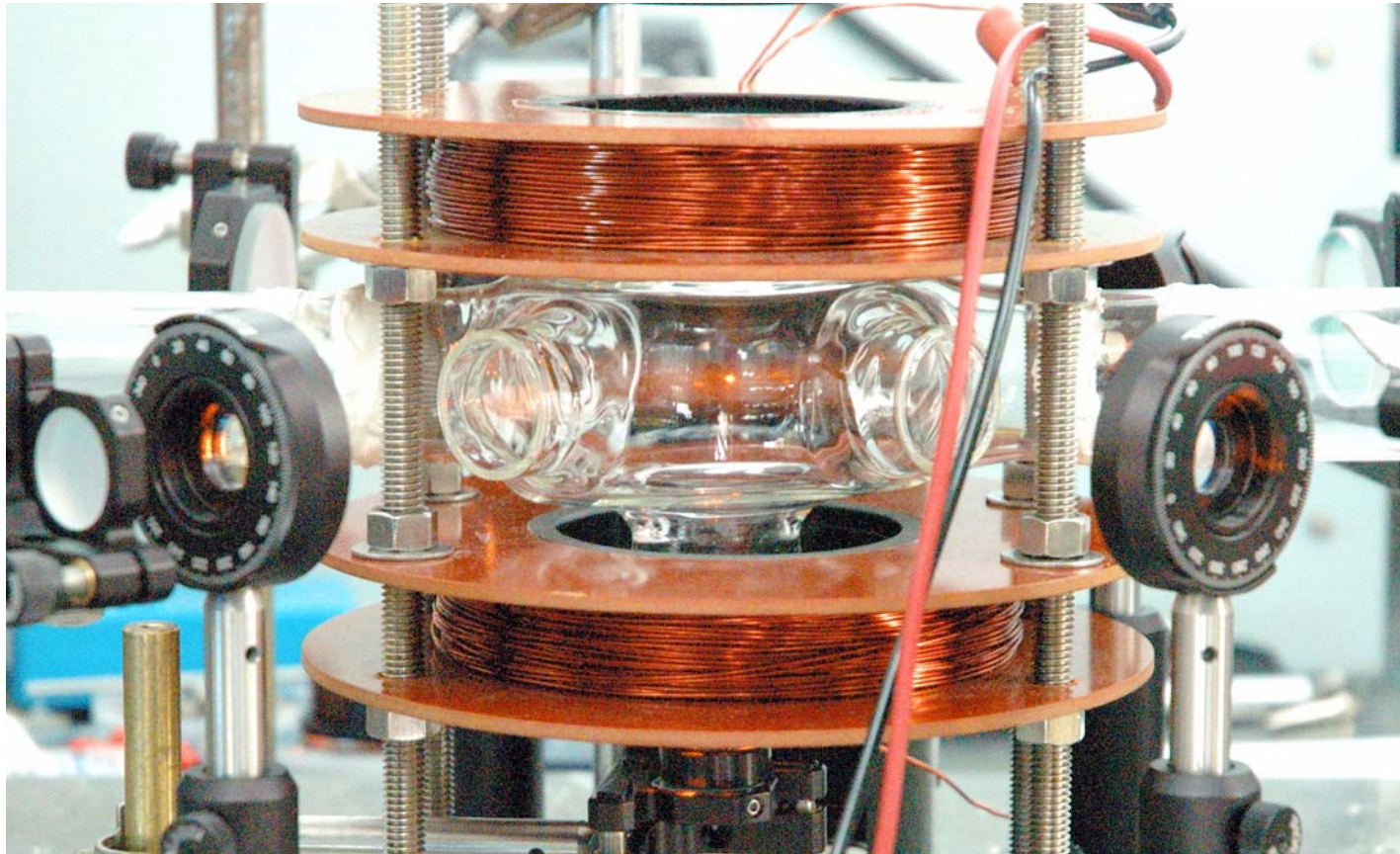
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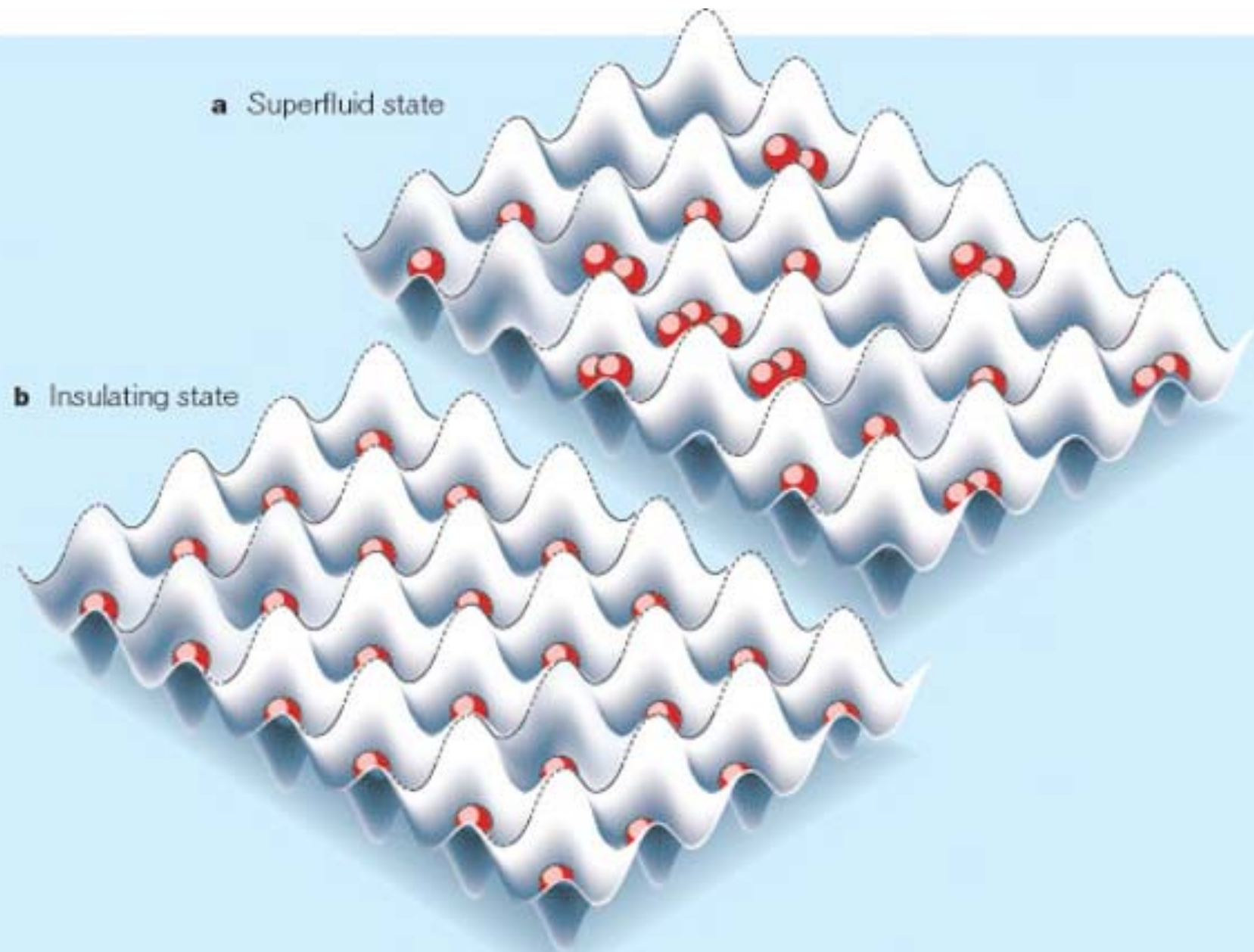
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Quantum criticality and dyonic black holes



Trap for ultracold ^{87}Rb atoms



M. Greiner, O. Mandel, T. Esslinger, T. W. Hänsch, and I. Bloch, *Nature* **415**, 39 (2002).

The Bose-Einstein condensate in a periodic potential

$$|G\rangle = \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle$$

Lowest energy state for many atoms

$$\begin{aligned} |BEC\rangle &= |G\rangle |G\rangle |G\rangle \\ &= \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle \\ &\quad + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle + \dots 27 \text{ terms} \end{aligned}$$

Large fluctuations in number of atoms in each potential well
 – *superfluidity* (atoms can “flow” without dissipation)

Breaking up the Bose-Einstein condensate

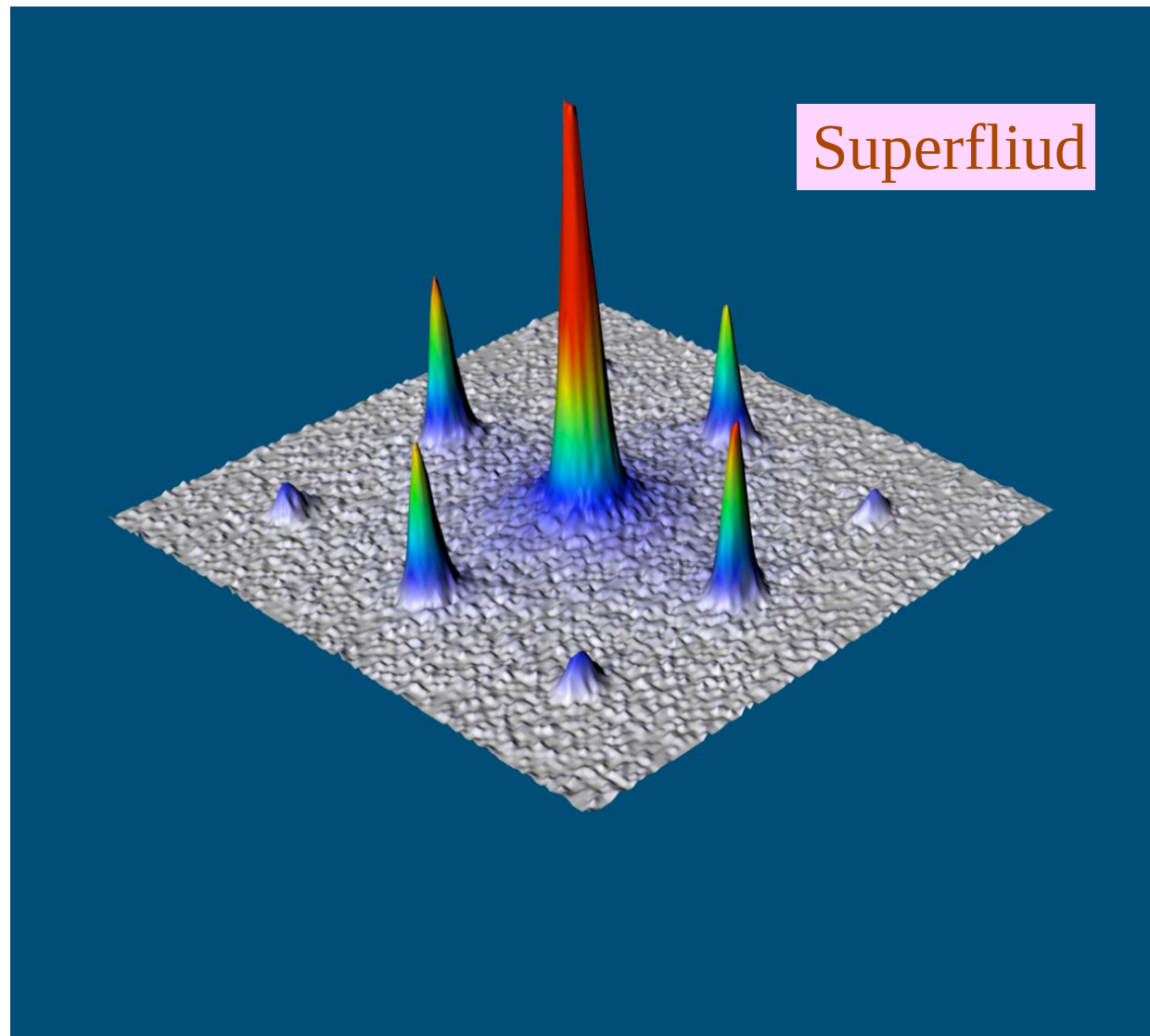
$$|G\rangle = \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right\rangle + \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \\ \hline \end{array} \right| \left| \begin{array}{|c|} \hline \bigcirc \\ \hline \end{array} \right\rangle$$

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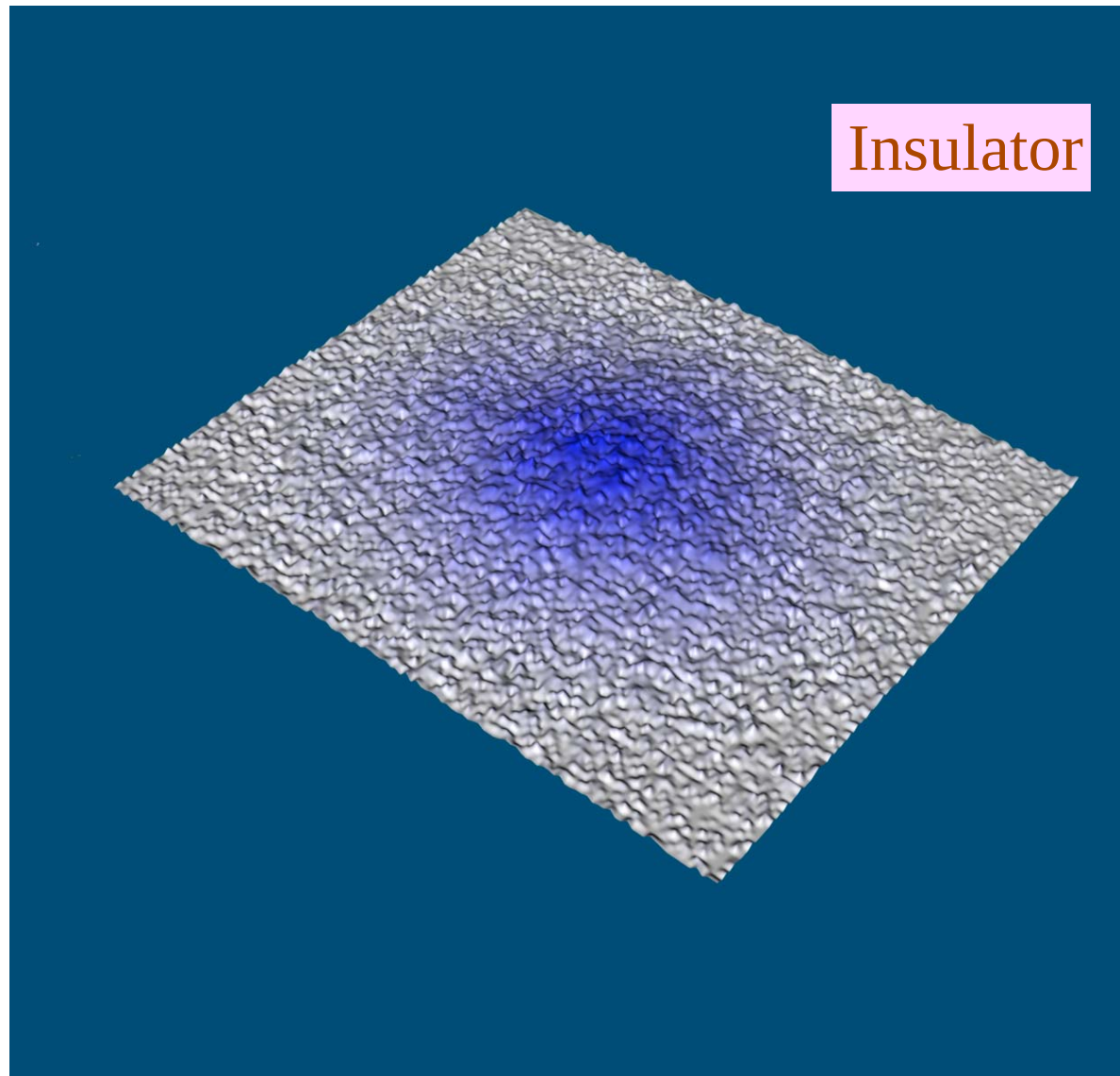
By tuning repulsive interactions between the atoms, states with multiple atoms in a potential well can be suppressed. The lowest energy state is then a *Mott insulator* – it has negligible number fluctuations, and atoms cannot “flow”

Velocity distribution of ^{87}Rb atoms



M. Greiner, O. Mandel, T. Esslinger, T. W. Hänsch, and I. Bloch, *Nature* **415**, 39 (2002).

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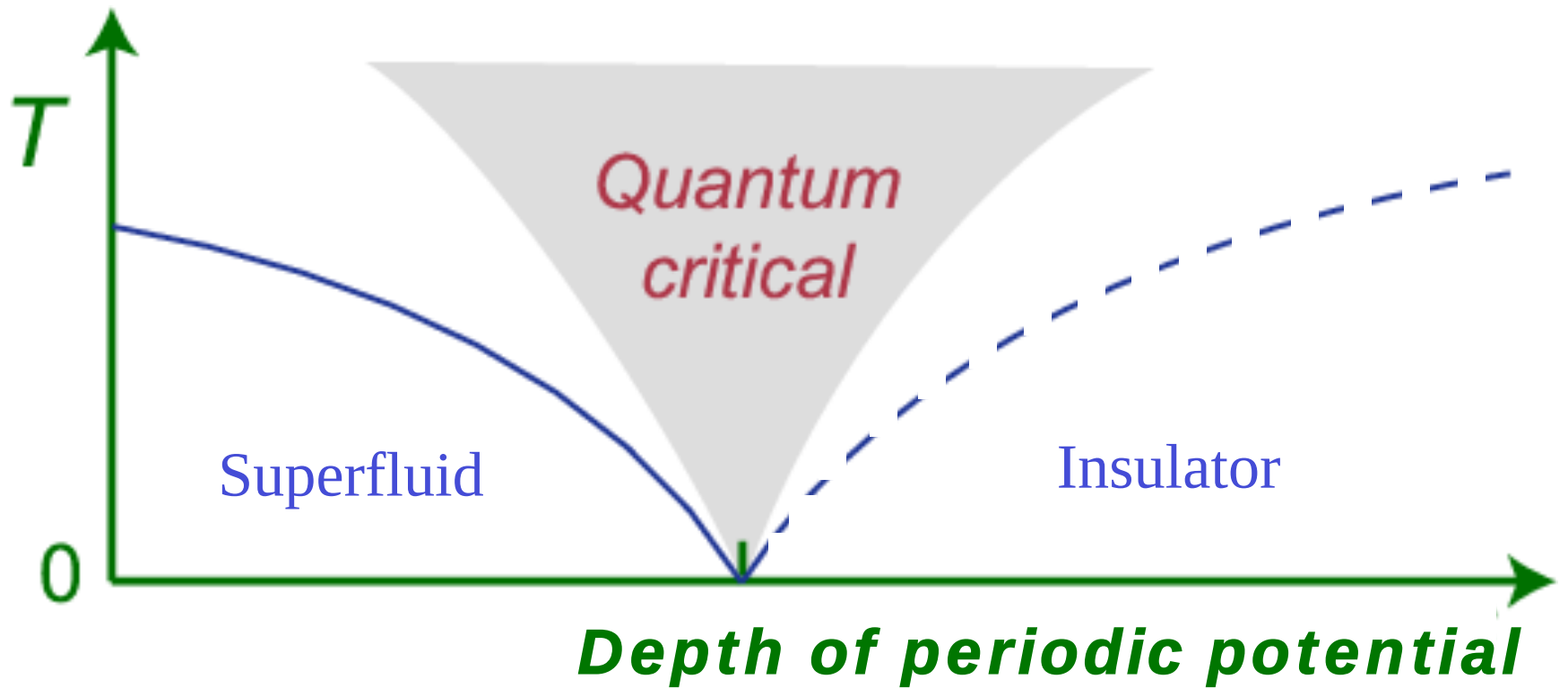
Superfluid

Insulator

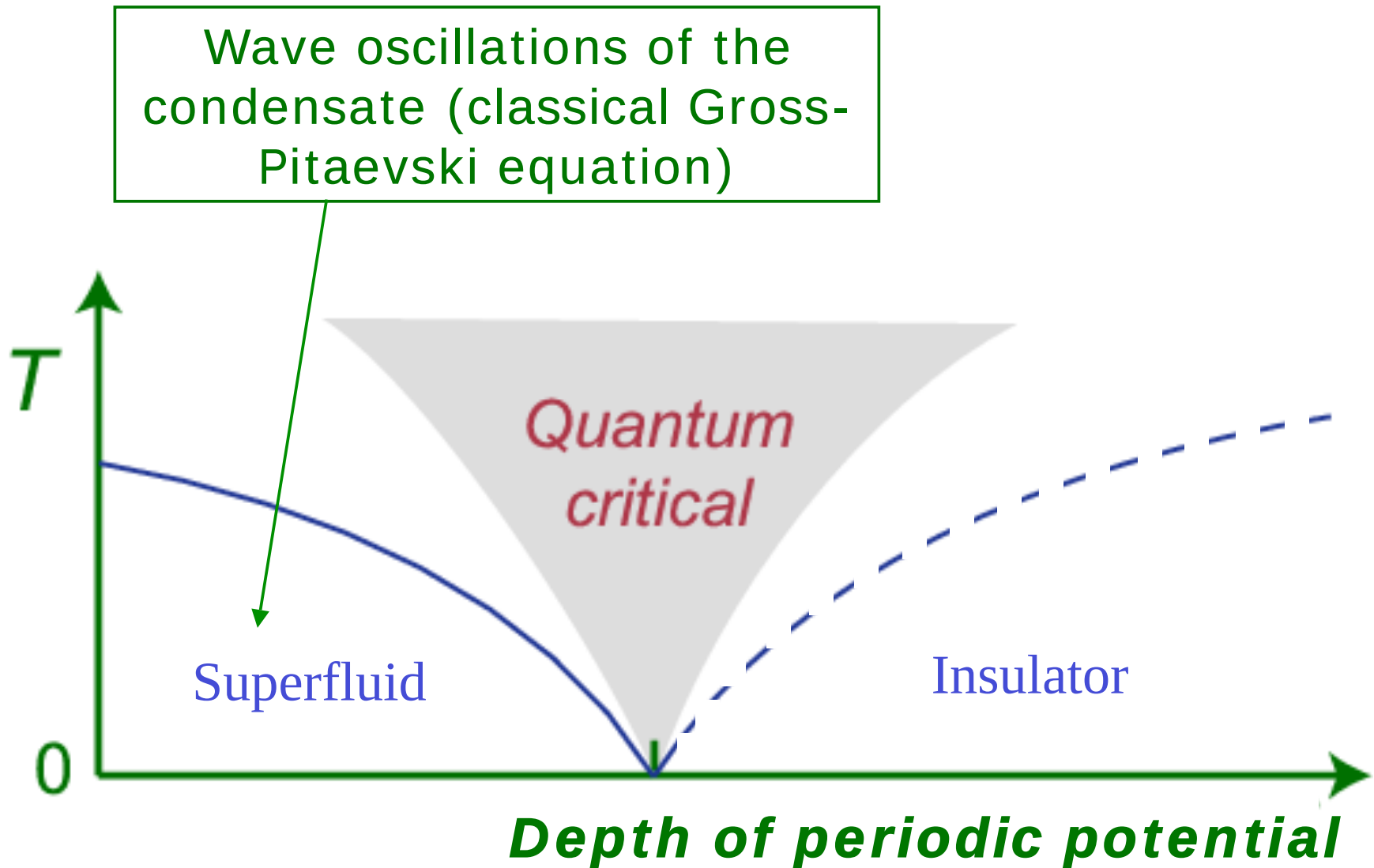
Depth of periodic potential



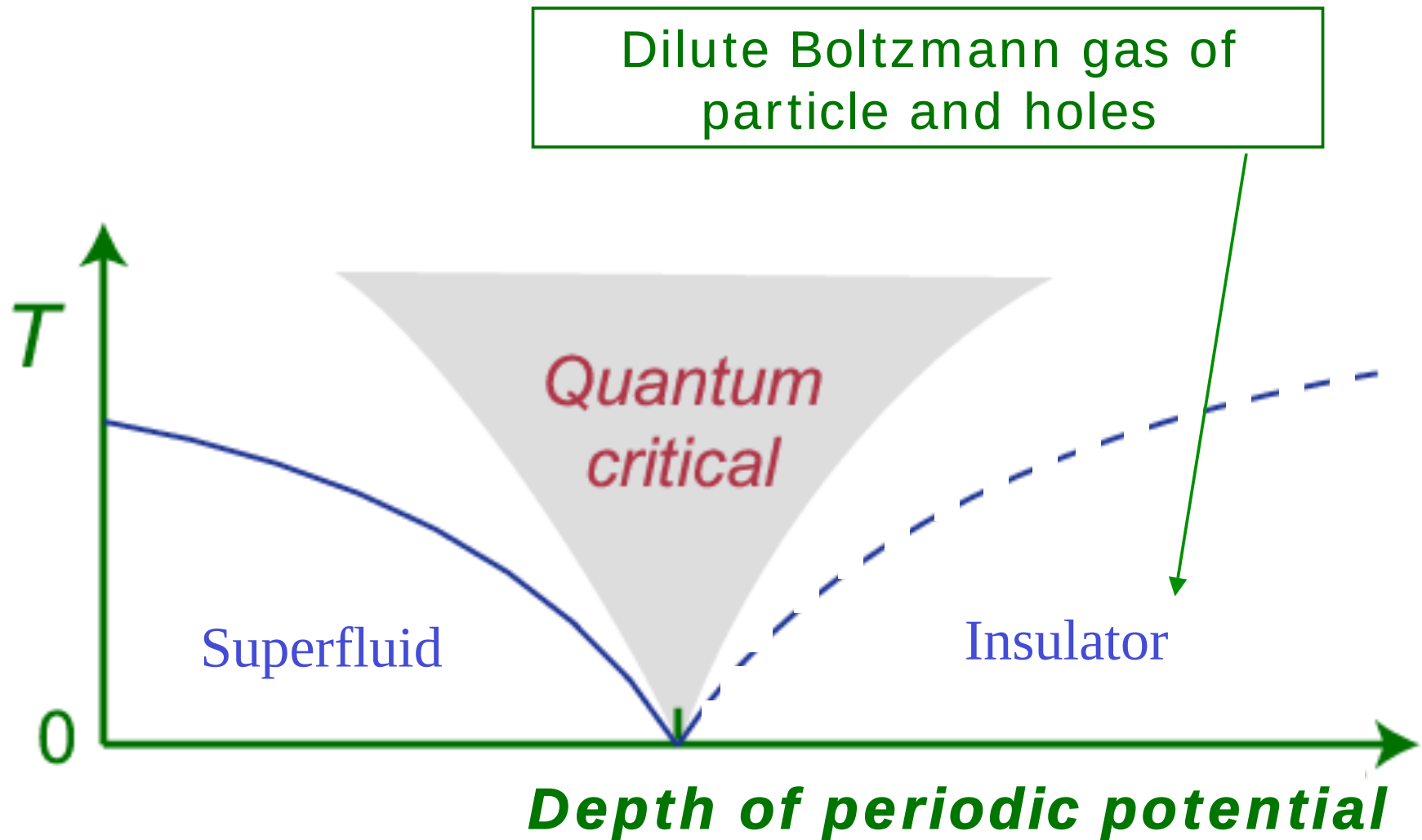
Non-zero temperature phase diagram



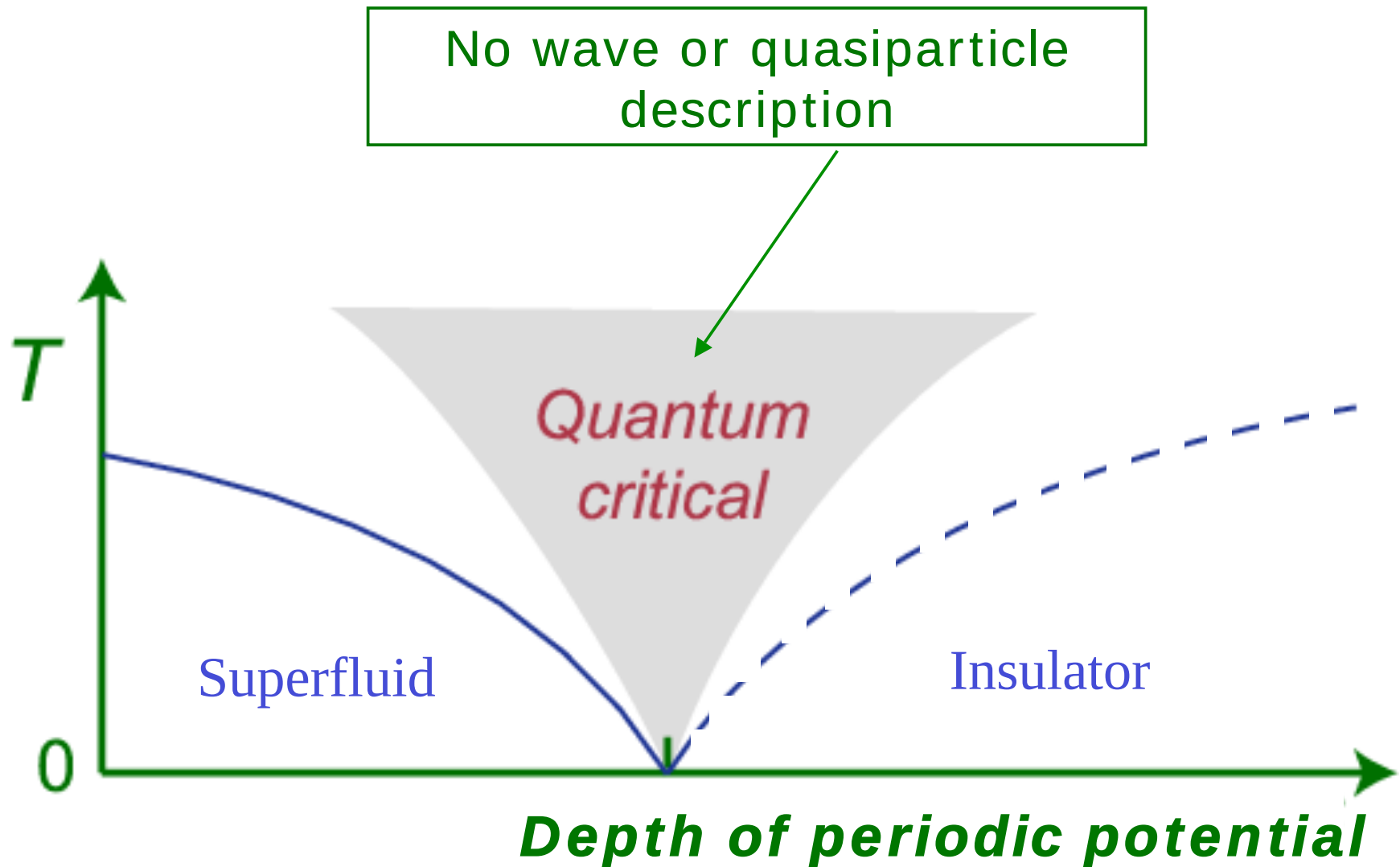
Non-zero temperature phase diagram



Non-zero temperature phase diagram



Non-zero temperature phase diagram



Resistivity of Bi films

Conductivity σ

$$\sigma_{\text{Superconductor}}(T \rightarrow 0) = \infty$$

$$\sigma_{\text{Insulator}}(T \rightarrow 0) = 0$$

$$\sigma_{\text{Quantum critical point}}(T \rightarrow 0) \approx \frac{4e^2}{h}$$

D. B. Haviland, Y. Liu, and A. M. Goldman,
Phys. Rev. Lett. **62**, 2180 (1989)

M. P. A. Fisher, *Phys. Rev. Lett.* **65**, 923 (1990)

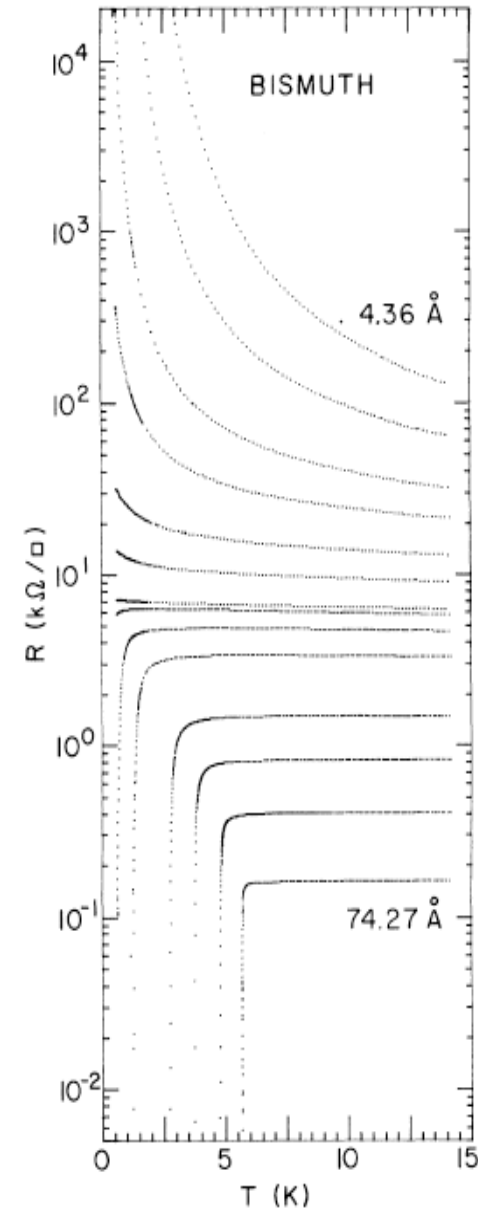
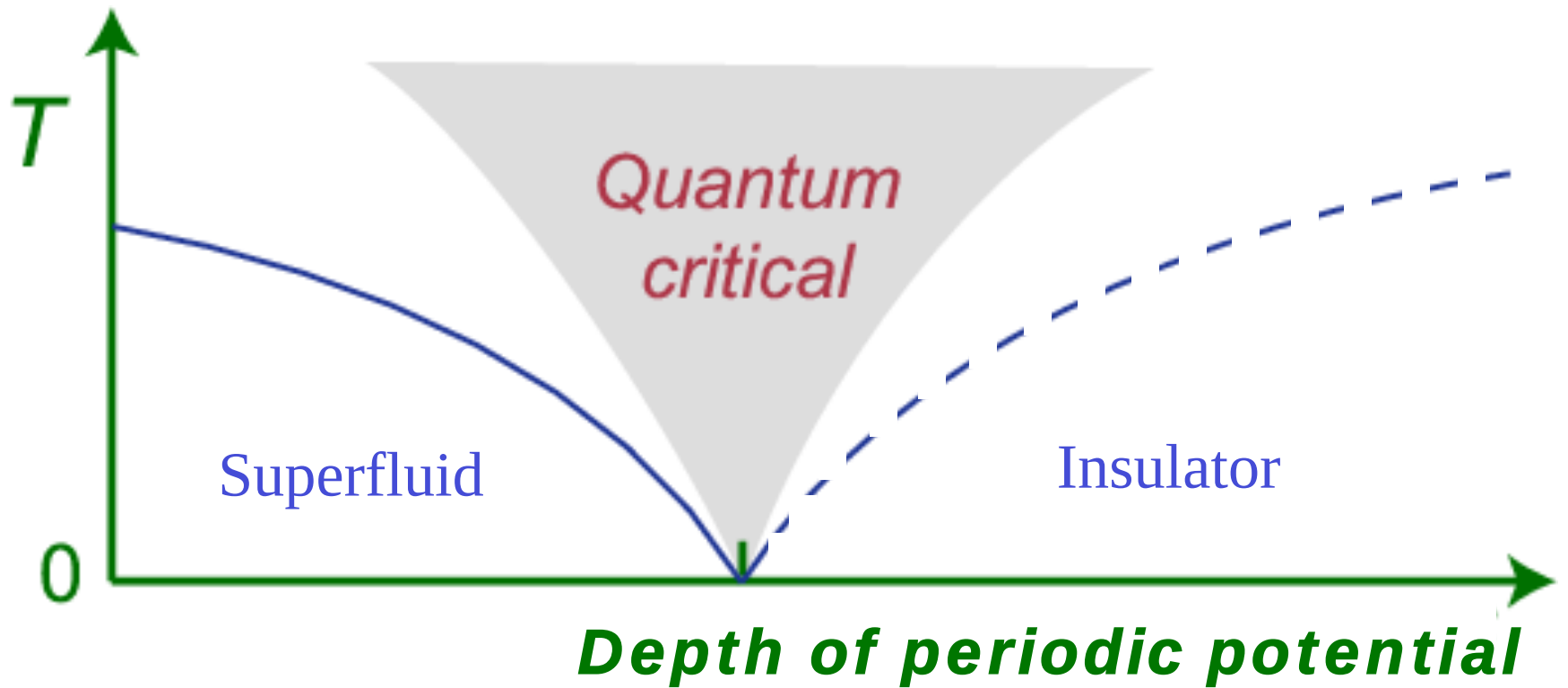
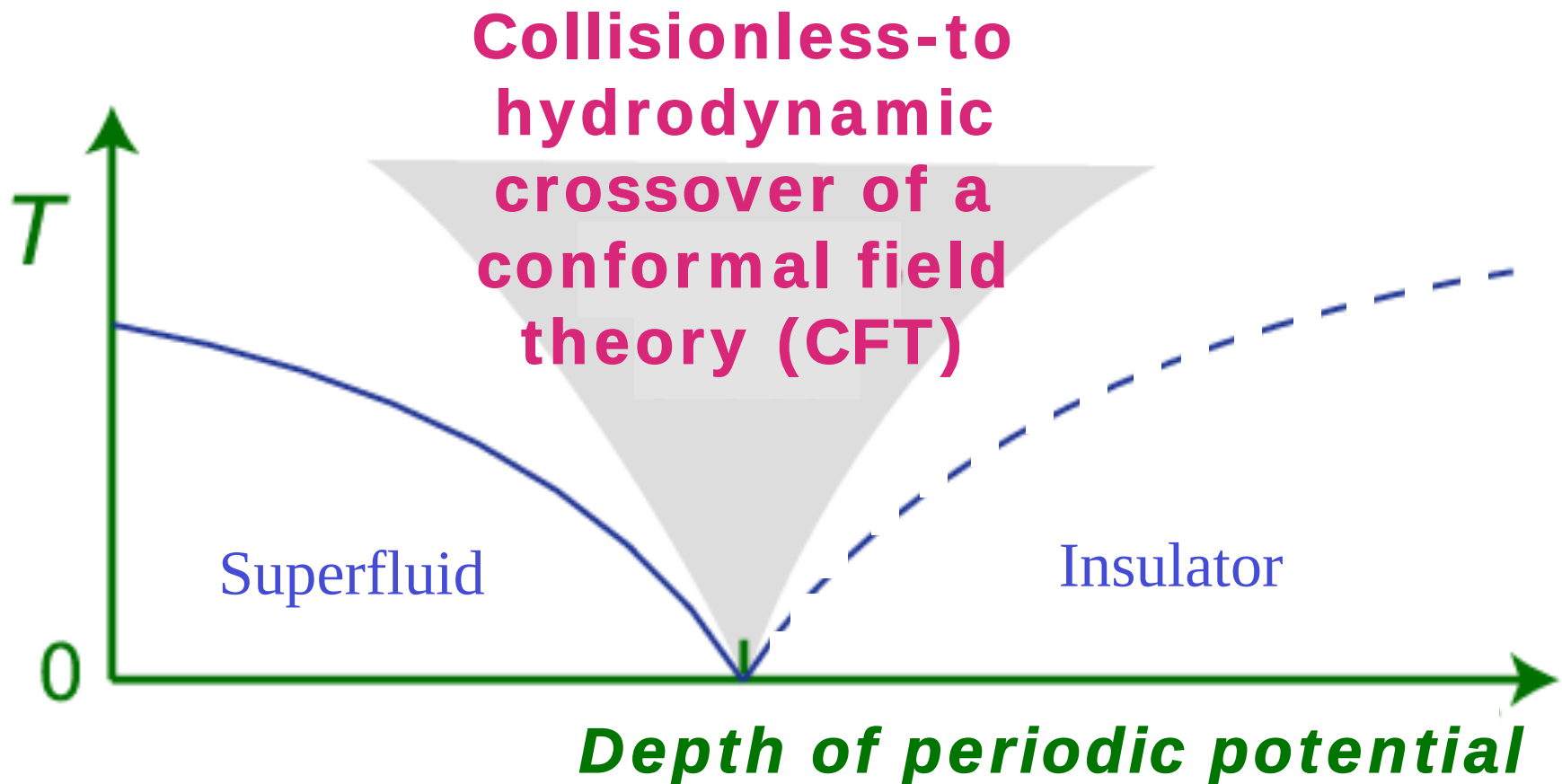


FIG. 1. Evolution of the temperature dependence of the sheet resistance $R(T)$ with thickness for a Bi film deposited onto Ge. Fewer than half of the traces actually acquired are shown. Film thicknesses shown range from 4.36 to 74.27 Å.

Non-zero temperature phase diagram



Non-zero temperature phase diagram



K. Damle and S. Sachdev, *Phys. Rev. B* **56**, 8714 (1997).

Hydrodynamics of a conformal field theory (CFT)

The scattering cross-section of the thermal excitations is universal and so transport coefficients are universally determined by $k_B T$

Charge diffusion constant $D_c = \Theta \frac{c^2}{k_B T}$

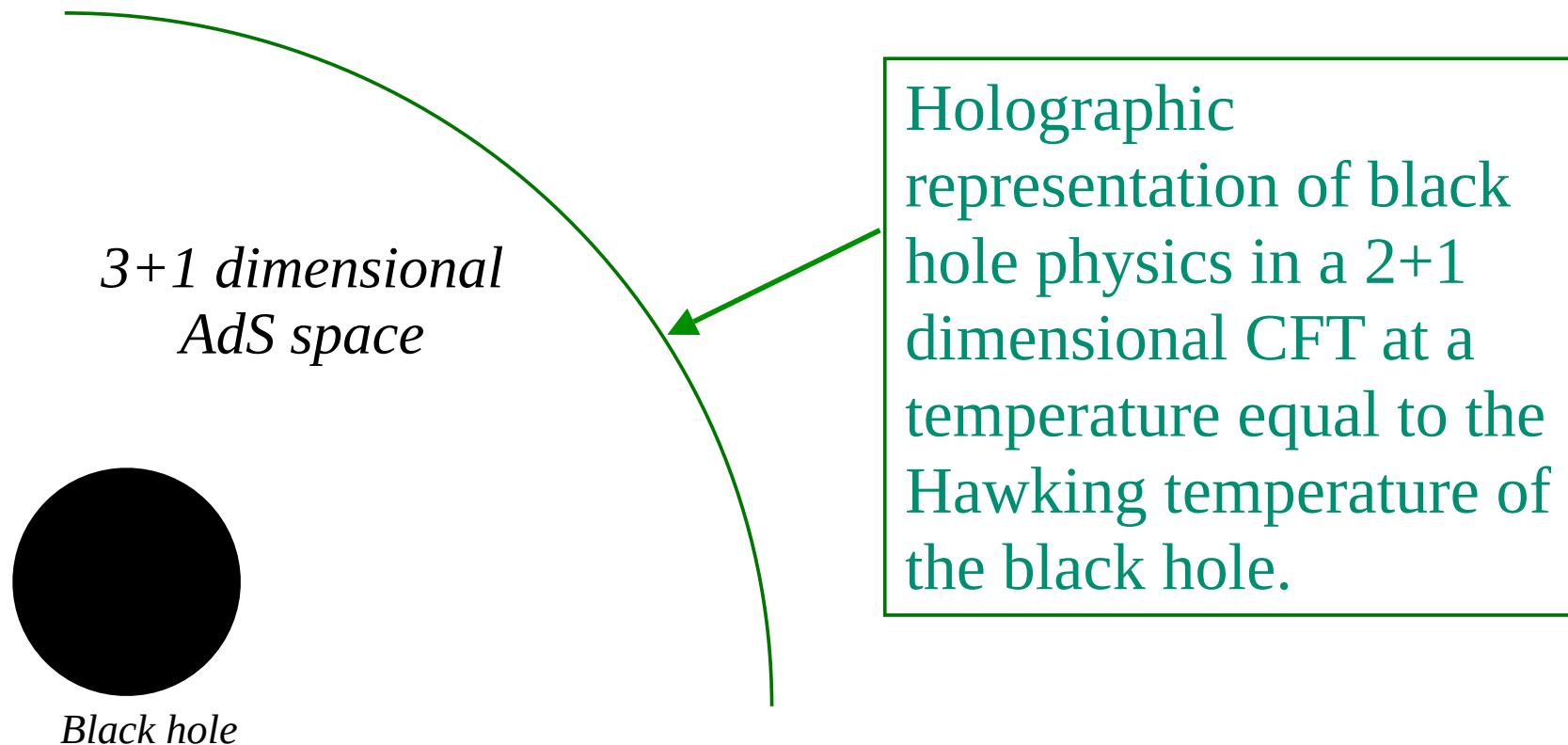
Conductivity $\sigma_Q = D_c \chi = \overline{\Theta} \frac{4e^2}{h}$

Hydrodynamics of a conformal field theory (CFT)

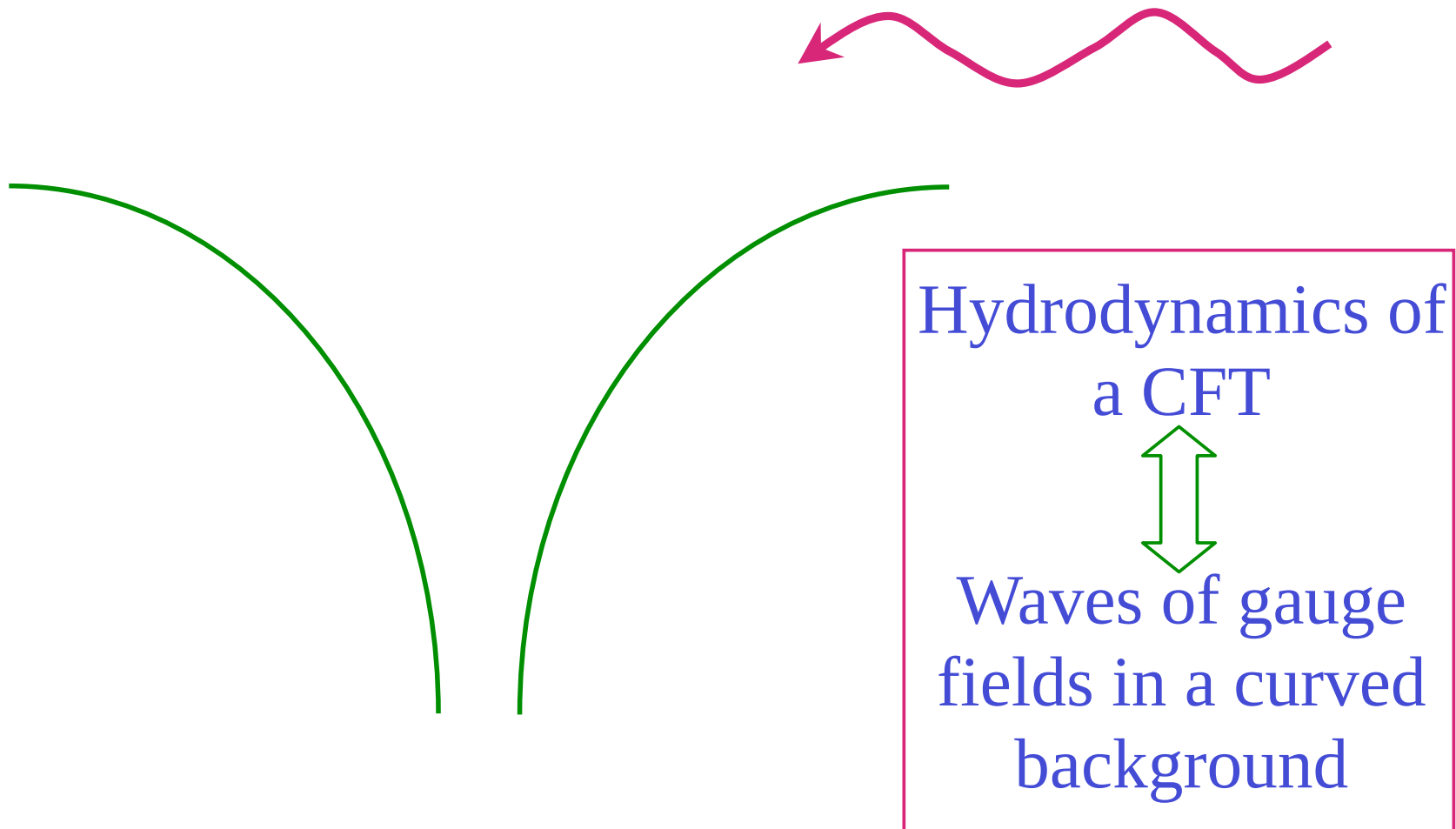
The AdS/CFT correspondence (Maldacena, Polyakov) relates the hydrodynamics of CFTs to the quantum gravity theory of the horizon of a black hole in Anti-de Sitter space.

Hydrodynamics of a conformal field theory (CFT)

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Hydrodynamics of a conformal field theory (CFT)



Hydrodynamics of a conformal field theory (CFT)

For the (unique) CFT with a $SU(N)$ gauge field and 16 supercharges, we know the exact diffusion constant associated with a global $SO(8)$ symmetry:

Spin diffusion constant $D_c = \frac{3}{4\pi} \frac{c^2}{k_B T}$

Spin conductivity $\sigma_Q = \frac{N^{3/2}}{3\sqrt{2}\pi}$

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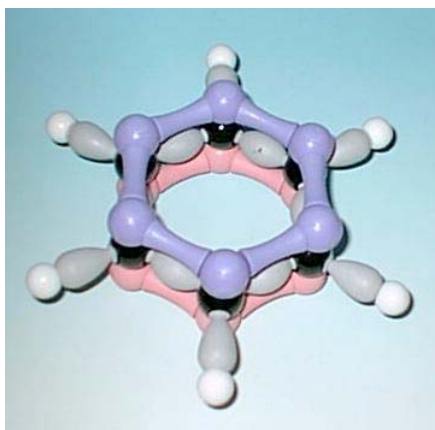
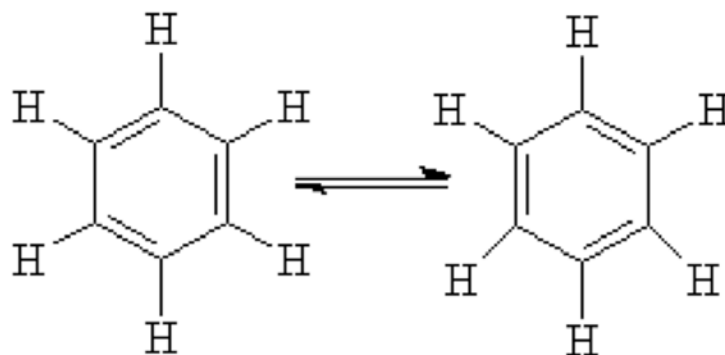
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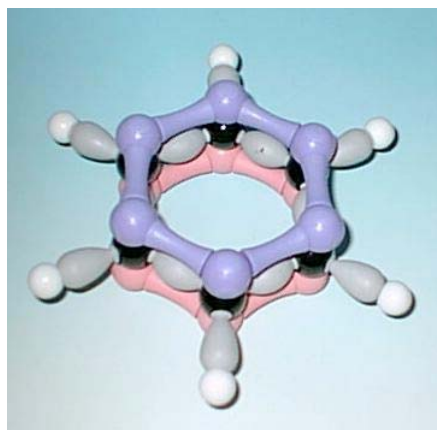
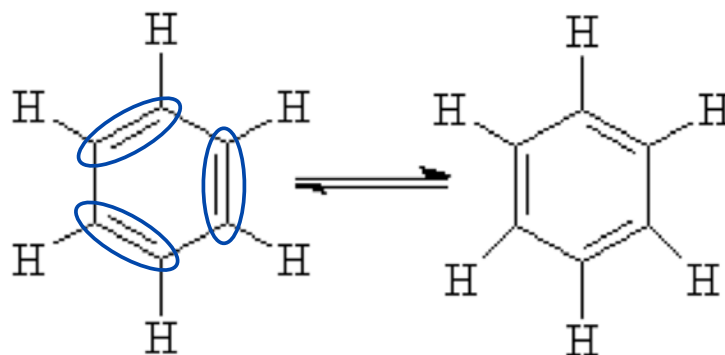
Quantum criticality and dyonic black holes

Valence bonds in benzene



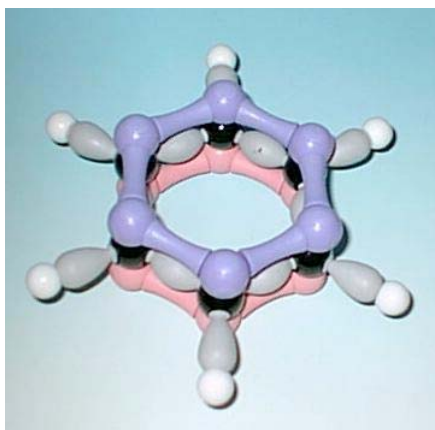
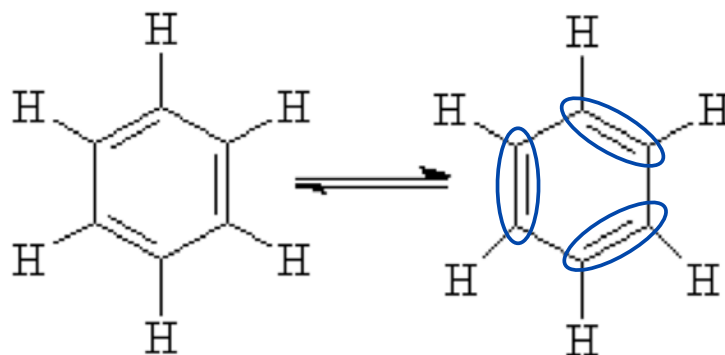
Resonance in benzene leads to a symmetric configuration of valence bonds
(*F. Kekulé, L. Pauling*)

Valence bonds in benzene



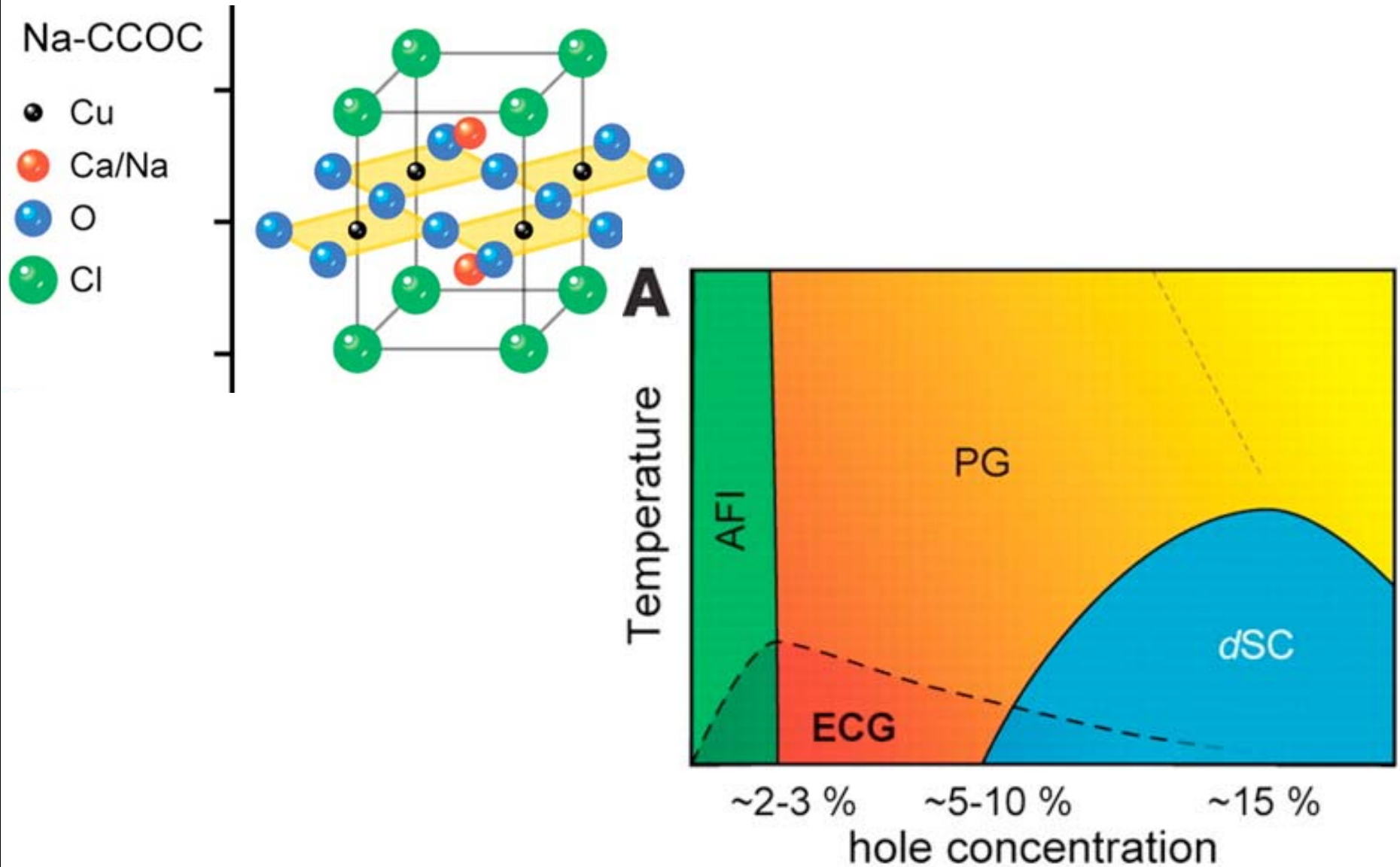
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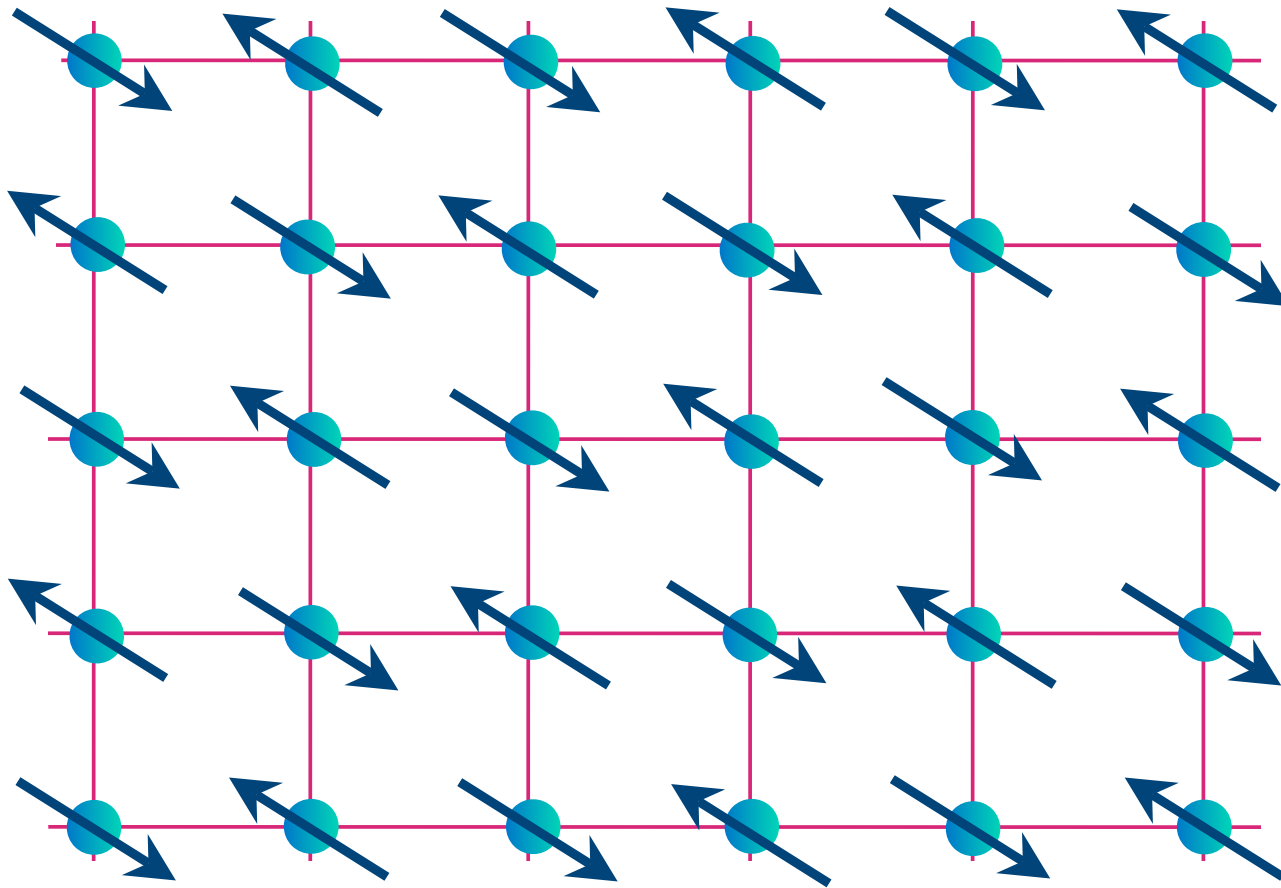


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Temperature-doping phase diagram of the cuprate superconductors

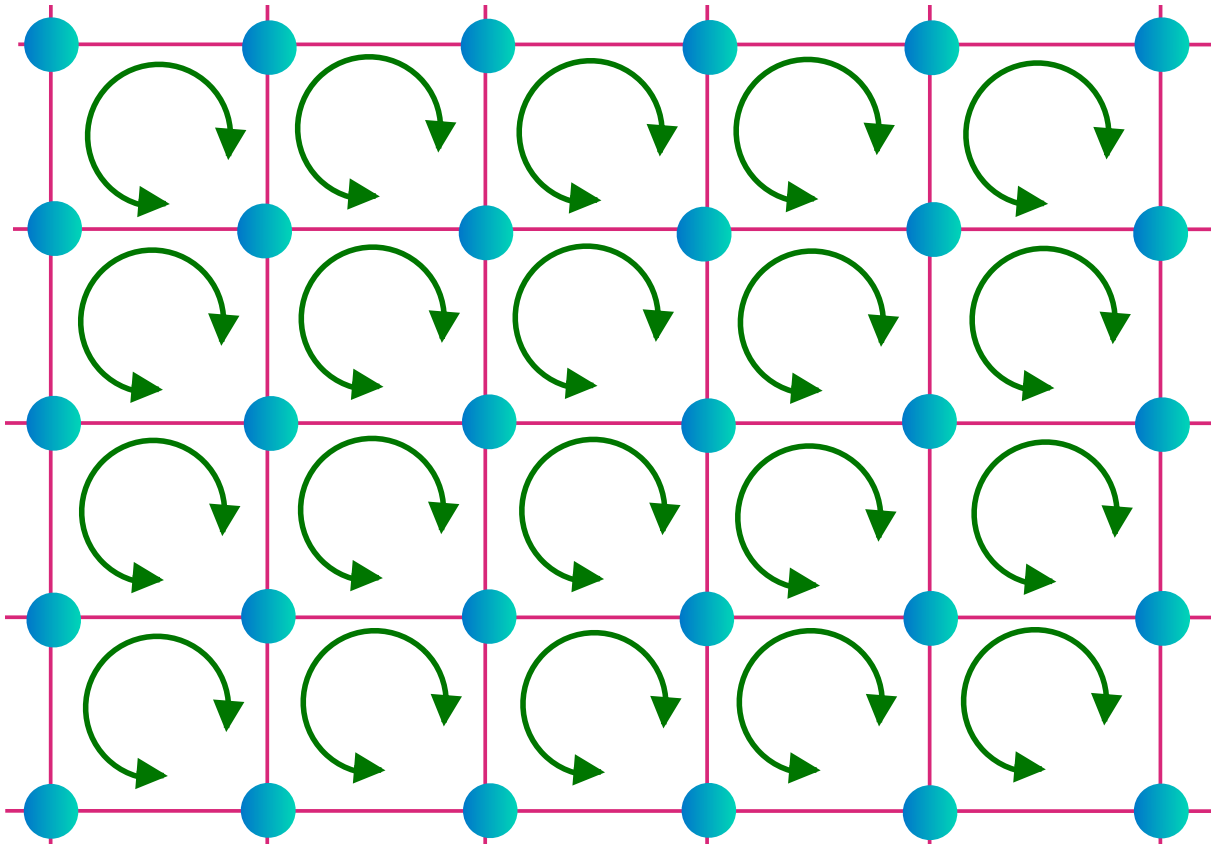


Antiferromagnetic (Neel) order in the insulator



$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j \quad ; \quad \vec{S}_i \Rightarrow \text{spin operator with } S = 1/2$$

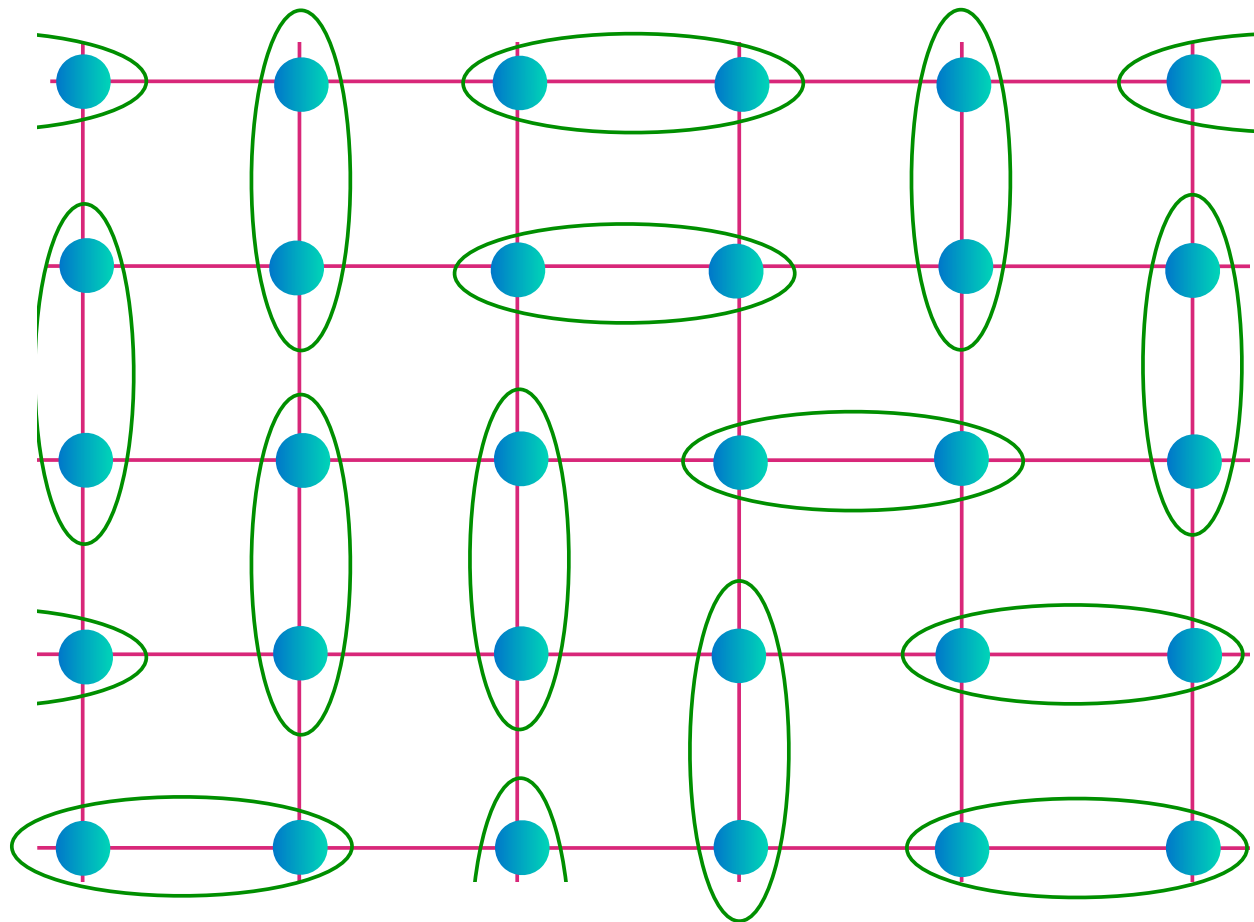
Induce formation of valence bonds by *e.g.*
ring-exchange interactions



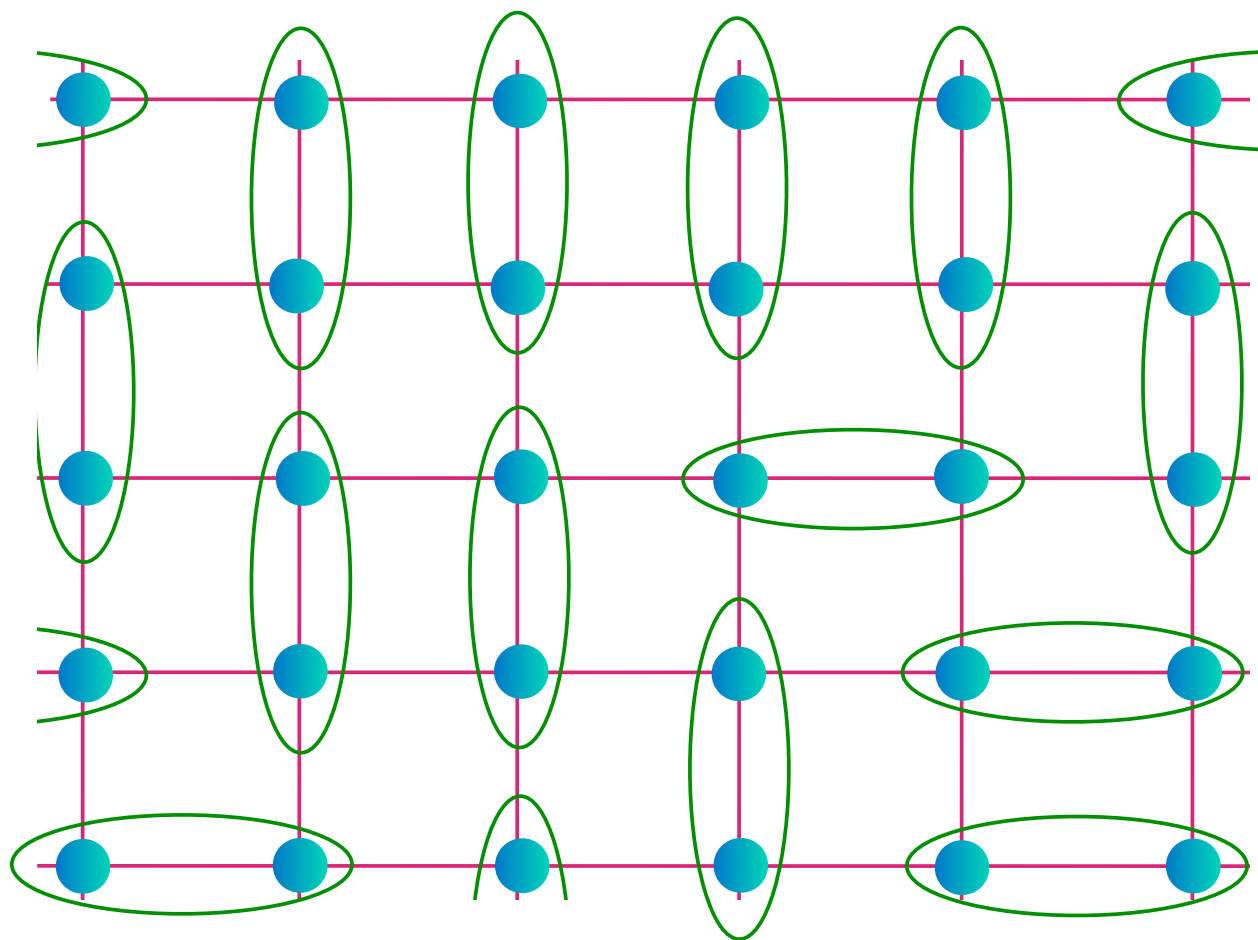
$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} \text{four spin exchange}$$

A. W. Sandvik, cond-mat/0611343

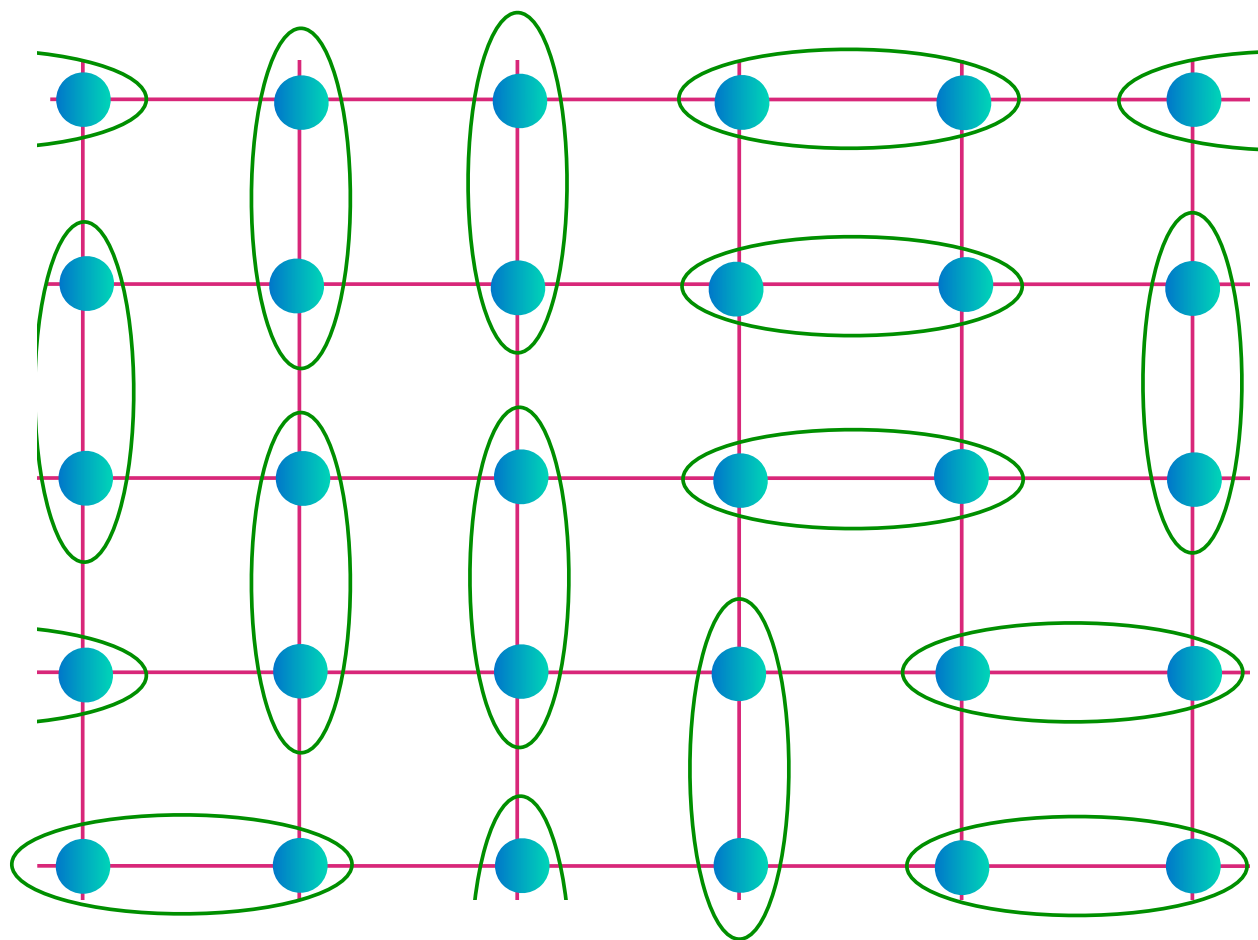
As in H_2 and benzene, each electron wants to pair up with another electron and form a valence bond



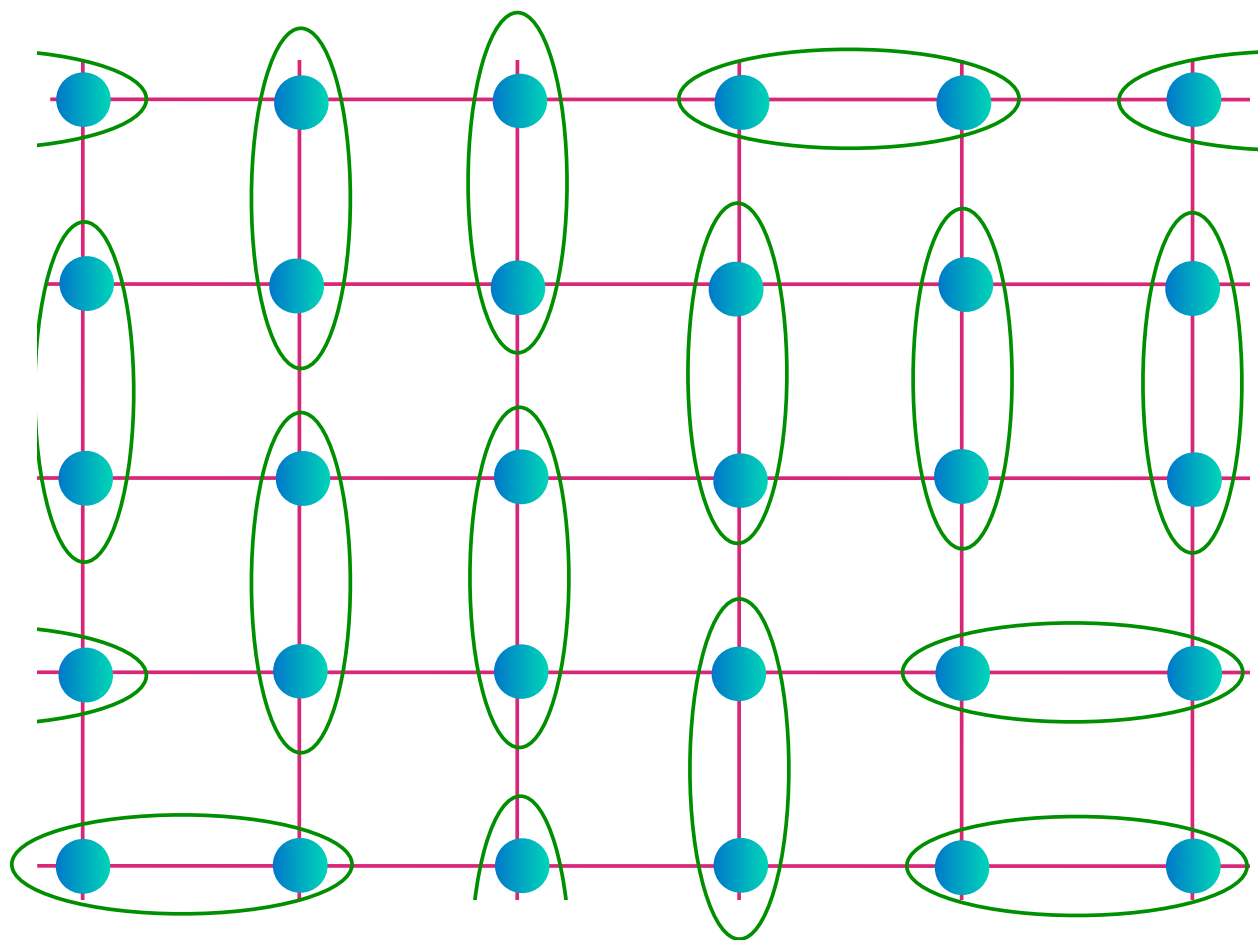
$$\text{[Two blue dots in a green oval]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$



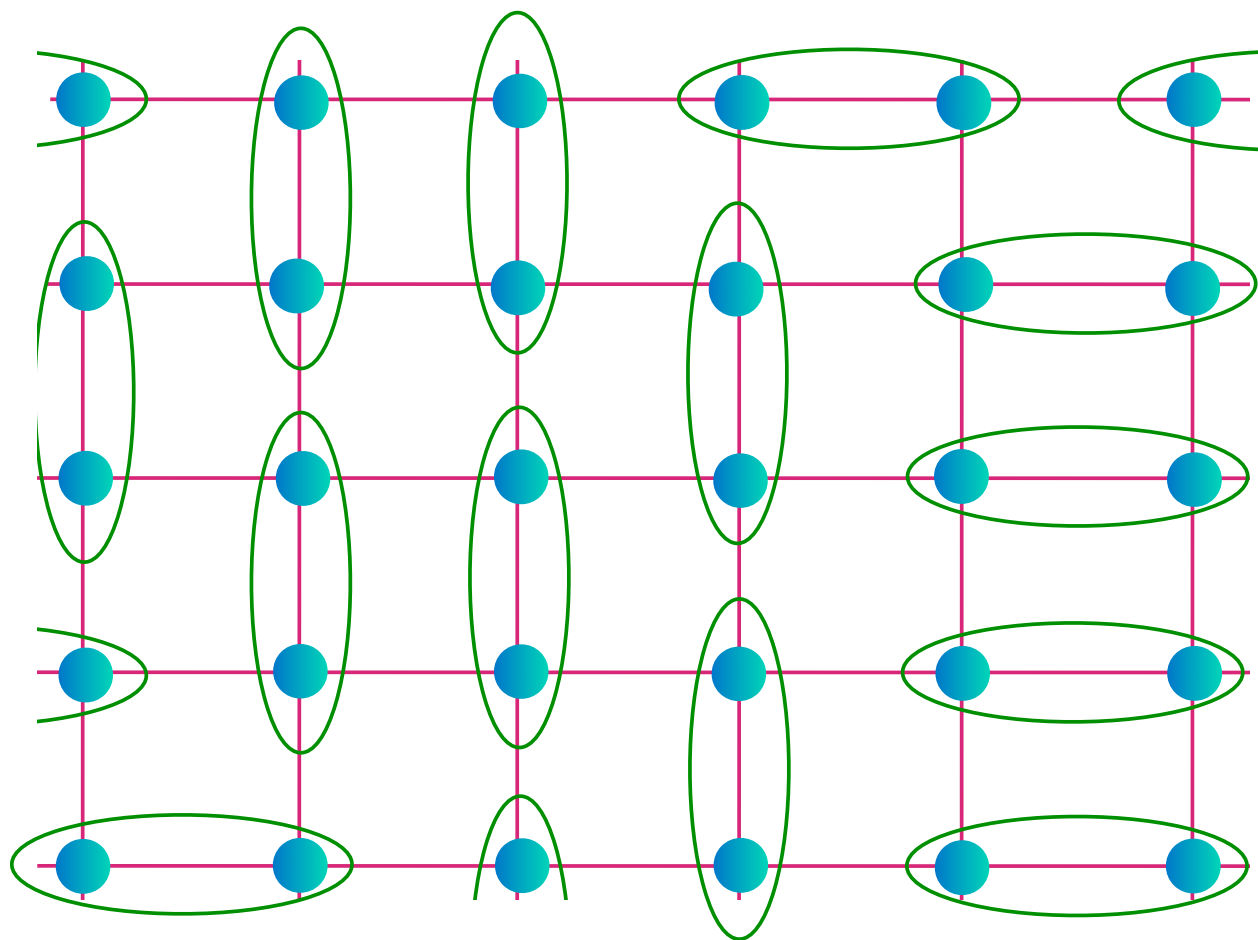
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$$\frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$

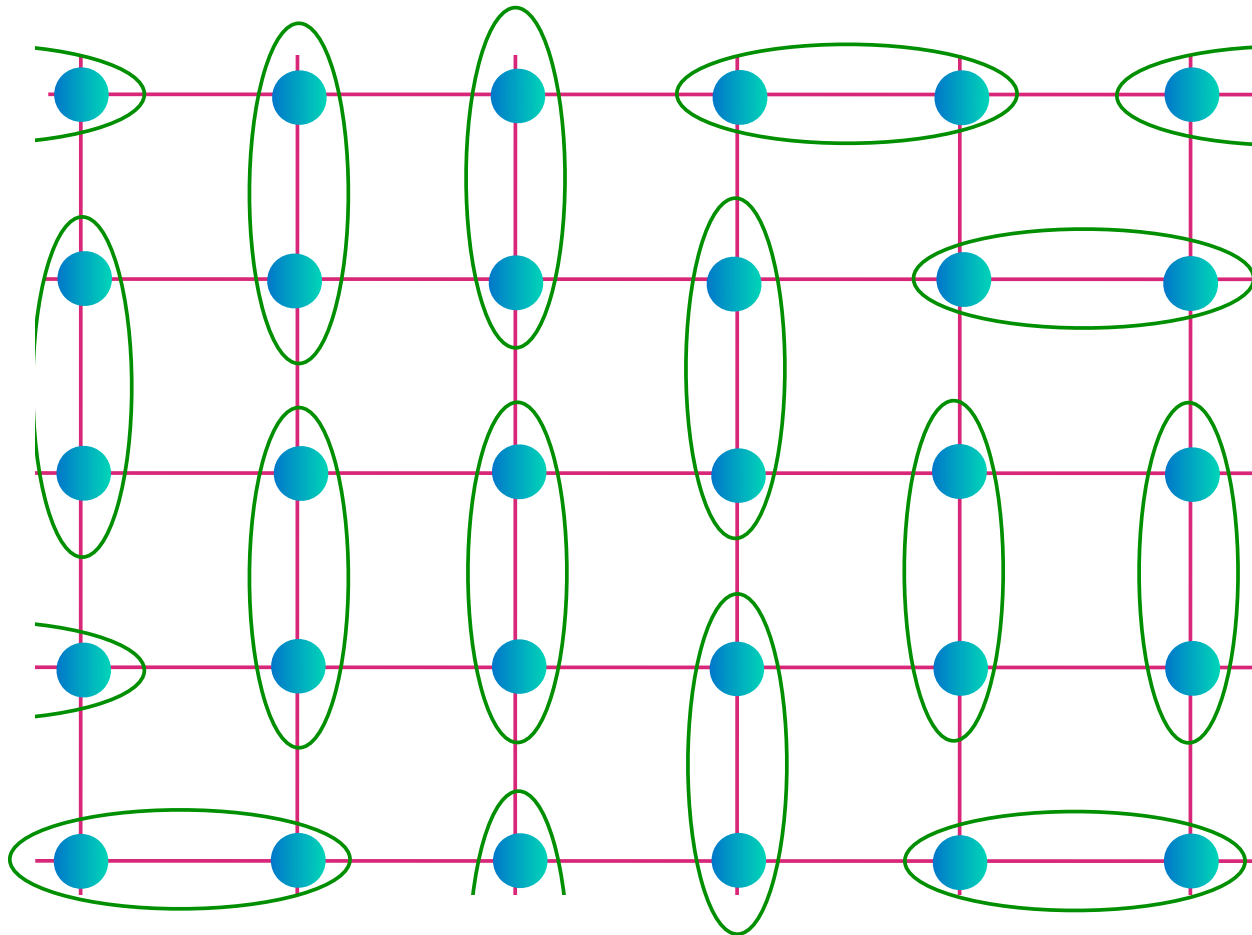


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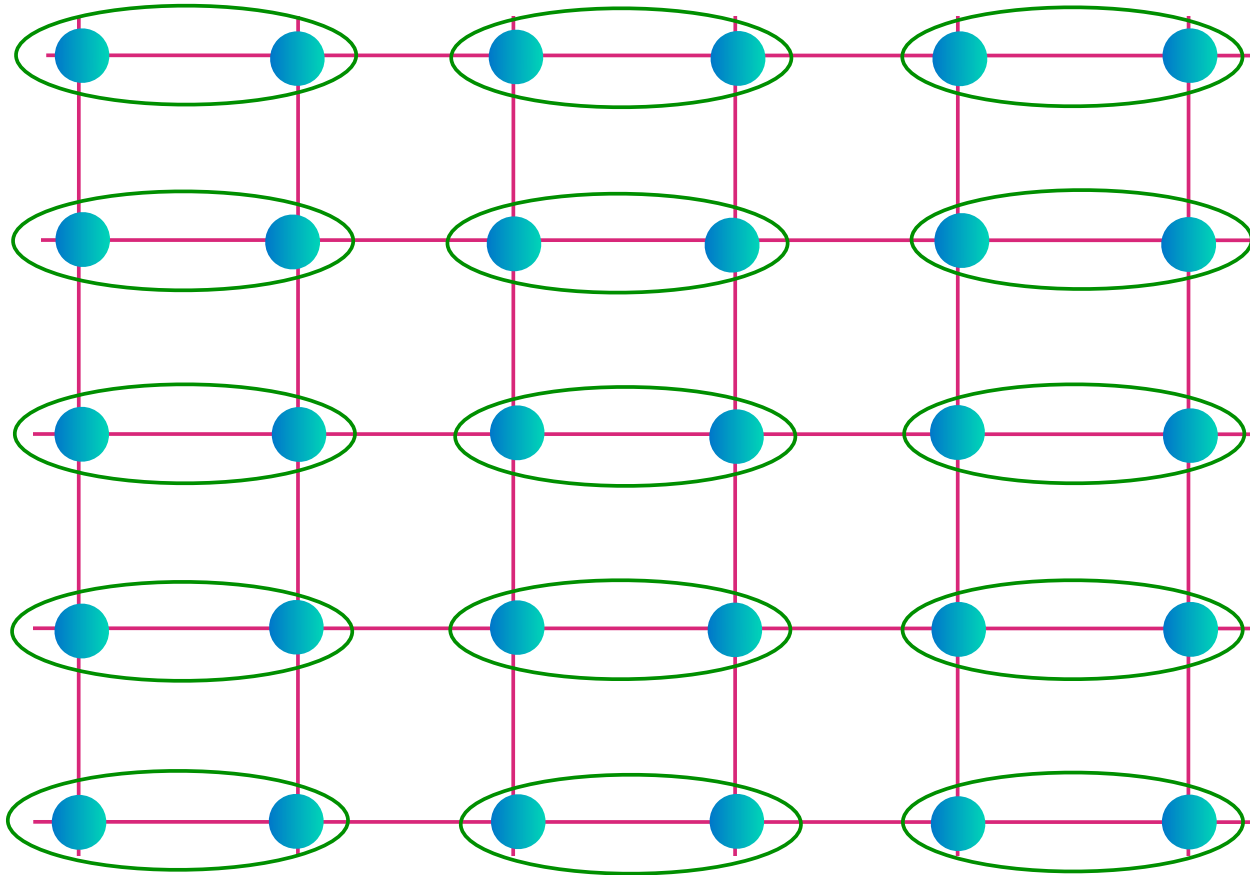
$$\frac{1}{\sqrt{2}} \left(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right)$$

Entangled liquid of valence bonds (Resonating valence bonds – RVB)



$$\begin{array}{|c|c|} \hline \bullet & \bullet \\ \hline \end{array} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Valence bond solid (VBS)

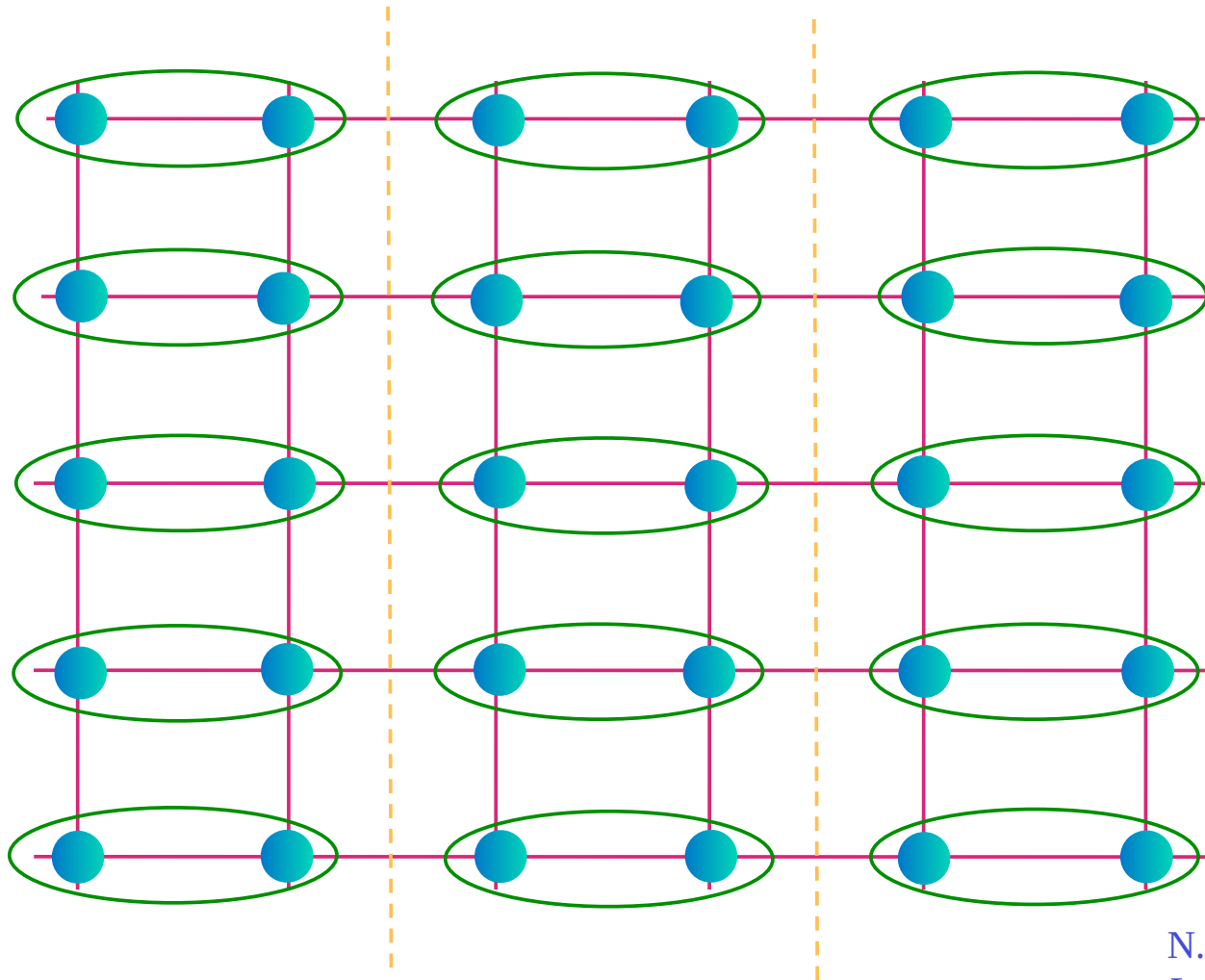


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N. Read and S. Sachdev, *Phys. Rev. Lett.* **62**, 1694 (1989).

R. Moessner and S. L. Sondhi, *Phys. Rev. B* **63**, 224401 (2001).

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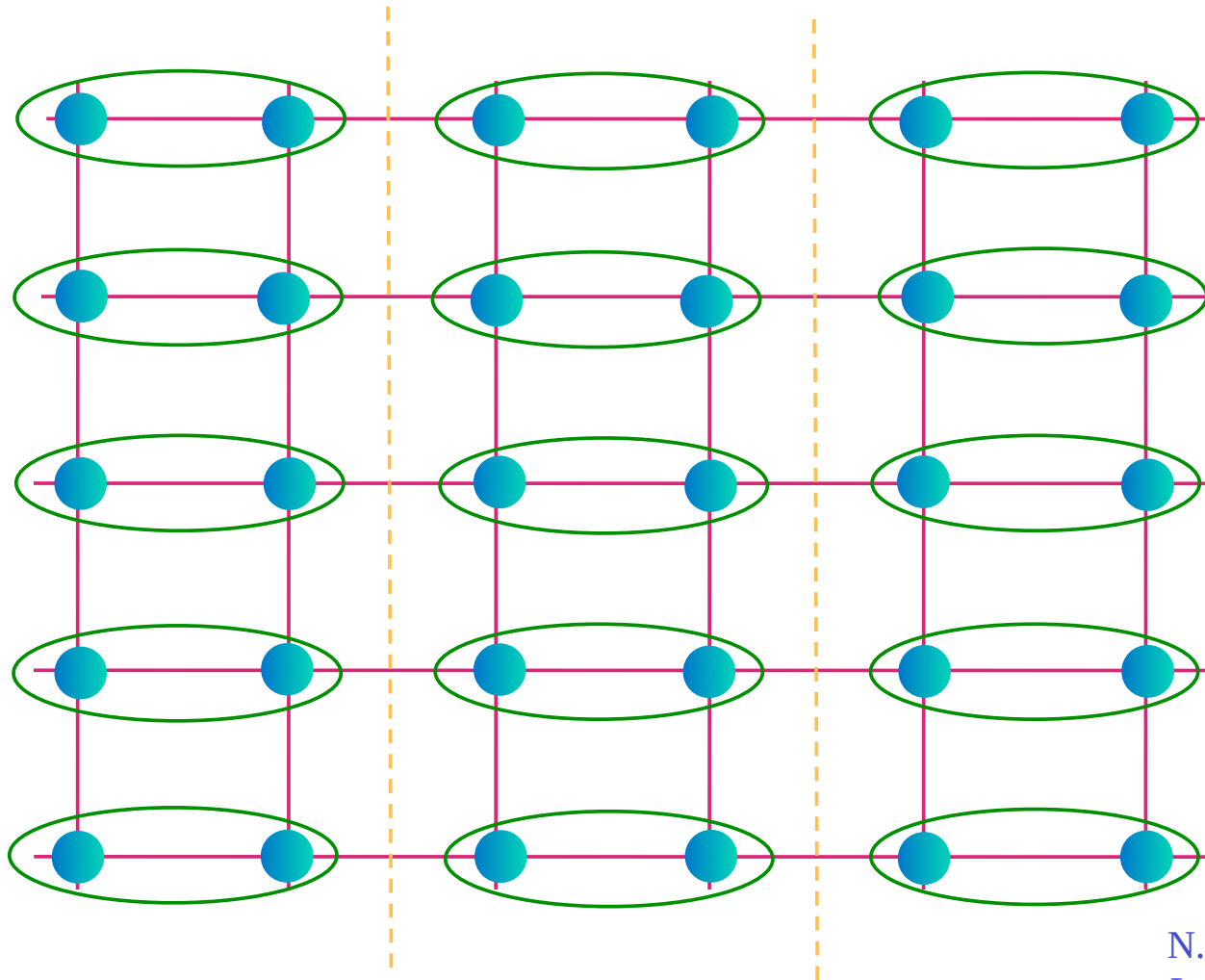
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More possibilities for entanglement with nearby states



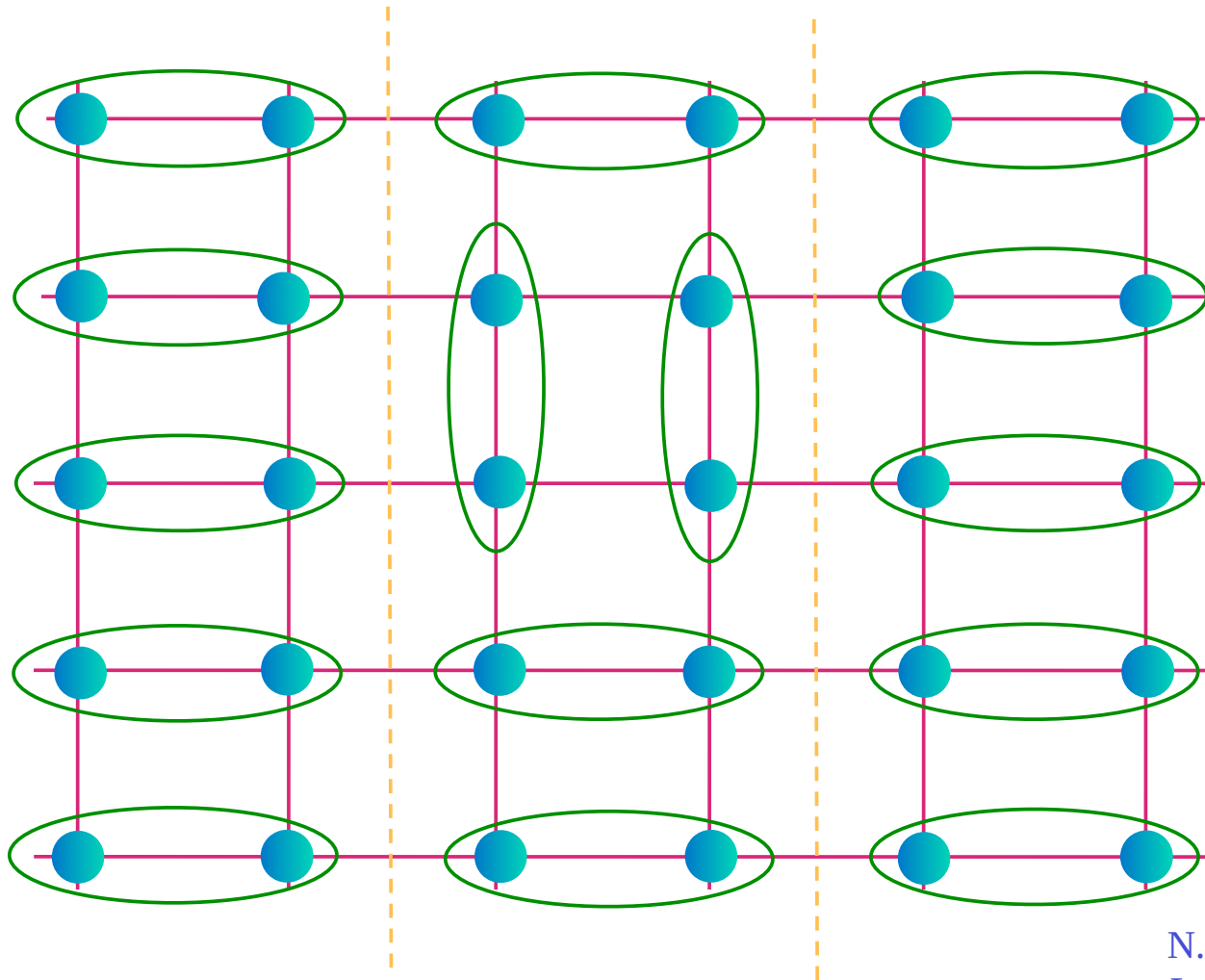
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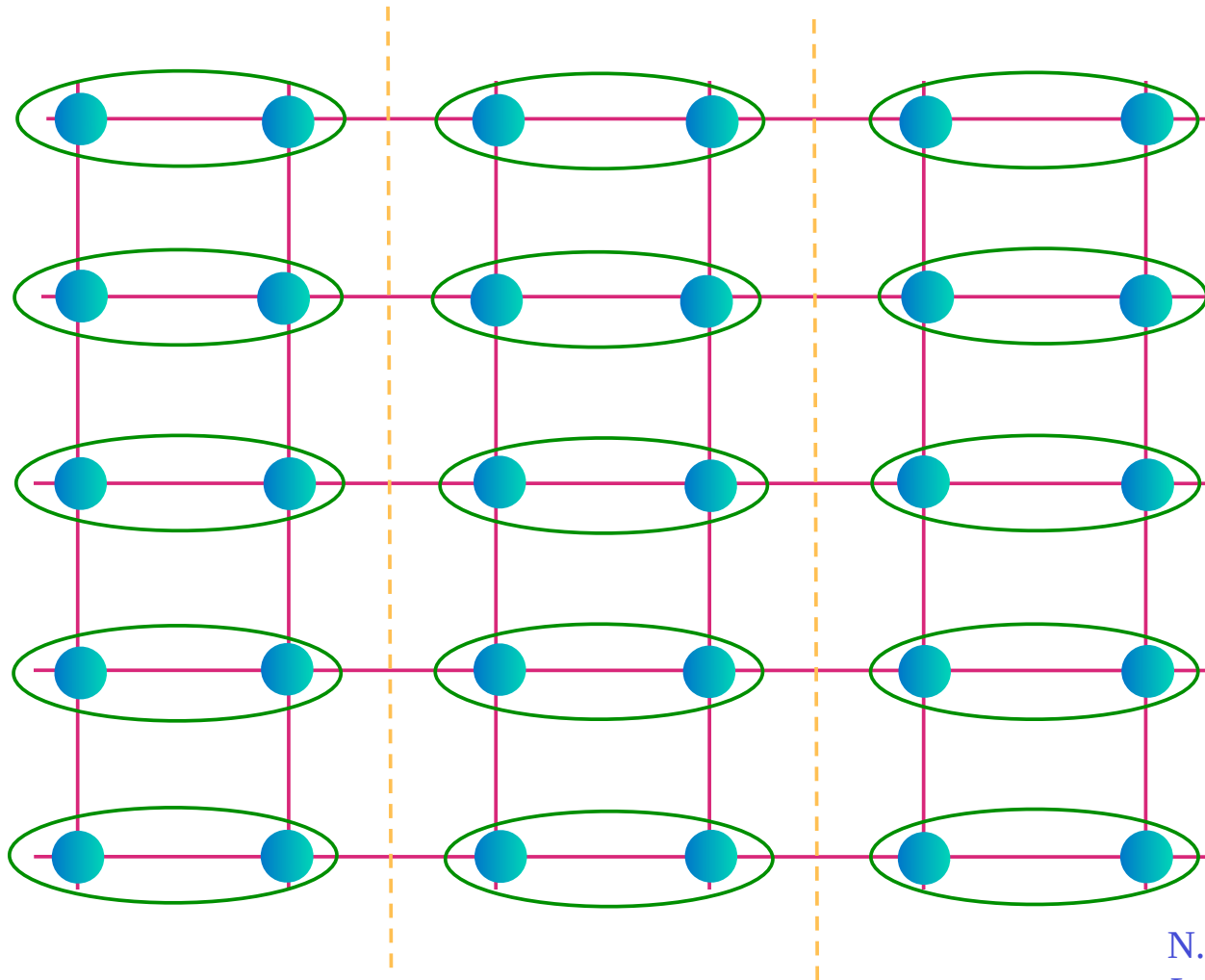
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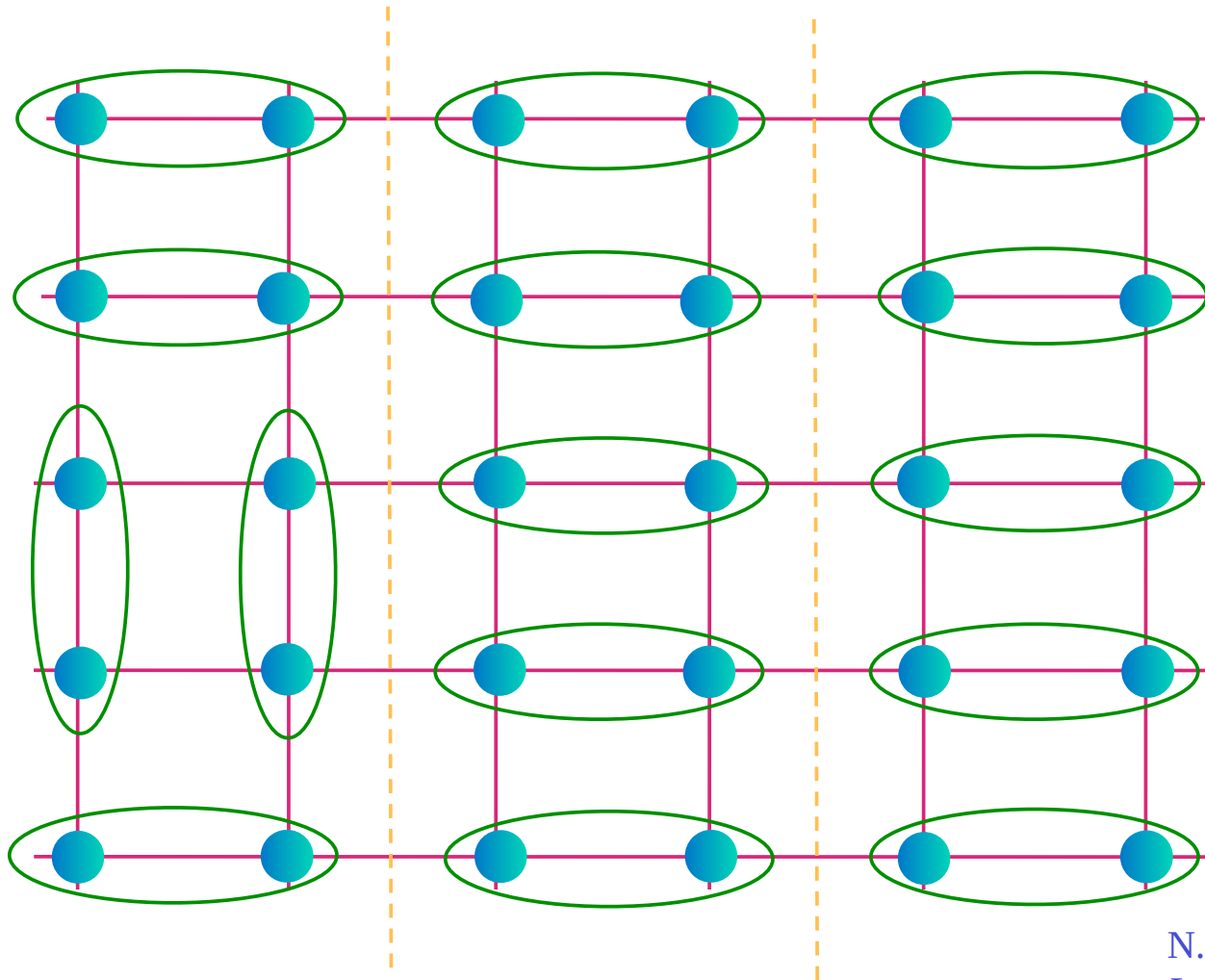
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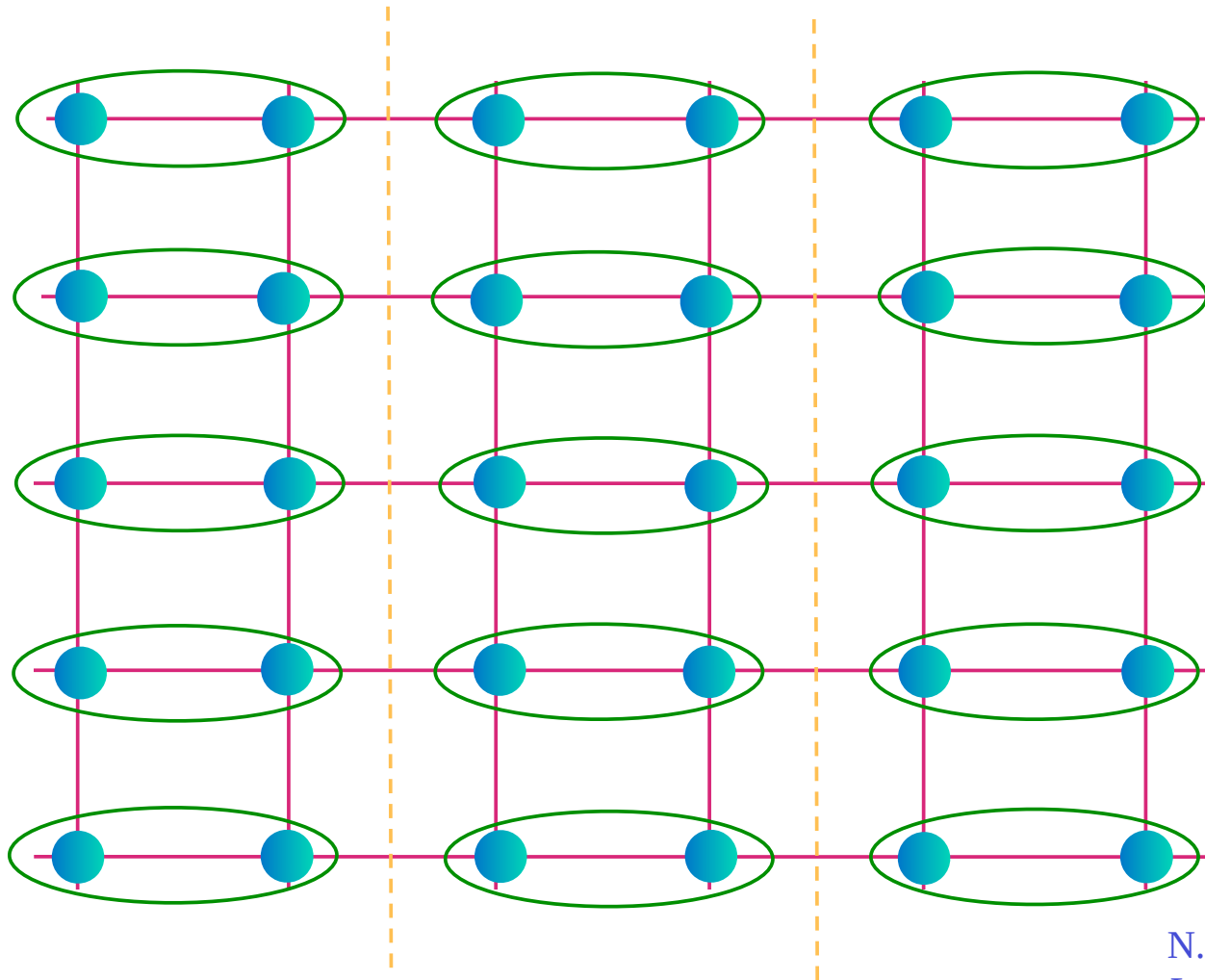
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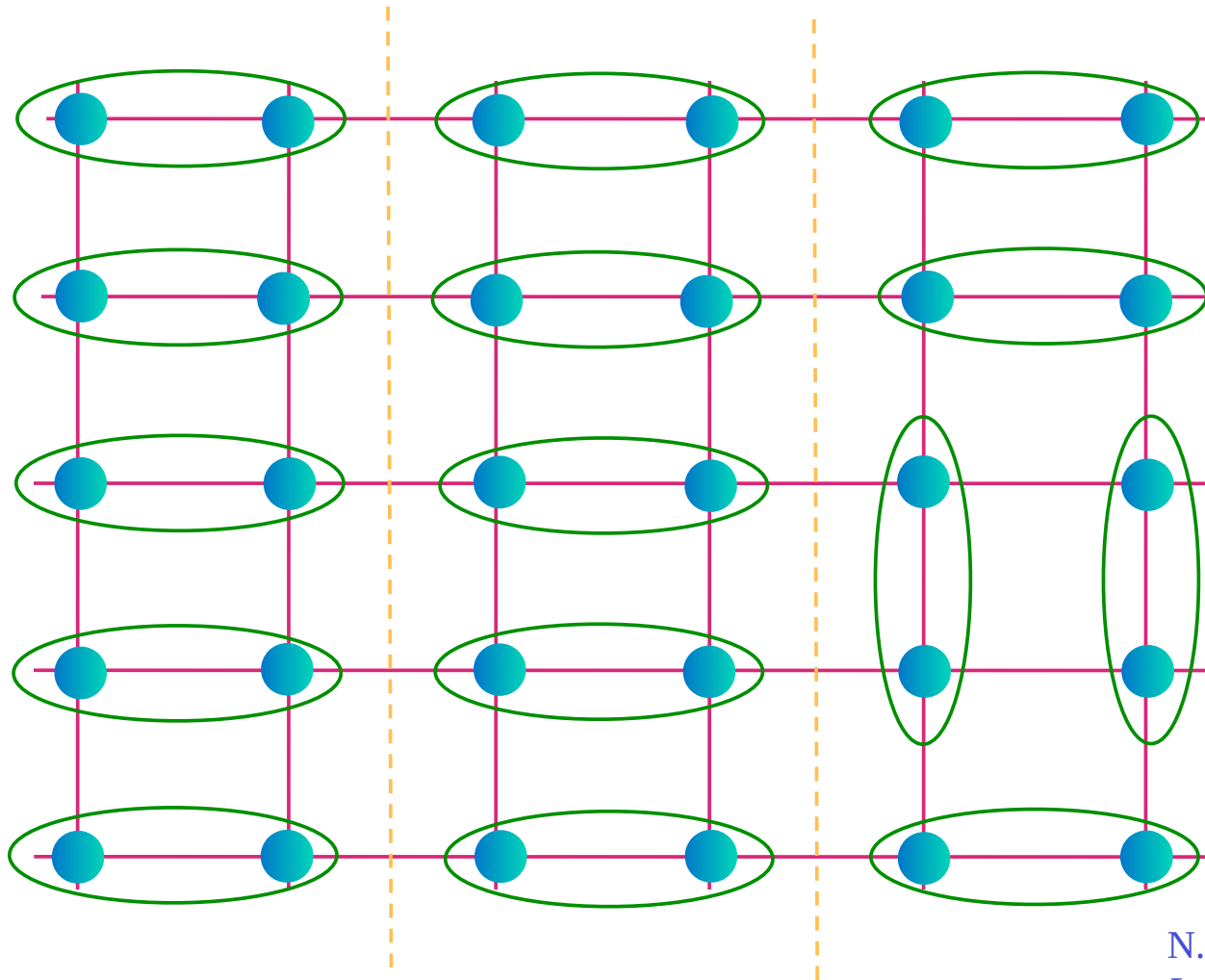
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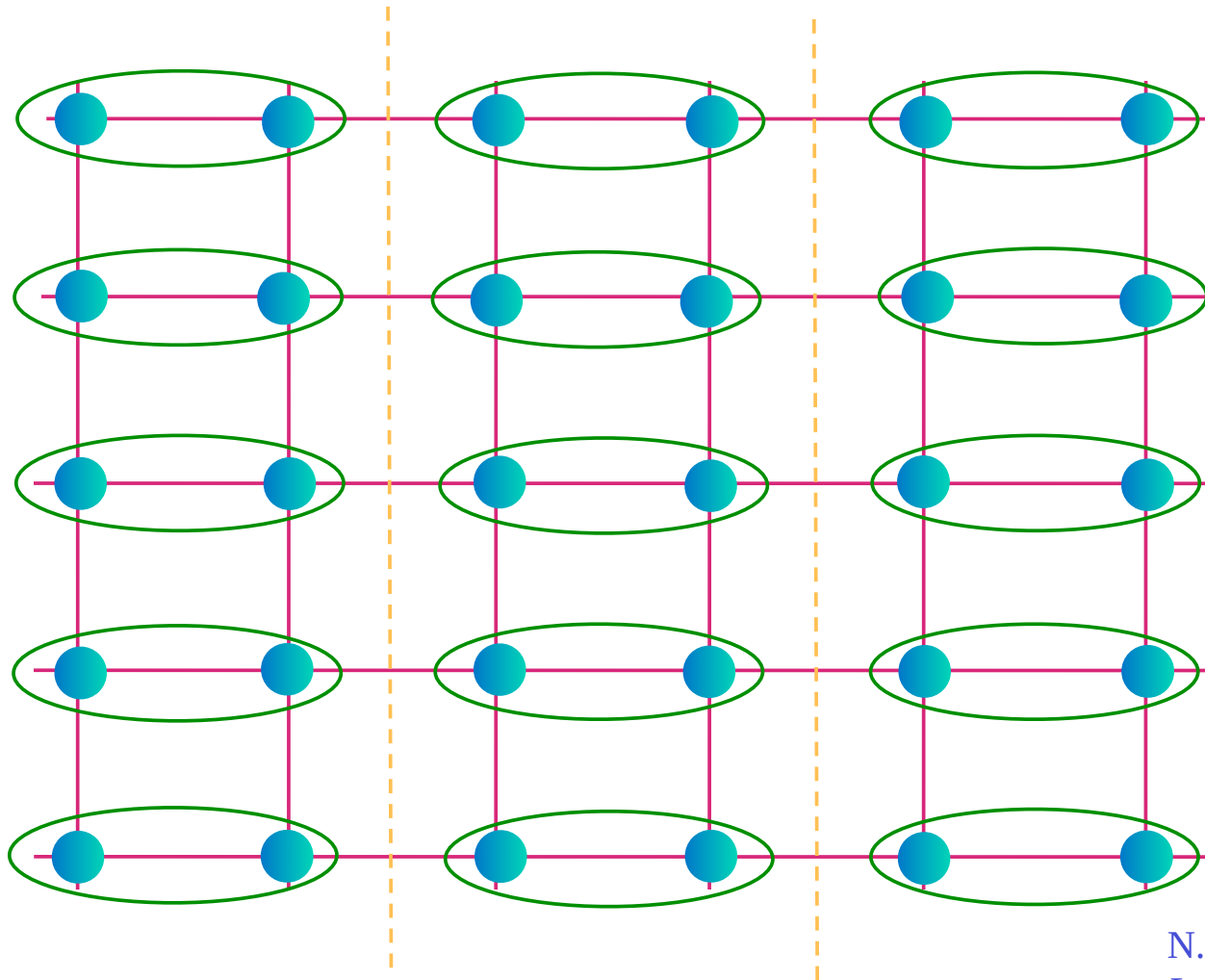
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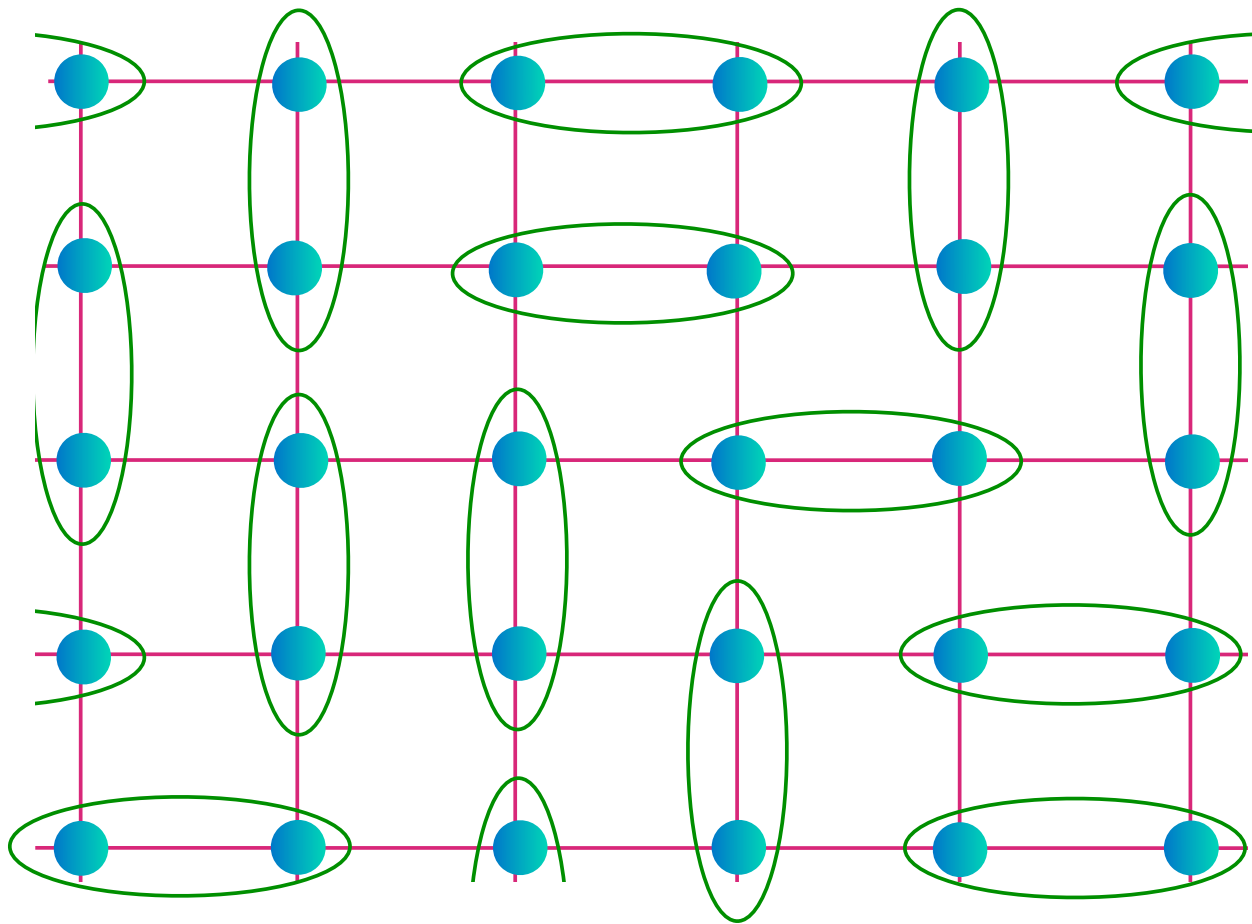


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N. Read and S. Sachdev, *Phys. Rev. Lett.* **62**, 1694 (1989).

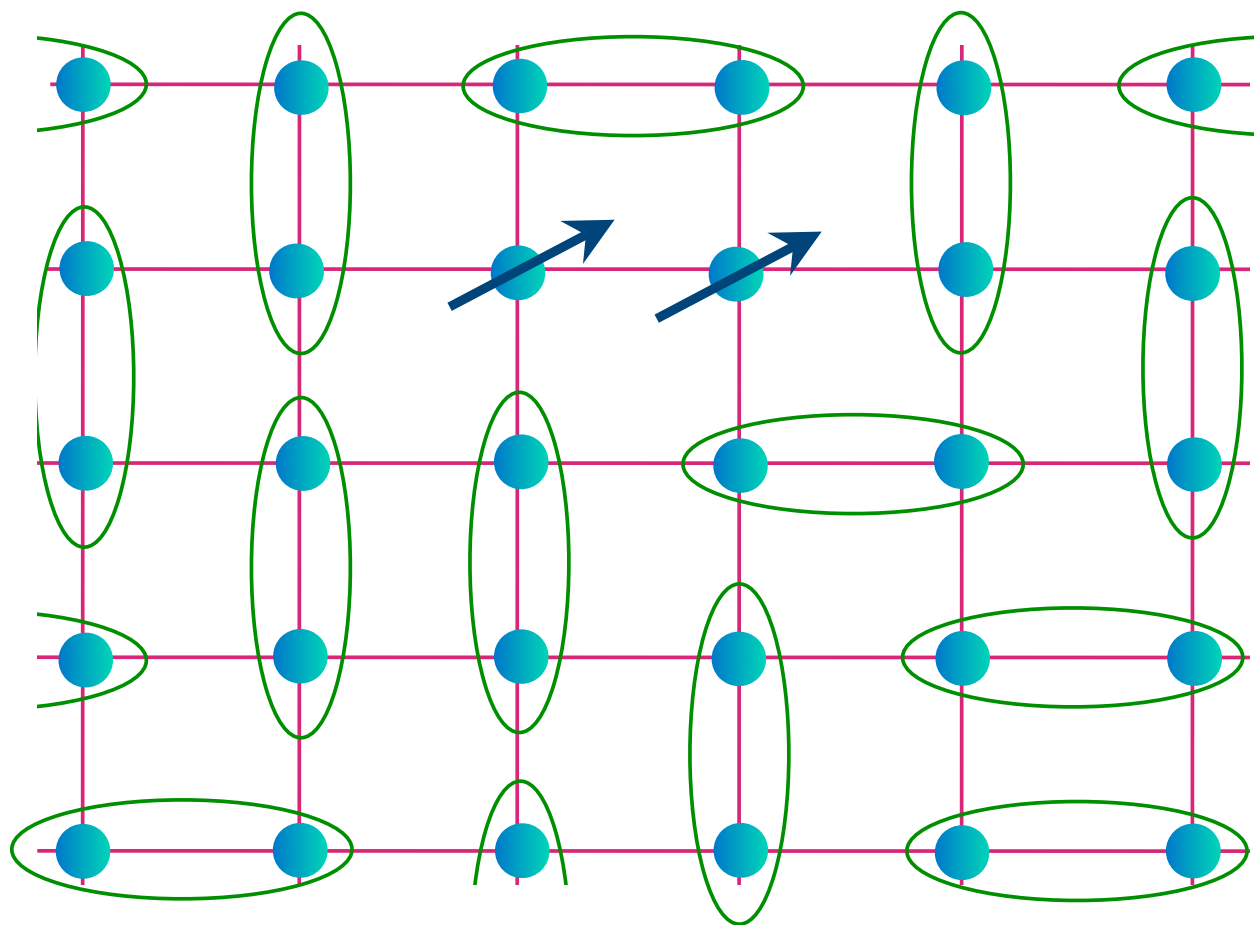
R. Moessner and S. L. Sondhi, *Phys. Rev. B* **63**, 224401 (2001).

Excitations of the RVB liquid



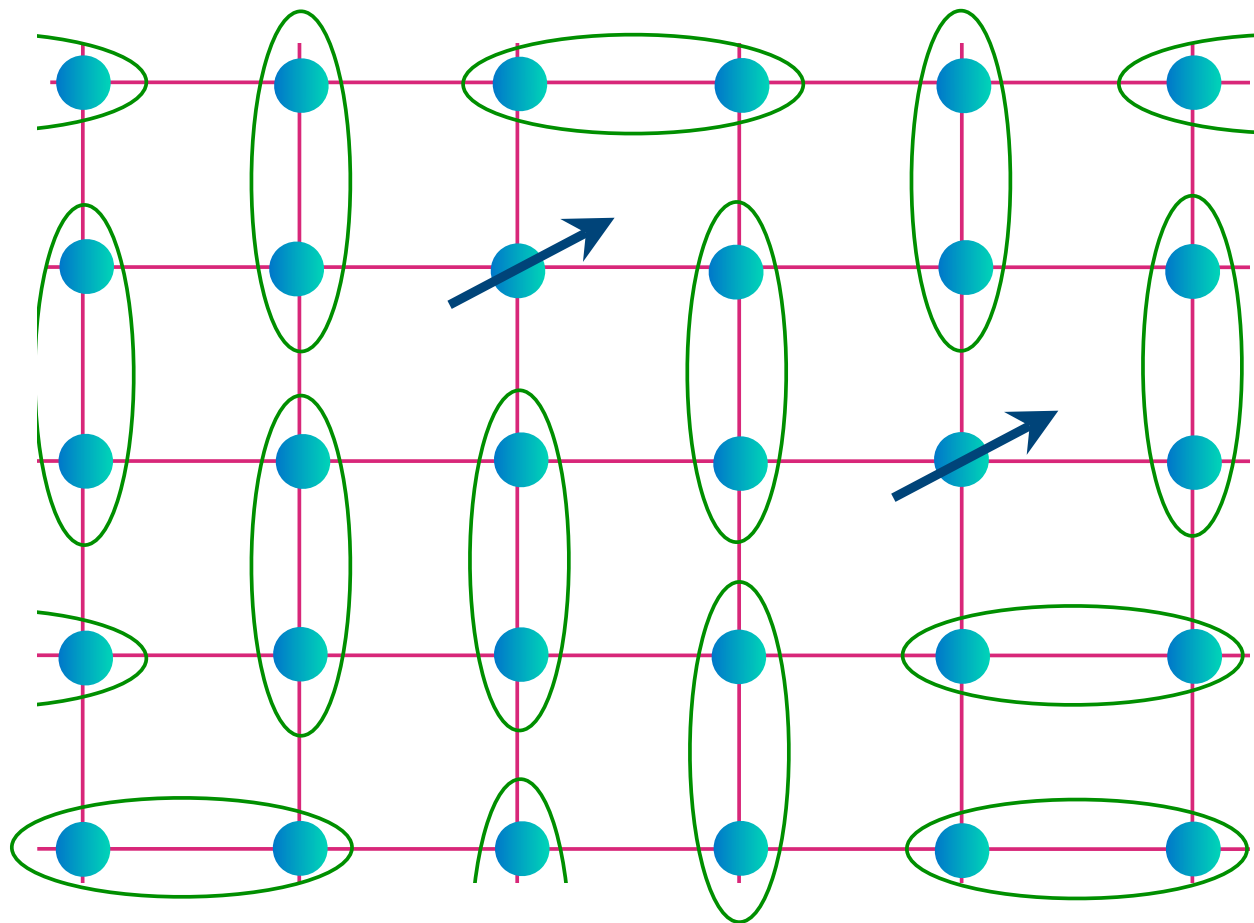
$$\text{[Diagram of two dots in an oval]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



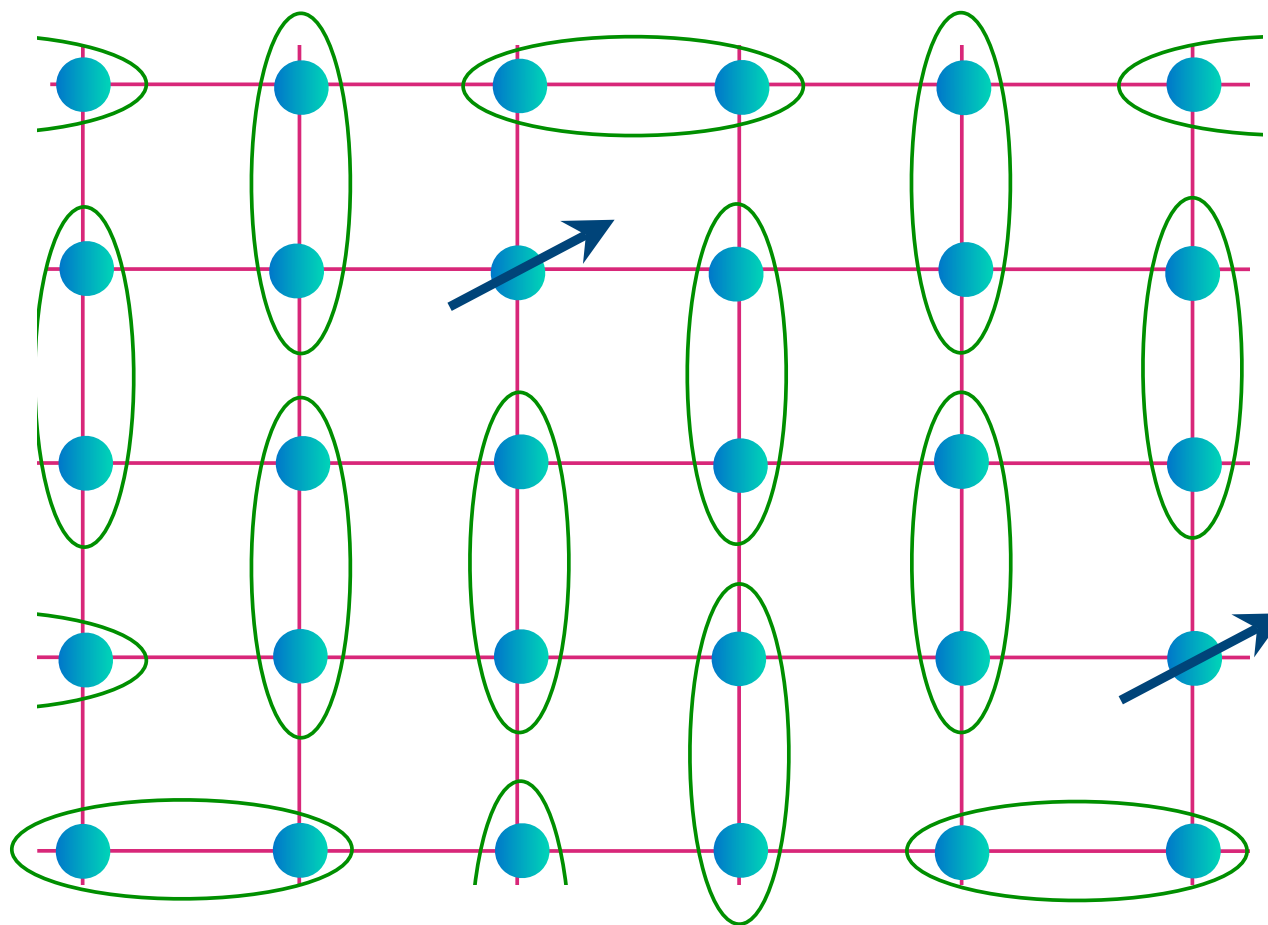
$$\text{[Diagram of two dots in an oval]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



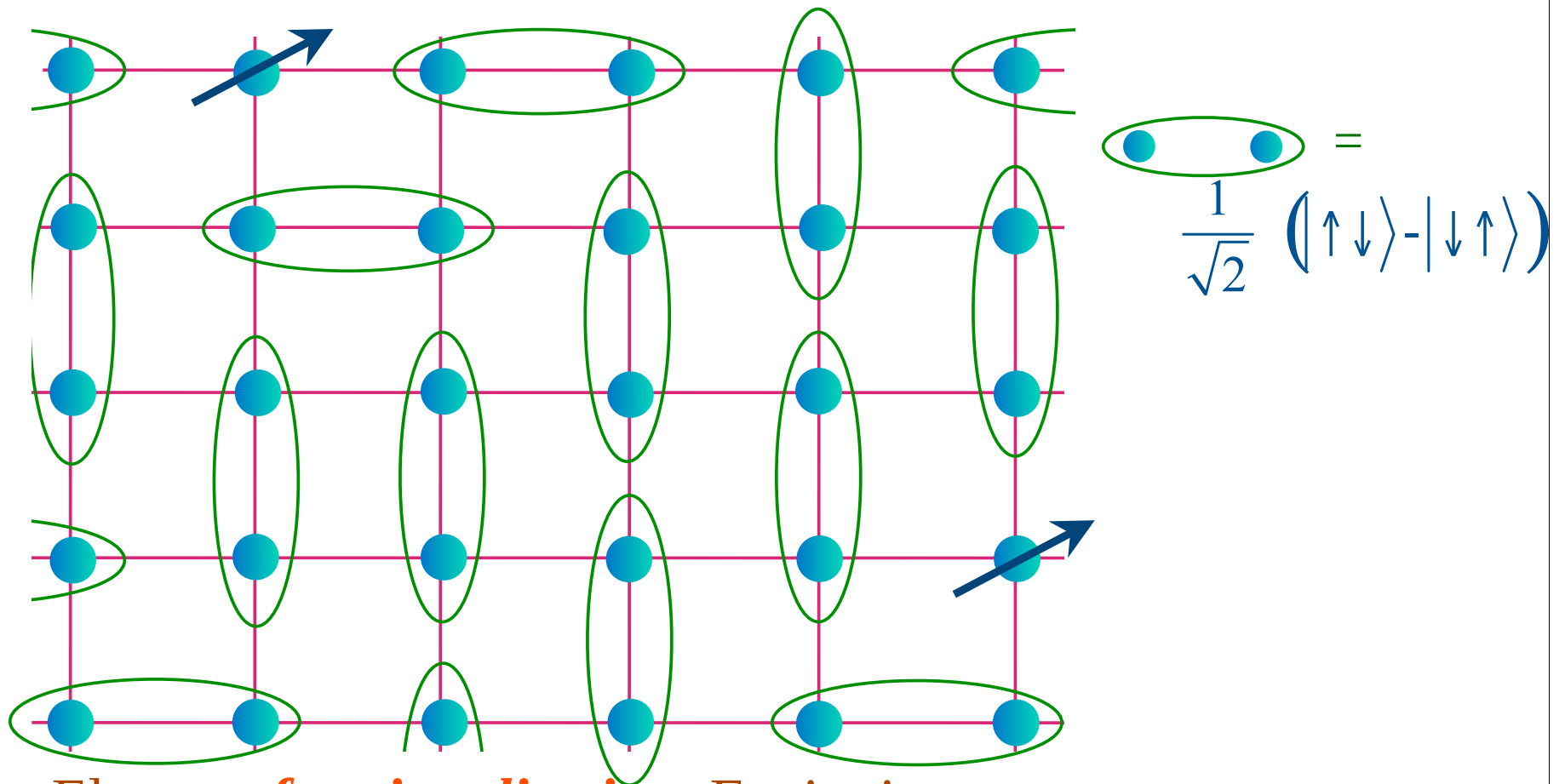
$$\text{[Diagram of two dots in a green oval]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



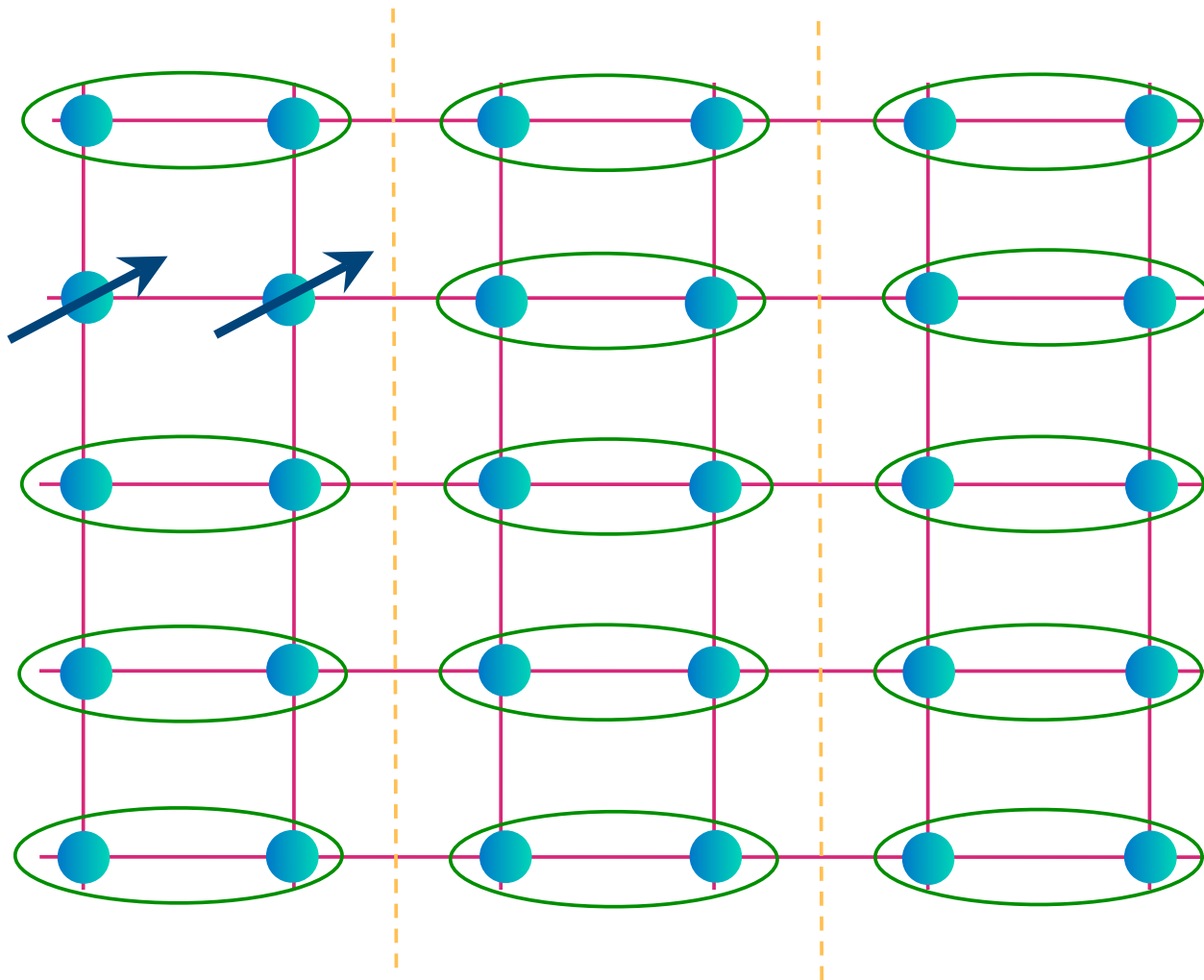
$$\text{[Diagram of two dots in an oval]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the RVB liquid



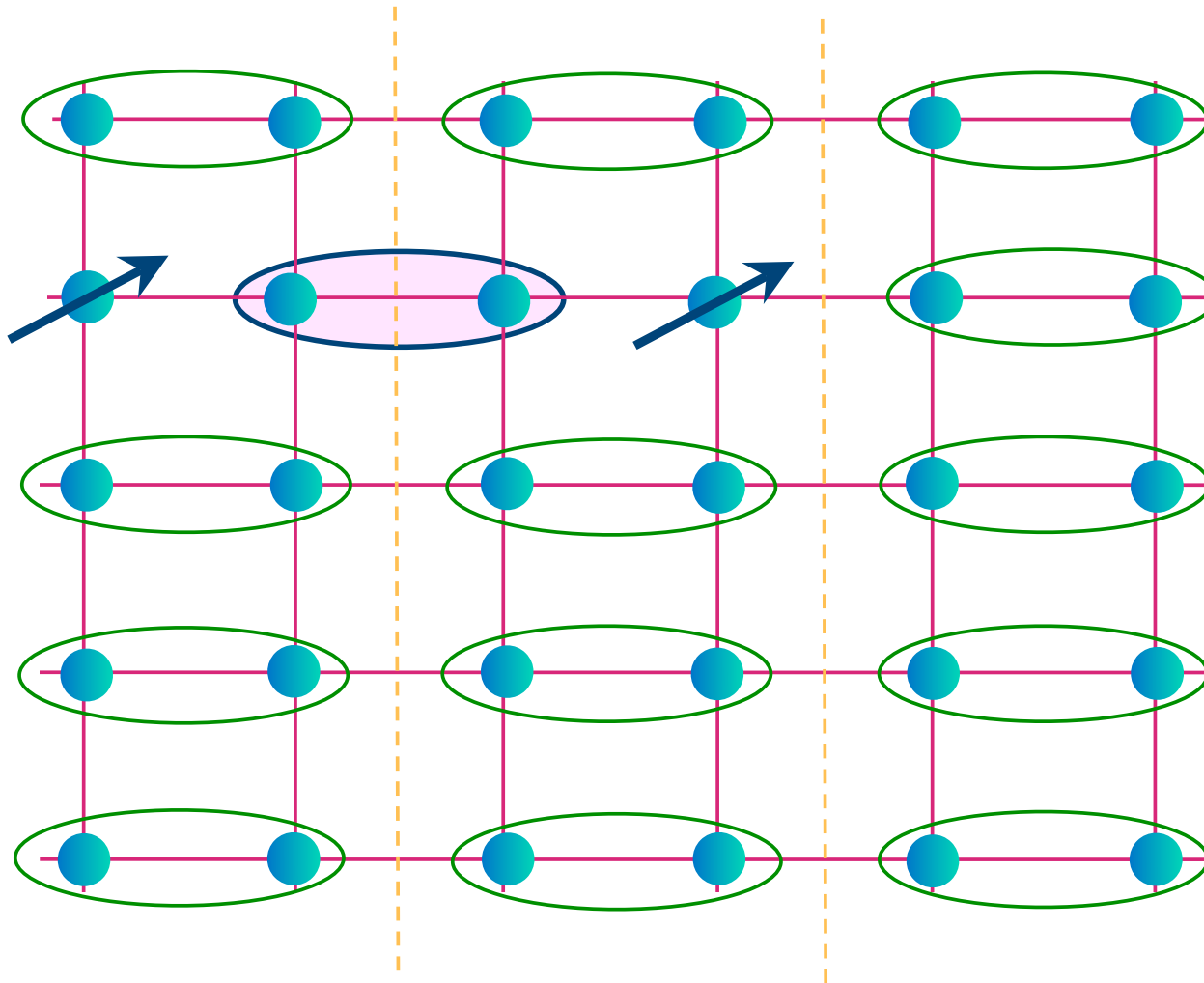
Electron **fractionalization**: Excitations carry spin $S=1/2$ but no charge

Excitations of the VBS



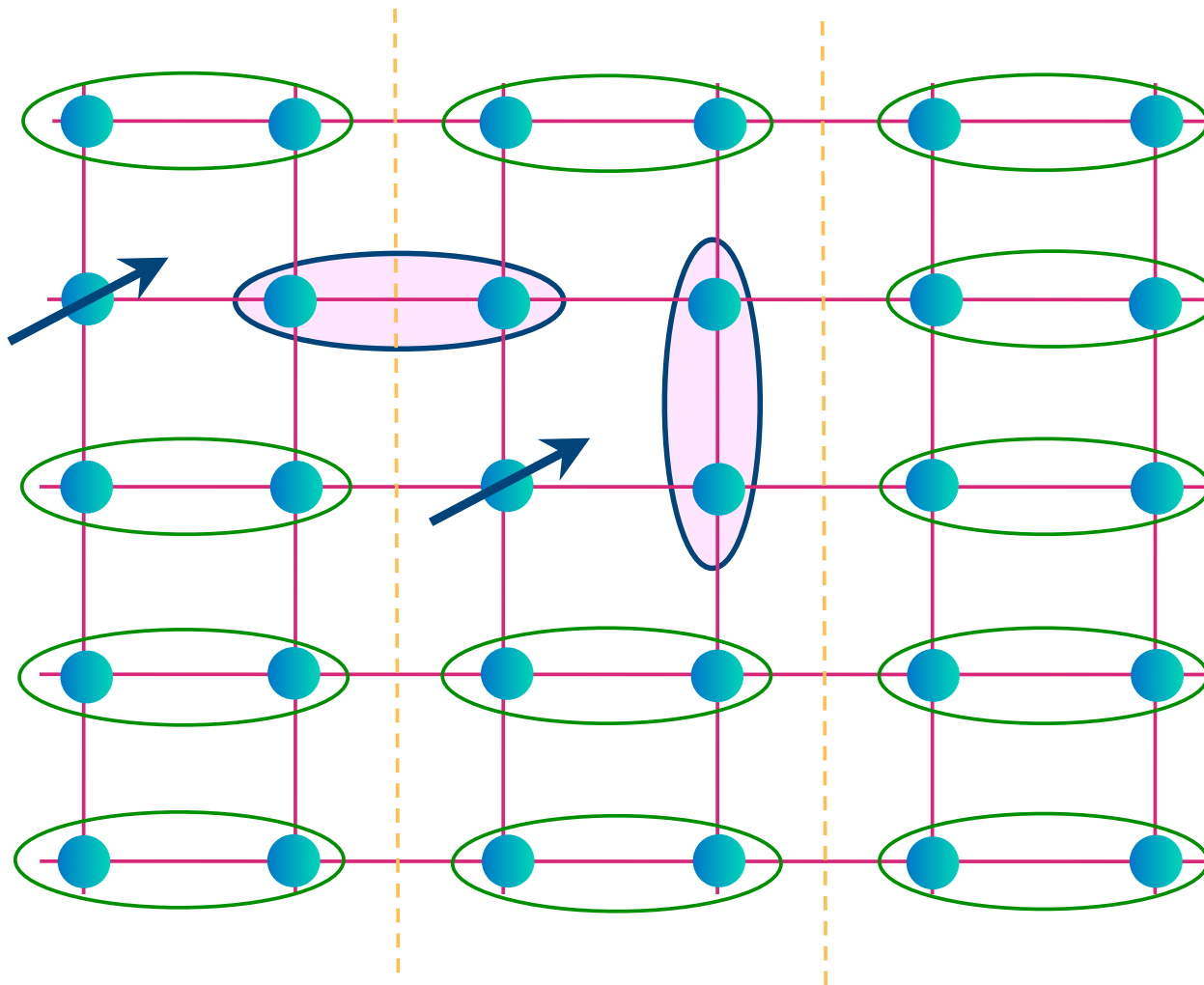
$$\begin{array}{|c|c|} \hline \bullet & \bullet \\ \hline \end{array} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the VBS



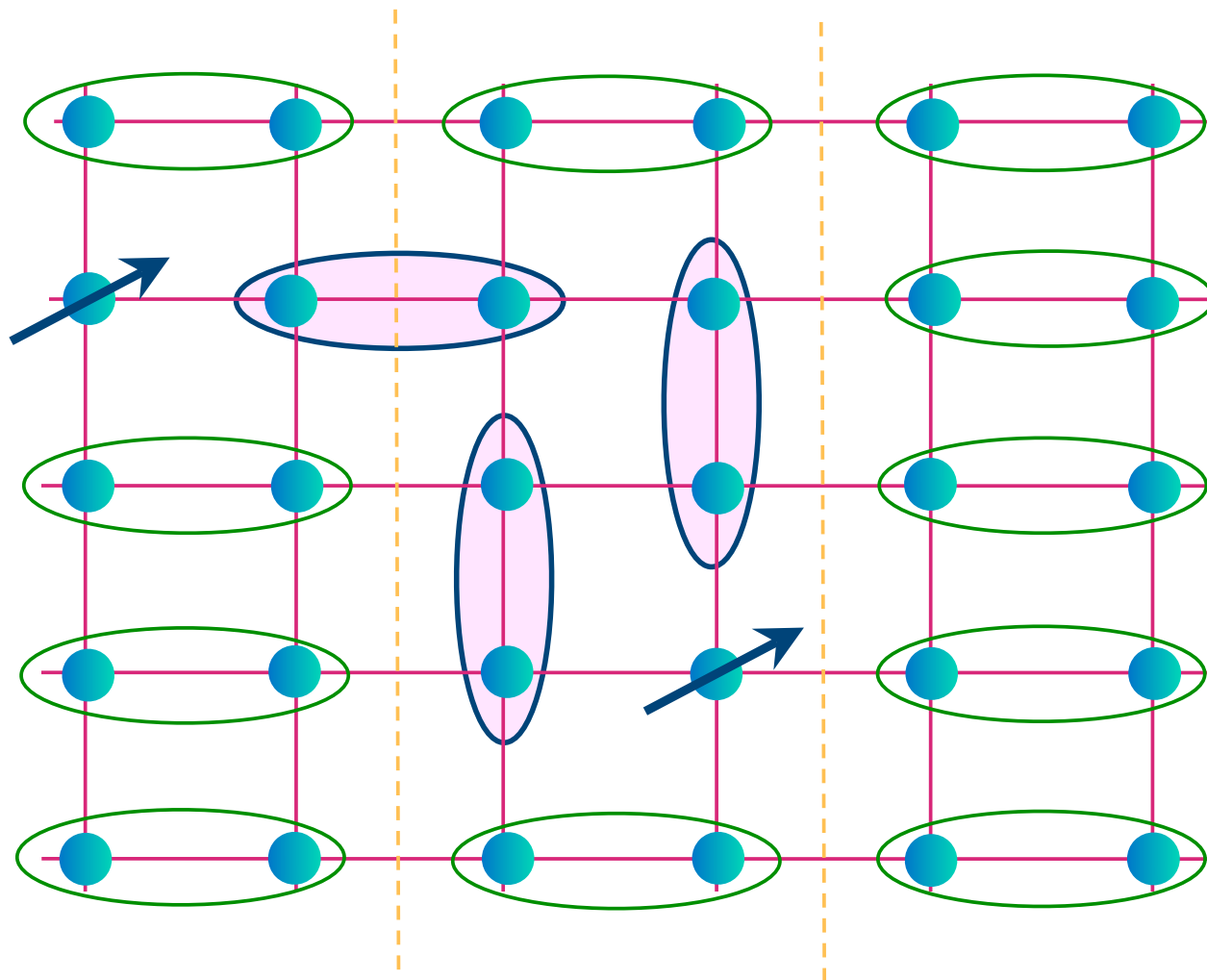
$$\begin{array}{|c|c|} \hline \bullet & \bullet \\ \hline \end{array} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the VBS



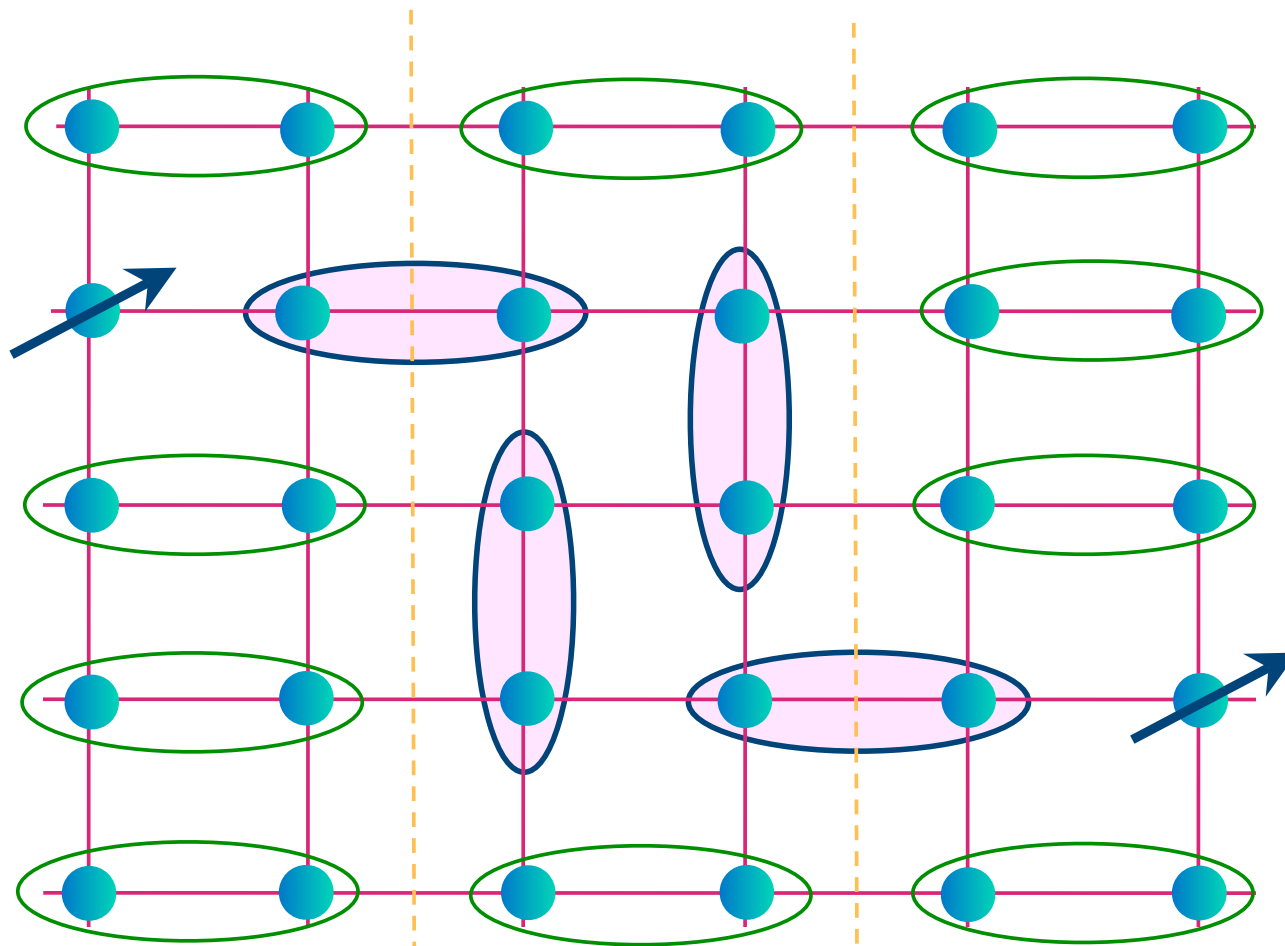
$$\begin{array}{|c|c|} \hline \bullet & \bullet \\ \hline \end{array} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Excitations of the VBS



$$\text{green oval} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

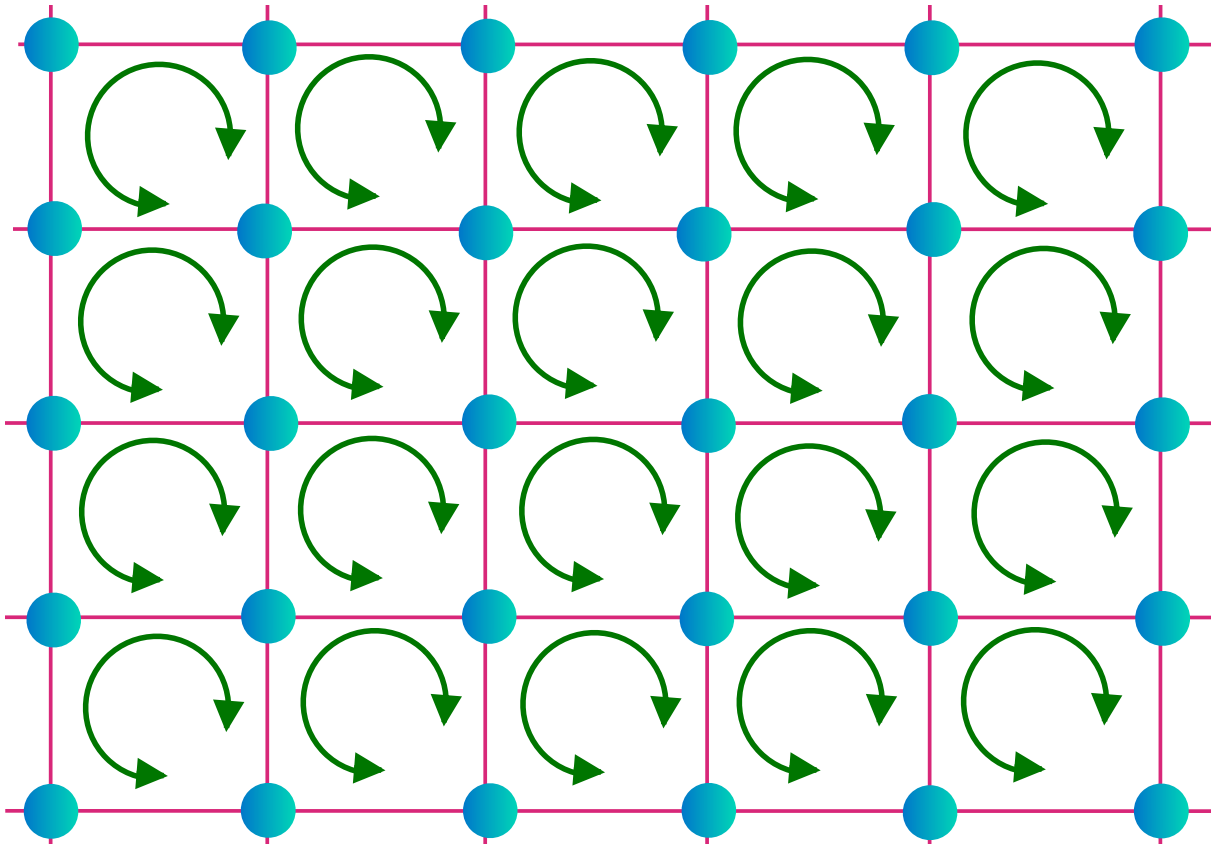
Excitations of the VBS



$$\text{[Diagram of two dots in a green oval]} = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

Free spins are unable to move apart: no fractionalization, but confinement

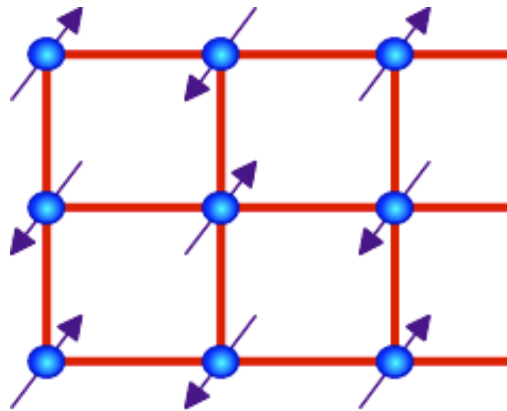
Phase diagram of square lattice antiferromagnet



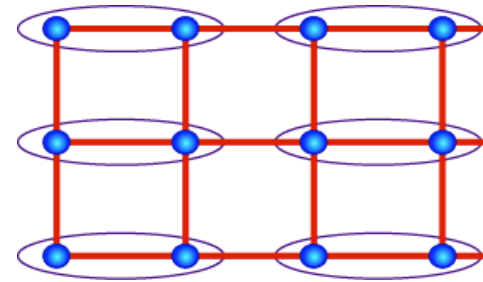
$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} \text{four spin exchange}$$

A. W. Sandvik, cond-mat/0611343

Phase diagram of square lattice antiferromagnet



Neel order



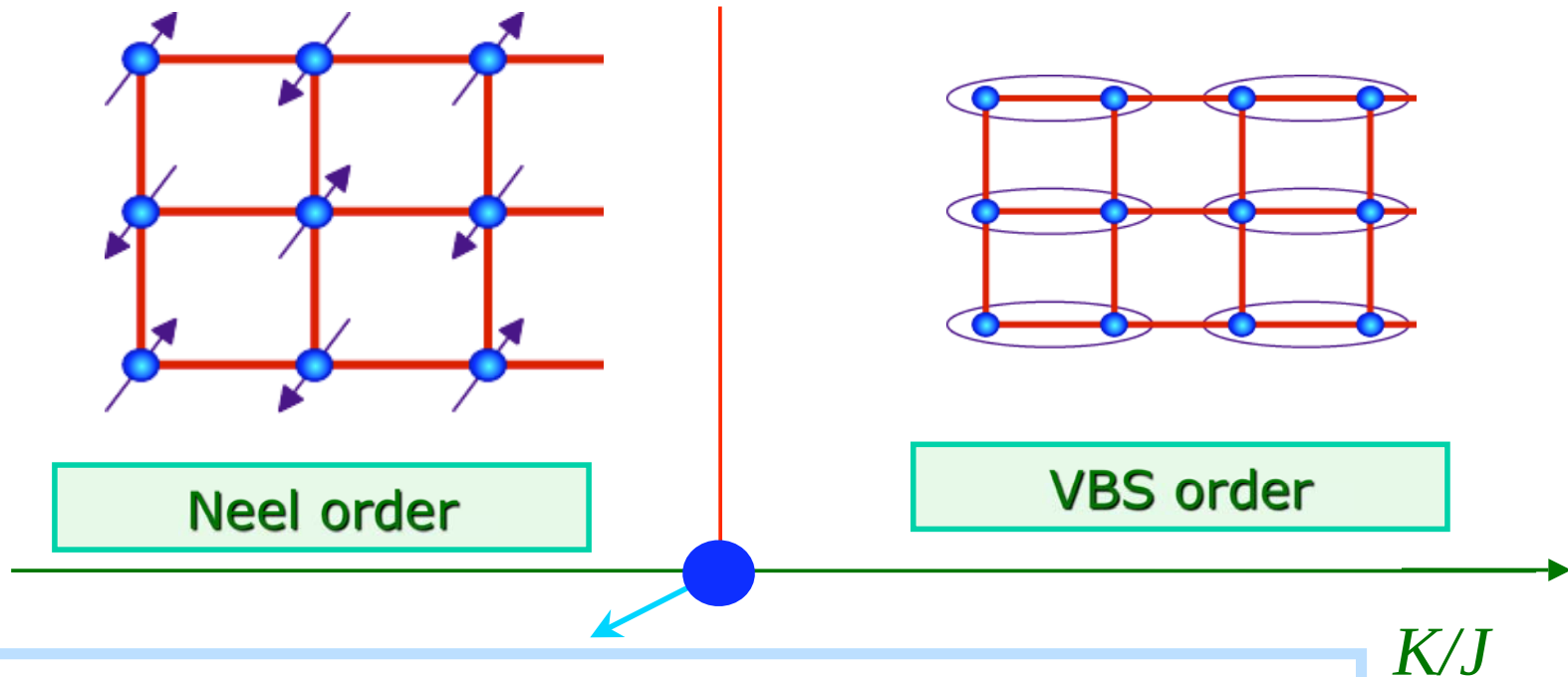
VBS order

K/J

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} \text{four spin exchange}$$

T. Senthil, A. Vishwanath, L. Balents, S. Sachdev and M.P.A. Fisher, *Science* **303**, 1490 (2004).

Phase diagram of square lattice antiferromagnet

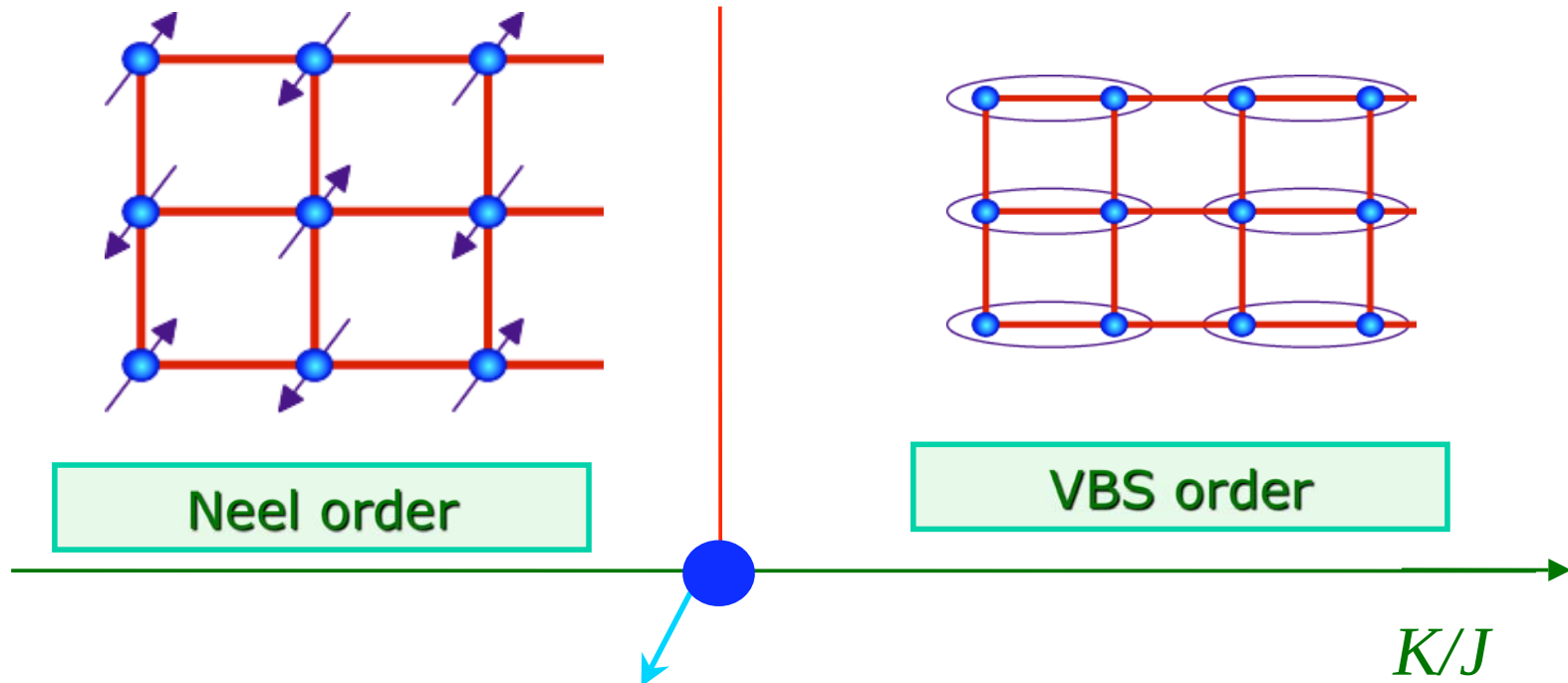


RVB physics appears at the quantum critical point which has fractionalized excitations: “deconfined criticality”

$$H = J \sum_{\langle ij \rangle} \vec{S}_i \cdot \vec{S}_j + K \sum_{\square} \text{four spin exchange}$$

T. Senthil, A. Vishwanath, L. Balents, S. Sachdev and M.P.A. Fisher, *Science* **303**, 1490 (2004).

Phase diagram of square lattice antiferromagnet



Second-order critical point described by

$$\mathcal{S}_{\text{critical}} = \int d^2x d\tau \left[|(\partial_\mu - iA_\mu)z_\alpha|^2 + r |z_\alpha|^2 + \frac{u}{2} (|z_\alpha|^2)^2 + \frac{1}{4e^2} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2 \right]$$

at its critical point $r = r_c$, where z_α are the neutral $S = 1/2$ spinons and A_μ is a *non-compact* $U(1)$ gauge field.

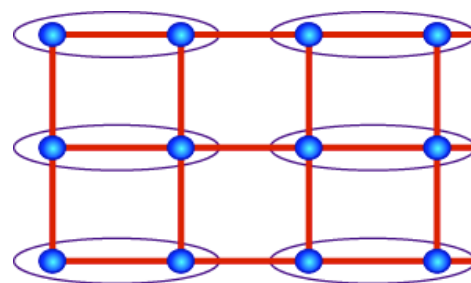
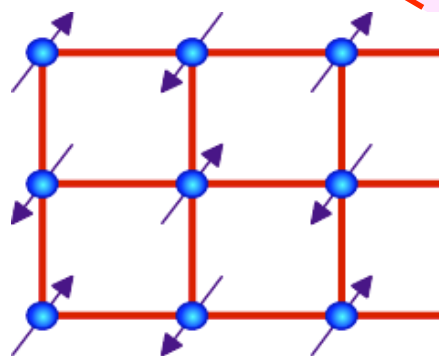
T. Senthil, A. Vishwanath, L. Balents, S. Sachdev and M.P.A. Fisher, *Science* **303**, 1490 (2004).

Temperature, T

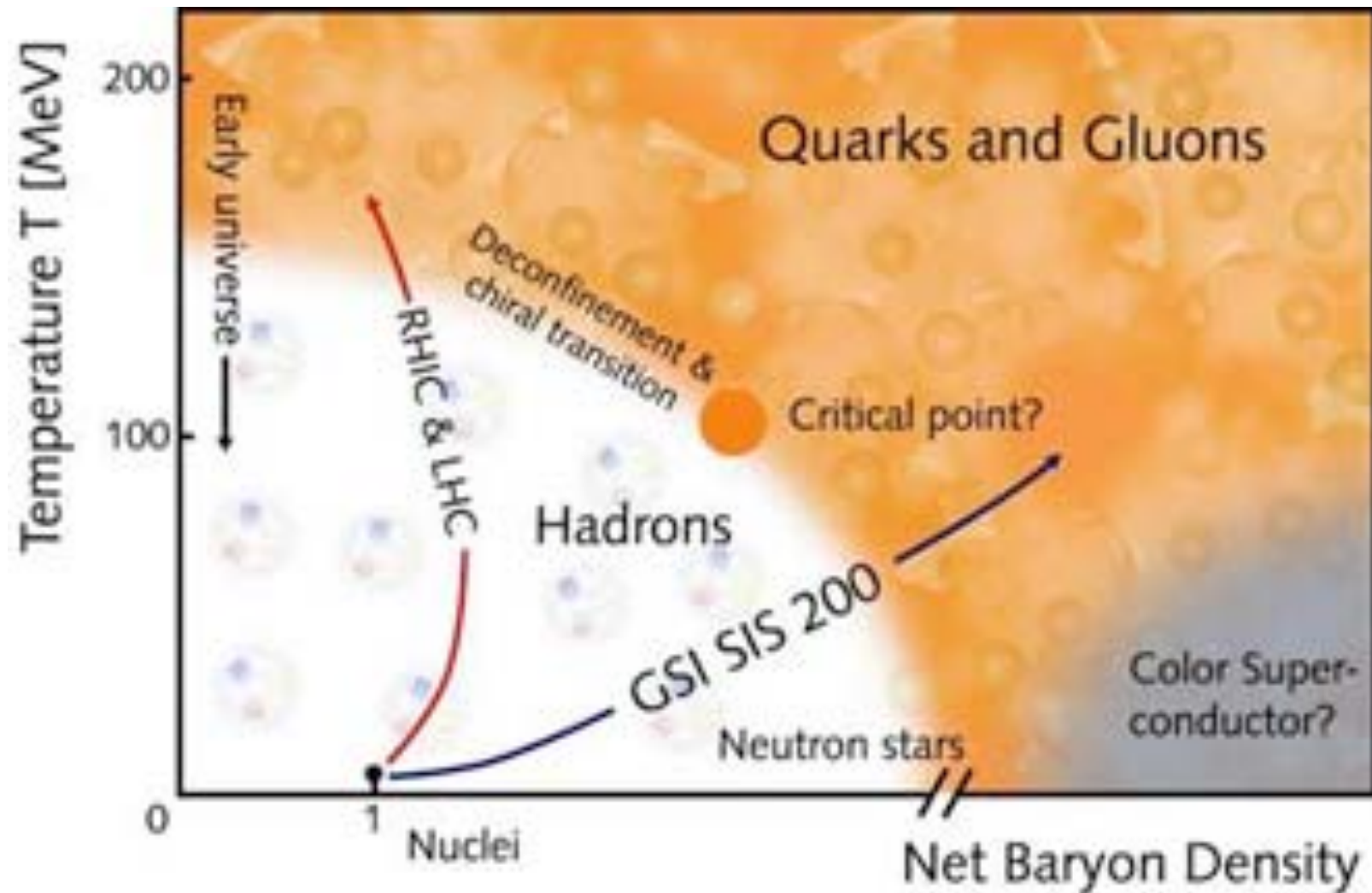
Quantum criticality of
fractionalized
excitations

0

K/J



Phases of nuclear matter



Outline

1. Superfluid-insulator quantum transitions

Experiments on ultracold atoms

2. Theory of quantum-critical transport

Collisionless- t_0 -hydrodynamic crossover of conformal field theories

3. Entanglement of valence bonds

Deconfined criticality in antiferromagnets

4. Nernst effect in the cuprate superconductors

Quantum criticality and dyonic black holes

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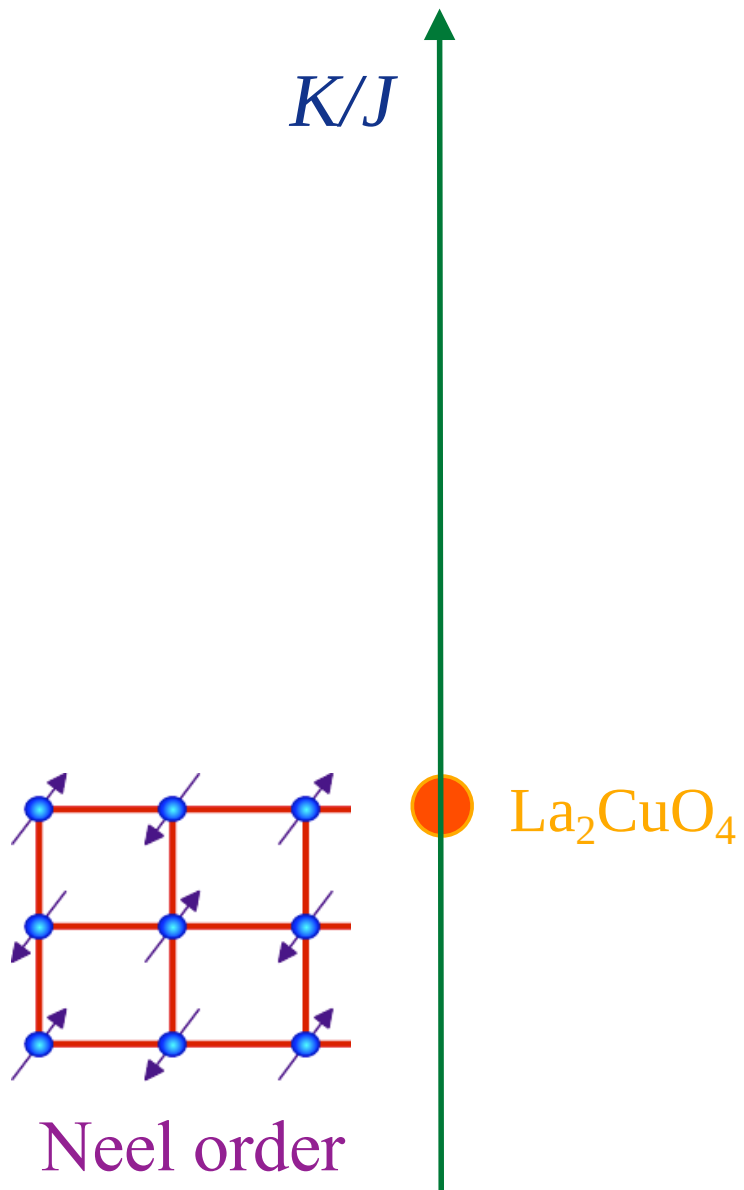
3. Entanglement of valence bonds

Deconfined criticality in antiferromagnets

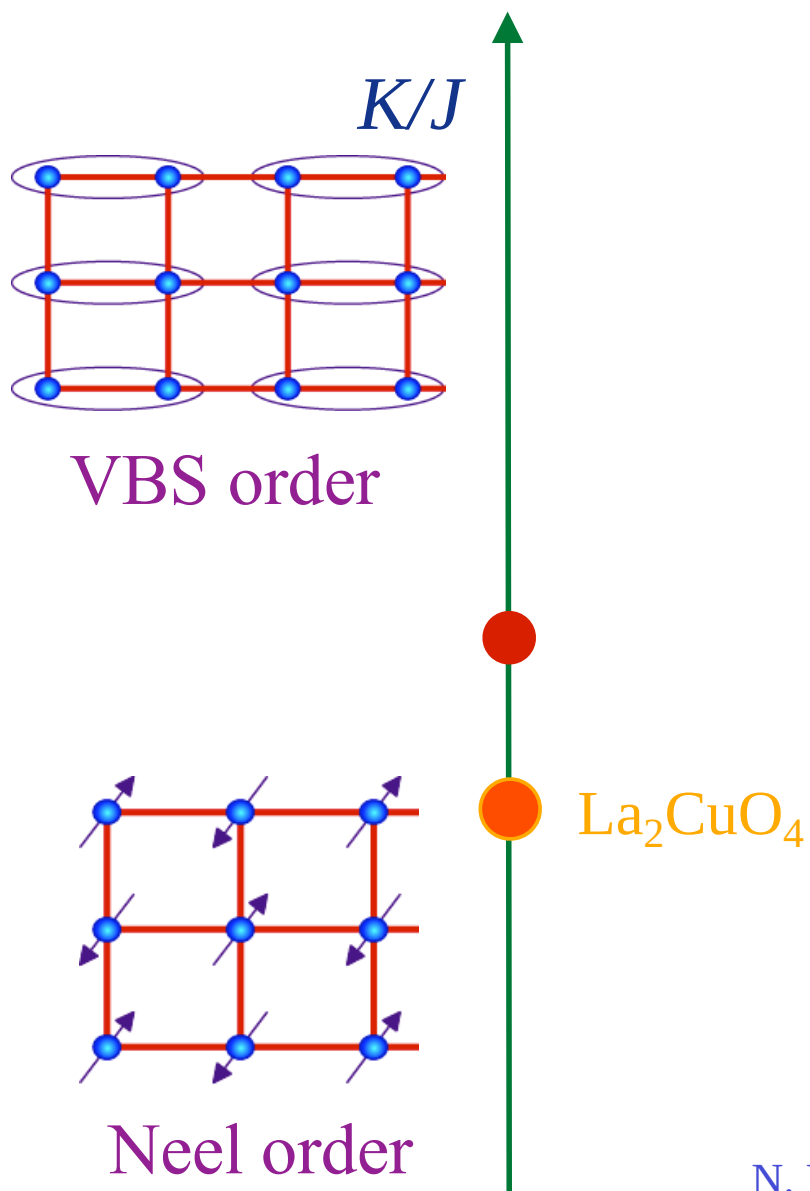
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Quantum criticality and dyonic black holes

Phase diagram of doped antiferromagnets



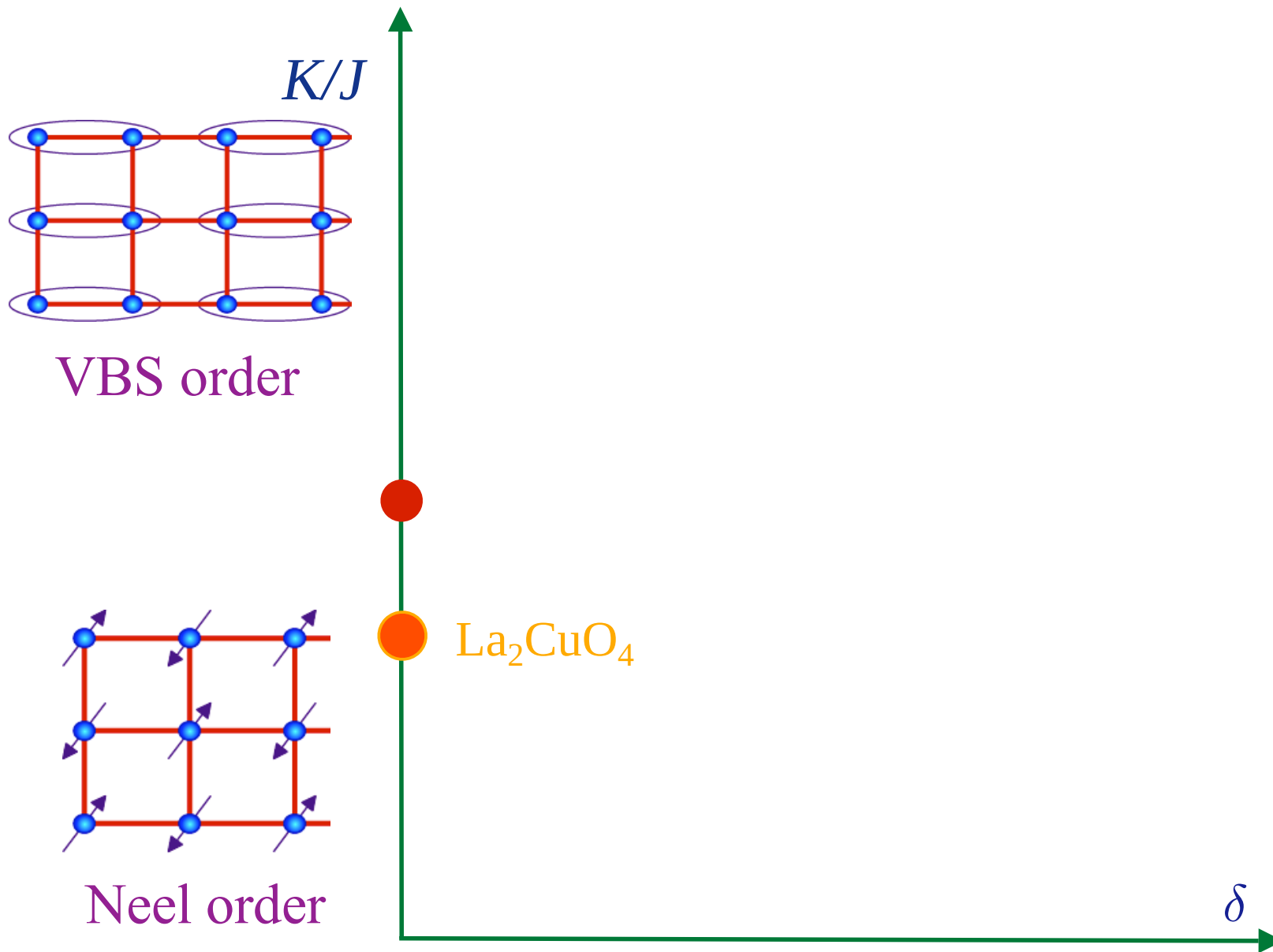
Phase diagram of doped antiferromagnets

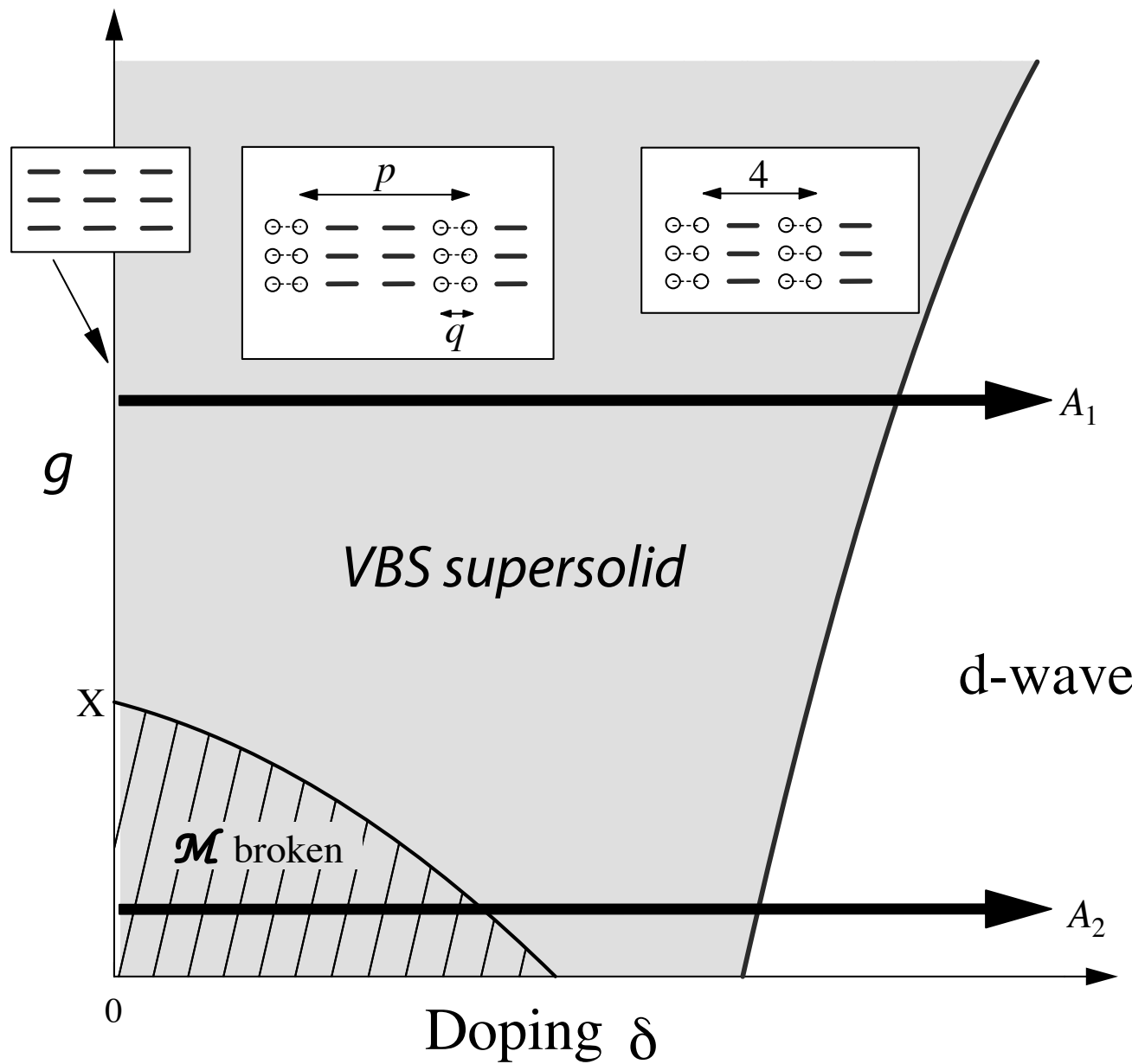


N. Read and S. Sachdev, *Phys. Rev. Lett.* **62**, 1694 (1989).

T. Senthil, A. Vishwanath, L. Balents, S. Sachdev and M.P.A. Fisher, *Science* **303**, 1490 (2004).

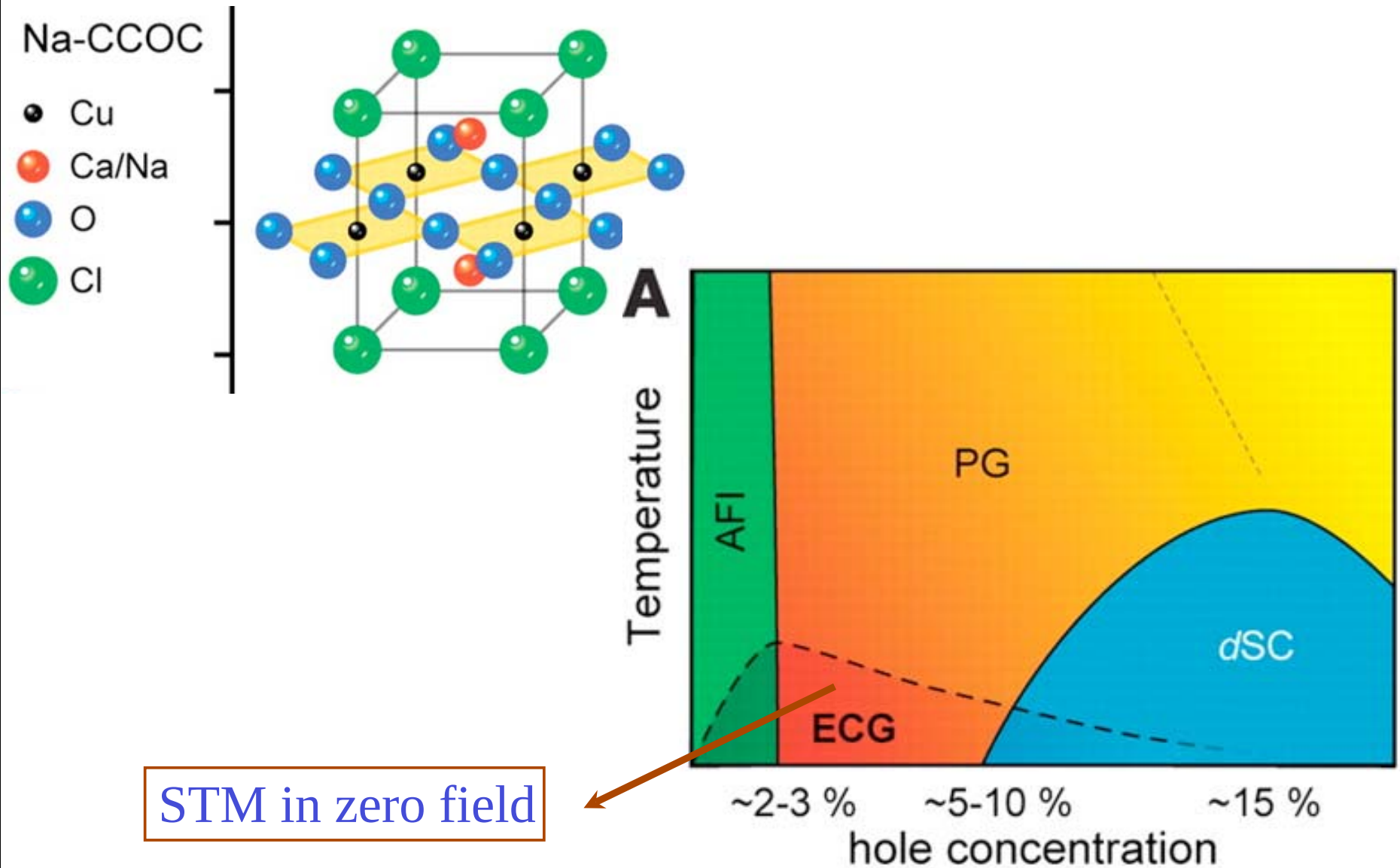
Phase diagram of doped antiferromagnets

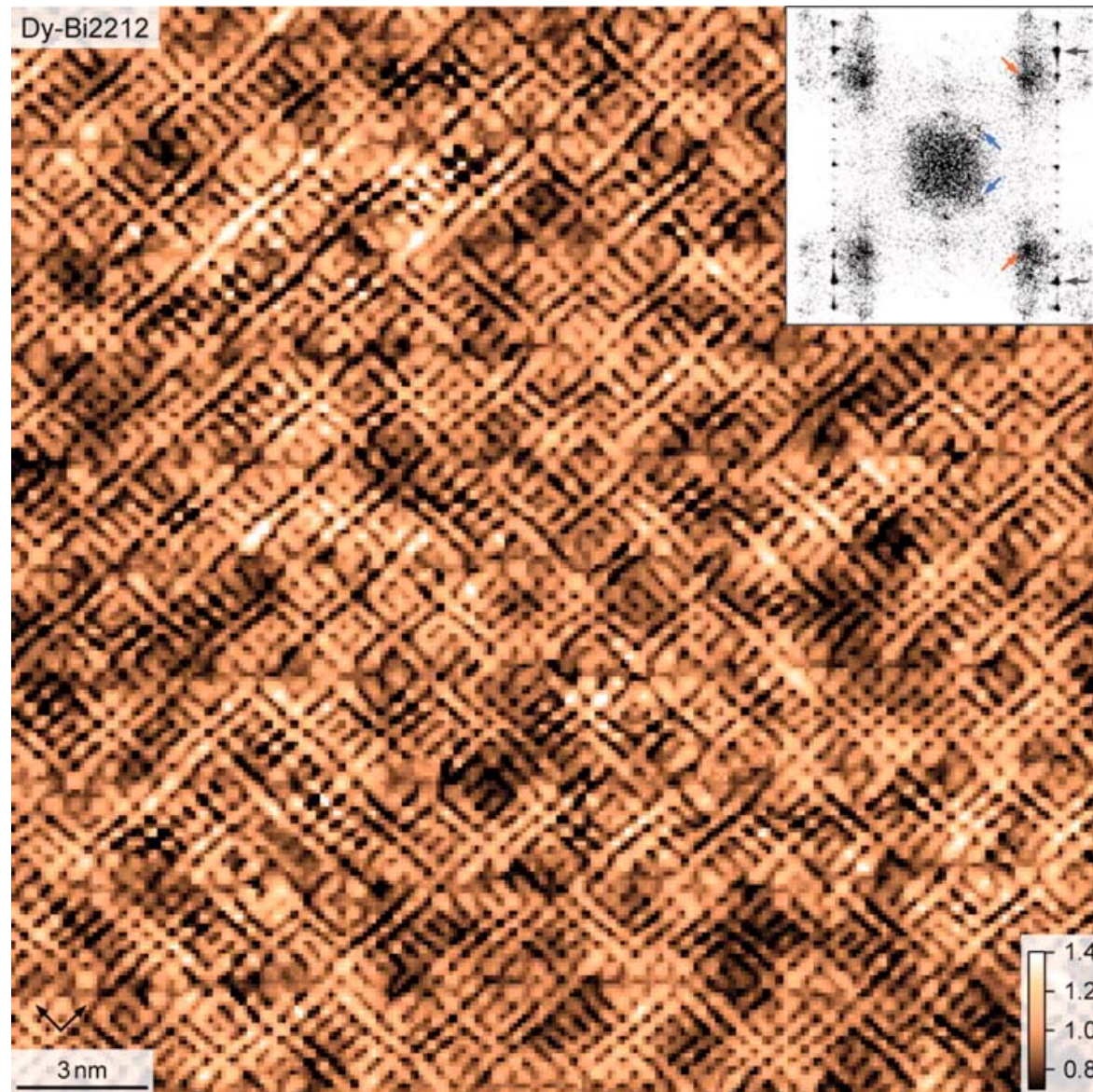




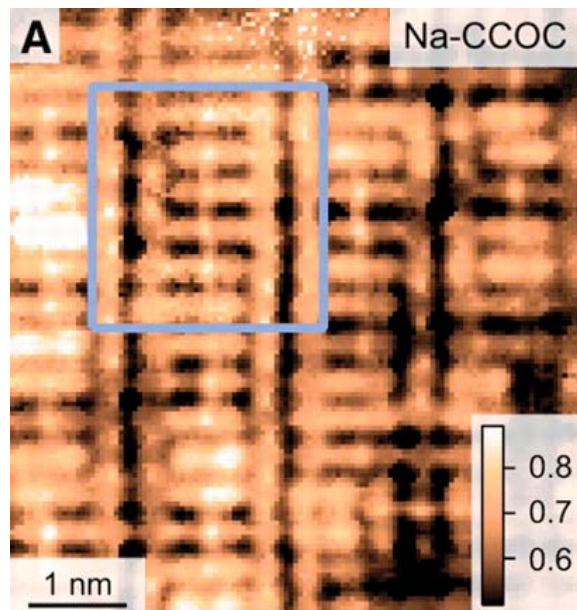
M. Vojta and S. Sachdev, *Phys. Rev. Lett.* **83**, 3916 (1999)

Temperature-doping phase diagram of the cuprate superconductors

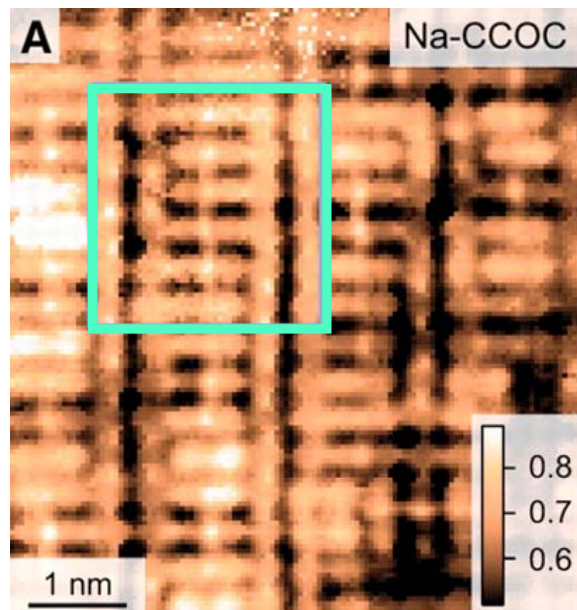




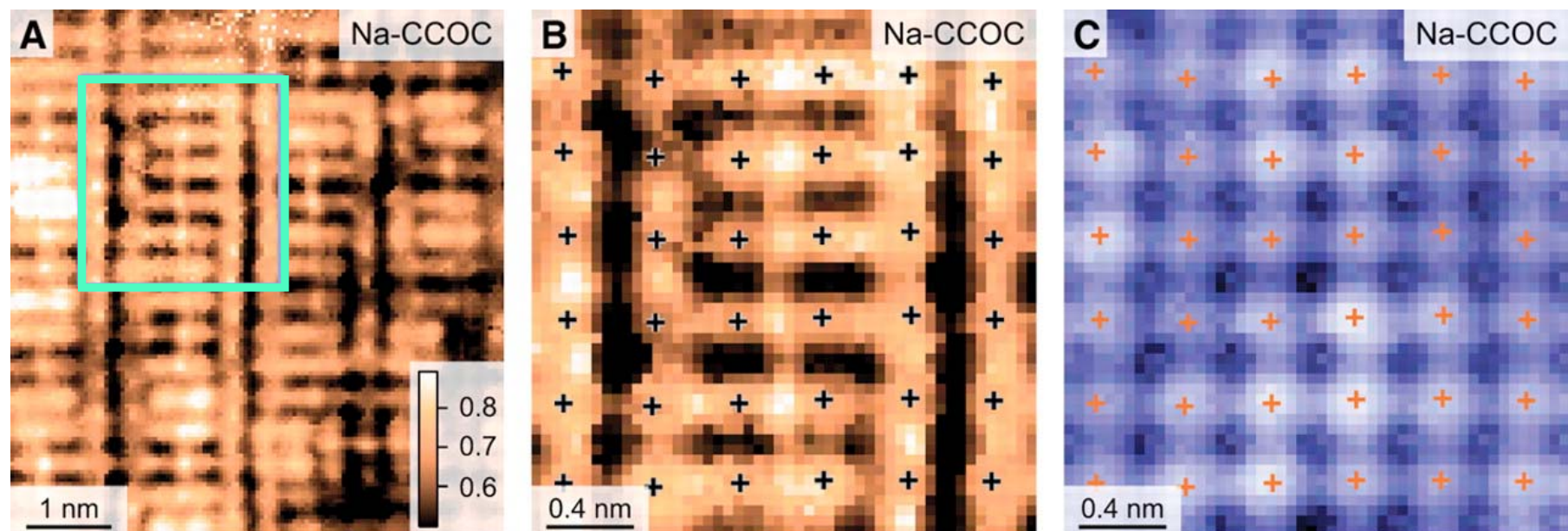
Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)



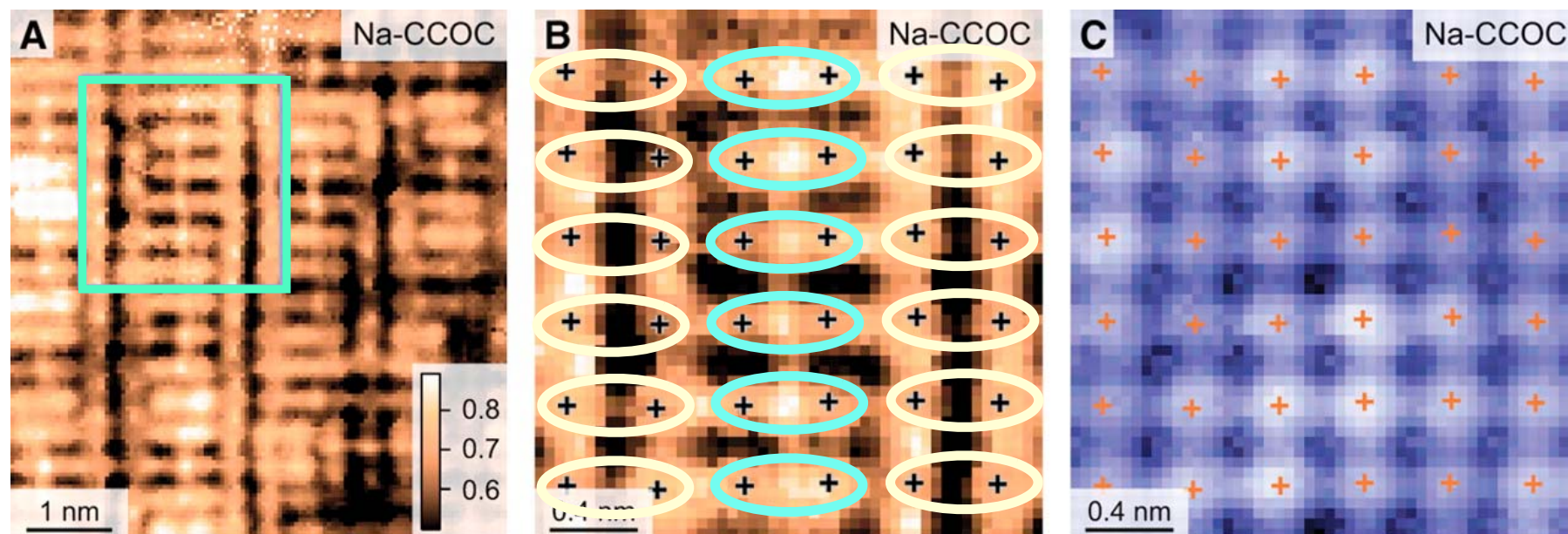
Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)



Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)



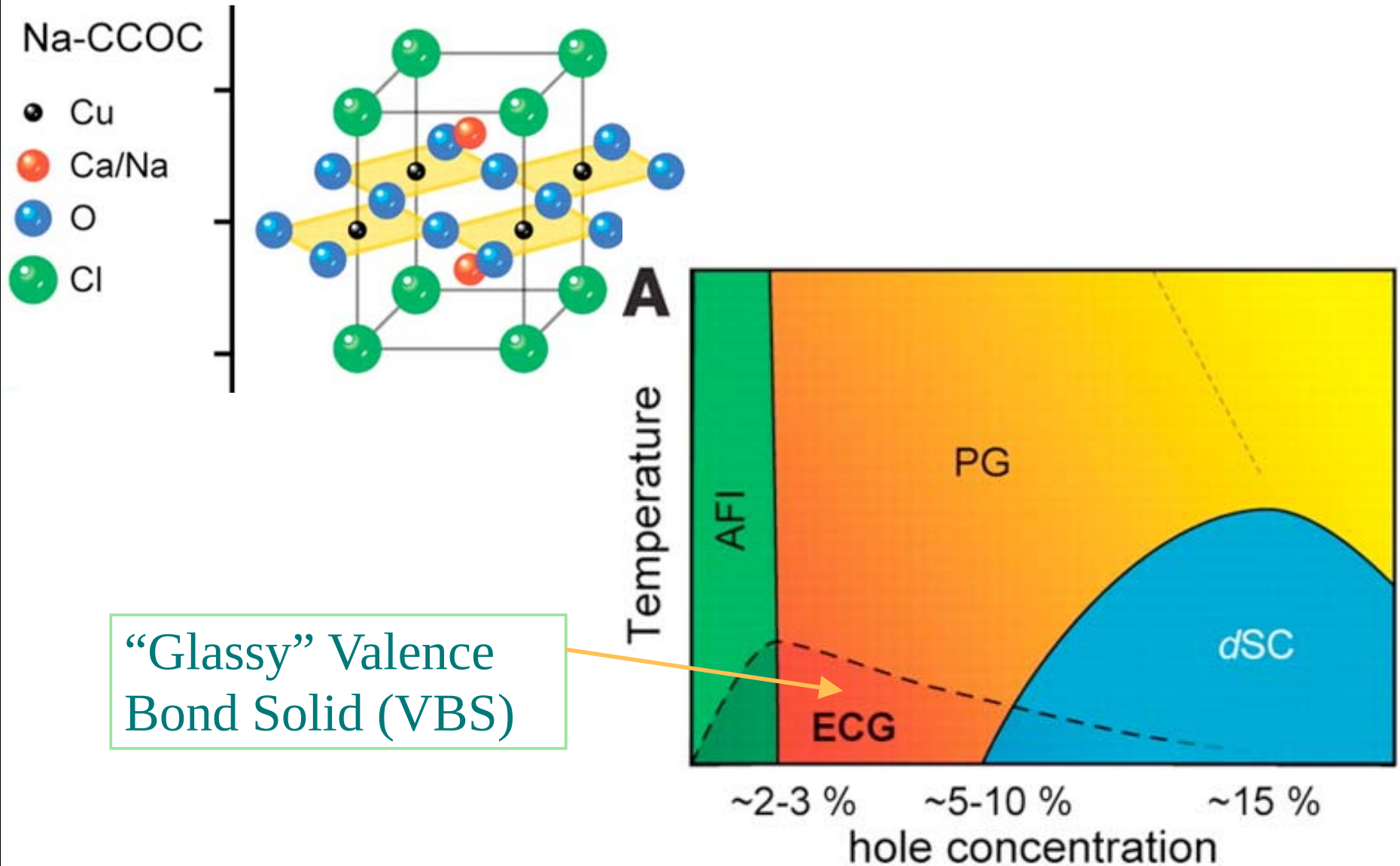
Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)



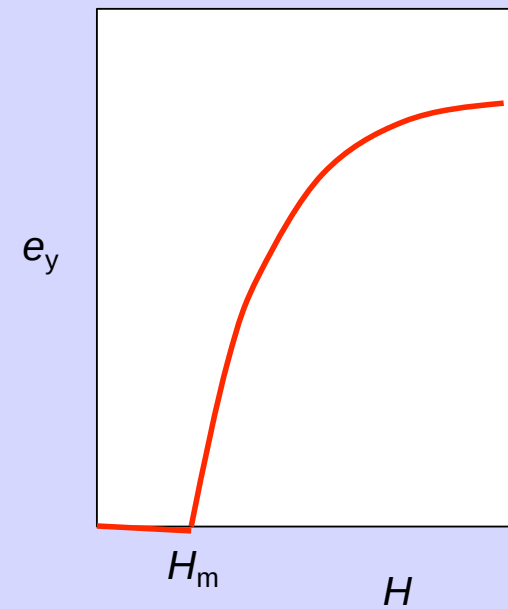
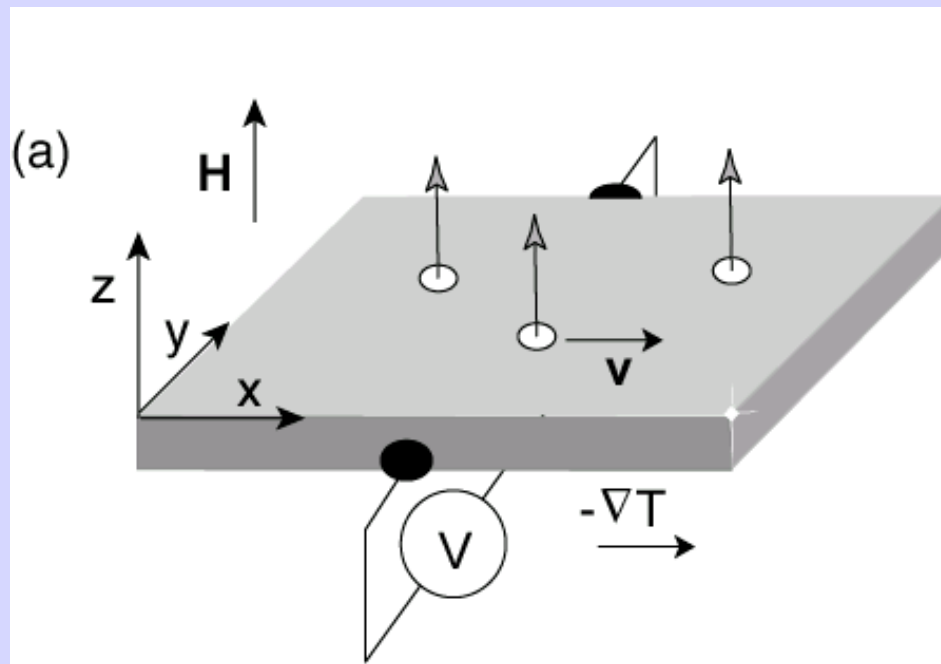
“Glassy” Valence Bond Solid (VBS)

Y. Kohsaka, C. Taylor, K. Fujita, A. Schmidt, C. Lupien, T. Hanaguri, M. Azuma, M. Takano, H. Eisaki, H. Takagi, S. Uchida, and J. C. Davis, *Science* **315**, 1380 (2007)

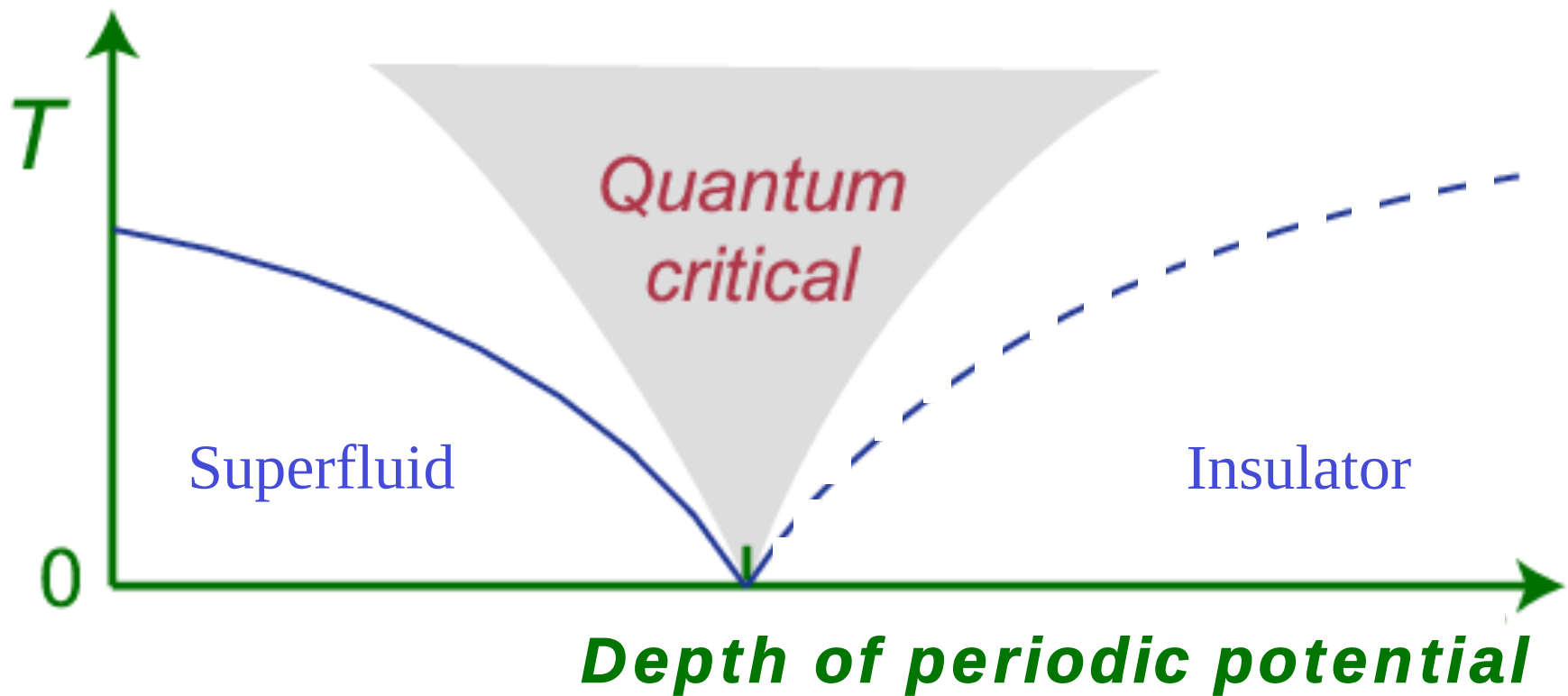
Temperature-doping phase diagram of the cuprate superconductors



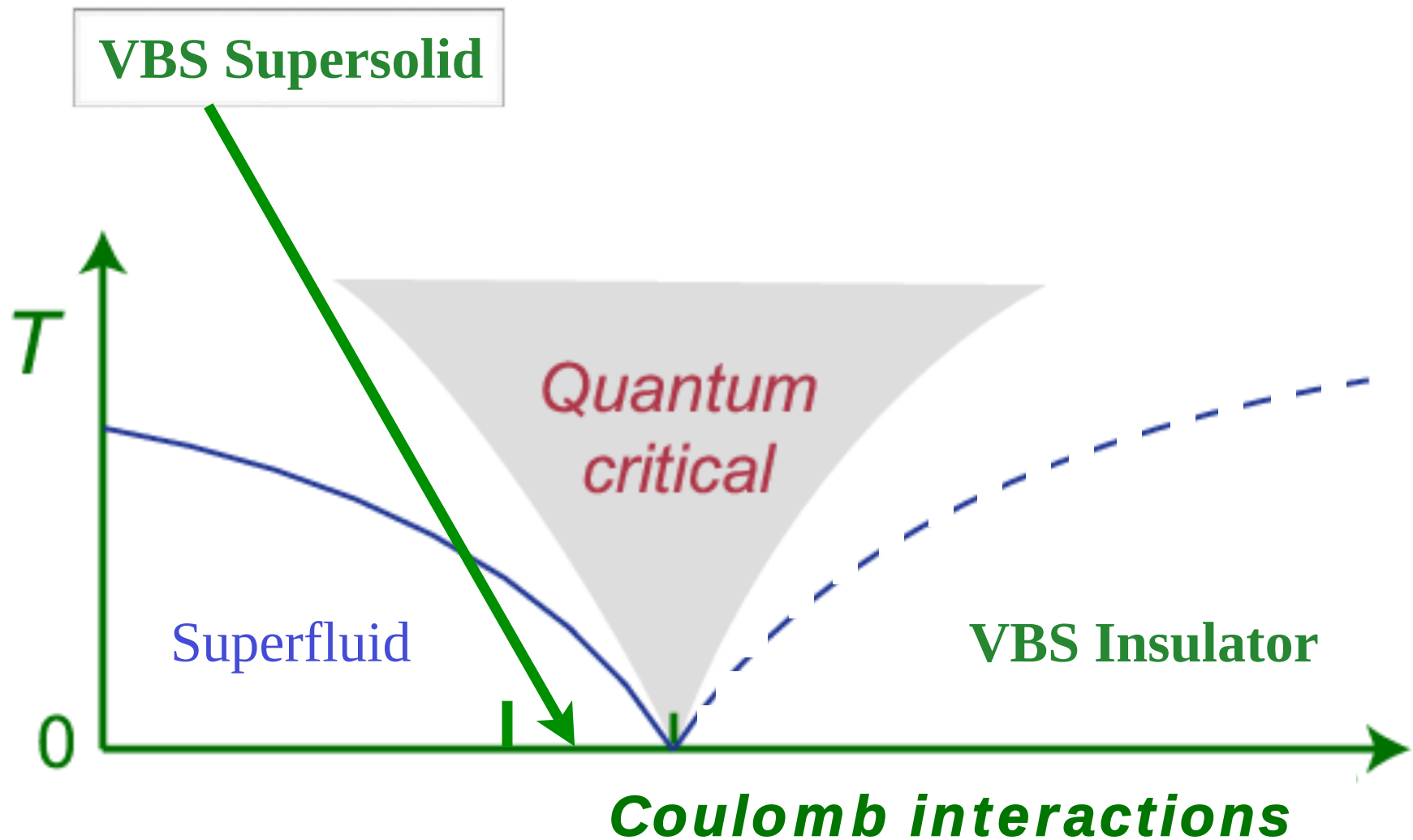
Nernst experiment



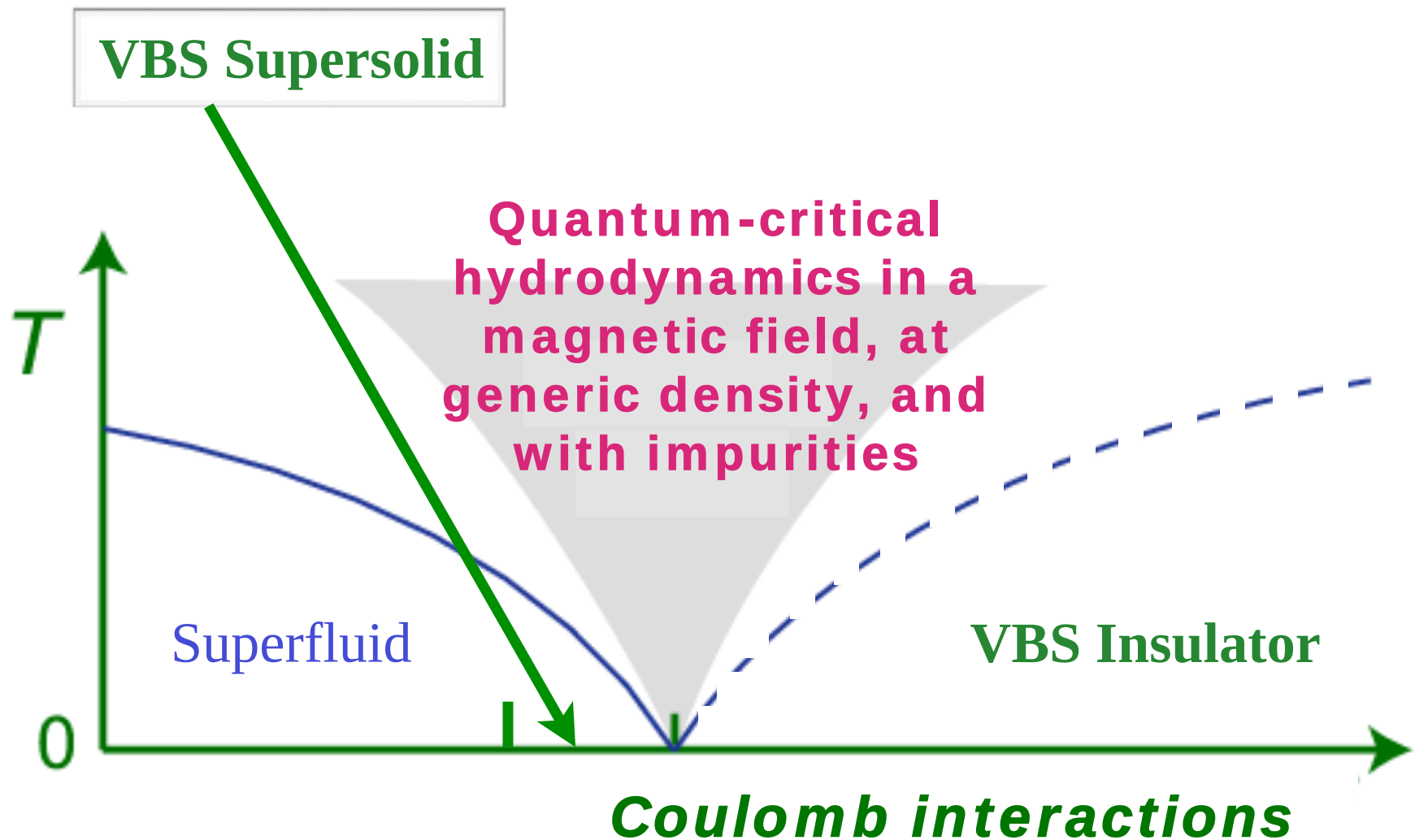
Non-zero temperature phase diagram



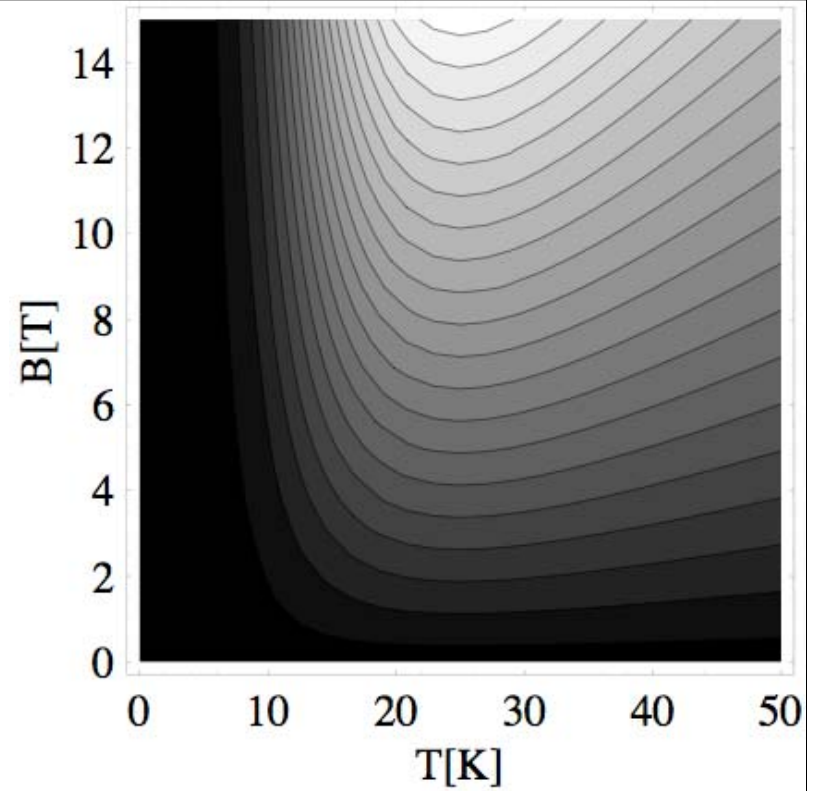
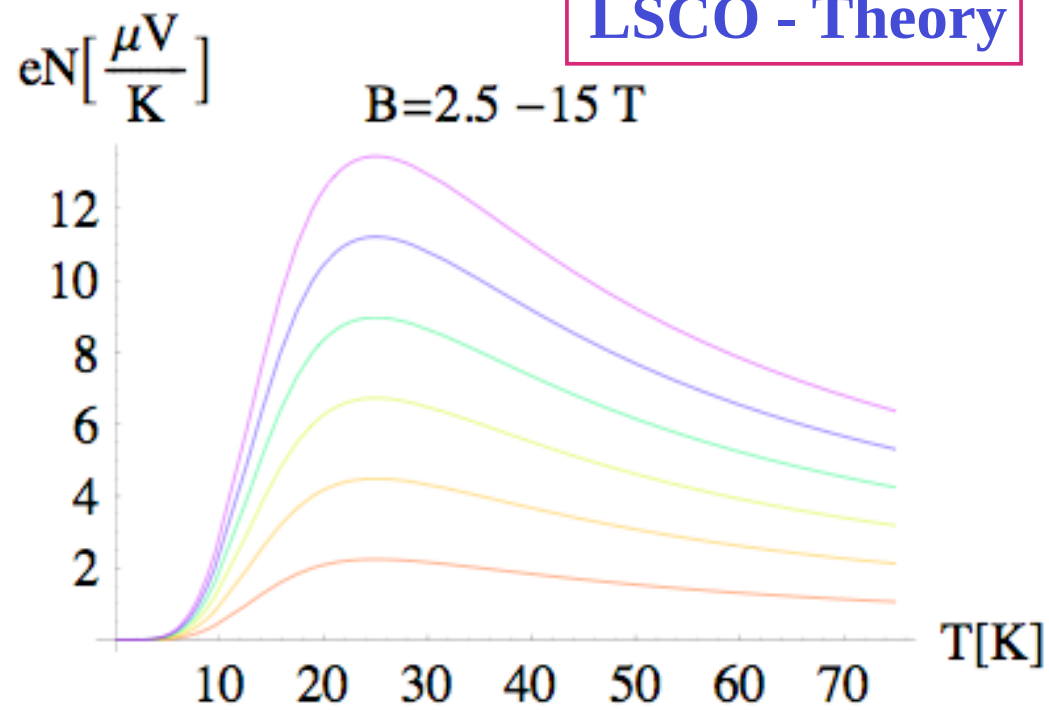
Non-zero temperature phase diagram



Non-zero temperature phase diagram

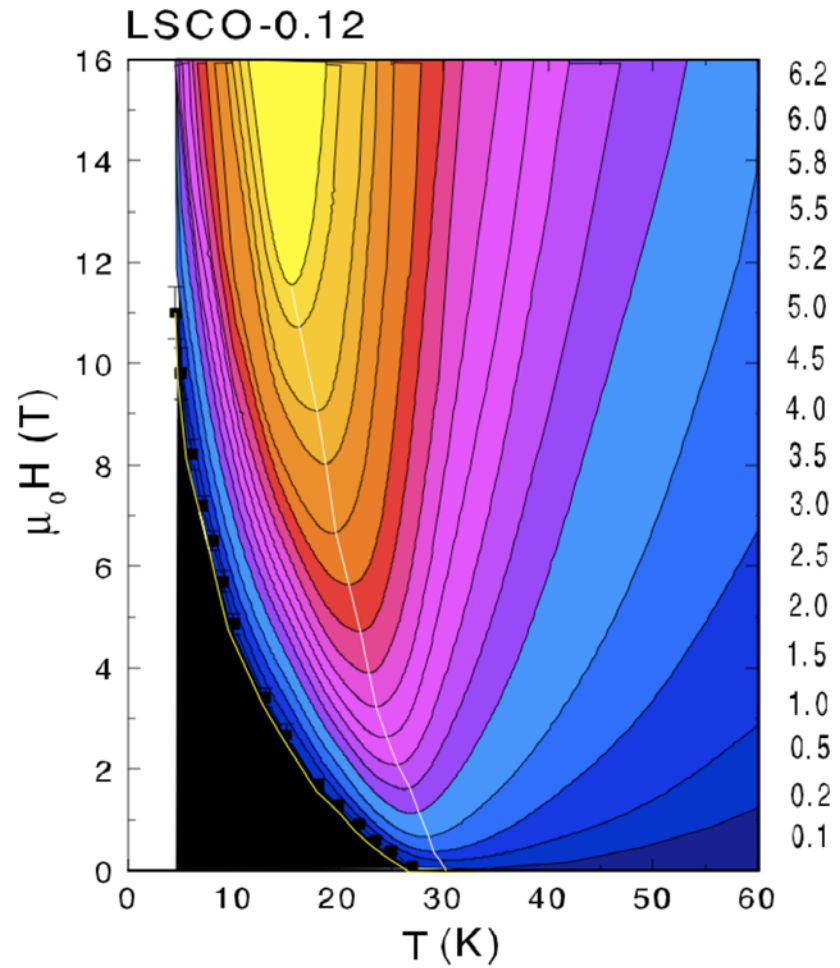
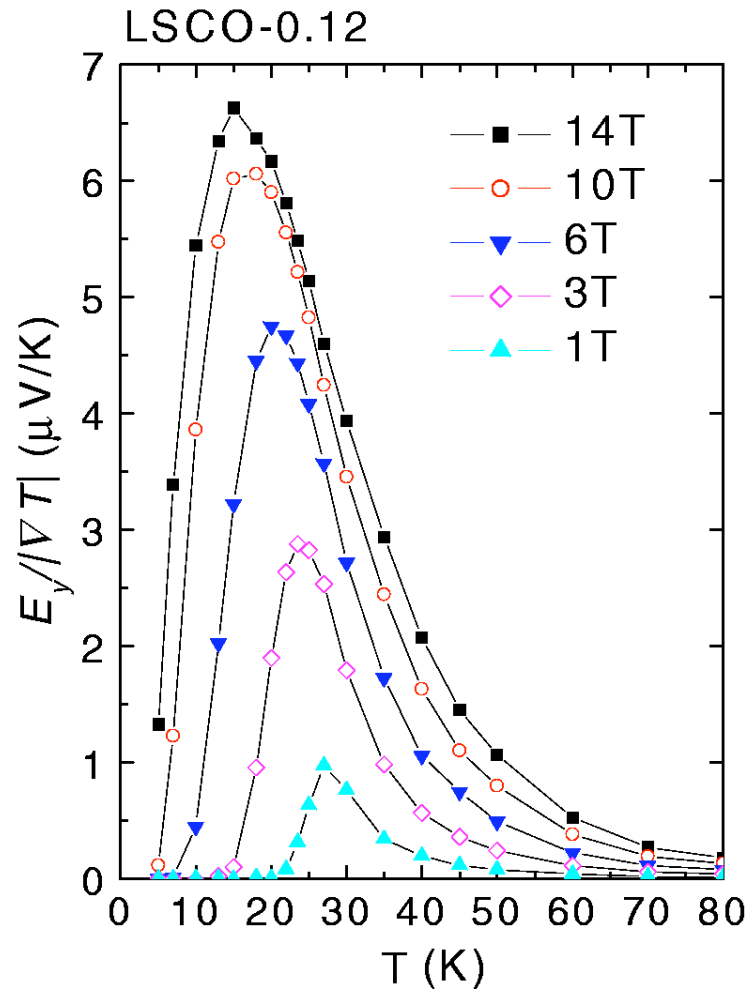


LSCO - Theory



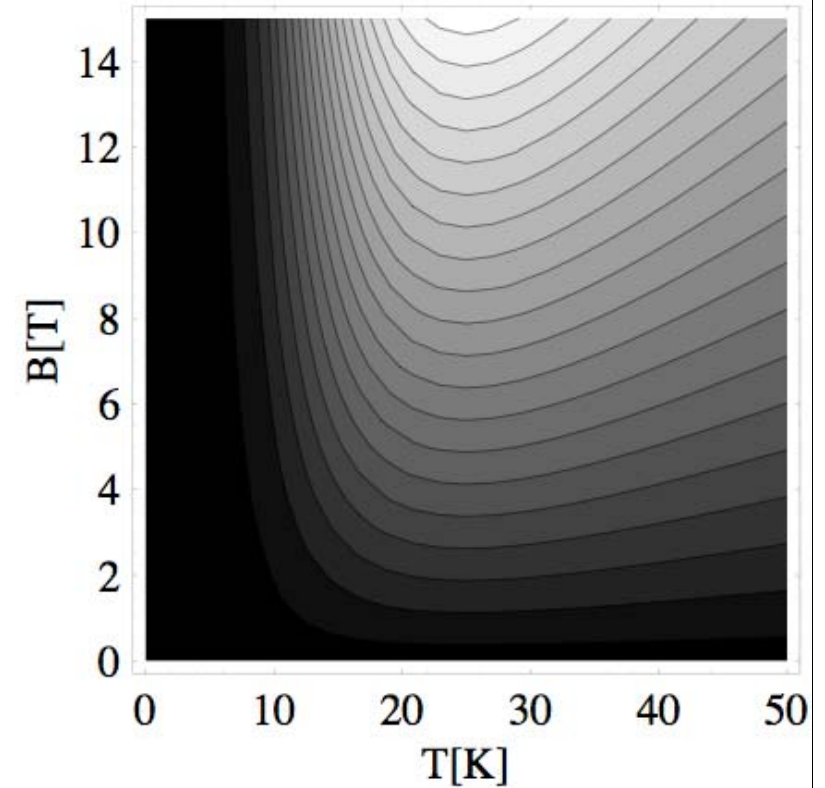
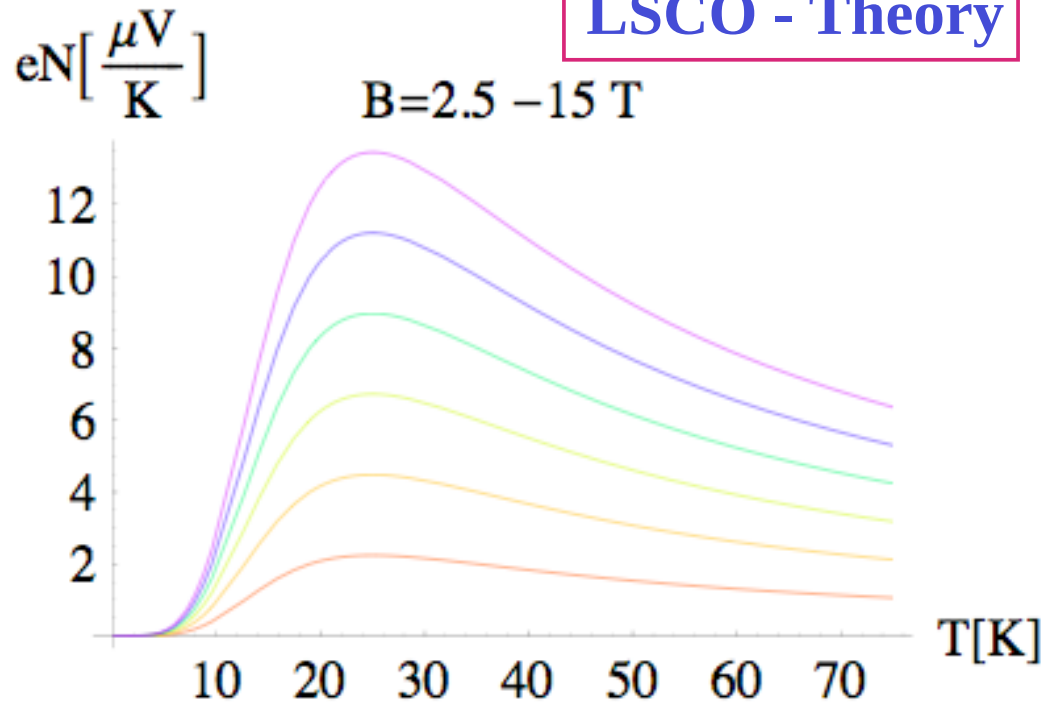
S.A. Hartnoll, P.K. Kovtun, M. Müller, and S. Sachdev, arXiv:0706.3215

LSCO - Experiments



N. P. Ong *et al.*

LSCO - Theory



Only input parameters

$$\hbar v = 47 \text{ meV } \text{\AA}$$

$$\tau_{\text{imp}} \approx 10^{-12} \text{ s}$$

Output

$$\omega_c = 6.2 \text{ GHz} \cdot \frac{B}{1 \text{ T}} \left(\frac{35 \text{ K}}{T} \right)^3$$

Similar to velocity estimates by
A.V. Balatsky and Z-X. Shen, *Science* **284**, 1137 (1999).

S.A. Hartnoll, P.K. Kovtun, M. Müller, and S. Sachdev, arXiv:0706.3215

To the solvable supersymmetric, Yang-Mills theory CFT, we add

- A chemical potential μ
- A magnetic field B

After the AdS/CFT mapping, we obtain the Einstein-Maxwell theory of a black hole with

- An electric charge
- A magnetic charge

The exact results are found to be in *precise* accord with *all* hydrodynamic results presented earlier

Conclusions

- Studies of new materials and trapped ultracold atoms are yielding new quantum phases, with novel forms of quantum entanglement.
- Some materials are of technological importance: e.g. high temperature superconductors.
- Exact solutions via black hole mapping have yielded first exact results for transport co-efficients in interacting many-body systems, and were valuable in determining general structure of hydrodynamics.
- Theory of VBS order and Nernst effect in cuprates.
- Tabletop “laboratories for the entire universe”: quantum mechanics of black holes, quark-gluon plasma, neutrons stars, and big-bang physics.