

Quantum phase transitions in condensed matter

Session L04: Lars Onsager Prize
APS March Meeting, Los Angeles

Subir Sachdev
March 7, 2018



Talk online: sachdev.physics.harvard.edu



Thanks to students and postdocs, and many other collaborators

- [Jinwu Ye](#), Associate Professor, Department of Physics and Astronomy, Mississippi State University
Thesis: [Some Examples of Quantum Phase Transitions](#)
- [T. Senthil](#), Professor, Department of Physics, Massachusetts Institute of Technology.
Thesis: [Quantum Phase Transitions in Random Spin Systems](#)
- [Kedar Damle](#), Department of Theoretical Physics, Tata Institute of Fundamental Research, Mumbai, India.
Thesis: [Turning on the Heat: Non-zero Temperature Dynamical Properties of Quantum Many-body Systems](#)
- [Chiranjeeb Buragohain](#), Microsoft Research.
Thesis: [Dynamical Properties of Quantum Antiferromagnets in One and Two Dimensions](#)
- [Ying Zhang](#), Finisterre Capital, London.
Thesis: [Competing Orders in the Cuprate Superconductors](#)
- [Anatoli Polkovnikov](#), Associate Professor, Boston University.
Thesis: [Manifestation of Quantum Fluctuations in Strongly Correlated Systems](#)
- [Stephen Powell](#), Assistant Professor, University of Nottingham
Thesis: [Quantum phases and transitions of many-body systems realized using cold atomic gases](#)
- [Adrian Del Maestro](#), Associate Professor, University of Vermont
Thesis: [The superconductor-metal quantum phase transition in ultra-narrow wires](#)
- [Emily Dunkel](#) (with [David Coker](#), Boston University), NASA Jet Propulsion Laboratory
Thesis: [Quantum Phenomena in Condensed Phase Systems](#)
- [Yang Qi](#), Institute for Advanced Studies, Tsinghua University
Thesis: [Spin and Charge Fluctuations in Strongly Correlated Systems](#).
- [Rudro Rana Biswas](#), Assistant Professor, Purdue University
Thesis: [Explorations in Dirac Fermions and Spin Liquids](#).
- [Eun Gook Moon](#), Assistant Professor, Korea Advanced Institute of Science and Technology
Thesis: [Superfluidity in Strongly Correlated Systems](#)
- [Max Metlitski](#), Assistant Professor, Department of Physics, Massachusetts Institute of Technology
Thesis: [Aspects of Critical Behavior of Two Dimensional Electron Systems](#)
- [Yejin Huh](#), Applied Scientist at Apple
Thesis: [Quantum Phase Transitions in d-wave Superconductors and Antiferromagnetic Kagome Lattices](#)
- [Susanne Pielawa](#), Lyft, Munich
Thesis: [Metastable Phases and Dynamics of Low-Dimensional Strongly-Correlated Atomic Quantum Gases](#)
- [Debanjan Chowdhury](#), Moore Foundation Postdoctoral Fellow, MIT
Thesis: [Interplay of Broken Symmetries and Quantum Criticality in Correlated Electronic Systems](#)
- [Junhyun Lee](#), Postdoctoral fellow, University of Maryland
Thesis: [Novel quantum phase transitions in low-dimensional systems](#)
- [Andrew Lucas](#), Postdoctoral fellow, Stanford University
Thesis: [Transport and hydrodynamics in holography, strange metals and graphene](#)
- [Shubhayu Chatterjee](#), Harvard University
- [Aavishkar Patel](#), Harvard University
- [Wenbo Fu](#), Harvard University
- [Seth Whitsitt](#), Harvard University
- Alex Thomson, Harvard University
- [Julia Steinberg](#), Harvard University

Students



Thanks to students and postdocs, and many other collaborators

- [Pierre Le Doussal](#), Directeur de Recherche de Classe Exceptionnelle, Laboratoire de Physique Théorique de l'École Normale Supérieure, Paris, France.
- [Rodolfo Jalabert](#), Professeur à l'Université Louis Pasteur, Institut de Physique et Chimie des Matériaux de Strasbourg, France.
- [Andrey Chubukov](#), William I. and Bianca M. Fine Chair in Theoretical Physics, University of Minnesota, Minneapolis.
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- [Predrag Nikolic](#), Associate Professor, George Mason University
- [Ribhu Kaul](#), Associate Professor, University of Kentucky
- [Markus Müller](#), Scientist, Paul Scherrer Institute, Switzerland.
- [Lars Fritz](#), Assistant Professor, University of Utrecht
- [Michael Levin](#), Associate Professor, University of Chicago
- [Cenke Xu](#), Associate Professor, University of California, Santa Barbara
- [Sean Hartnoll](#), Associate Professor, Stanford University
- [Erez Berg](#), Associate Professor, University of Chicago
- [Liang Fu](#), Lawrence C. (1944) and Sarah W. Biedenharn Career Development Associate Professor of Physics, Massachusetts Institute of Technology
- [Liza Huijse](#), Software Engineer at [Karius, Inc.](#)
- [Chris Laumann](#), Assistant Professor, Boston University
- [Matthias Punk](#), Faculty, LMU Munich
- [Philipp Strack](#), ZEISS Group
- [Brian Swingle](#), Assistant Professor, University of Maryland
- [Dmitry Abanin](#), Professor of Physics, University of Geneva
- [Ling-Yan \(Janet\) Hung](#), Professor of Physics, Fudan University, Shanghai
- [Jay Sau](#), Assistant Professor, University of Maryland
- [Sarang Gopalakrishnan](#), Postdoctoral Fellow, Caltech
- [Andrea Allais](#), Cruise Automation, San Francisco
- [Johannes Bauer](#), SCL Group, London
- [Paul Chesler](#), Harvard University
- [Andreas Eberlein](#), Harvard University
- [William Witczak-Krempa](#), Assistant Professor, University of Montreal
- [Richard Davison](#), Harvard University
- [Chong Wang](#), Harvard University
- [Mathias Scheurer](#), Harvard University.

Postdocs



Continuous quantum transitions

(A)

Broken symmetry

No broken symmetry



Continuous quantum transitions

(A)

Broken symmetry

No broken symmetry

Exact solution by
Onsager of the
 $D=2$ Ising model

Onsager Prizes:

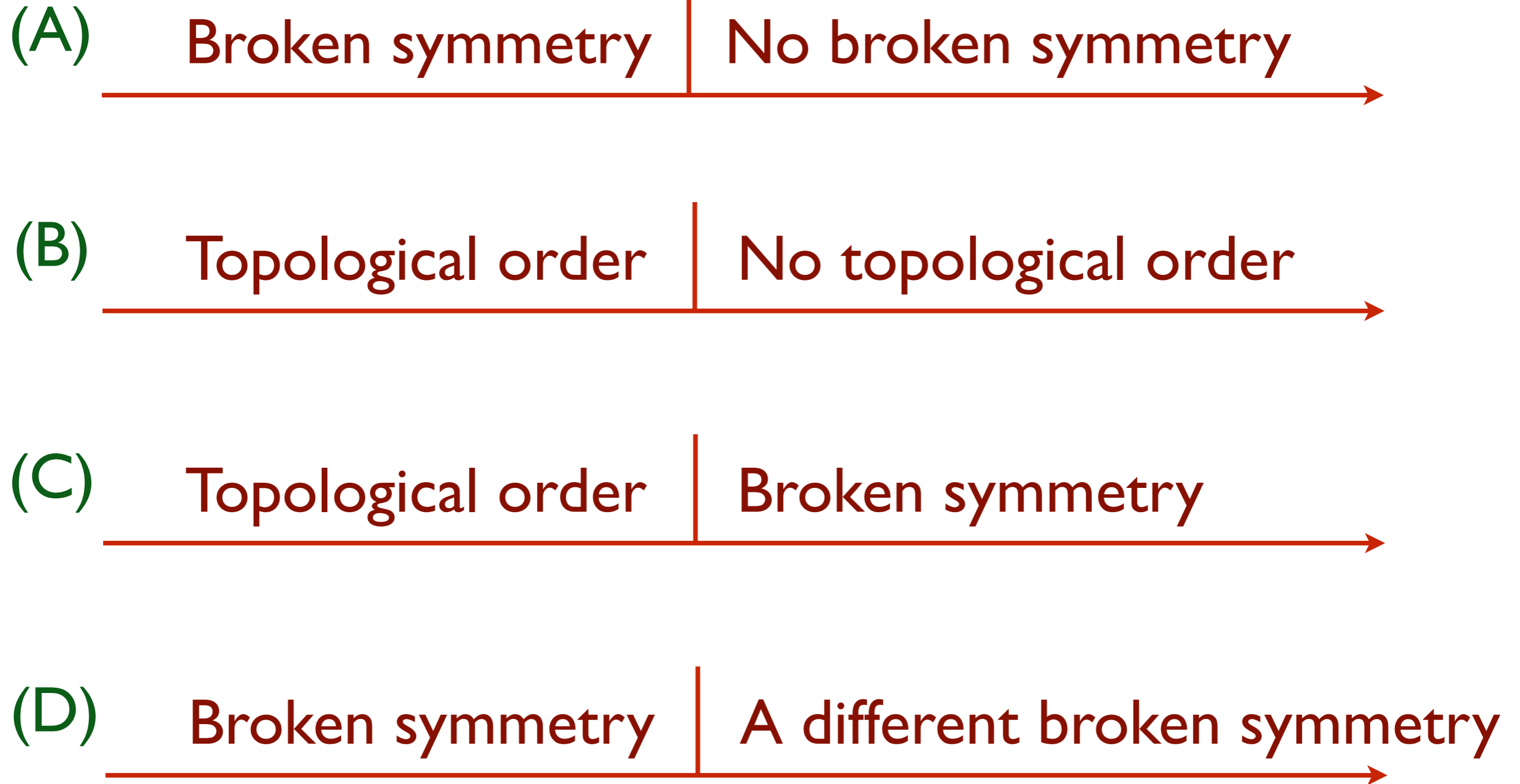
M.E. Fisher

L. Kadanoff

A. I. Larkin

V. L. Pokrovsky

Continuous quantum transitions

- (A) Broken symmetry | No broken symmetry
- (B) Topological order | No topological order
- (C) Topological order | Broken symmetry
- (D) Broken symmetry | A different broken symmetry
- 

Continuous quantum transitions

Theory with emergent gauge fields: Higgs/
confining phases and phase transitions

try

-
- (B) Topological order | No topological order
- (C) Topological order | Broken symmetry
- (D) Broken symmetry | A different broken symmetry

Continuous quantum transitions

(A) Broken symmetry | No broken symmetry



A horizontal red line with an arrow pointing to the right. A vertical red line is positioned in the middle of the horizontal line, dividing it into two equal segments. The text "Broken symmetry" is centered in the left segment, and "No broken symmetry" is centered in the right segment. The entire diagram is enclosed in a blue rounded rectangular border.

(B) Topological order | No topological order



A horizontal red line with an arrow pointing to the right. A vertical red line is positioned in the middle of the horizontal line, dividing it into two equal segments. The text "Topological order" is centered in the left segment, and "No topological order" is centered in the right segment.

(C) Topological order | Broken symmetry



A horizontal red line with an arrow pointing to the right. A vertical red line is positioned in the middle of the horizontal line, dividing it into two equal segments. The text "Topological order" is centered in the left segment, and "Broken symmetry" is centered in the right segment.

(D) Broken symmetry | A different broken symmetry



A horizontal red line with an arrow pointing to the right. A vertical red line is positioned in the middle of the horizontal line, dividing it into two equal segments. The text "Broken symmetry" is centered in the left segment, and "A different broken symmetry" is centered in the right segment.

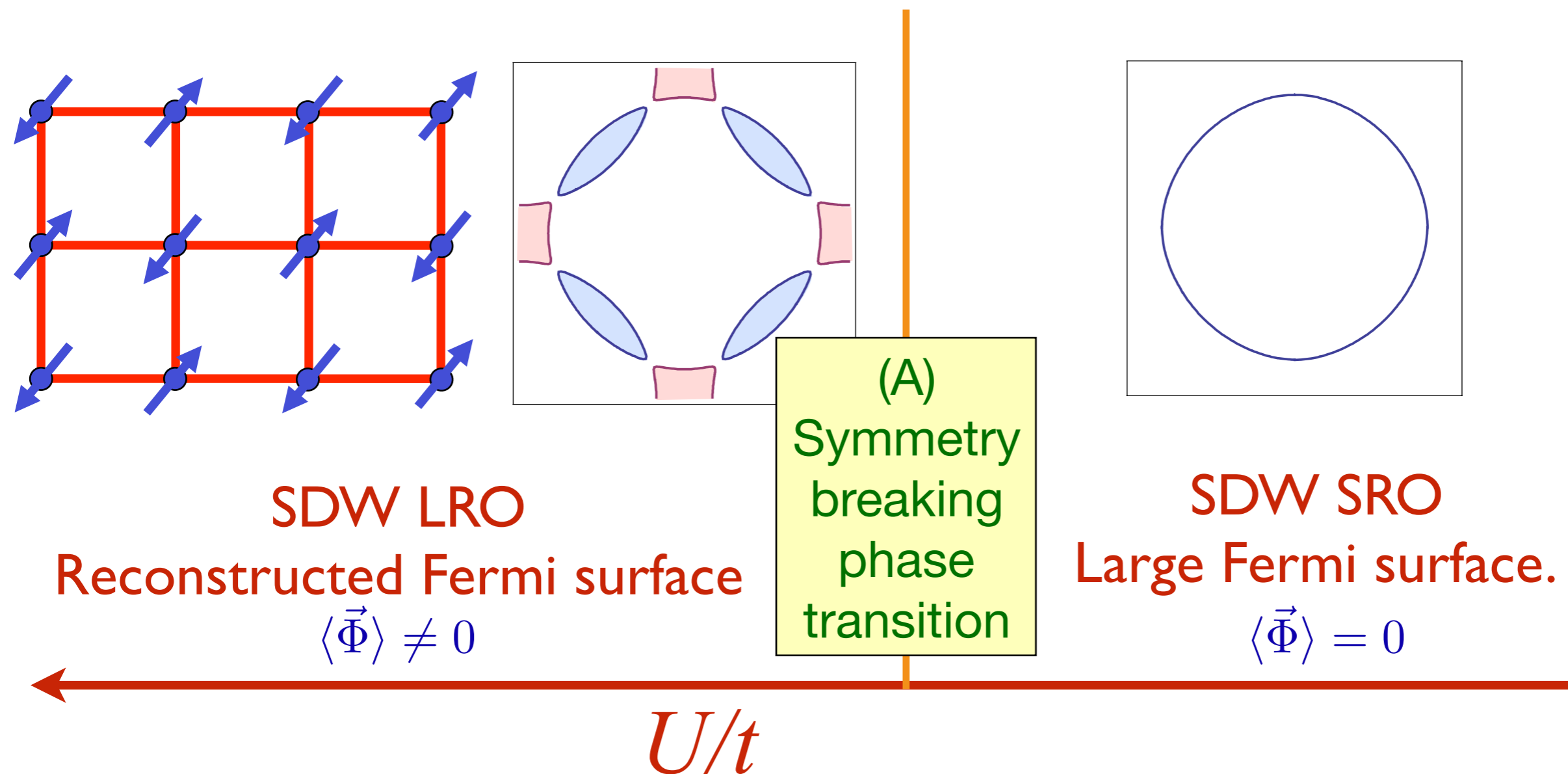
Antiferromagnetism in the Hubbard Model

$$H = - \sum_{i < j} t_{ij} c_{i\alpha}^\dagger c_{j\alpha} + U \sum_i \left(n_{i\uparrow} - \frac{1}{2} \right) \left(n_{i\downarrow} - \frac{1}{2} \right) - \mu \sum_i c_{i\alpha}^\dagger c_{i\alpha}$$

$t_{ij} \rightarrow$ "hopping". $U \rightarrow$ local repulsion, $\mu \rightarrow$ chemical potential

Mean-field theory with a spin density wave (SDW)

$$\text{order parameter } \vec{\Phi}_i = (-1)^{i_x+i_y} \langle c_{i\alpha}^\dagger \vec{\sigma}_{\alpha\beta} c_{i\beta} \rangle / 2$$

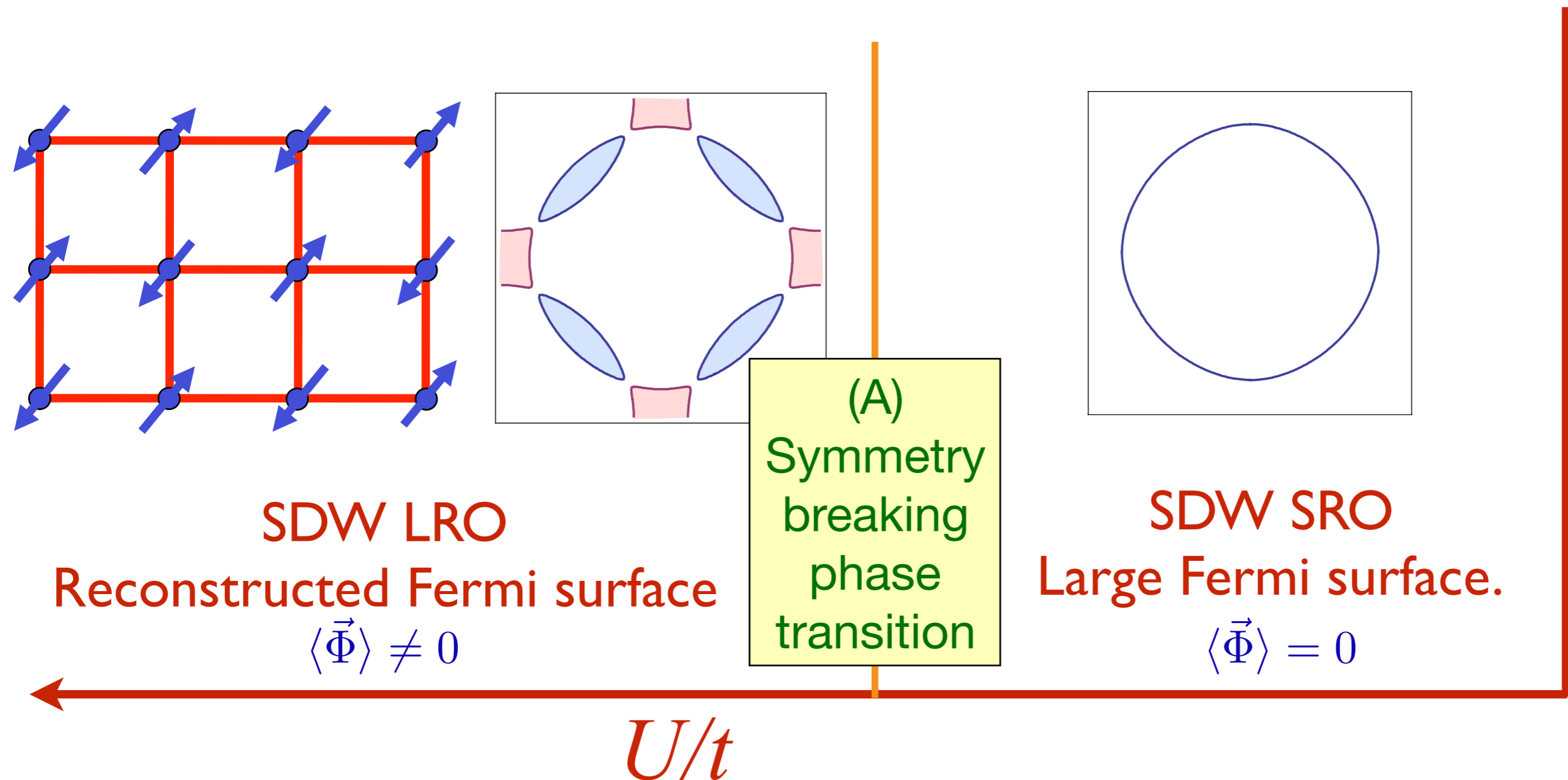


Antiferromagnetism in the Hubbard Model

$$H = - \sum_{i < j} t_{ij} c_{i\alpha}^\dagger c_{j\alpha} + U \sum_i \left(n_{i\uparrow} - \frac{1}{2} \right) \left(n_{i\downarrow} - \frac{1}{2} \right) - \mu \sum_i c_{i\alpha}^\dagger c_{i\alpha}$$

$t_{ij} \rightarrow$ "hopping". $U \rightarrow$ local repulsion, $\mu \rightarrow$ chemical potential

Both states have Luttinger volume Fermi surfaces



Continuous quantum transitions

(A) Broken symmetry | No broken symmetry

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Broken symmetry" and the right side contains "No broken symmetry".

(B) Topological order | No topological order

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Topological order" and the right side contains "No topological order". The entire diagram is enclosed in a blue rounded rectangular border.

(C) Topological order | Broken symmetry

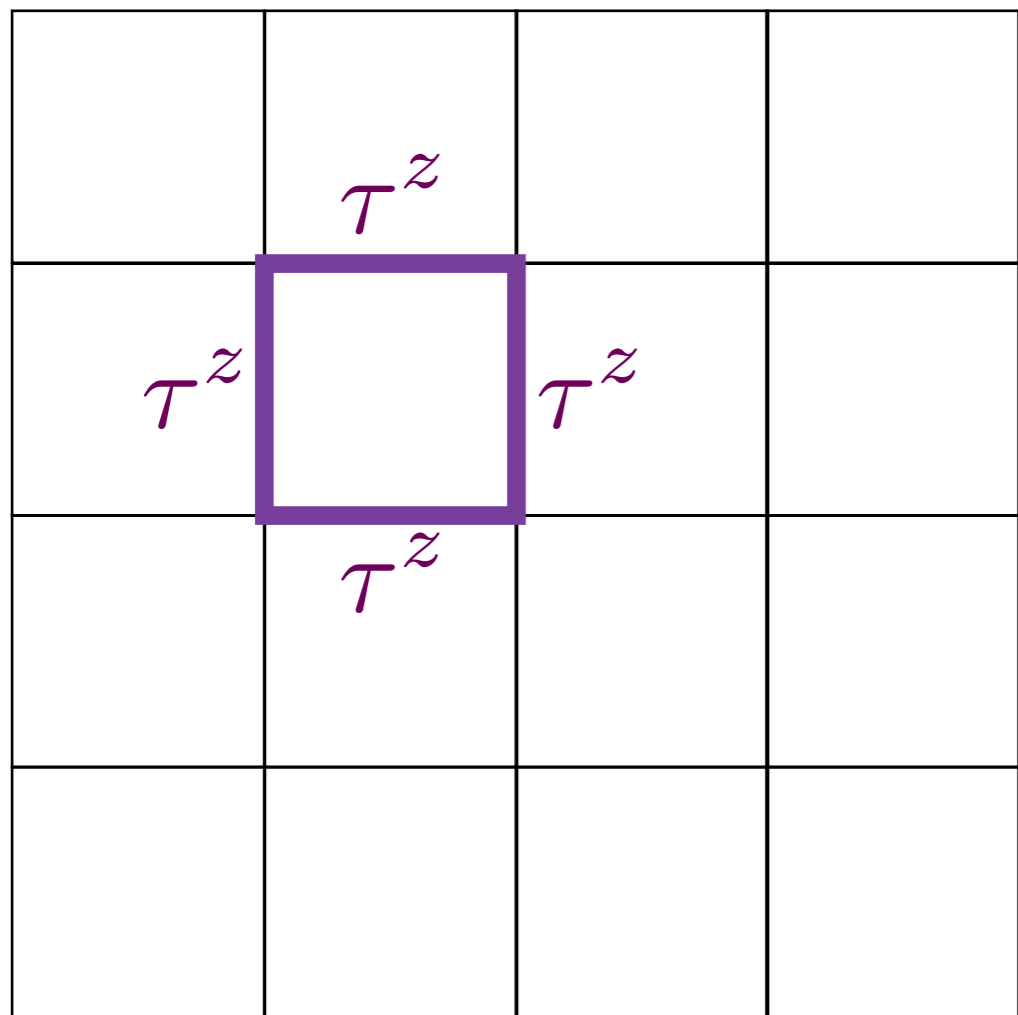
A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Topological order" and the right side contains "Broken symmetry".

(D) Broken symmetry | A different broken symmetry

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Broken symmetry" and the right side contains "A different broken symmetry".

Z_2 lattice gauge theory

(Wegner, 1971
Onsager Prize, 2015)



$$H = - \sum_{\square} \tau^z \tau^z \tau^z \tau^z - g \sum_i \tau^x$$

$$G_i = \begin{array}{c|c} \tau^x & \tau^x \\ \hline & \tau^x \end{array}$$

Gauss's Law: $[H, G_i] = 0$, $G_i = 1$

\mathbb{Z}_2 lattice gauge theory

Deconfined phase.
 \mathbb{Z}_2 flux expelled.
 \mathbb{Z}_2 (toric code)
topological order.

(B)
Topological
phase
transition

Confined phase.
 \mathbb{Z}_2 flux proliferates.
No topological order.

$$H = - \sum_{\square} \tau^z \tau^z \tau^z \tau^z - g \sum_i \tau^x$$

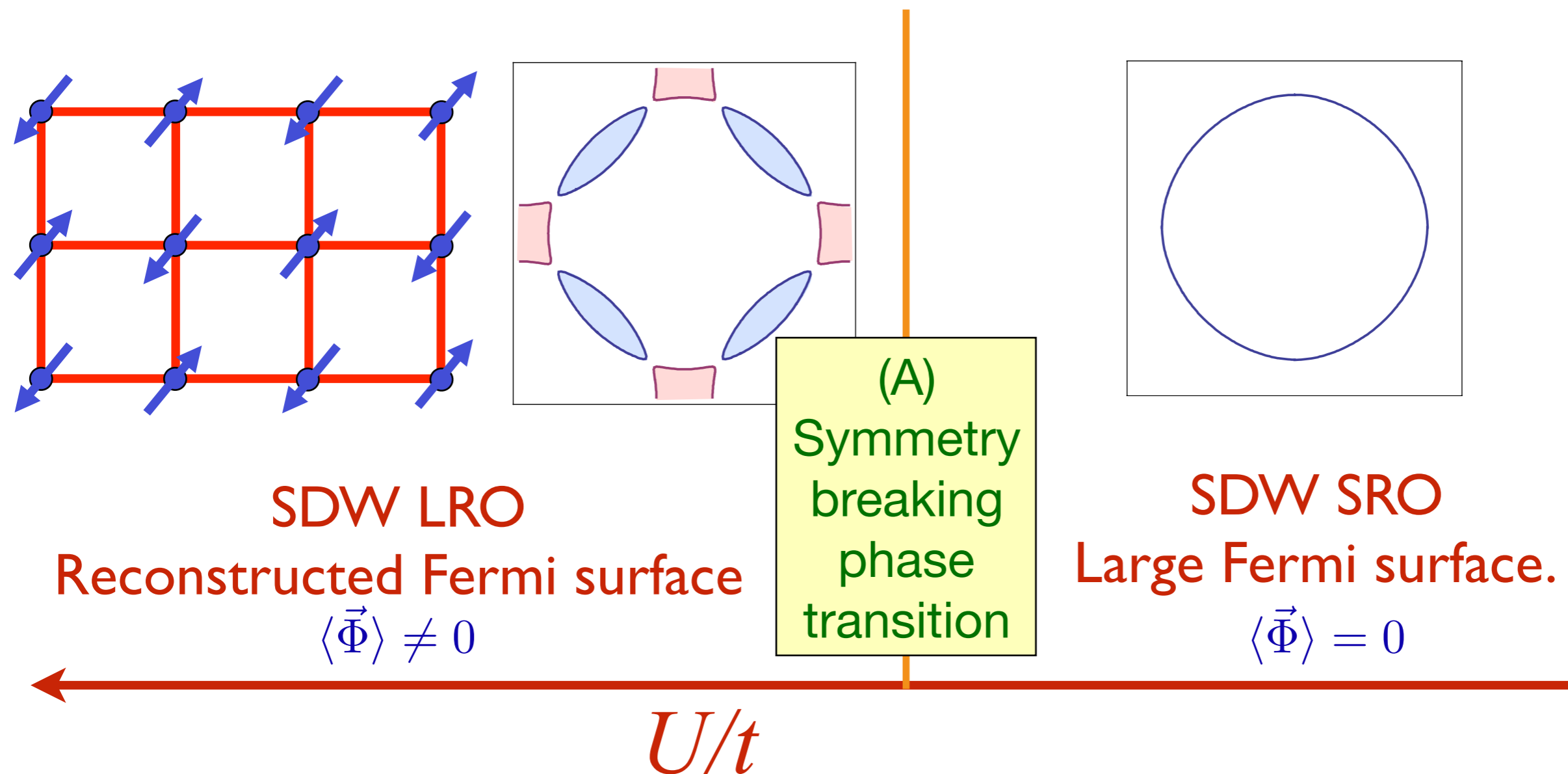
Antiferromagnetism in the Hubbard Model

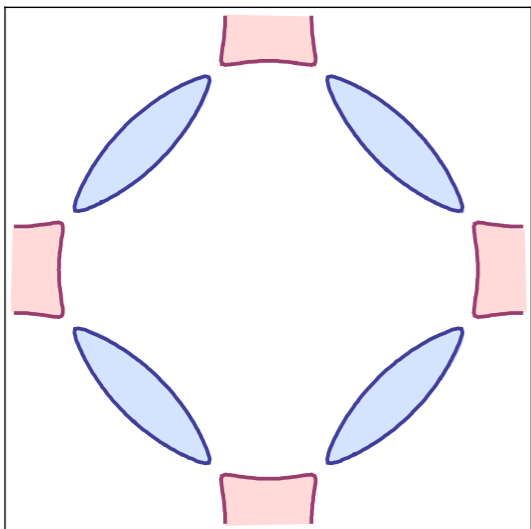
$$H = - \sum_{i < j} t_{ij} c_{i\alpha}^\dagger c_{j\alpha} + U \sum_i \left(n_{i\uparrow} - \frac{1}{2} \right) \left(n_{i\downarrow} - \frac{1}{2} \right) - \mu \sum_i c_{i\alpha}^\dagger c_{i\alpha}$$

$t_{ij} \rightarrow$ "hopping". $U \rightarrow$ local repulsion, $\mu \rightarrow$ chemical potential

Mean-field theory with a spin density wave (SDW)

$$\text{order parameter } \vec{\Phi}_i = (-1)^{i_x+i_y} \langle c_{i\alpha}^\dagger \vec{\sigma}_{\alpha\beta} c_{i\beta} \rangle / 2$$



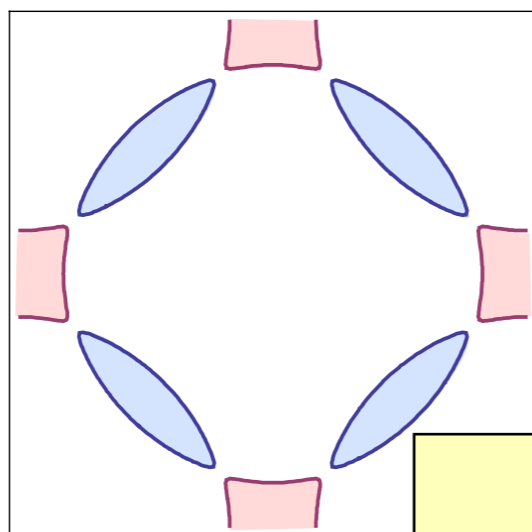
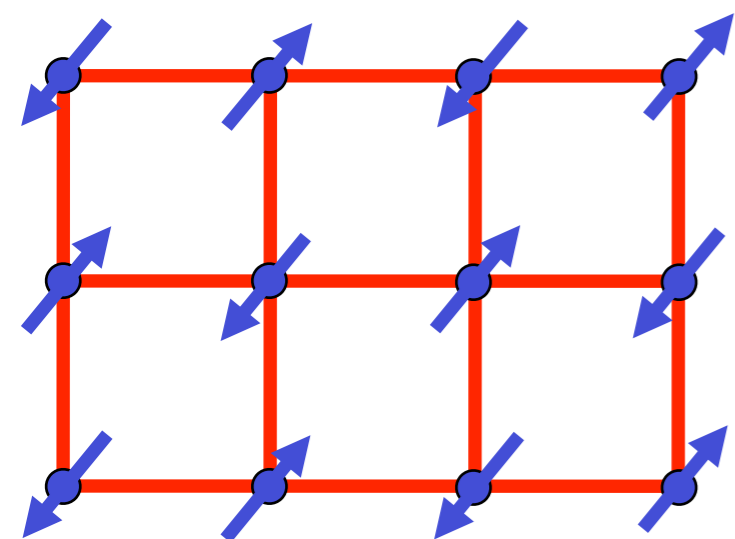


SDW SRO

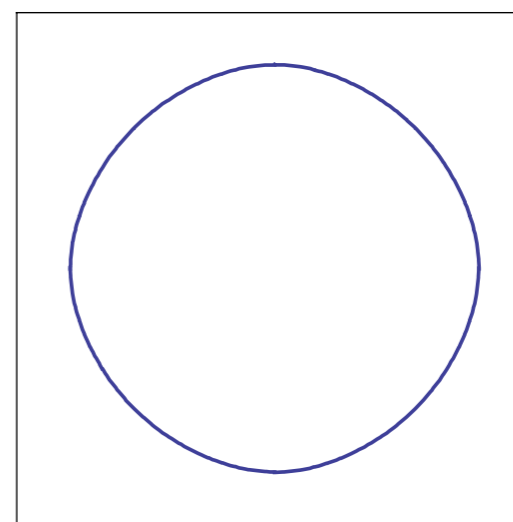
Z_2 or $U(1)$ topological order.
Reconstructed Fermi surface.

$$\langle \vec{\Phi} \rangle = 0$$

(B) Topological
phase transition



(A)
Symmetry
breaking
phase
transition



SDW LRO
Reconstructed Fermi surface

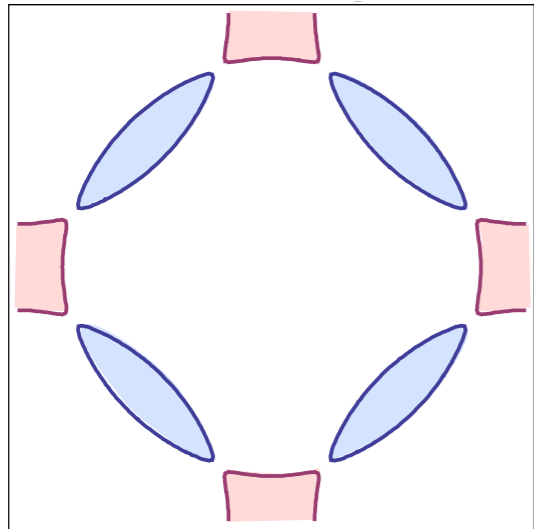
$$\langle \vec{\Phi} \rangle \neq 0$$

SDW SRO
Large Fermi surface.

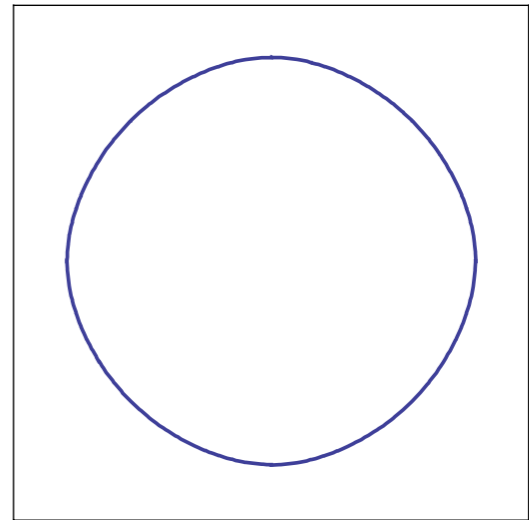
$$\langle \vec{\Phi} \rangle = 0$$

U/t

g



(B)
Topological
phase
transition



SDW SRO

Reconstructed Fermi surface
with non-Luttinger volume.

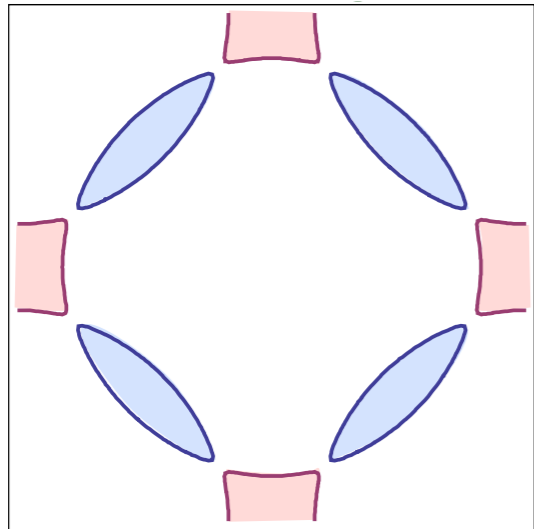
Z_2 vortices or hedgehogs expelled.

Z_2 or $U(1)$ topological order.

SDW SRO

Large Fermi surface
with Luttinger volume.
No topological order

Metallic states with non-Luttinger volume
Fermi surfaces must have topological order



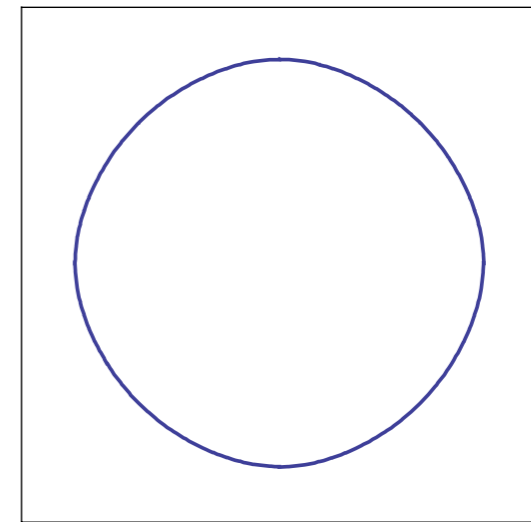
SDW SRO

Reconstructed Fermi surface
with non-Luttinger volume.

Z_2 vortices or hedgehogs expelled.

Z_2 or $U(1)$ topological order.

(B)
Topological
phase
transition;
phases of a
theory
with an emergent
 $SU(2)$ gauge field.



SDW SRO

Large Fermi surface
with Luttinger volume.
No topological order

Metallic states with non-Luttinger volume
Fermi surfaces must have topological order

T. Senthil, M. Vojta, and S. Sachdev, PRB **69**, 035111 (2004)

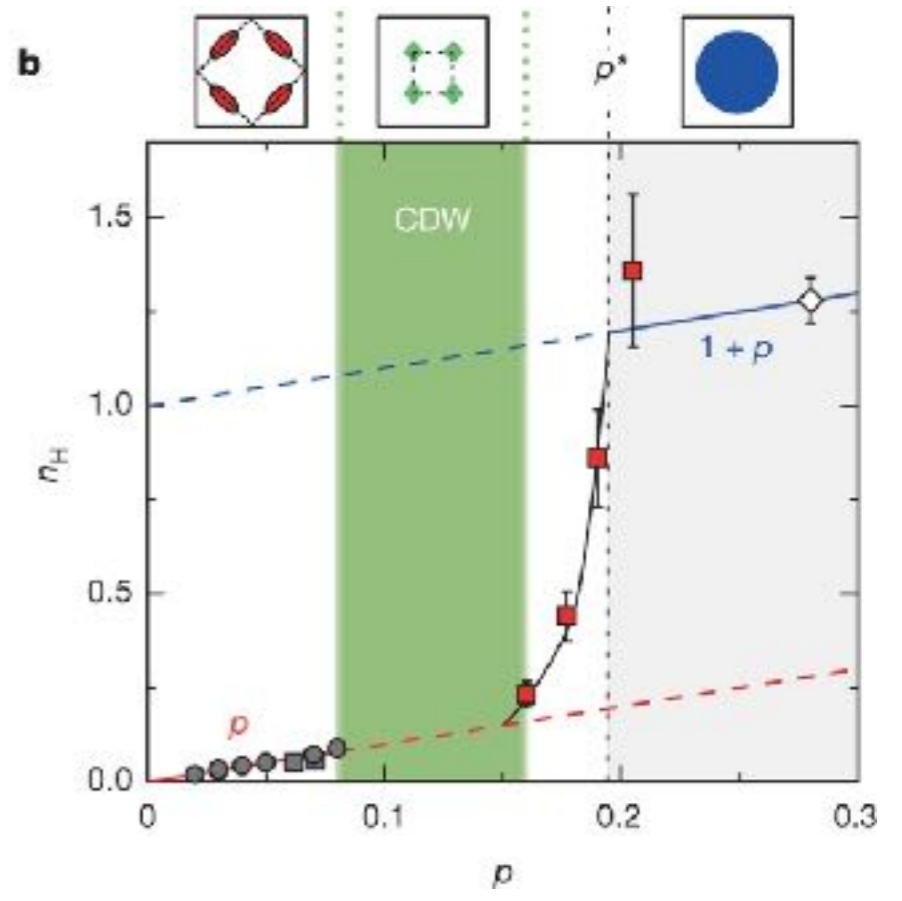
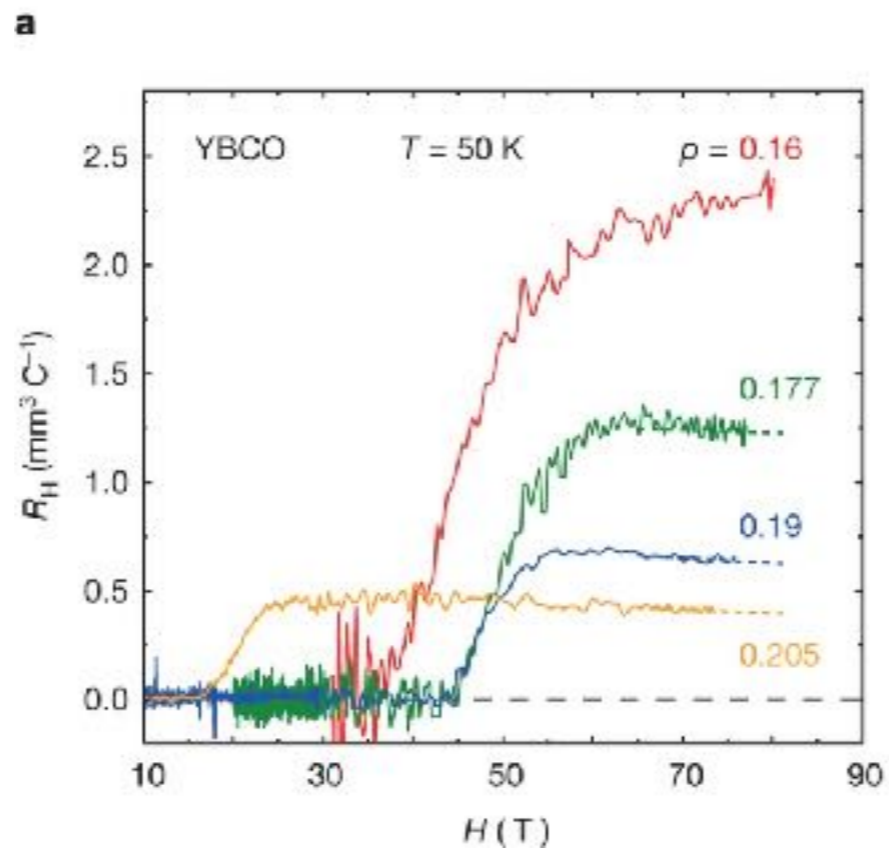
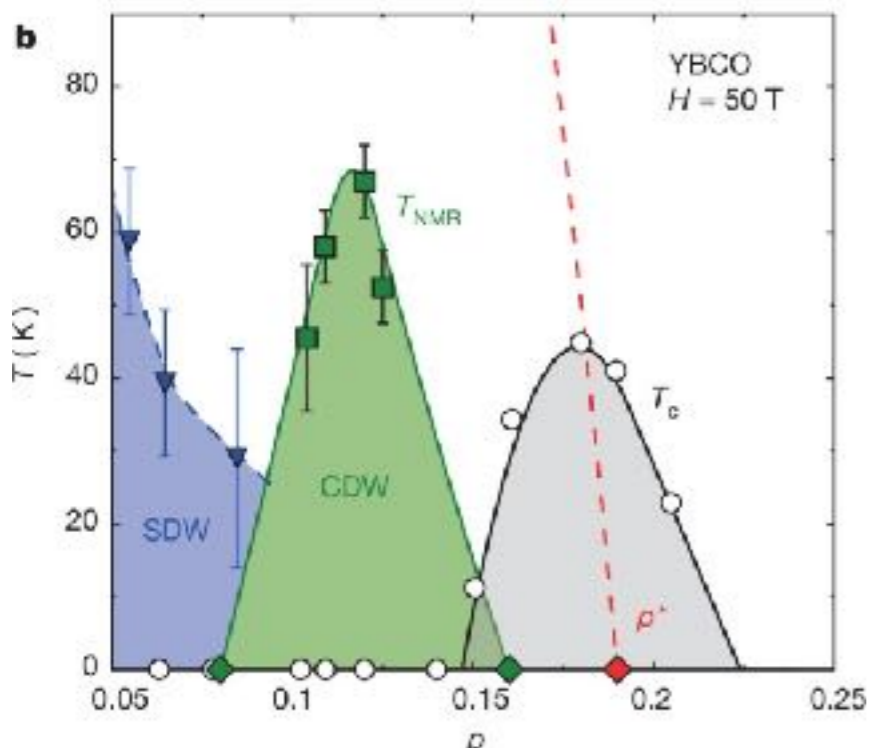
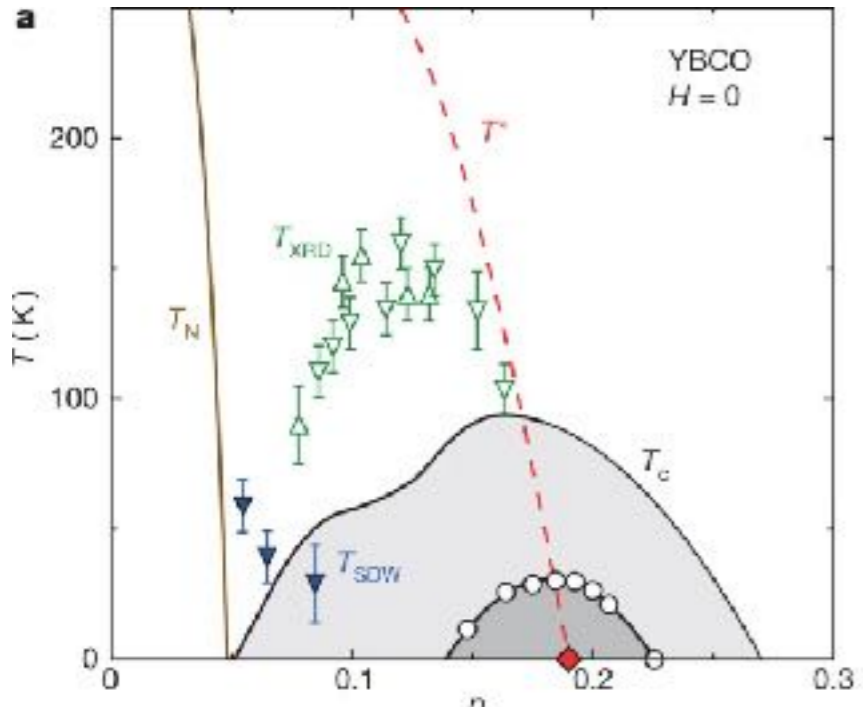
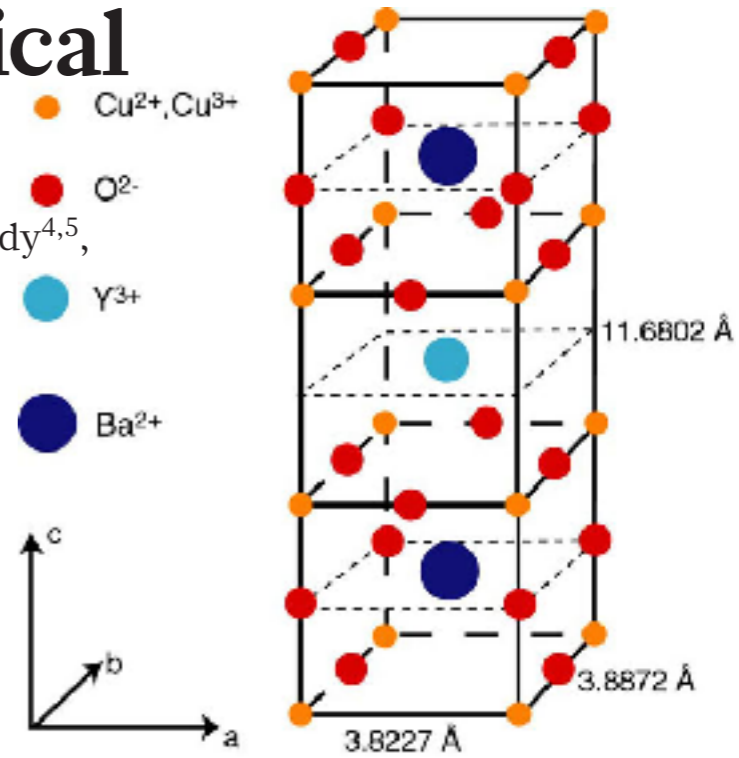
S. Sachdev, M.A. Metlitski, Y. Qi, and C. Xu, PRB **80**, 155129 (2009)

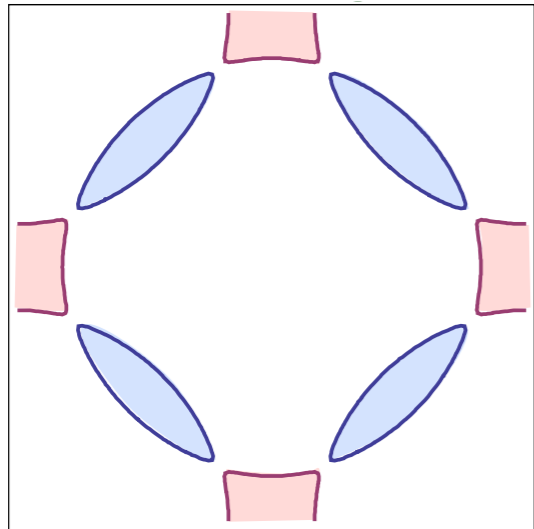
g

Change of carrier density at the pseudogap critical point of a cuprate superconductor

S. Badoux¹, W. Tabis^{2,3}, F. Laliberté², G. Grissonnanche¹, B. Vignolle², D. Vignolles², J. Béard², D. A. Bonn^{4,5}, W. N. Hardy^{4,5}, R. Liang^{4,5}, N. Doiron-Leyraud¹, Louis Taillefer^{1,5} & Cyril Proust^{2,5}

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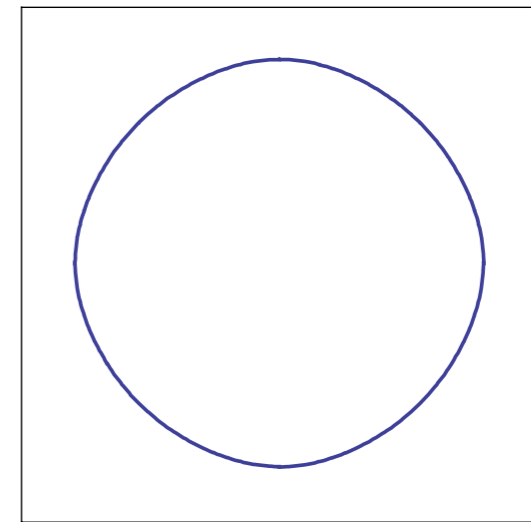
SDW SRO

Reconstructed Fermi surface
with non-Luttinger volume.

Z_2 vortices or hedgehogs expelled.

Z_2 or $U(1)$ topological order.

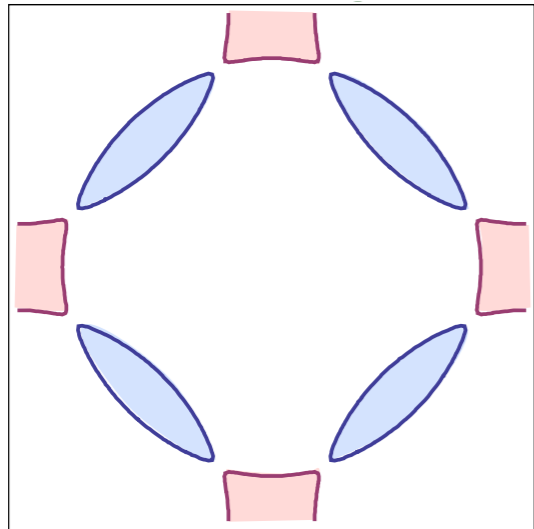
(B)
Topological
phase
transition;
phases of a
theory
with an emergent
 $SU(2)$ gauge field.



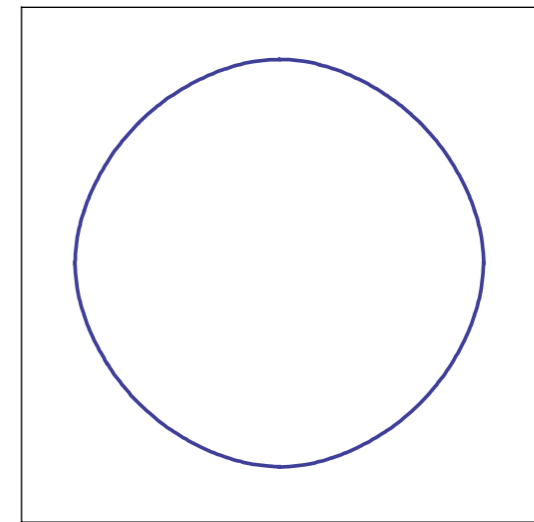
SDW SRO

Large Fermi surface
with Luttinger volume.
No topological order

Can model the doping dependence of the Hall
effect in the hole-doped cuprates



(B)
 Topological
 phase
 transition;
 phases of a
 theory
 with an emergent
 SU(2) gauge field.



SDW SRO

Reconstructed Fermi surface
 with non-Luttinger volume.

Z₂ vortices or hedgehogs expelled.

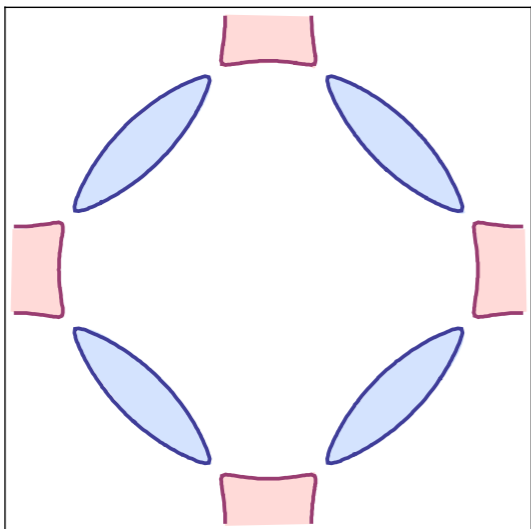
Z₂ or U(1) topological order.

SDW SRO

Large Fermi surface
 with Luttinger volume.
 No topological order



SU(2) gauge theory fits the real and imaginary parts of the electron Green's function computed by multi-site DMFT

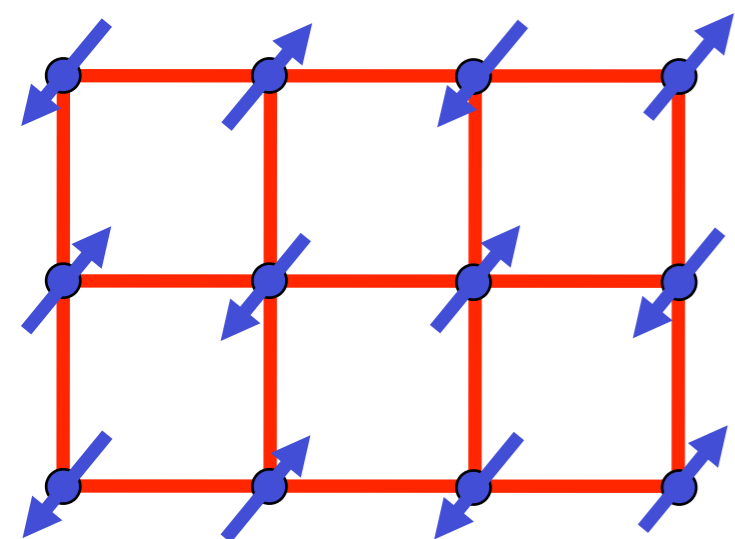


SDW SRO

Z_2 or $U(1)$ topological order.
Reconstructed Fermi surface.

$$\langle \vec{\Phi} \rangle = 0$$

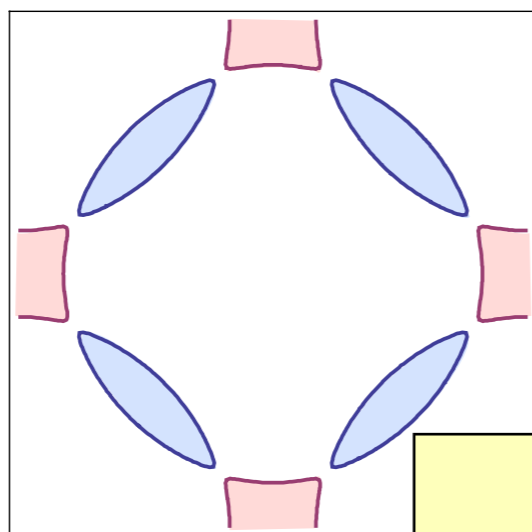
(B) Topological
phase transition



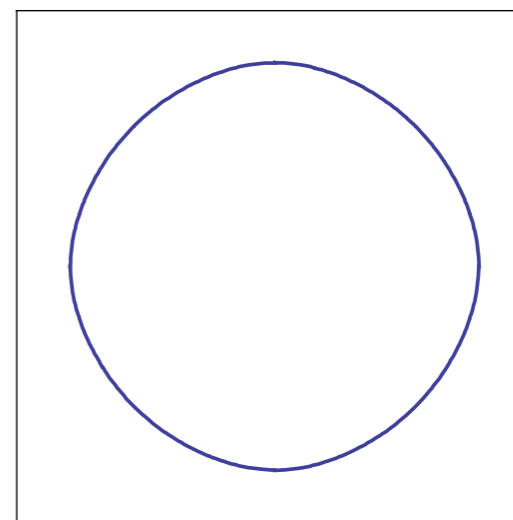
SDW LRO

Reconstructed Fermi surface

$$\langle \vec{\Phi} \rangle \neq 0$$



(A)
Symmetry
breaking
phase
transition



SDW SRO

Large Fermi surface.

$$\langle \vec{\Phi} \rangle = 0$$

g

U/t

SDW SRO

Z_2 or $U(1)$ topological order.
Reconstructed Fermi surface.

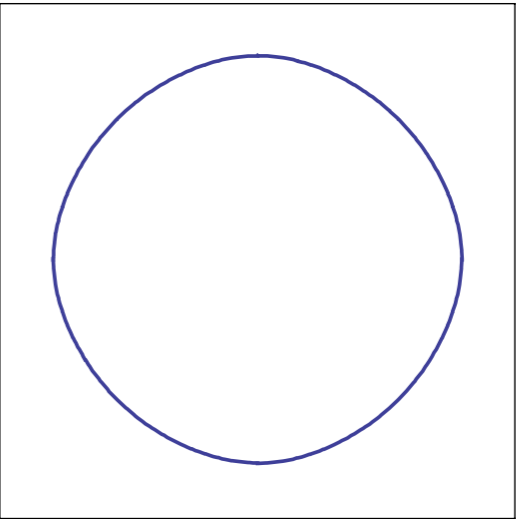
$$\langle \vec{\Phi} \rangle = 0$$

(C) Symmetry breaking and topological phase transition

(B) Topological phase transition

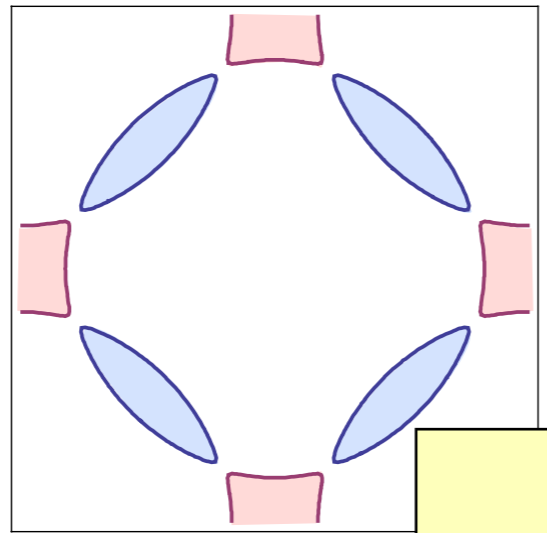
(A) Symmetry breaking phase transition

g



SDW SRO
Large Fermi surface.

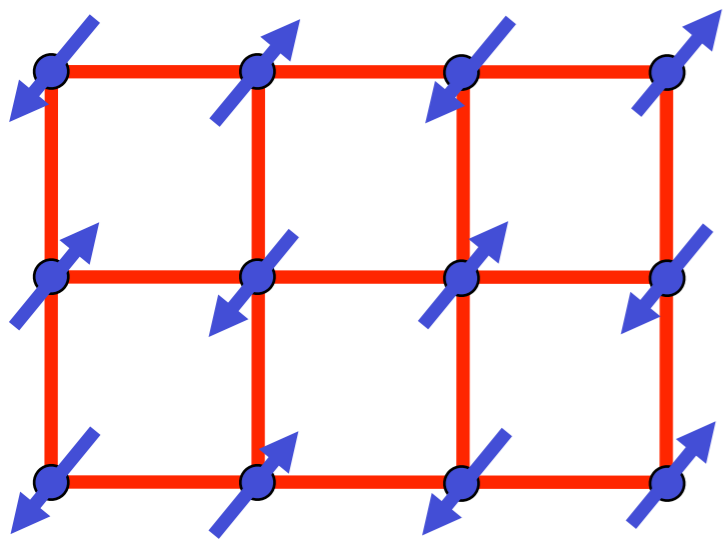
$$\langle \vec{\Phi} \rangle = 0$$



SDW LRO

Reconstructed Fermi surface

$$\langle \vec{\Phi} \rangle \neq 0$$



U/t



Continuous quantum transitions

(A) Broken symmetry | No broken symmetry

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Broken symmetry" and the right side contains "No broken symmetry".

(B) Topological order | No topological order

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Topological order" and the right side contains "No topological order".

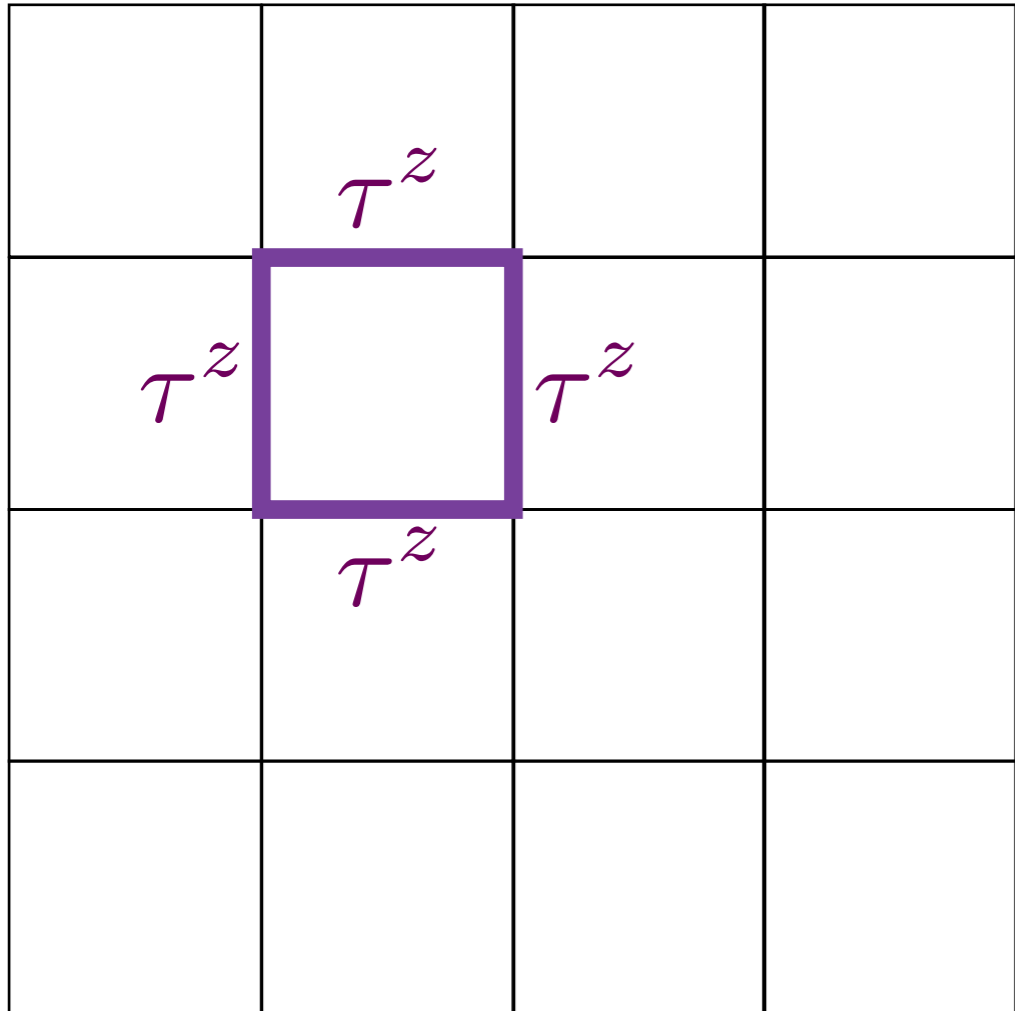
(C) Topological order | Broken symmetry

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Topological order" and the right side contains "Broken symmetry". The entire diagram is enclosed in a blue rounded rectangular border.

(D) Broken symmetry | A different broken symmetry

A horizontal red arrow pointing to the right, divided by a vertical red line. The left side contains the text "Broken symmetry" and the right side contains "A different broken symmetry".

Odd Z_2 lattice gauge theory



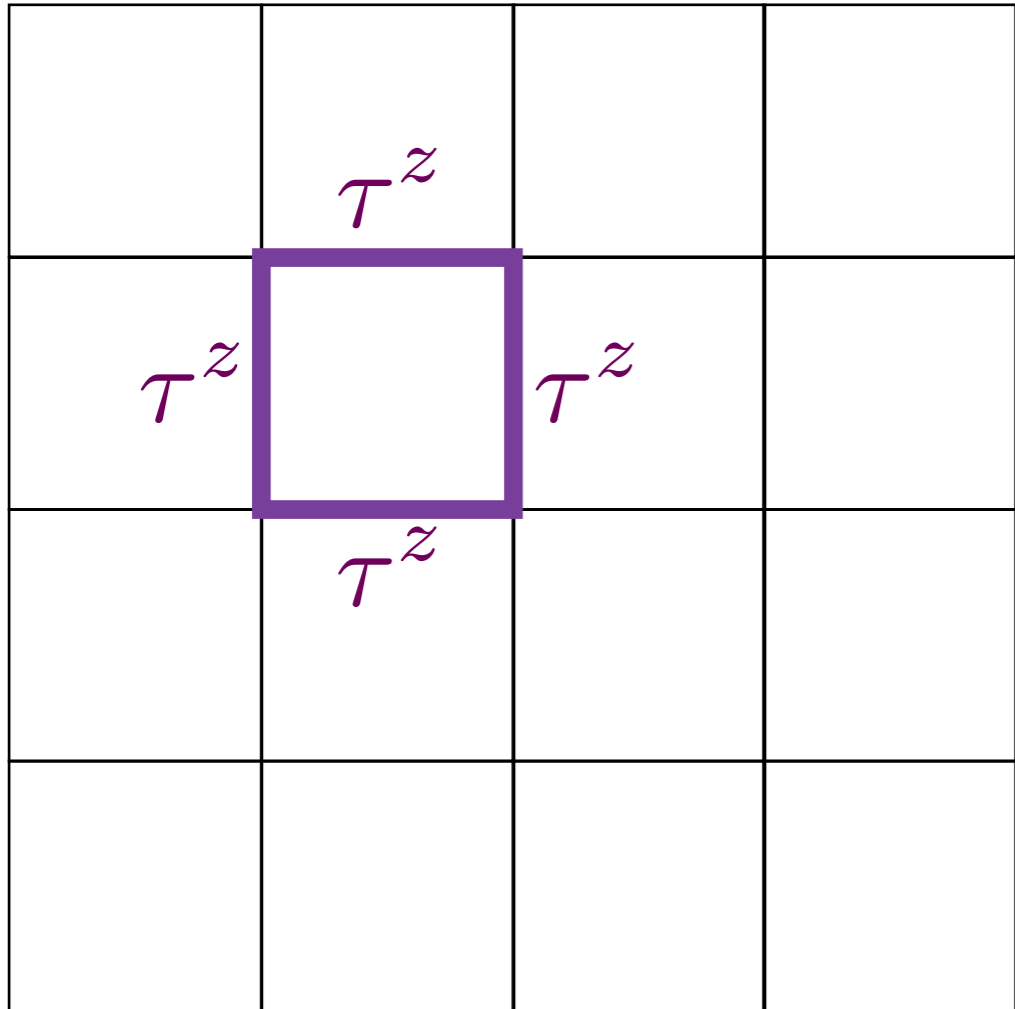
$$H = - \sum_{\square} \tau^z \tau^z \tau^z \tau^z - g \sum_i \tau^x$$

$$G_i = \begin{array}{c|c} & \tau^x \\ \hline \tau^x & \tau^x \\ \hline & \tau^x \end{array}$$

Gauss's Law with background electric charges:

$$[H, G_i] = 0 \quad , \quad G_i = -1$$

Odd Z_2 lattice gauge theory



$$H = - \sum_{\square} \tau^z \tau^z \tau^z \tau^z - g \sum_i \tau^x$$

$$G_i = \begin{array}{c|c} & \tau^x \\ \hline \tau^x & \tau^x \\ \hline & \tau^x \end{array}$$

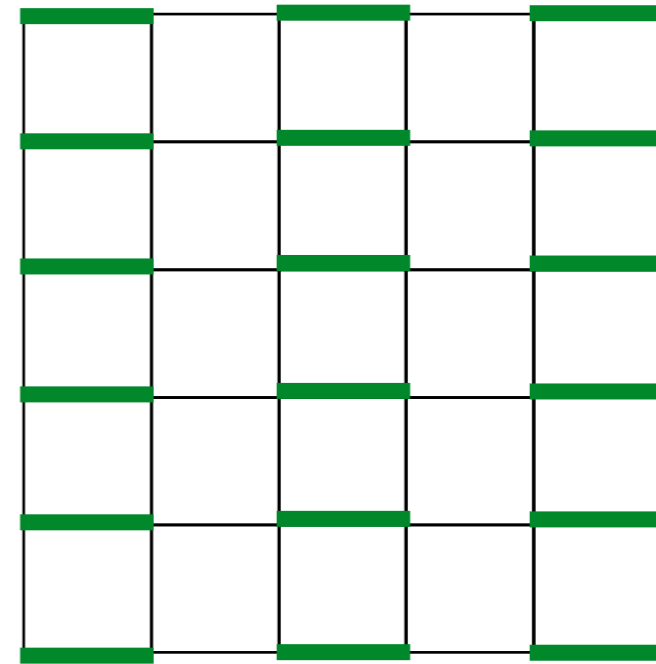
Gauss's Law with background electric charges:

$$[H, G_i] = 0 \quad , \quad G_i = -1$$

Odd Z_2 lattice gauge theory

Deconfined phase.
 Z_2 flux expelled.
 Z_2 (toric code)
topological order.

(C)
Symmetry-
breaking
and
topological
phase
transition



Confined phase.
 Z_2 flux proliferates.
No topological order.
Electric field lines lead
to symmetry breaking and
valence bond solid (VBS) order

$$H = - \sum_{\square} \tau^z \tau^z \tau^z \tau^z - g \sum_i \tau^x \quad G_i = -1 \quad g$$

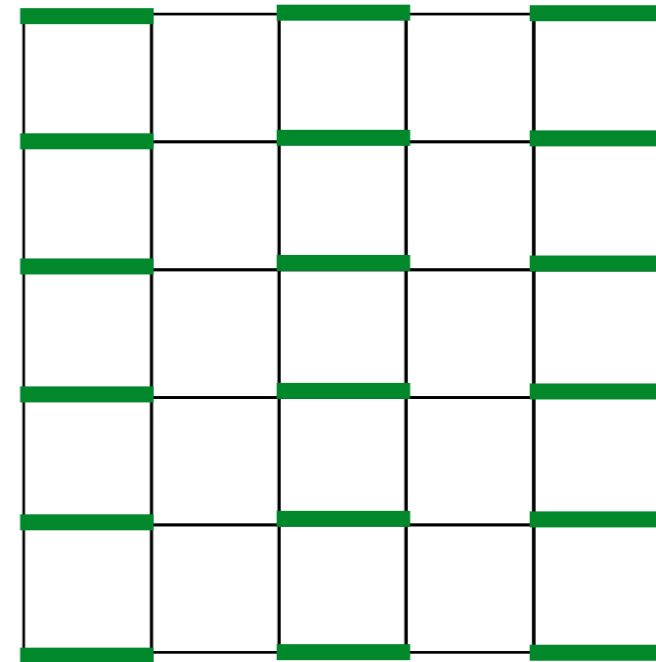
R. Jalabert and S. Sachdev, PRB **44**, 686 (1991)

T. Senthil, L. Balents, S. Sachdev, A. Vishwanath, and M. P.A. Fisher, PRB **70**, 144407 (2004)

Odd Z_2 lattice gauge theory

Deconfined phase.
 Z_2 flux expelled.
 Z_2 (toric code)
topological order.

(C)
Symmetry-
breaking
and
topological
phase
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phases of a
theory with an
emergent
U(1) gauge field



Confined phase.
 Z_2 flux proliferates.
No topological order.
Electric field lines lead
to symmetry breaking and
valence bond solid (VBS) order

$$H = - \sum_{\square} \tau^z \tau^z \tau^z \tau^z - g \sum_i \tau^x \quad G_i = -1 \quad g$$

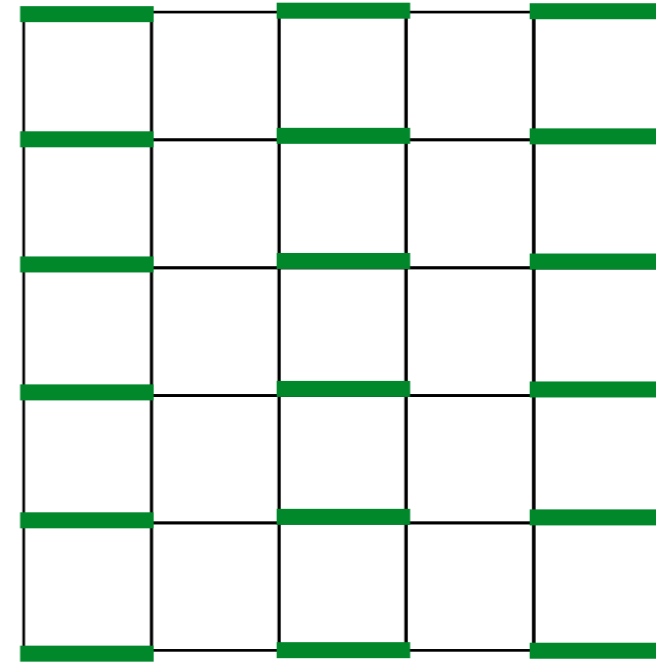
R. Jalabert and S. Sachdev, PRB **44**, 686 (1991)

T. Senthil, L. Balents, S. Sachdev, A. Vishwanath, and M. P.A. Fisher, PRB **70**, 144407 (2004)

Odd Z_2 lattice gauge theory

Deconfined phase.
 Z_2 flux expelled.
 Z_2 (toric code)
topological order.

(C)
Symmetry-
breaking
and
topological
phase
transition;
phases of a
theory with an
emergent
U(1) gauge field



Confined phase.
 Z_2 flux proliferates.
No topological order.
Electric field lines lead
to symmetry breaking and
valence bond solid (VBS) order

Similar phases and transitions in frustrated square lattice
antiferromagnets with spin $S=1/2$

g

R. Jalabert and S. Sachdev, PRB **44**, 686 (1991)

N. Read and S. Sachdev, PRL **66**, 1773 (1991)

Continuous quantum transitions

(A) Broken symmetry | No broken symmetry

A horizontal red arrow with a vertical red line in the middle. The left side of the arrow is labeled "Broken symmetry" and the right side is labeled "No broken symmetry".

(B) Topological order | No topological order

A horizontal red arrow with a vertical red line in the middle. The left side of the arrow is labeled "Topological order" and the right side is labeled "No topological order".

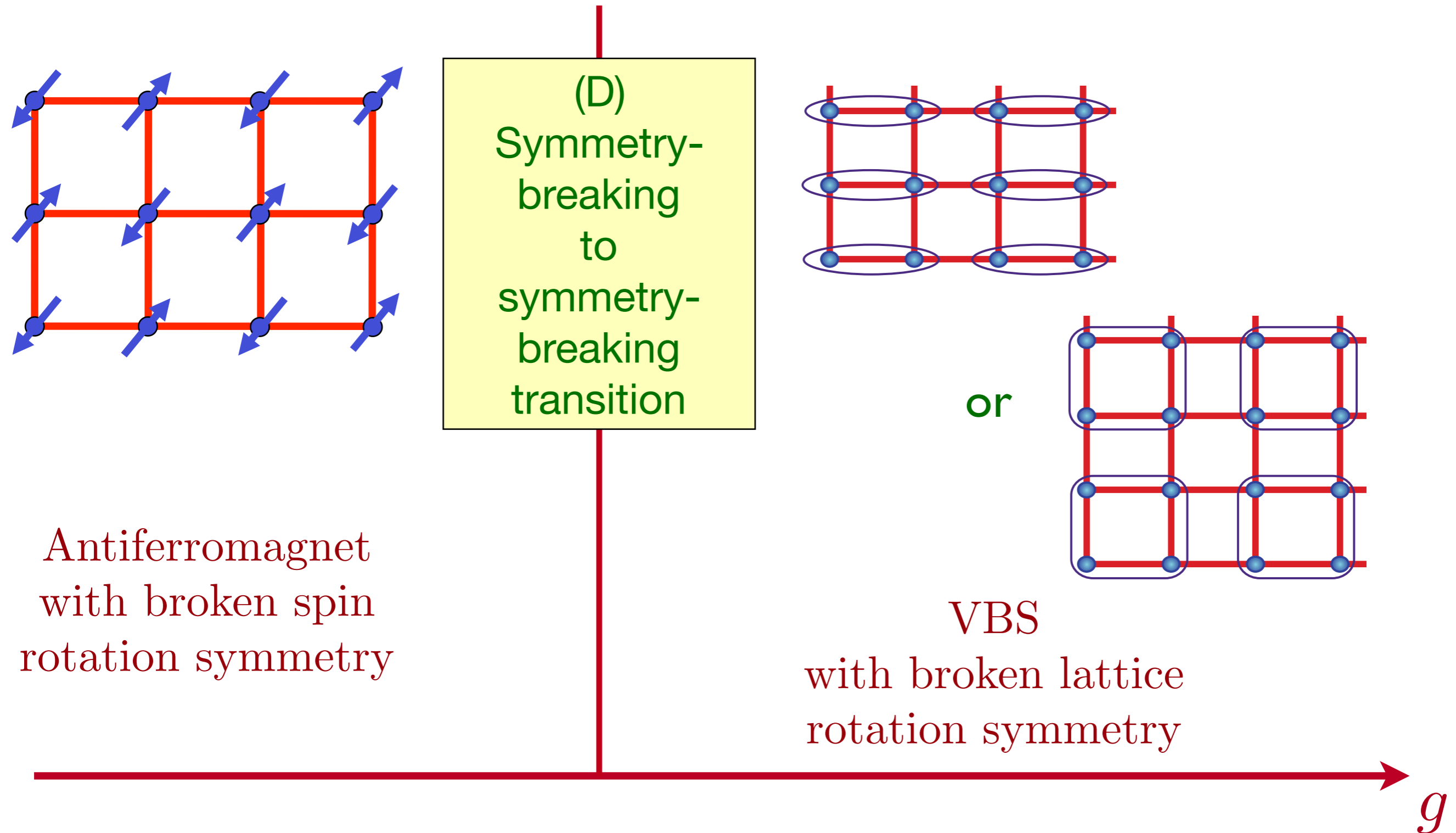
(C) Topological order | Broken symmetry

A horizontal red arrow with a vertical red line in the middle. The left side of the arrow is labeled "Topological order" and the right side is labeled "Broken symmetry".

(D) Broken symmetry | A different broken symmetry

A horizontal red arrow with a vertical red line in the middle. The left side of the arrow is labeled "Broken symmetry" and the right side is labeled "A different broken symmetry". The entire diagram is enclosed in a blue rounded rectangular border.

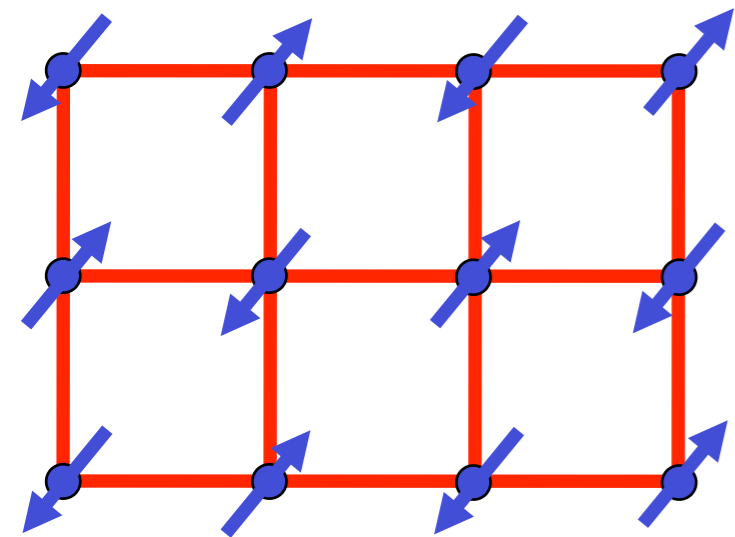
Quantum phases of a $S=1/2$ square lattice antiferromagnet



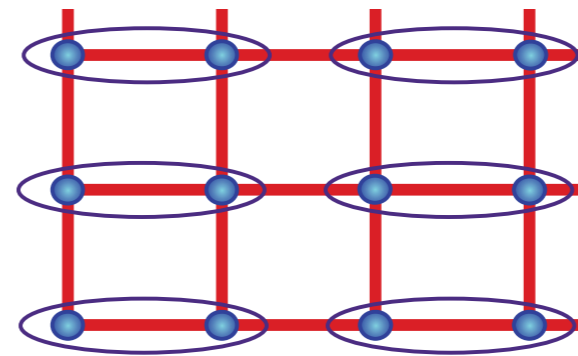
N. Read and S. Sachdev, PRL **62**, 1694 (1989)

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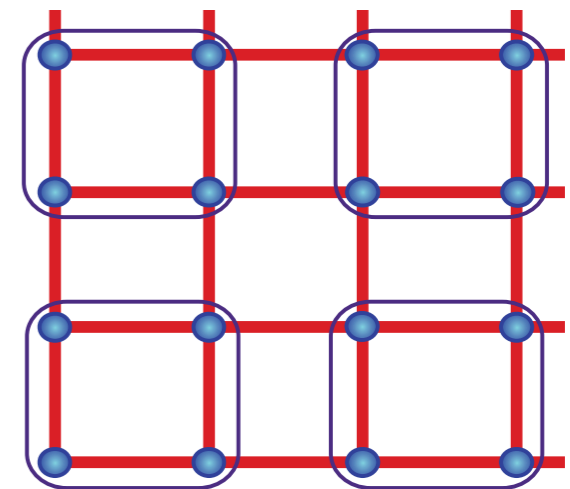
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(D)
Symmetry-breaking to symmetry-breaking transition; phases of a theory with an emergent U(1) gauge field



or



VBS

with broken lattice rotation symmetry

Antiferromagnet with broken spin rotation symmetry

Phases described by Higgs/confining phases of a U(1) gauge theory, which is deconfined only at criticality

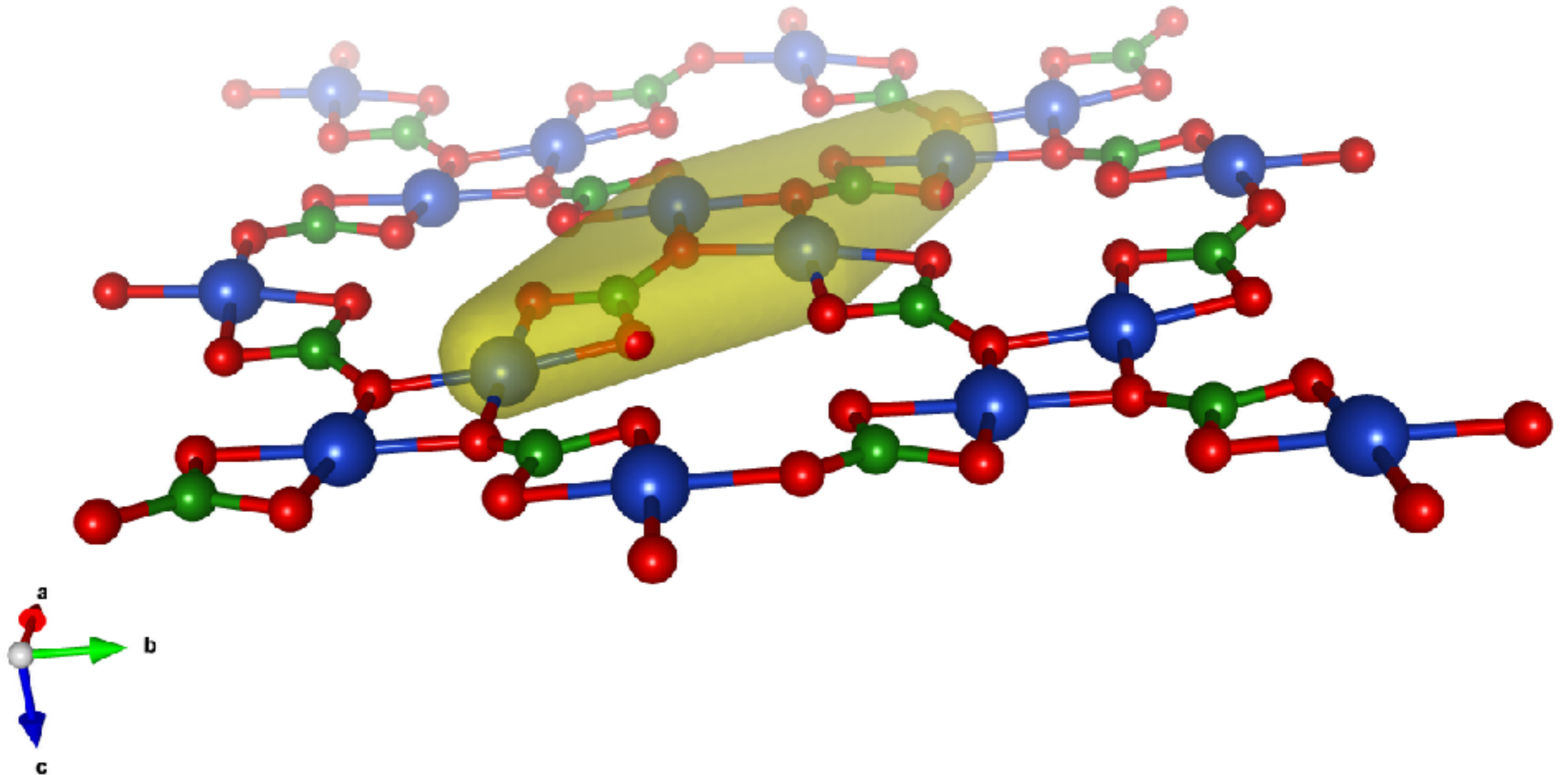
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4-spin plaquette singlet state in the Shastry–Sutherland compound $\text{SrCu}_2(\text{BO}_3)_2$

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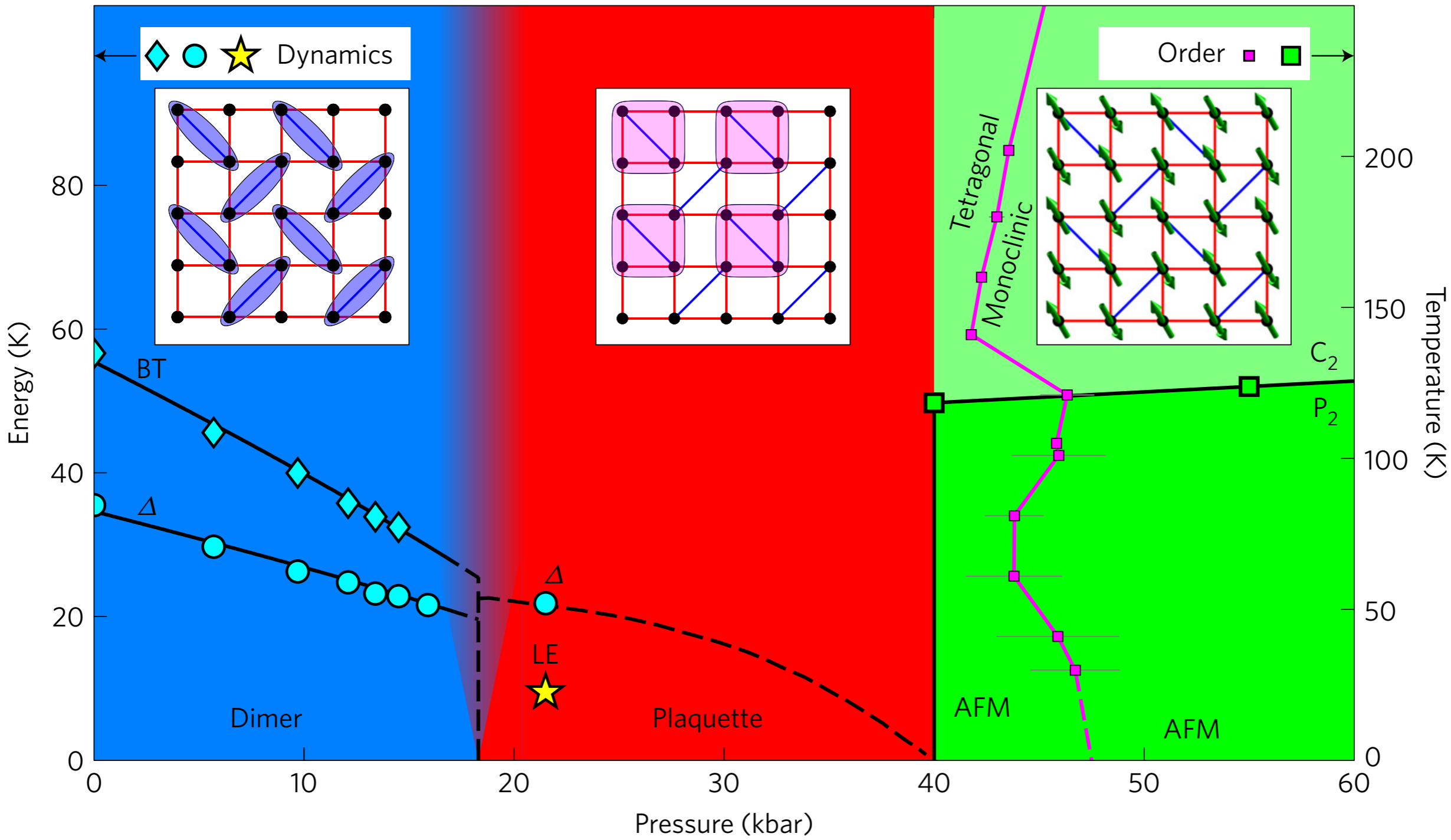
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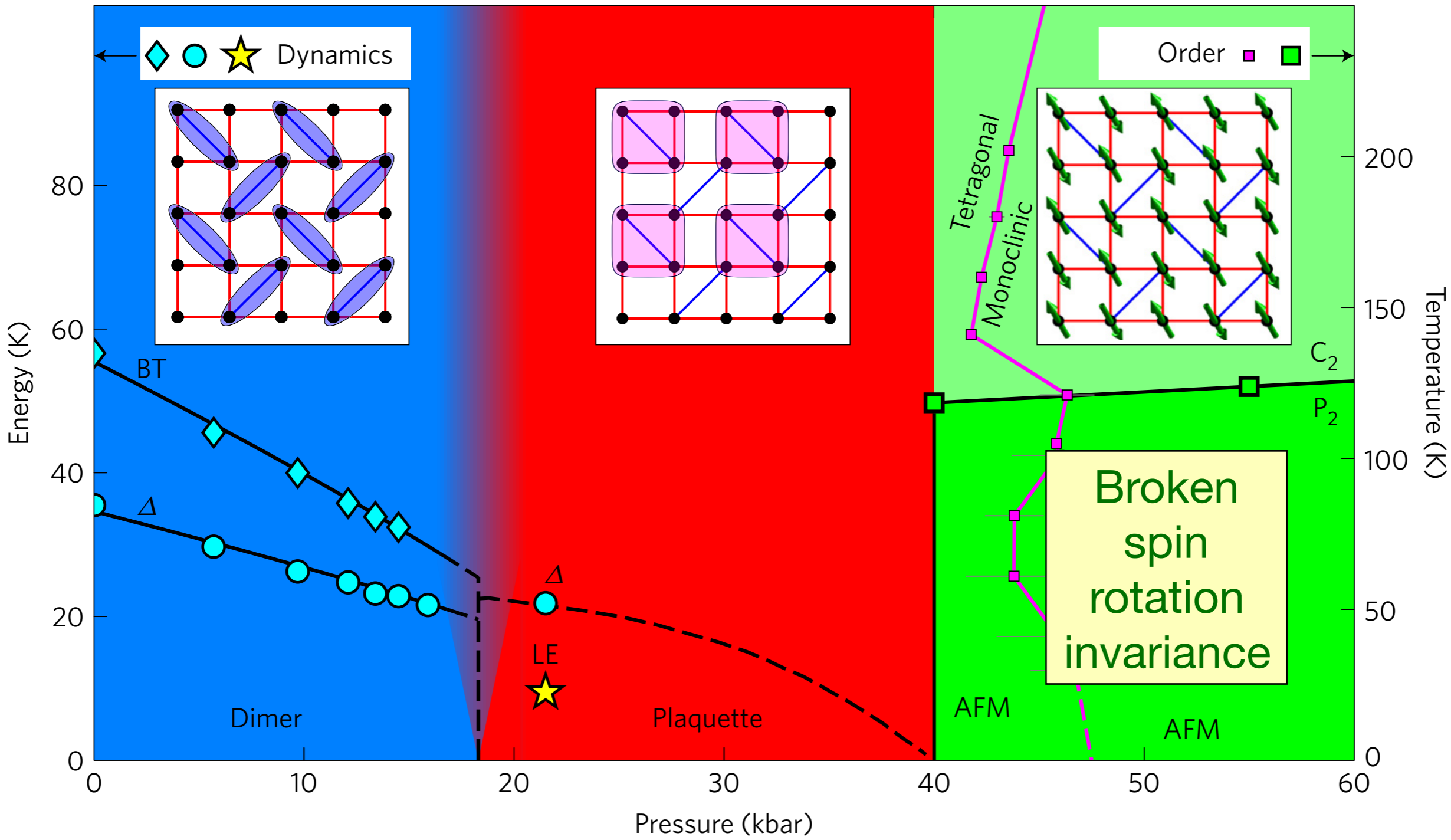
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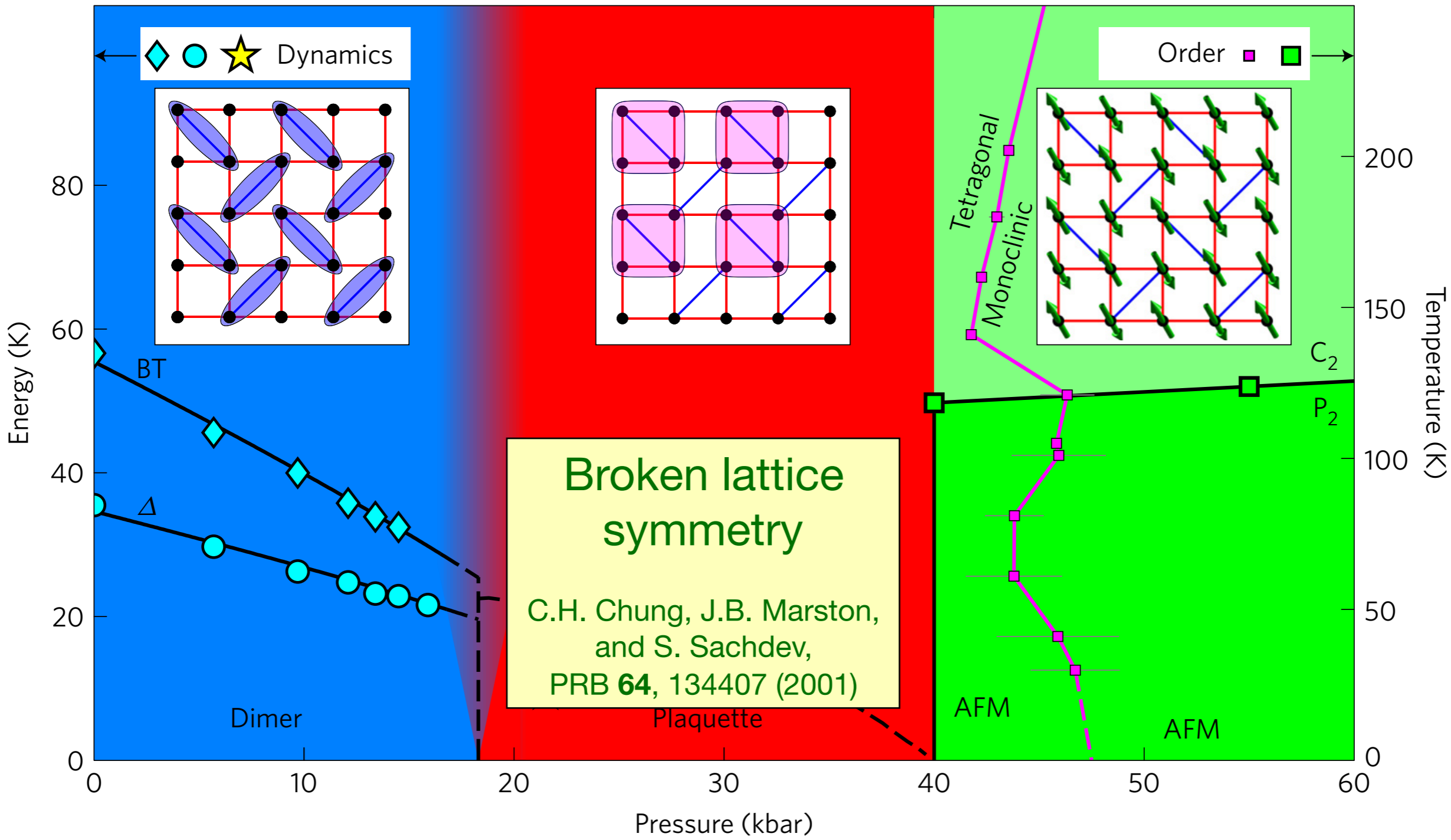
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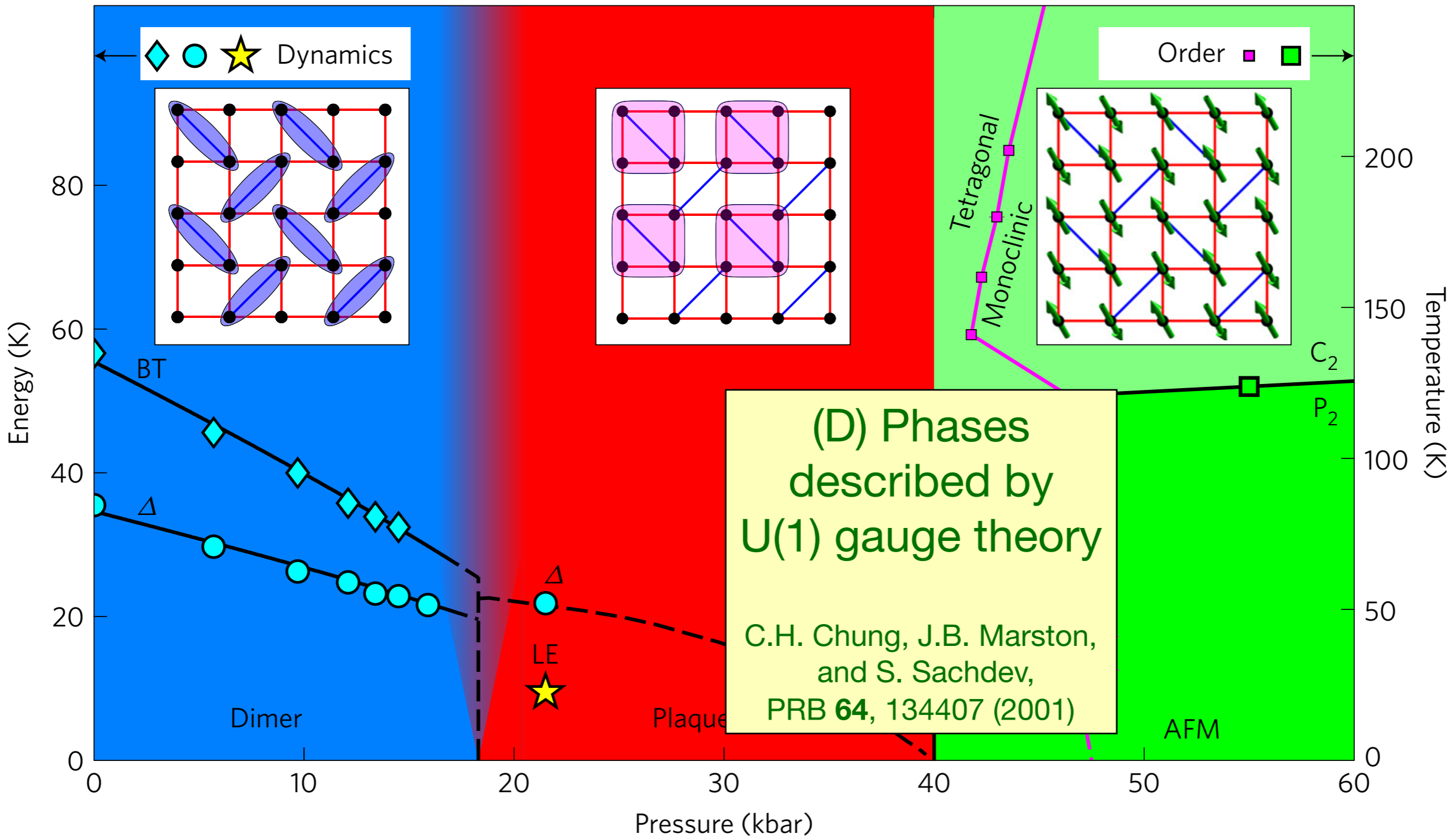
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Strongly-coupled quantum criticality

States of quantum matter with:

- No quasiparticle excitations.
- Strong interactions are at a universal RG fixed point, and this leads to the fastest possible ‘dephasing’ and ‘local thermalization’ in a time of order $\hbar/(k_B T)$.

S.Sachdev, *Quantum Phase Transitions*, 1991
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Onsager’s Ising criticality, other critical states in one spatial dimension, and quantum impurity models, do not have these properties

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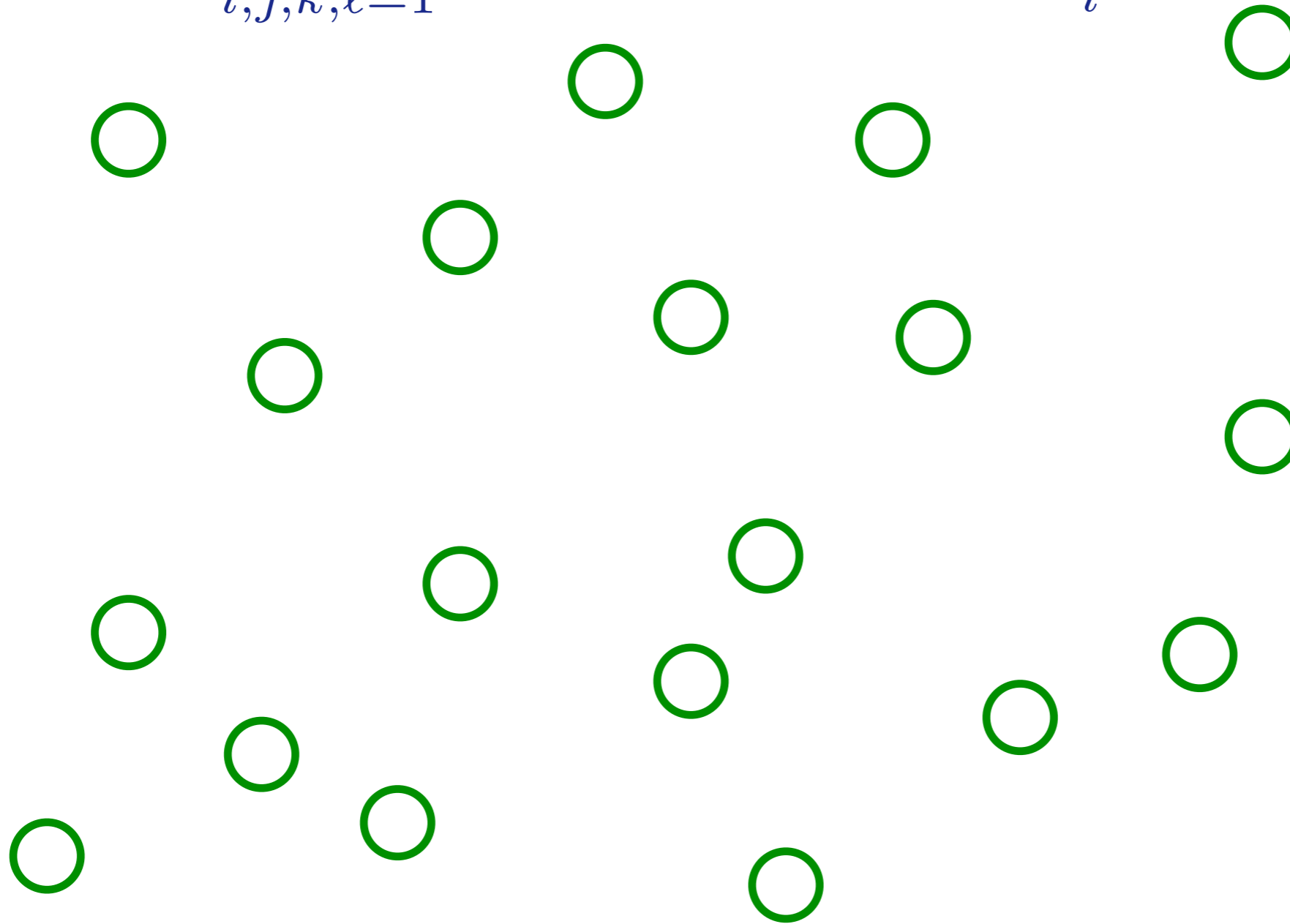
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‘Solvable’ models with these properties:
the SYK models

The Sachdev-Ye-Kitaev (SYK) model

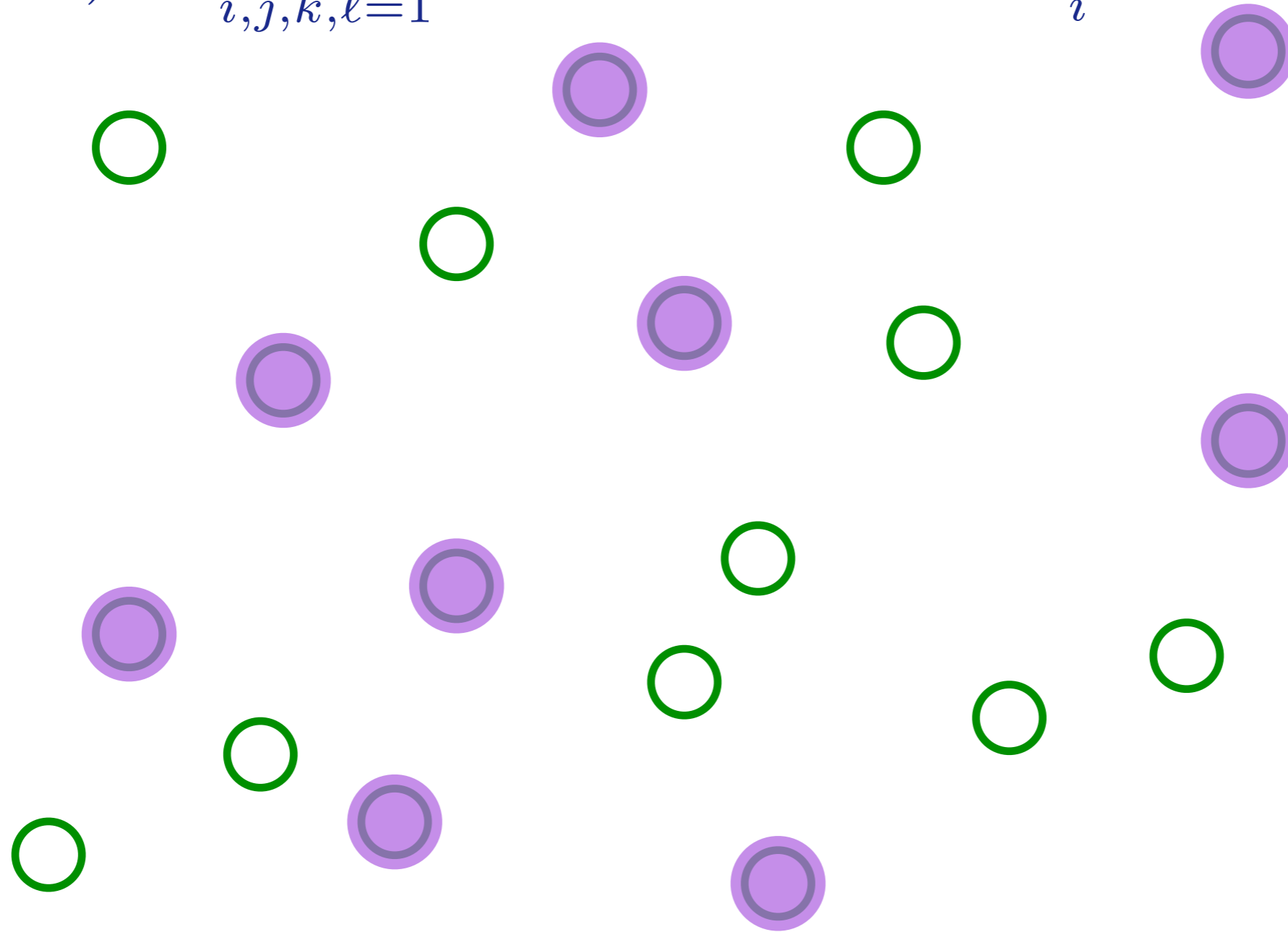
$$H = \frac{1}{(2N)^{3/2}} \sum_{i,j,k,l=1}^N U_{ij;kl} f_i^\dagger f_j^\dagger f_k f_l - \mu \sum_i f_i^\dagger f_i$$



Pick a set of random positions

The SYK model

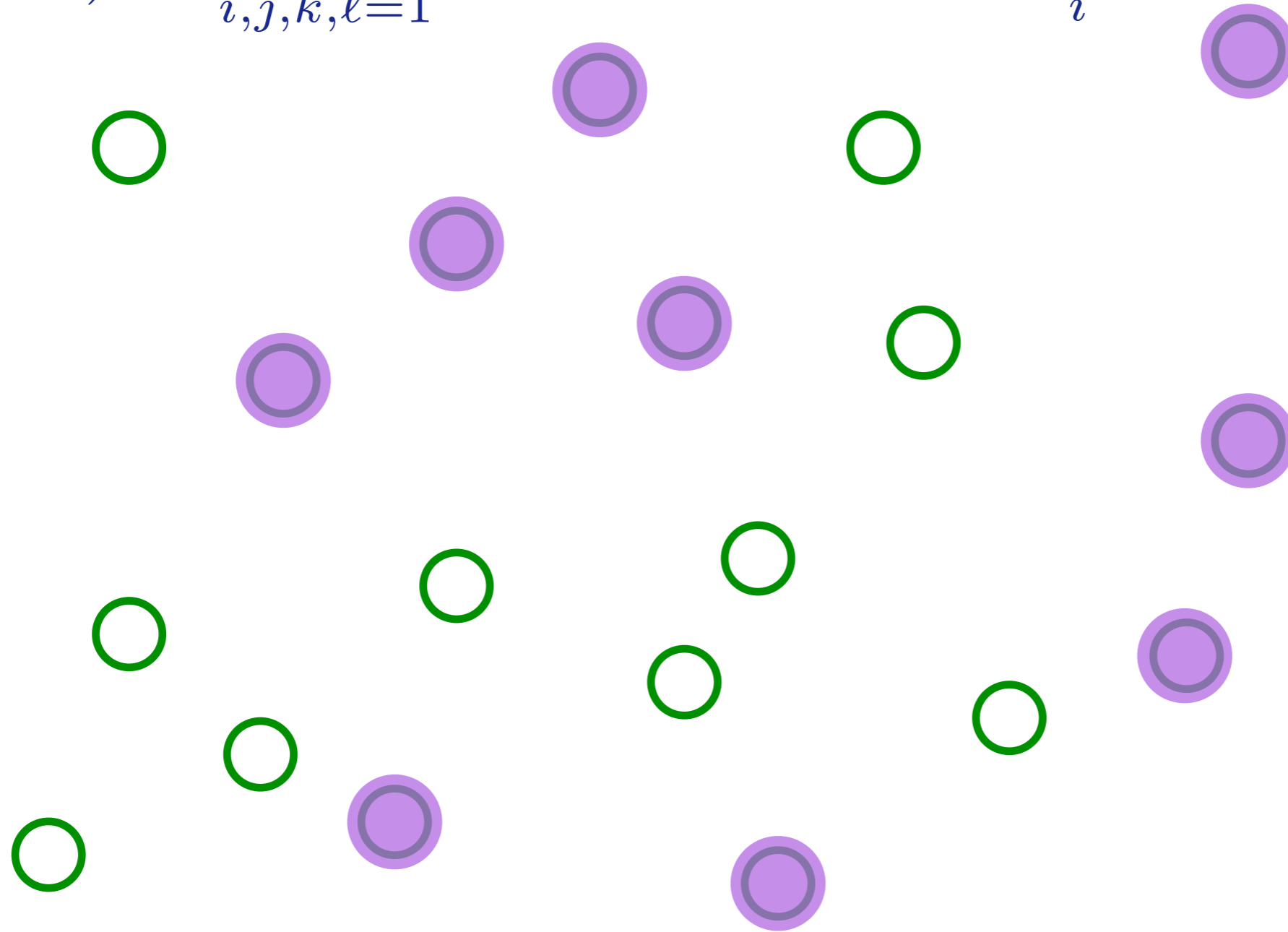
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Place electrons randomly on some sites

The SYK model

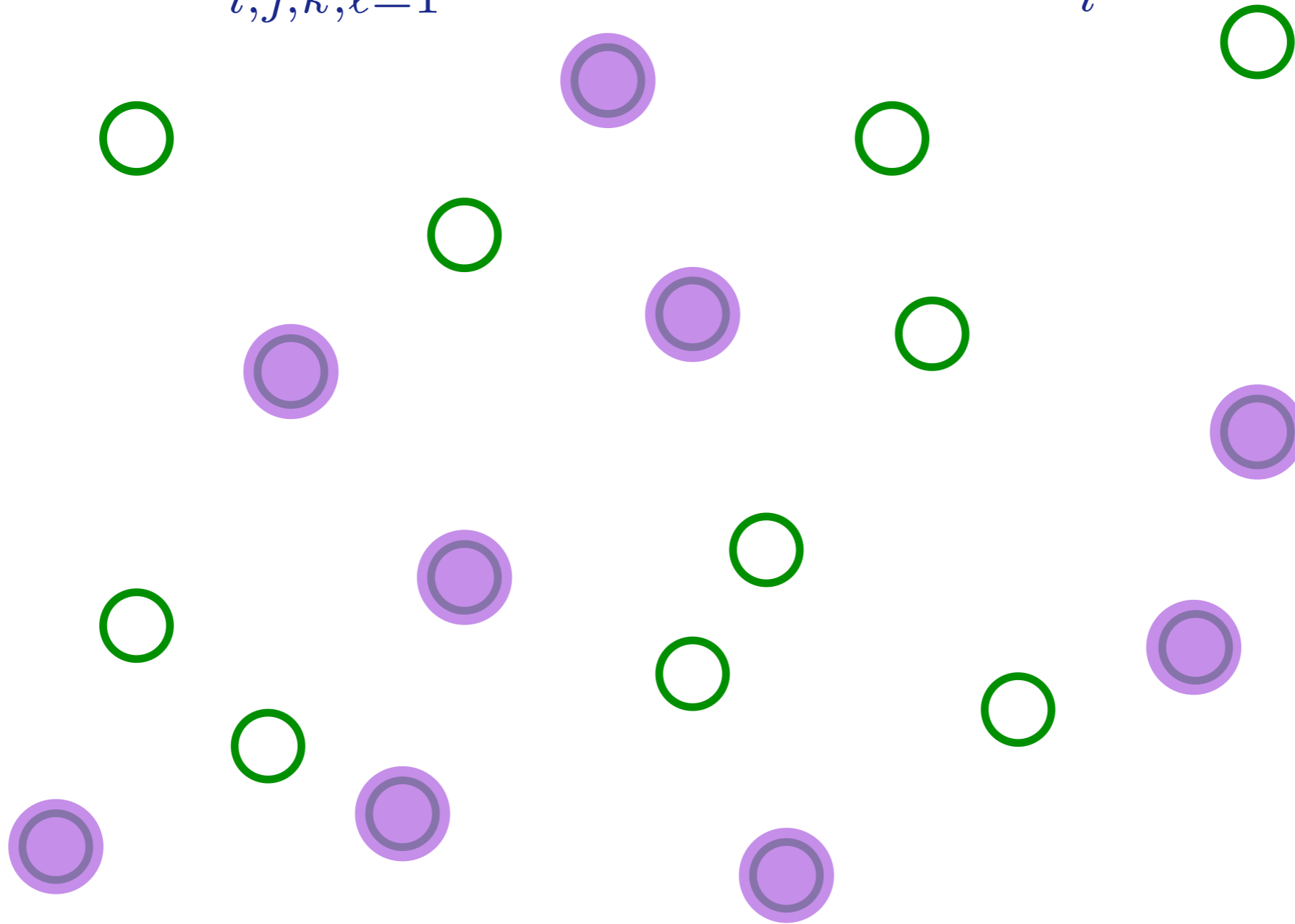
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Entangle electrons pairwise randomly

The SYK model

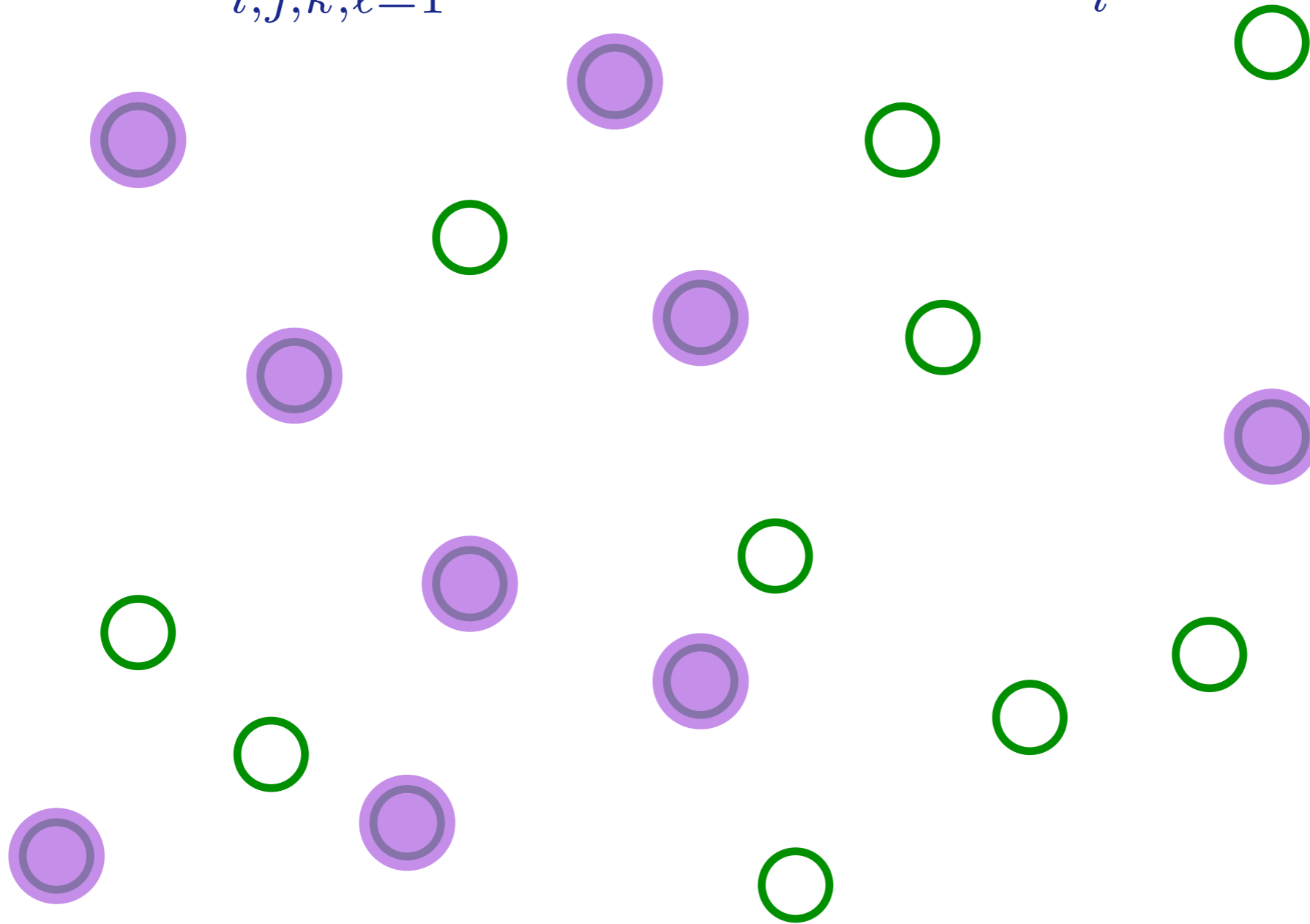
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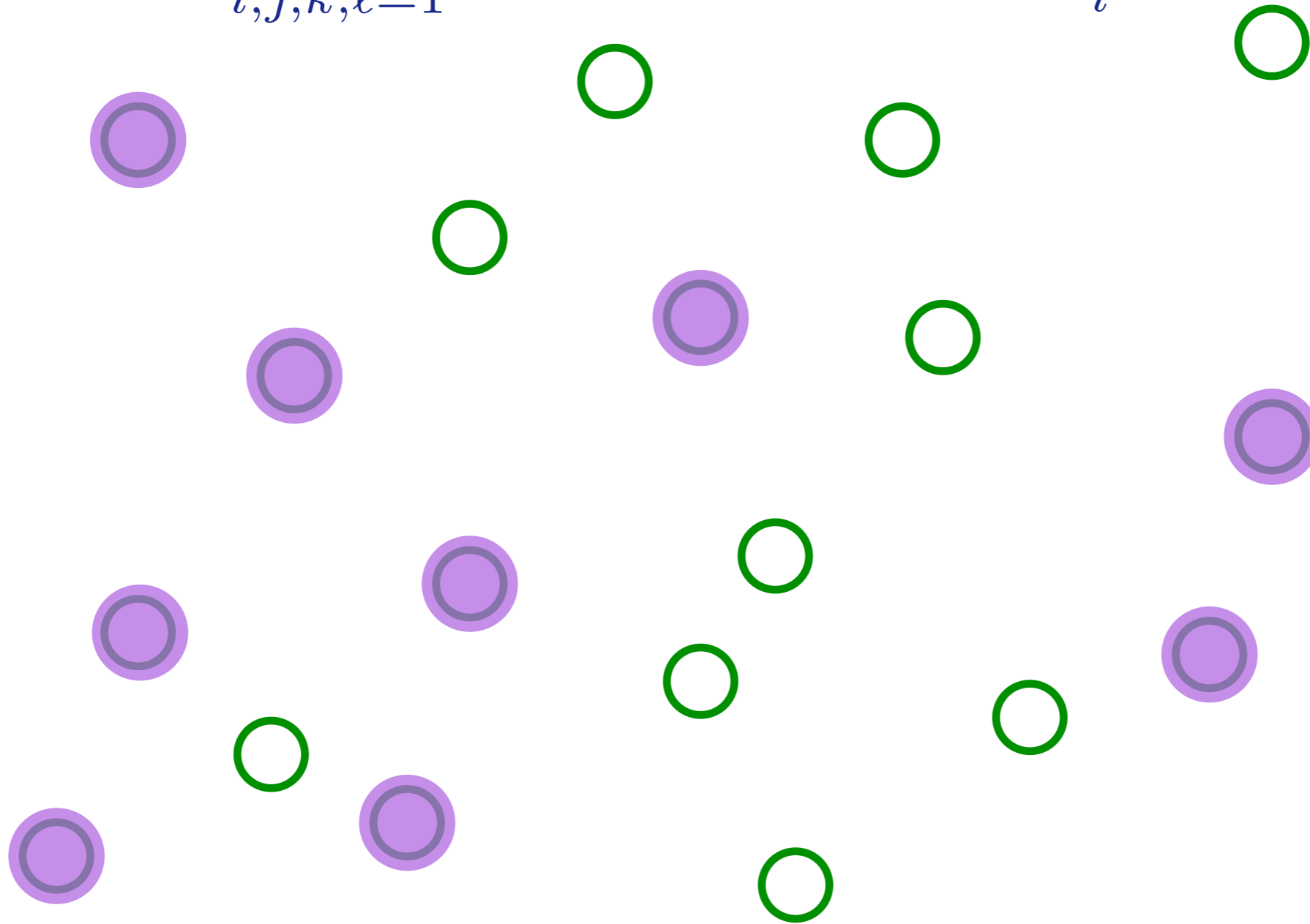
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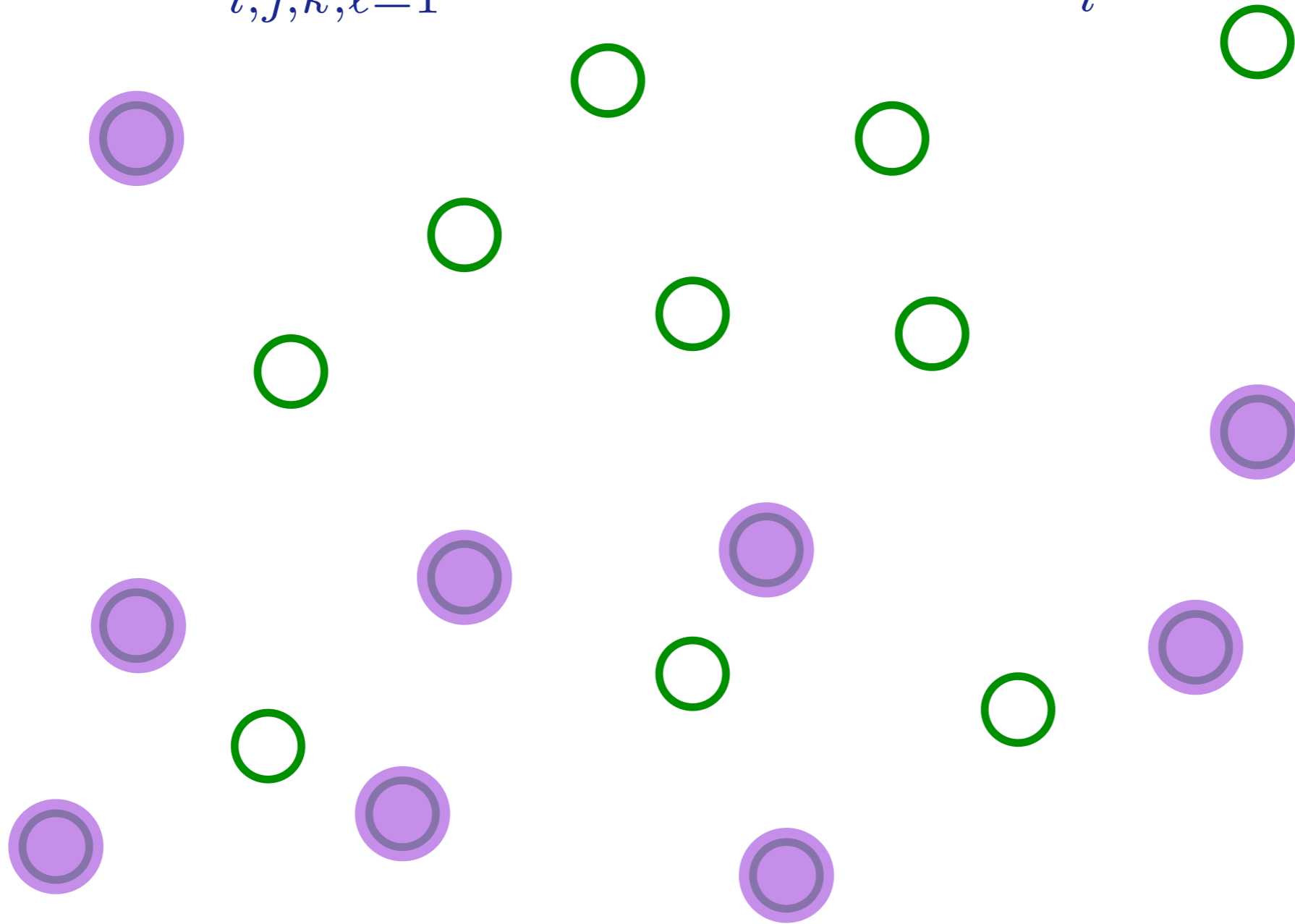
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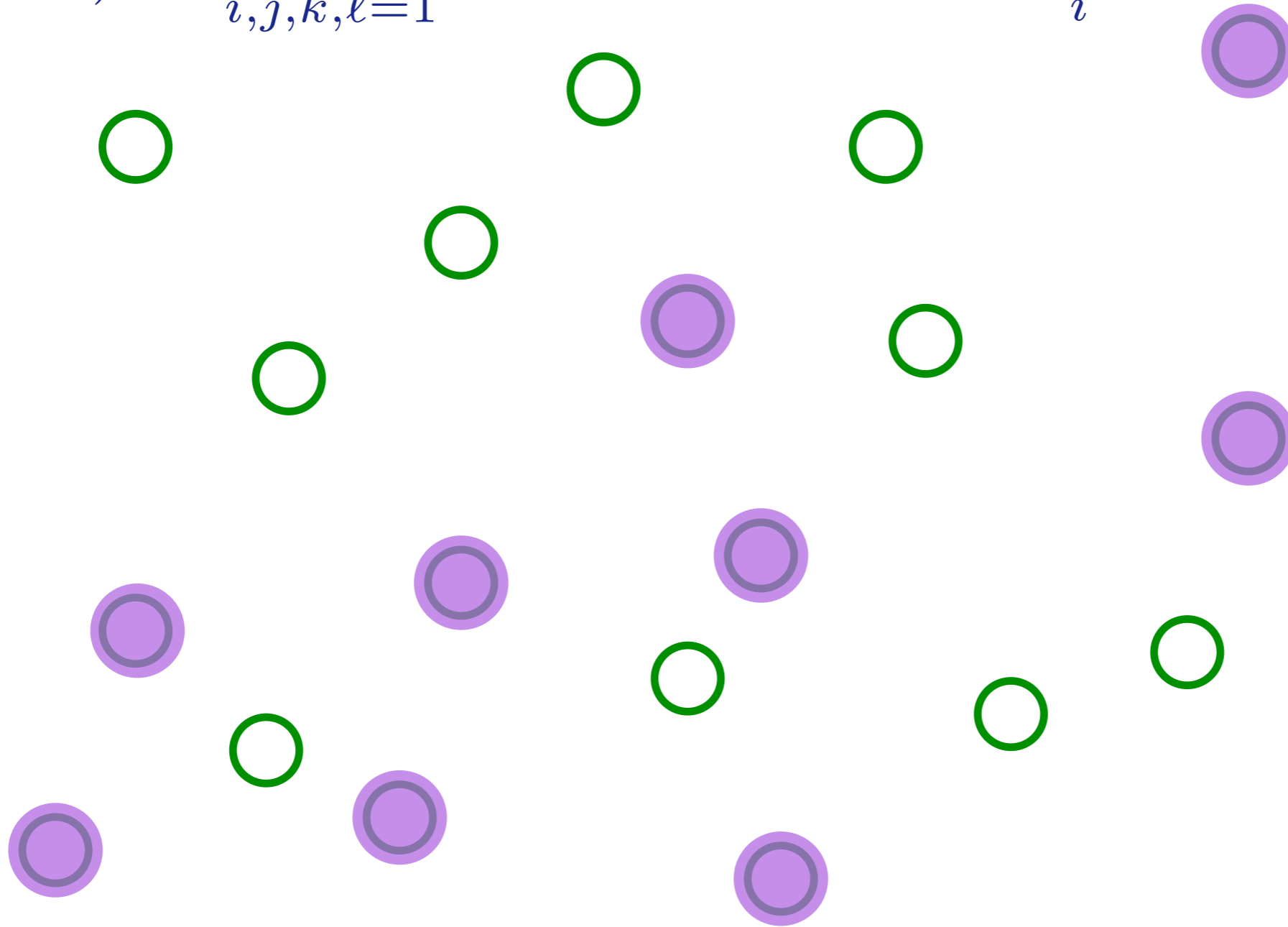
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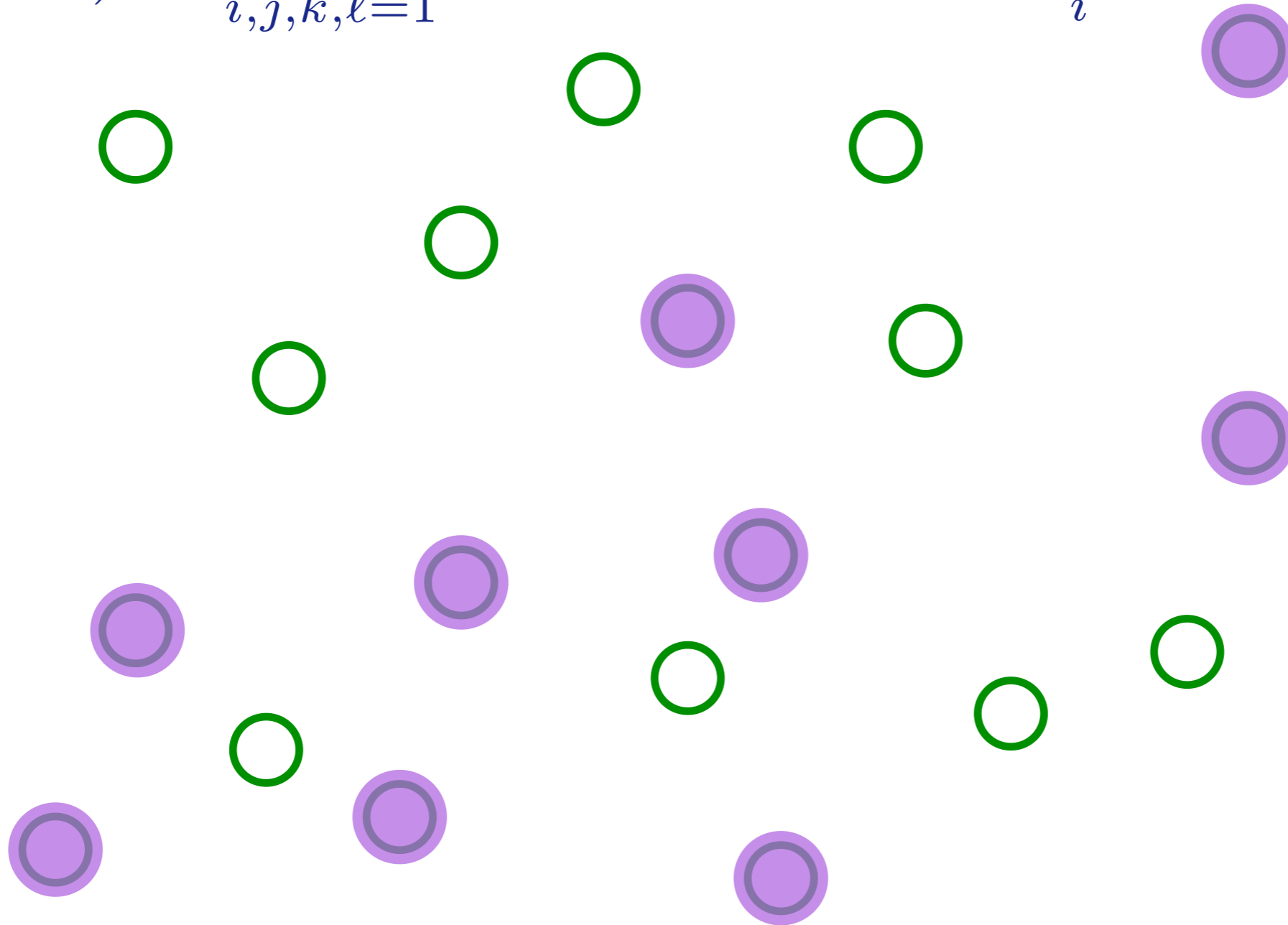
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This describes both a strange metal and a black hole!

The SYK model

There are 2^N many body levels with energy E , which do not admit a quasiparticle decomposition. Shown are all values of E for a single cluster of size $N = 12$. The $T \rightarrow 0$ state has an entropy $S_{GPS} = N s_0$ with

$$s_0 = \frac{G}{\pi} + \frac{\ln(2)}{4} = 0.464848\dots$$
$$< \ln 2$$

where G is Catalan's constant, for the half-filled case $Q = 1/2$.

GPS: A. Georges, O. Parcollet, and S. Sachdev, PRB **63**, 134406 (2001)

Many-body level spacing $\sim 2^{-N} = e^{-N \ln 2}$

Non-quasiparticle excitations with spacing $\sim e^{-N s_0}$

The SYK model

- Low energy, many-body density of states

$$\rho(E) \sim e^{Ns_0} \sinh(\sqrt{2(E - E_0)N\gamma})$$

D. Stanford and E. Witten, 1703.04612

A. M. Garcia-Garcia, J.J.M. Verbaarschot, 1701.06593

D. Bagrets, A. Altland, and A. Kamenev, 1607.00694

- Low temperature entropy $S = Ns_0 + N\gamma T + \dots$

- $T = 0$ fermion Green's function $G(\tau) \sim \tau^{-1/2}$ at large τ . (Fermi liquids with quasiparticles have $G(\tau) \sim 1/\tau$)

S. Sachdev and J. Ye, Phys. Rev. Lett. **70**, 3339 (1993)

- $T > 0$ Green's function has conformal invariance

$$G \sim (T / \sin(\pi k_B T \tau / \hbar))^{1/2}$$

A. Georges and O. Parcollet PRB **59**, 5341 (1999)

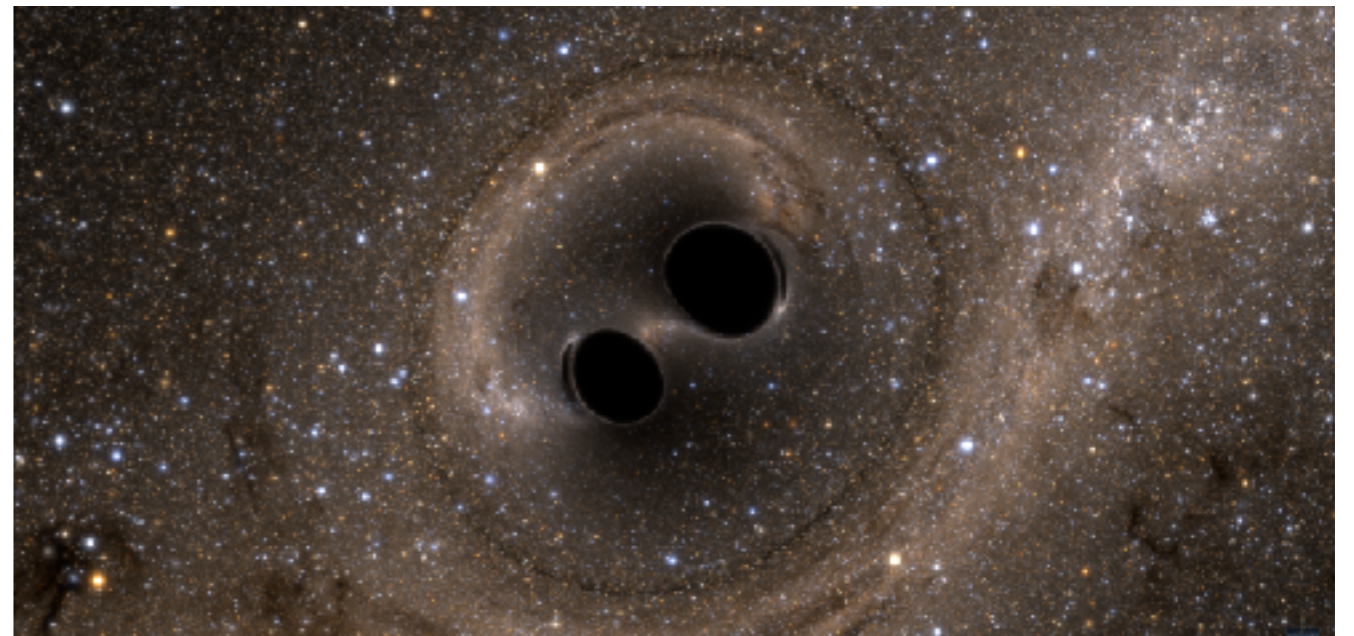
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A. Eberlein, V. Kasper, S. Sachdev, and J. Steinberg, arXiv:1706.07803

Black holes share the properties of strongly-coupled quantum criticality

- Black holes have an entropy and a temperature, $T_H = \hbar c^3 / (8\pi G M k_B)$.
- Black holes relax to thermal equilibrium in a time $\sim \hbar / (k_B T_H) = 8\pi G M / c^3$.

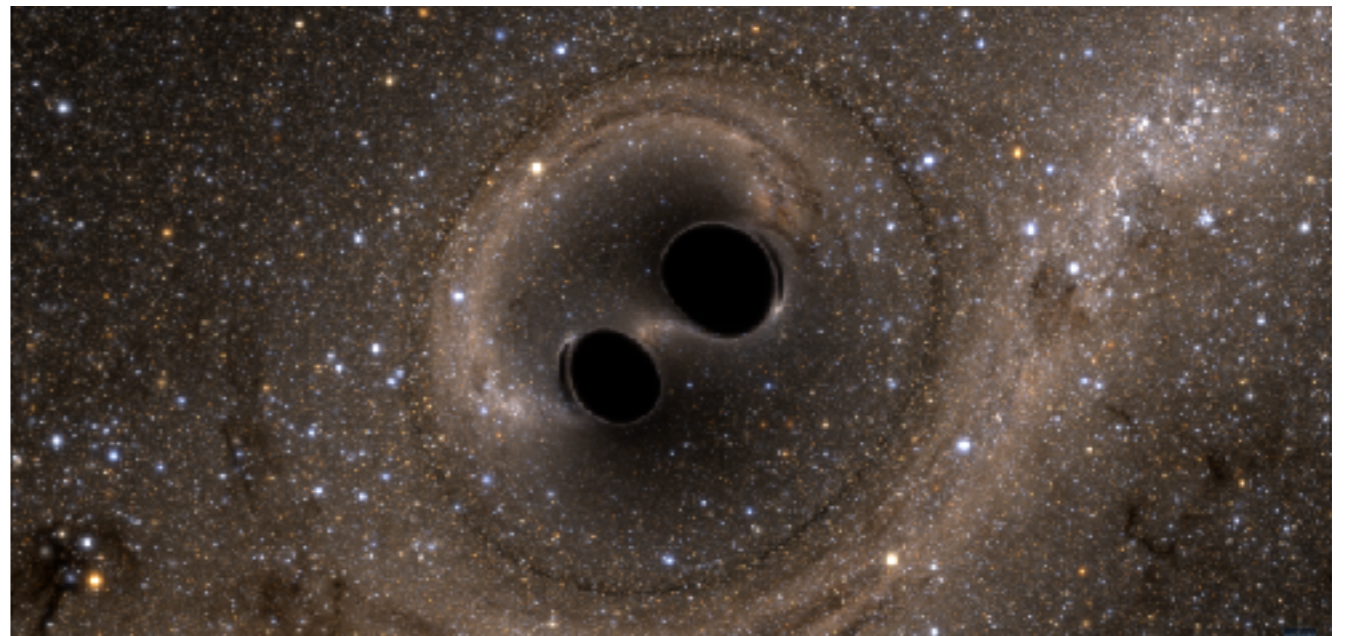
**Black
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- Black holes relax to thermal equilibrium in a time $\sim \hbar / (k_B T_H) = 8\pi G M / c^3$.
- The entropy of black holes is proportional to their surface area. So they can only be equivalent to quantum-critical systems in one lower dimension.

**Black
holes**



The SYK model

- Low energy, many-body density of states

$$\rho(E) \sim e^{Ns_0} \sinh(\sqrt{2(E - E_0)N\gamma})$$

D. Stanford and E. Witten, 1703.04612

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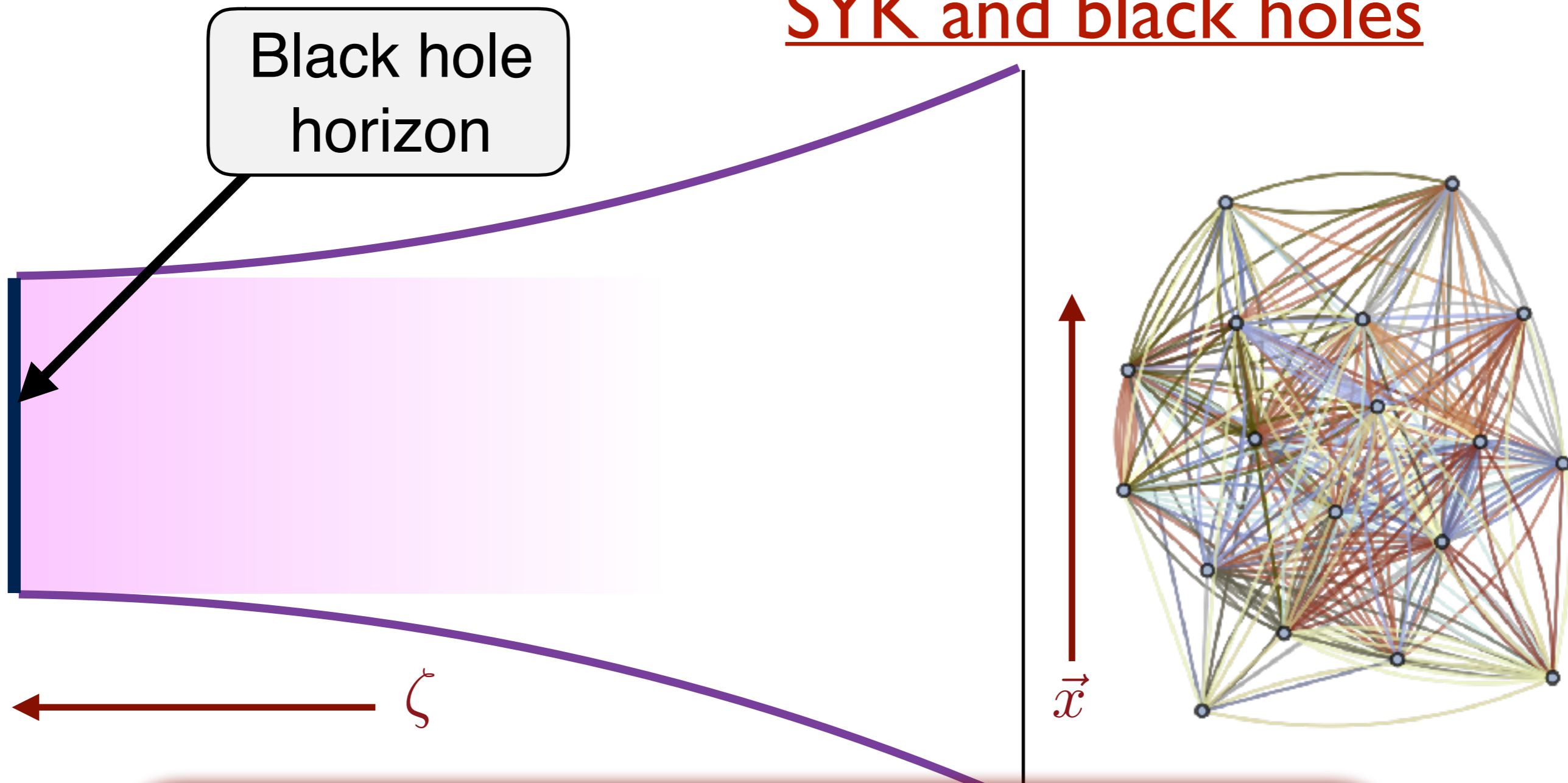
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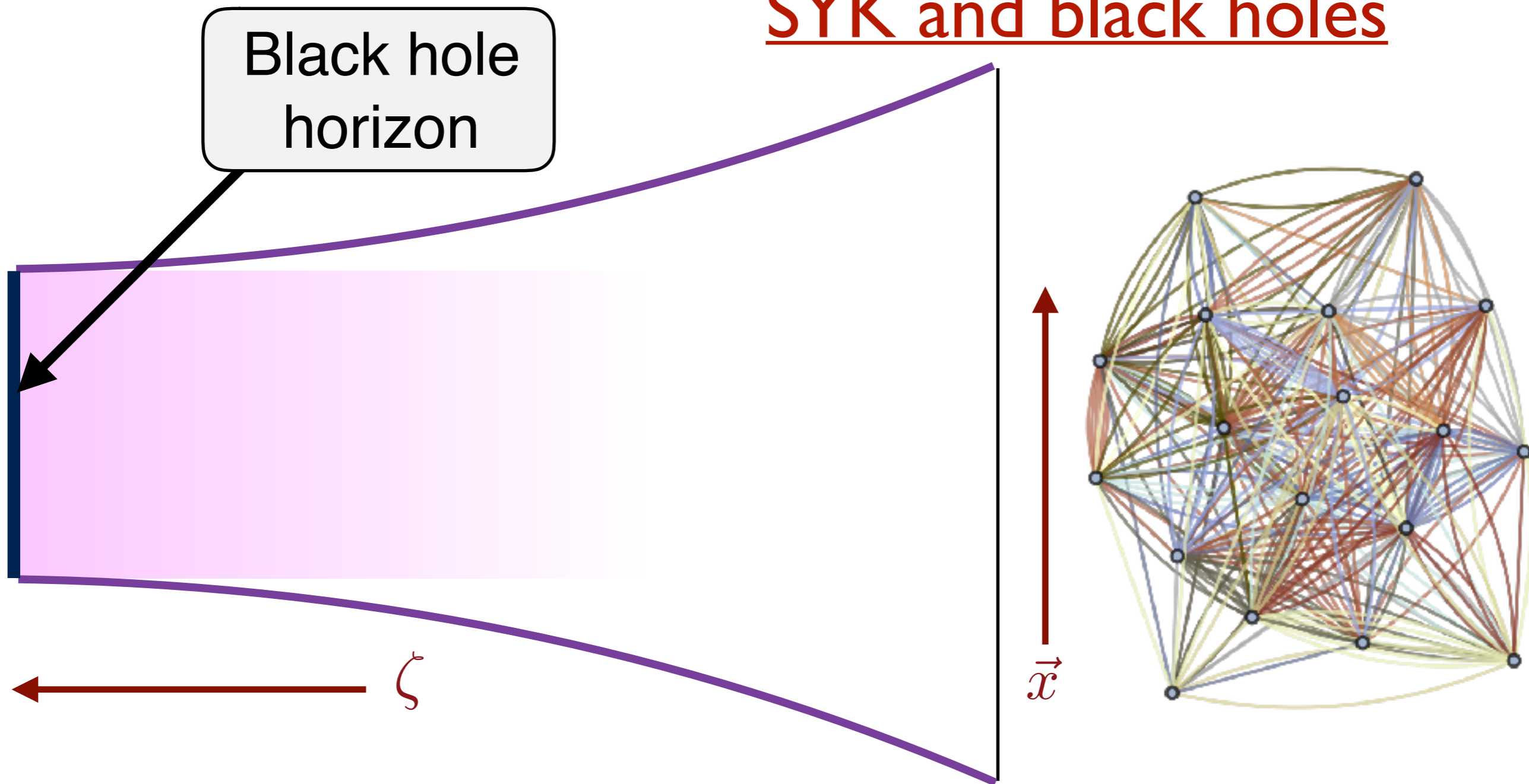
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SYK and black holes



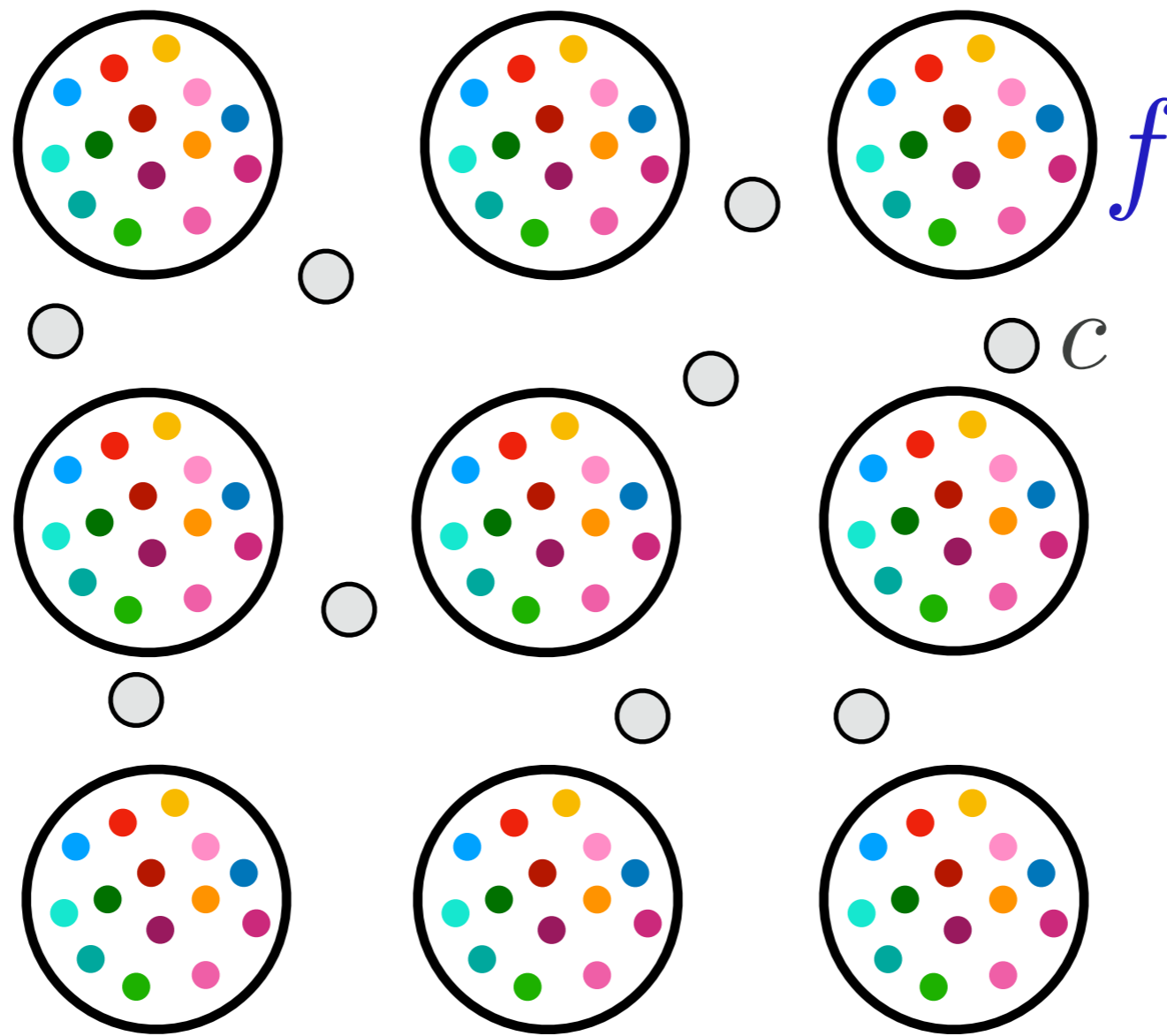
Black holes with a near-horizon AdS_2 geometry (described by quantum gravity in $1+1$ spacetime dimensions) match the properties of the $0+1$ dimensional SYK model in the previous slide: Ns_0 is the Bekenstein-Hawking entropy

SYK and black holes



Both the SYK model and the black hole saturate the lower bound on the Lyapunov time to quantum chaos: $\tau_L = \hbar / (2\pi k_B T)$.

SYK building blocks for strange metals



Yields solvable models with linear-in- T resistivity
(possibly ‘bad metals’ with $\rho > h/e^2$) and
linear-in- B magnetoresistance with B/T scaling.

Looking ahead:

- Experimental and theoretical studies of *metals* with bulk topological order and Fermi surfaces, possibly with non-Luttinger volumes.
- Experimental and theoretical studies of ‘strange’, ‘bad’, ‘incoherent’, ‘ultra-quantum’ metals: microscopic basis for SYK building blocks.