

Future of gapless UQM

UQM meeting
Caltech
September 20, 2022

Subir Sachdev

Talk online: sachdev.physics.harvard.edu

PHYSICS



HARVARD



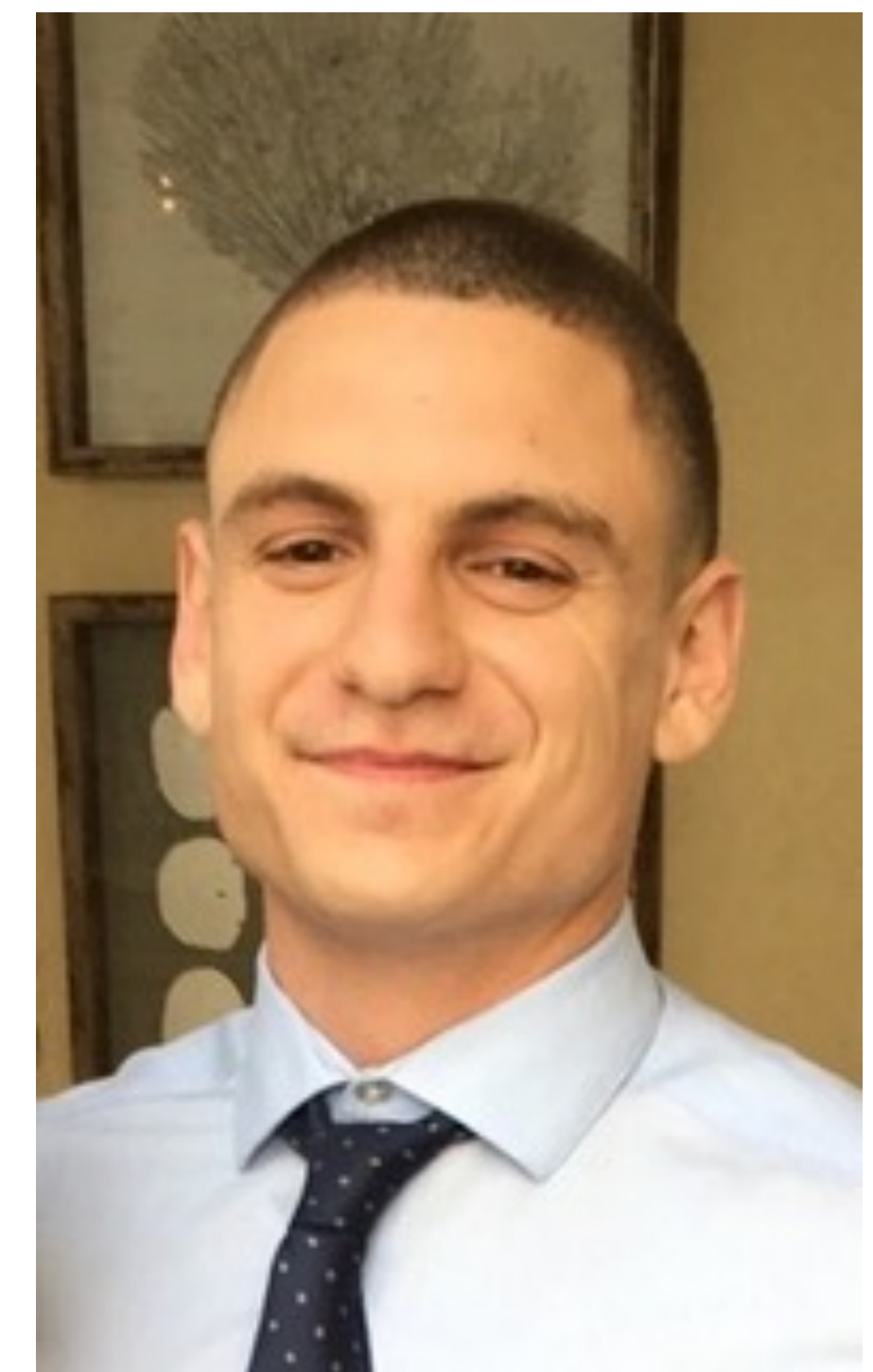
Aavishkar Patel

Flatiron Institute, NYC



Haoyu Guo

Harvard



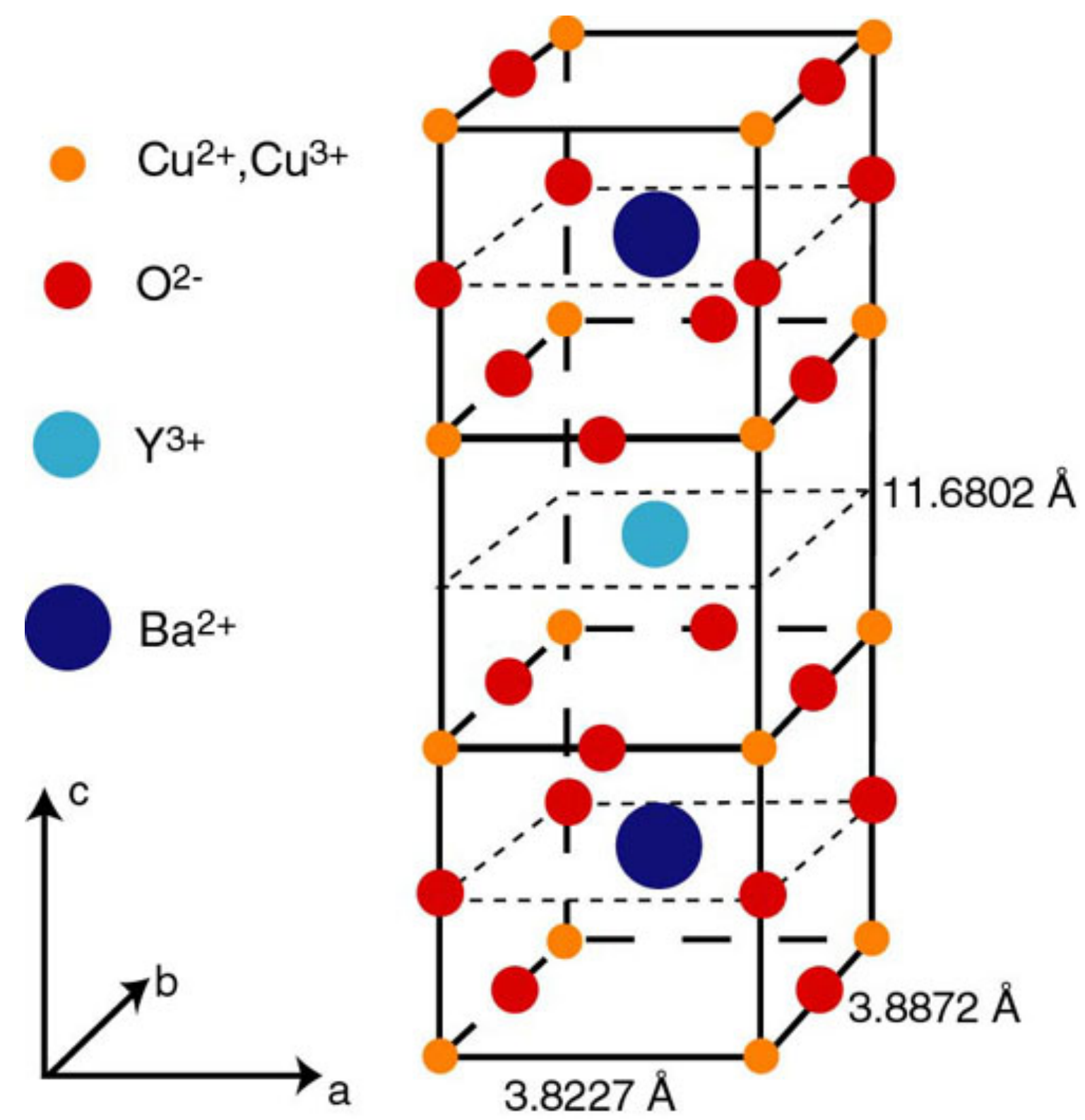
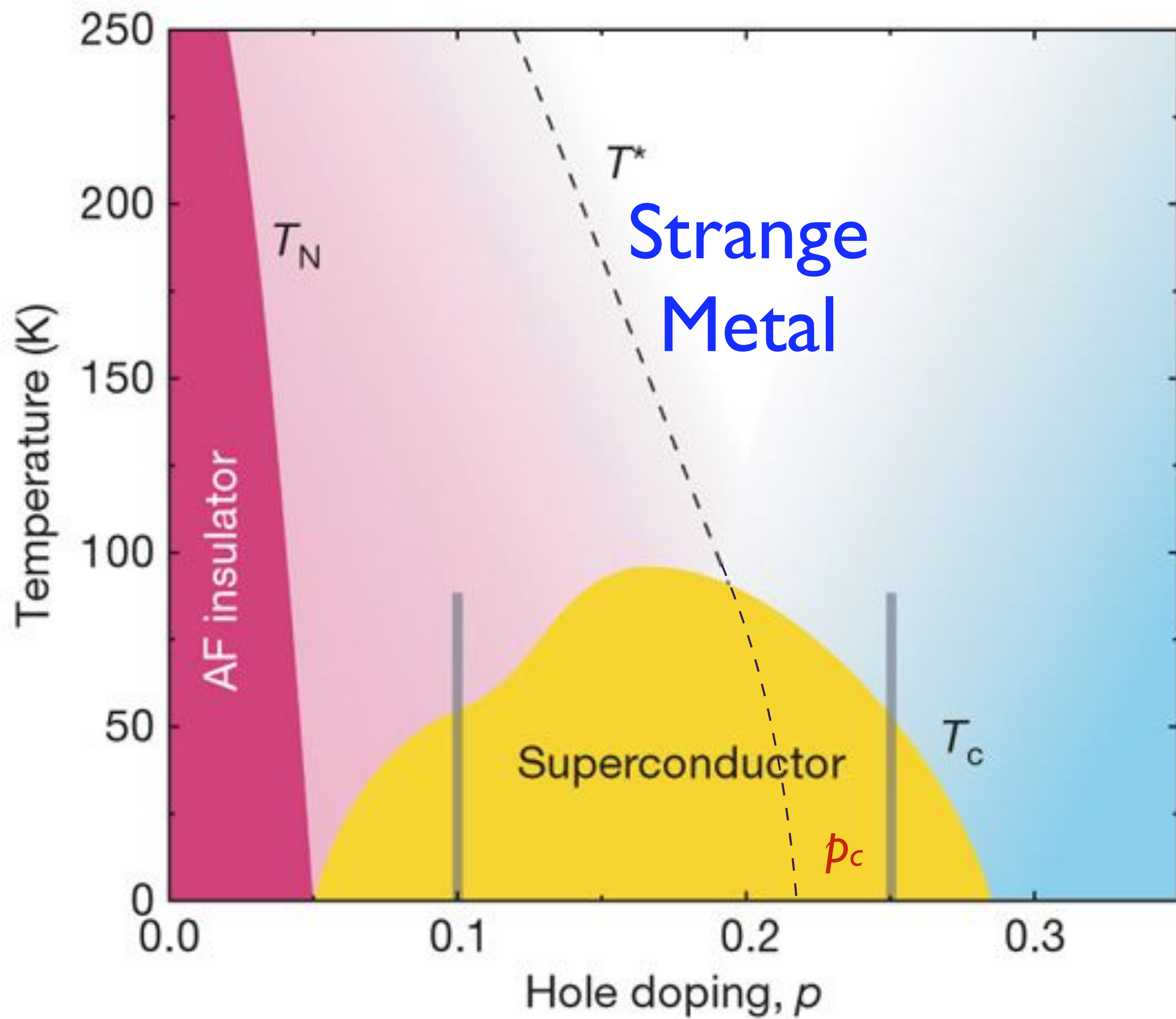
Ilya Esterlis

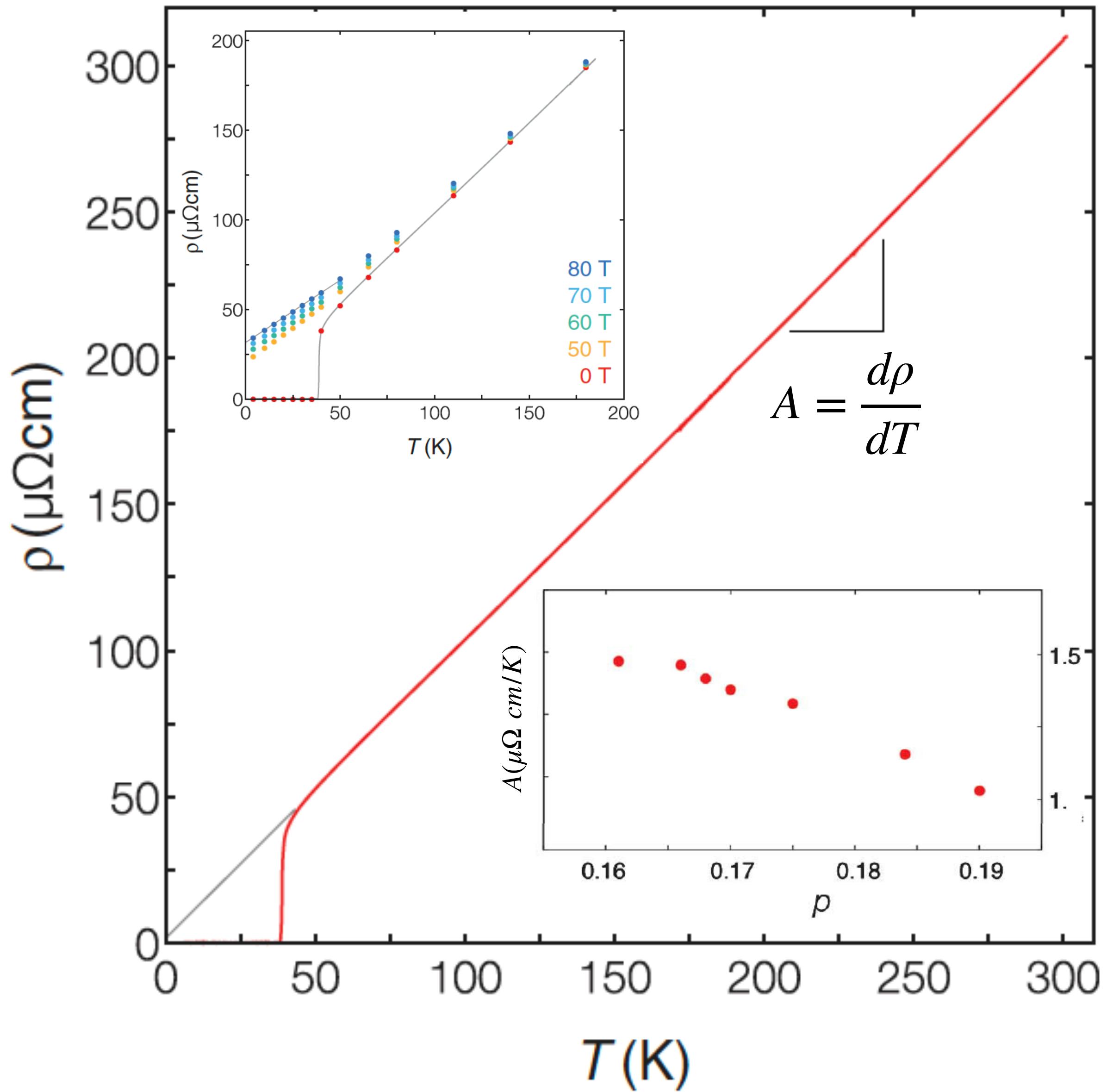
Harvard → Wisconsin

arXiv: 2103.08615, 2203.04990, 2207.08841

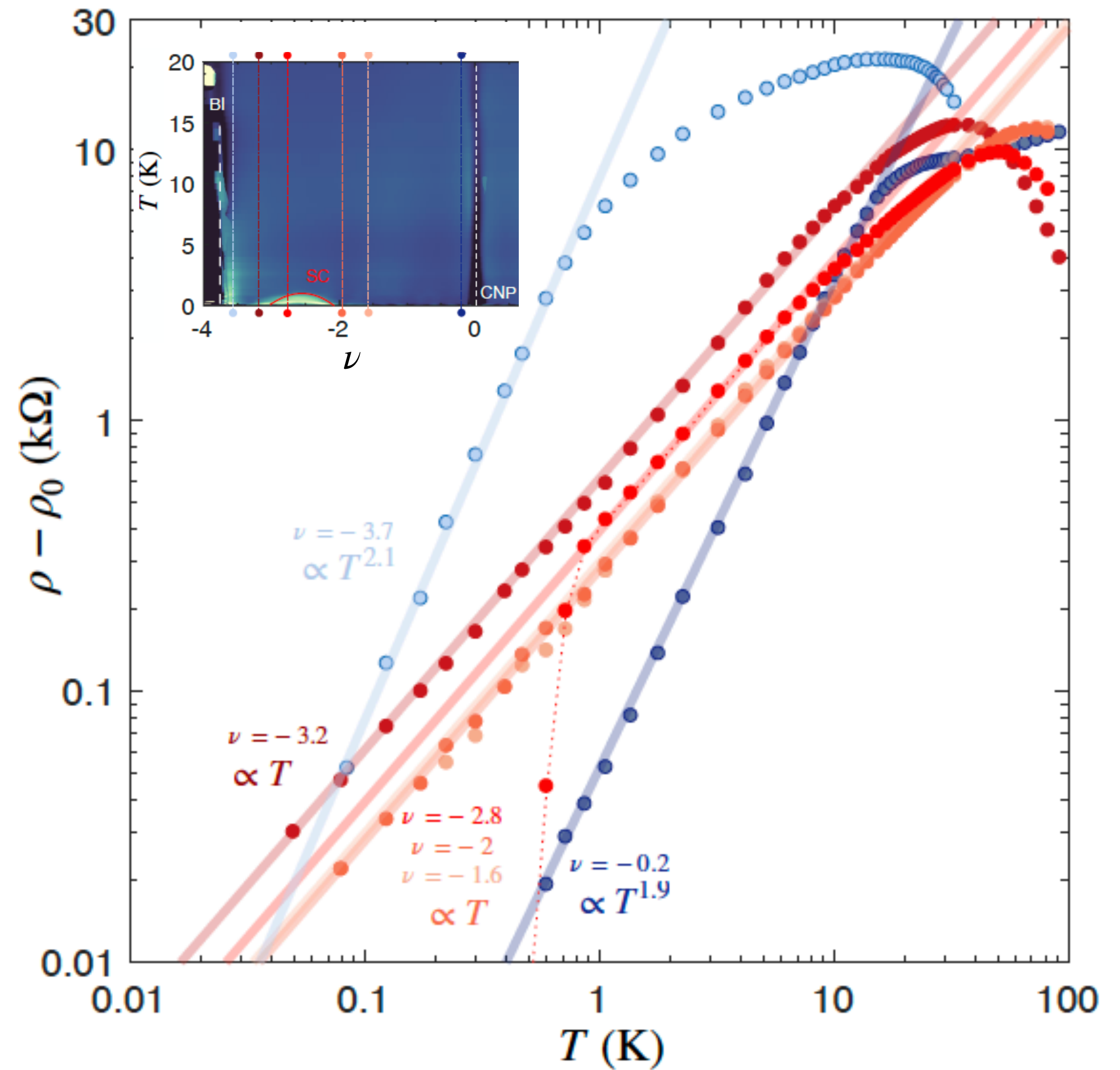
E. E. Aldape, T. Cookmeyer, Aavishkar A. Patel, and Ehud Altman, arXiv:2012.00763

**Strange
metals**





LSCO: Giraldo-Gallo et al. 2018



MATBG: Jaoui et al. 2021

Properties of a strange metal:

- Resistivity $\rho(T) = \rho_0 + AT + \dots$ as $T \rightarrow 0$
and $\rho(T) < h/e^2$ (in $d = 2$).
Metals with $\rho(T) > h/e^2$ are bad metals.

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- Specific heat $\sim T \ln(1/T)$ as $T \rightarrow 0$.

S.A. Hartnoll and A.P. MacKenzie, arXiv:2107.07802

Properties of a strange metal:

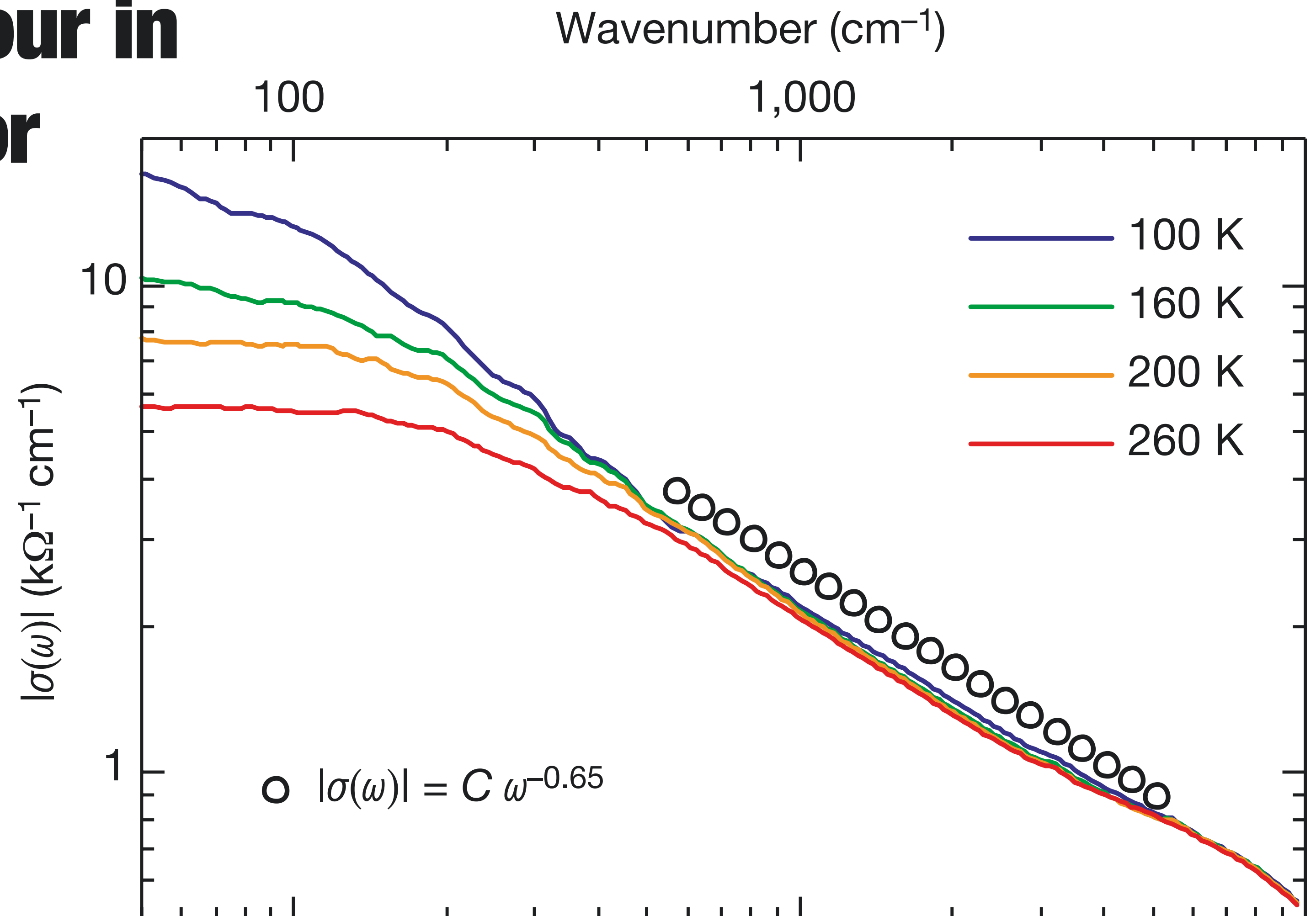
- Optical conductivity

Nature **425**, 271 (2003)

D. van der Marel^{1*}, H. J. A. Molegraaf^{1*}, J. Zaanen², Z. Nussinov^{2*},
F. Carbone^{1*}, A. Damascelli^{3*}, H. Eisaki^{3*}, M. Greven³, P. H. Kes² & M. Li²

Quantum critical behaviour in a high- T_c superconductor

But no $\hbar\omega/(k_B T)$ scaling.



- Optical conductivity

B. Michon,^{1,2,3} C. Berthod,³ C. W. Rischau,³ A. Ataei,⁴ L. Chen,⁴
S. Komiya,⁵ S. Ono,⁵ L. Taillefer,^{4,6} D. van der Marel,³ and A. Georges^{7,8,3}

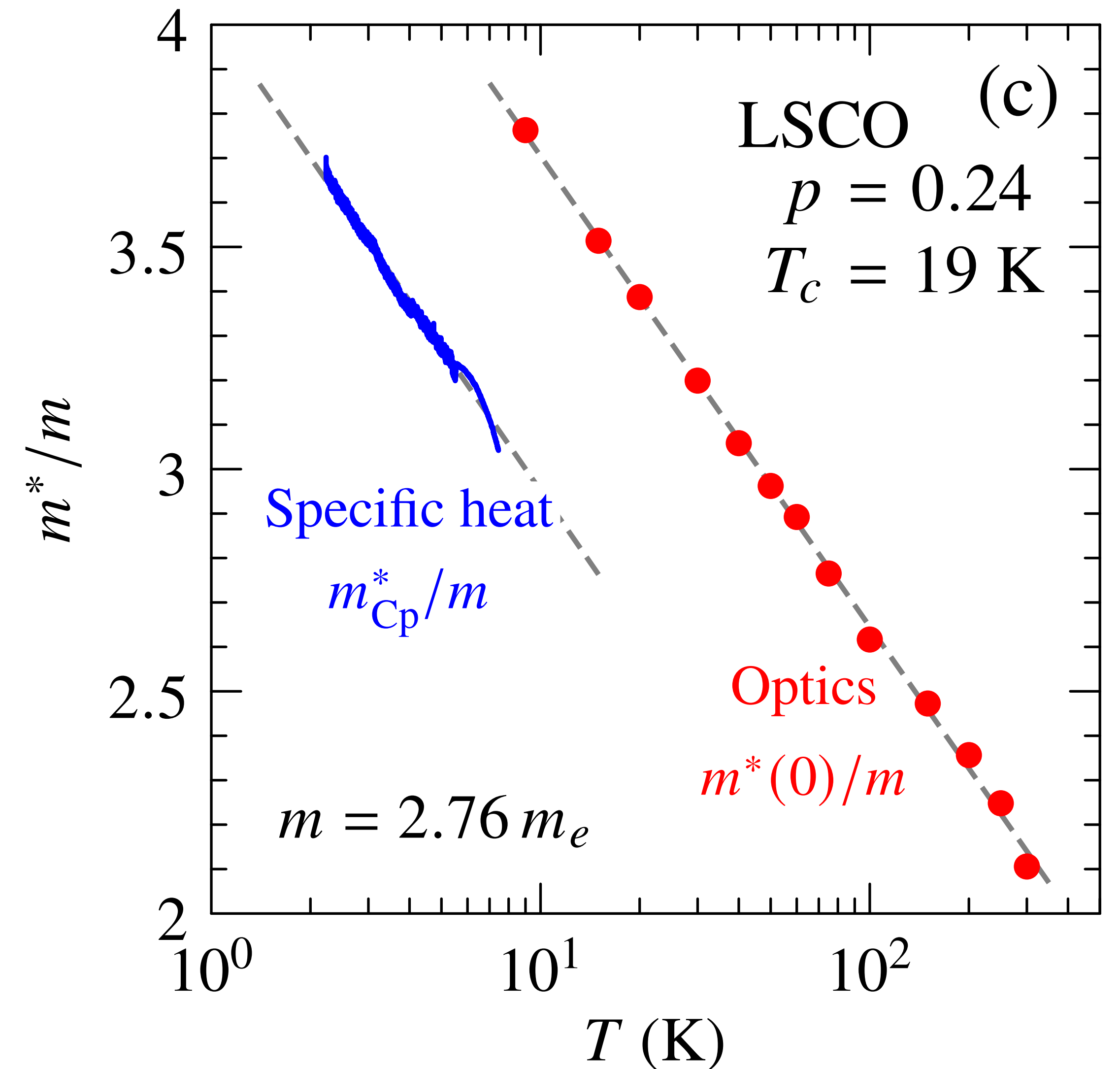
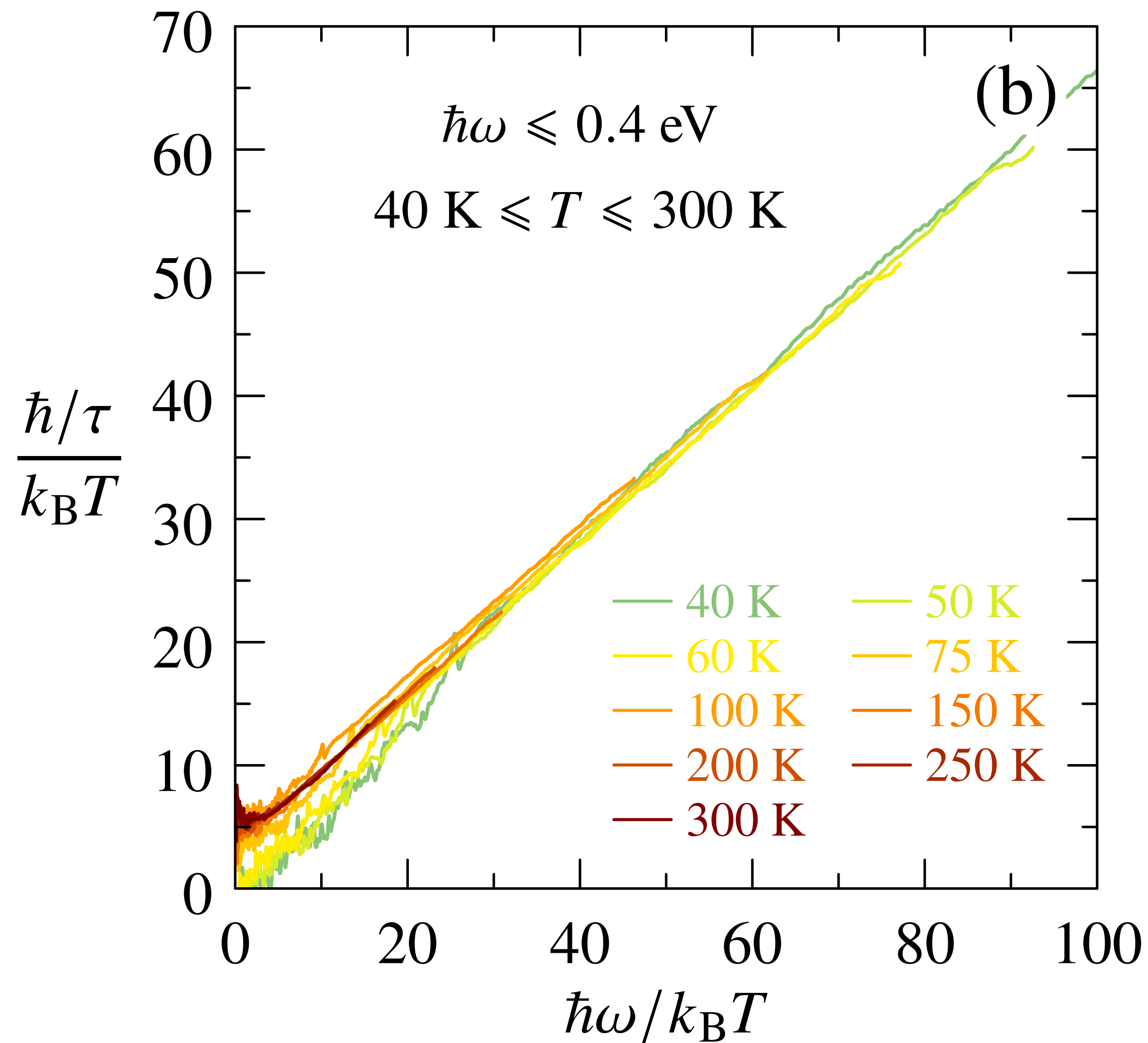
Planckian Behavior of Cuprate Superconductors: Reconciling the Scaling of Optical Conductivity with Resistivity and Specific Heat

$$\sigma(\omega) = \frac{K}{\frac{1}{\tau_{\text{trans}}(\omega)} - i\omega \frac{m^*(\omega)}{m}} \quad ; \quad \frac{1}{\tau_{\text{trans}}(\omega)} \sim |\omega| \Phi_{\sigma} \left(\frac{\hbar\omega}{k_B T} \right)$$

$$\text{Causality: } \frac{m^*(\omega)}{m} \sim \ln \left(\frac{\Lambda}{\text{Max}(\hbar\omega, k_B T)} \right)$$

● Optical conductivity

B. Michon,^{1,2,3} C. Berthod,³ C. W. Rischau,³ A. Ataei,⁴ L. Chen,⁴
S. Komiya,⁵ S. Ono,⁵ L. Taillefer,^{4,6} D. van der Marel,³ and A. Georges^{7,8,3}



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S.A. Hartnoll and A.P. MacKenzie, arXiv:2107.07802

- Optical conductivity

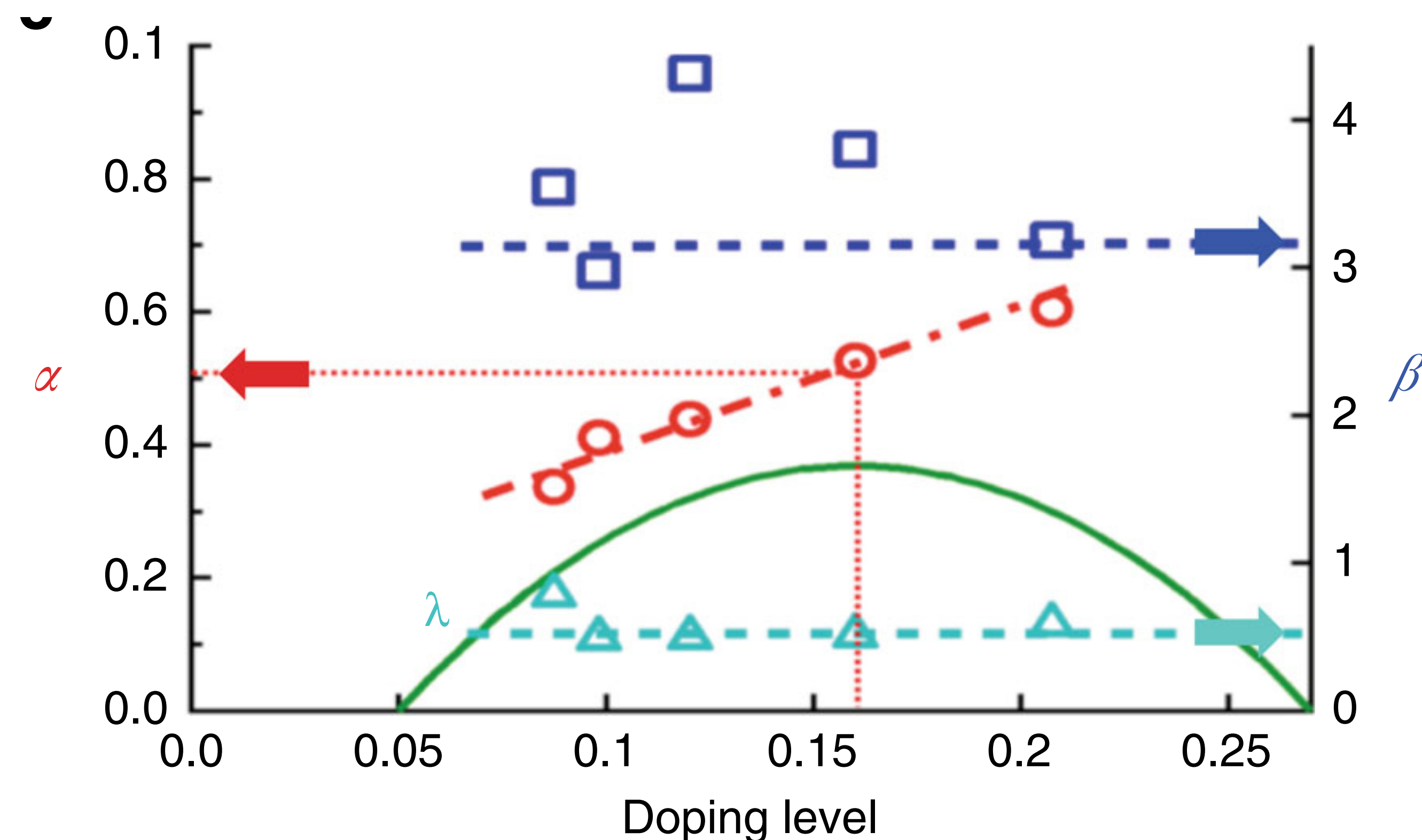
$$\sigma(\omega) = \frac{K}{\frac{1}{\tau_{\text{trans}}(\omega)} - i\omega \frac{m^*(\omega)}{m}} \quad ; \quad \frac{1}{\tau_{\text{trans}}(\omega)} \sim |\omega| \Phi_{\sigma} \left(\frac{\hbar\omega}{k_B T} \right)$$

B. Michon.....A. Georges, arXiv:2205.04030

● Photoemission

A unified form of low-energy nodal electronic interactions in hole-doped cuprate superconductors

T.J. Reber^{1,5*}, X. Zhou^{1*}, N.C. Plumb^{1,6}, S. Parham¹, J.A. Waugh¹, Y. Cao¹, Z. Sun^{1,7}, H. Li¹, Q. Wang¹, J.S. Wen², Z.J. Xu², G. Gu², Y. Yoshida³, H. Eisaki³, G.B. Arnold¹ & D.S. Dessau^{1,4*}



$$\Sigma''_{\text{PLL}}(\omega) = \Gamma_0 + \lambda \frac{[(\hbar\omega)^2 + (\beta k_B T)^2]^\alpha}{(\hbar\omega_N)^{2\alpha-1}}$$

Properties of a strange metal:

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S.A. Hartnoll and A.P. MacKenzie, arXiv:2107.07802

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$$\sigma(\omega) = \frac{K}{\frac{1}{\tau_{\text{trans}}(\omega)} - i\omega \frac{m^*(\omega)}{m}} \quad ; \quad \frac{1}{\tau_{\text{trans}}(\omega)} \sim |\omega| \Phi_{\sigma} \left(\frac{\hbar\omega}{k_B T} \right)$$

B. Michon.....A. Georges, arXiv:2205.04030

- Photoemission: nearly “marginal Fermi liquid” electron spectral density:

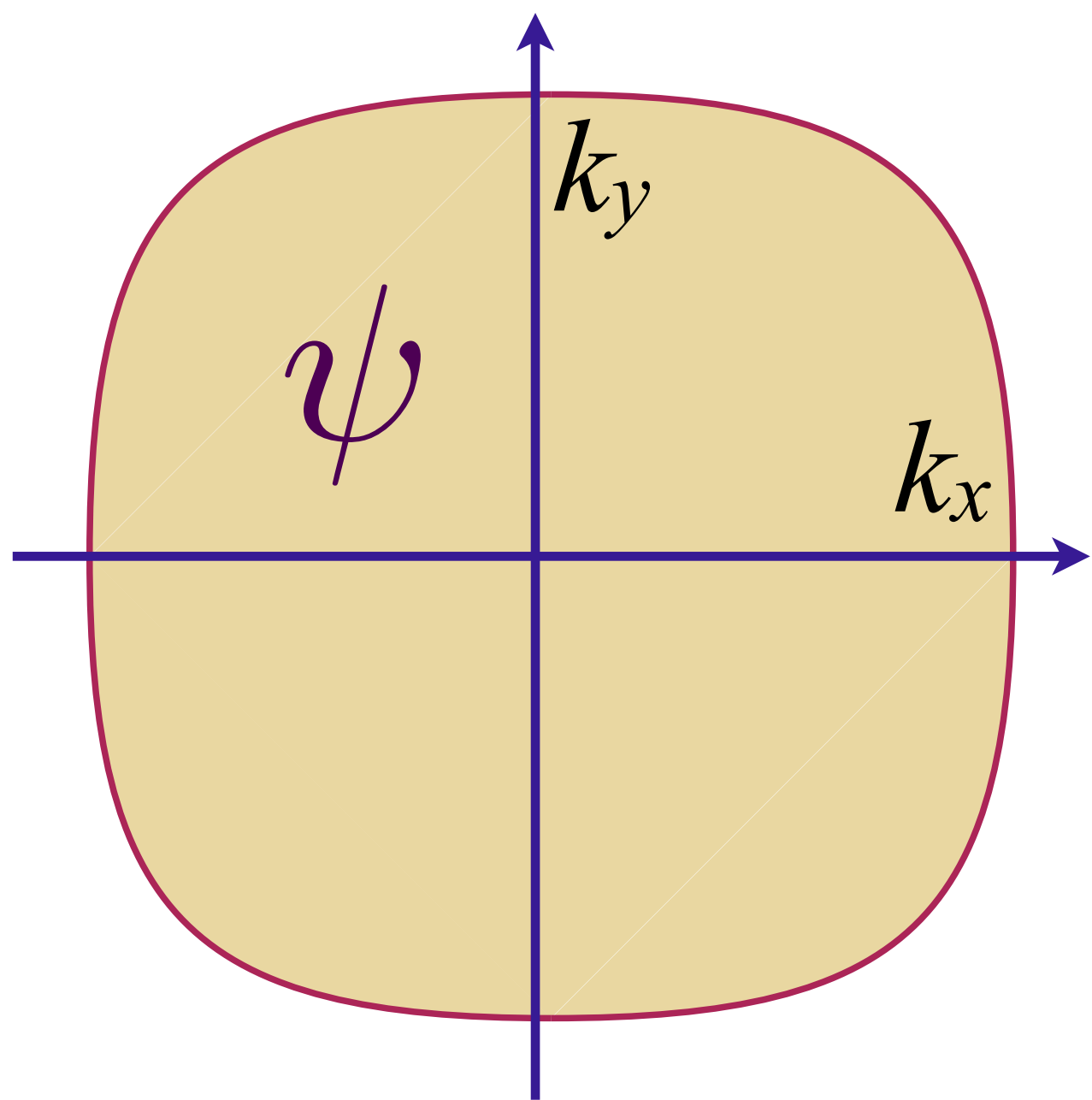
$$\text{Im}\Sigma(\omega) \sim |\omega|^{2\alpha} \Phi_{\Sigma} \left(\frac{\hbar\omega}{k_B T} \right) \quad \text{with } \alpha \approx 1/2$$

T.J. Reber....D. Dessau,

Nature Communications **10**, 5737 (2019)

**Fermions
and
critical bosons**

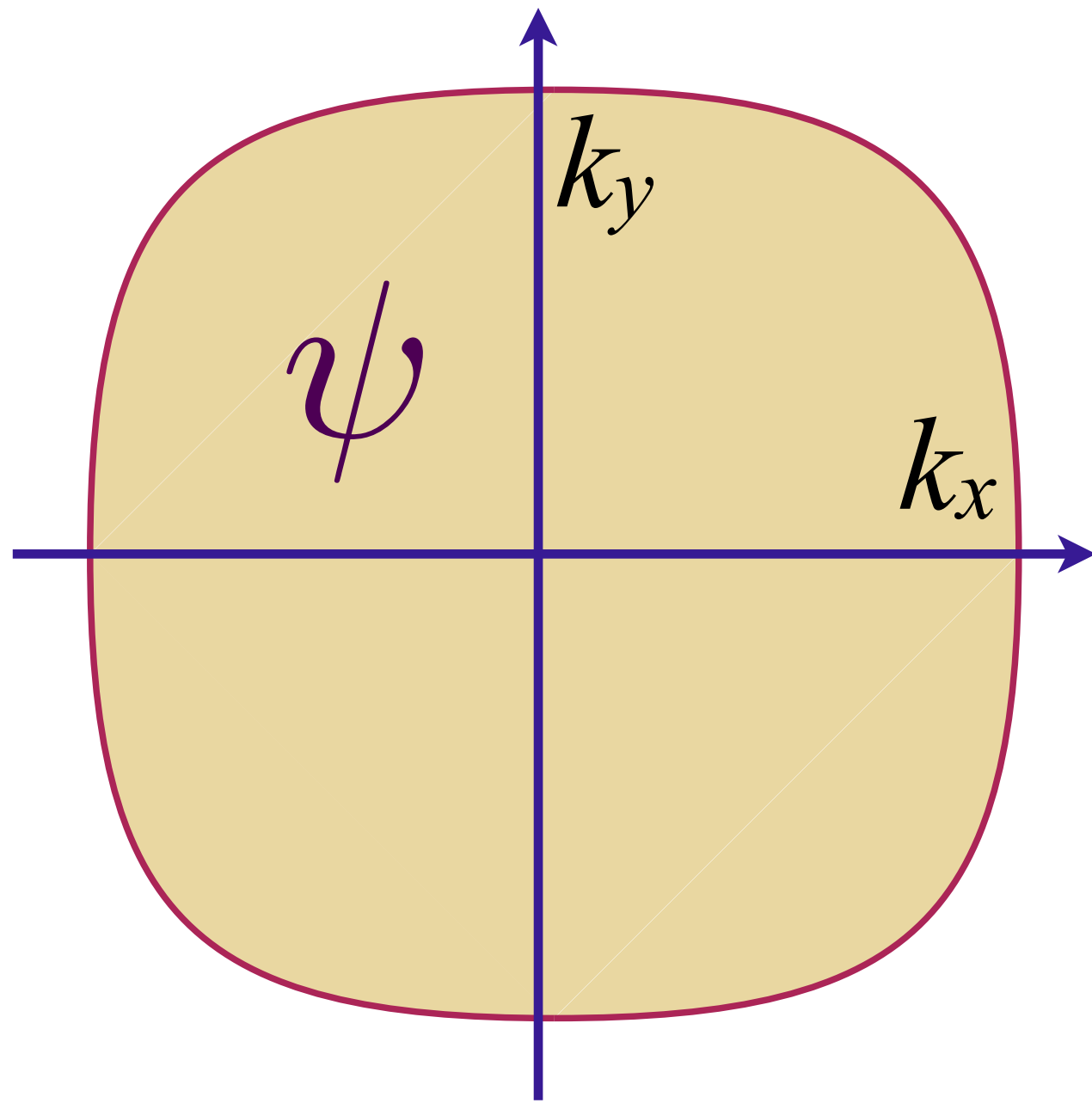
$$\mathcal{L}_\psi = \psi_{\mathbf{k}}^\dagger \left(\frac{\partial}{\partial \tau} + \varepsilon(\mathbf{k}) \right) \psi_{\mathbf{k}}$$



$$-J \psi^\dagger \psi^\dagger \psi \psi$$

Fermi surface coupled to a critical boson

$$\mathcal{L}_\psi = \psi_{\mathbf{k}}^\dagger \left(\frac{\partial}{\partial \tau} + \varepsilon(\mathbf{k}) \right) \psi_{\mathbf{k}}$$



a critical boson ϕ
e.g. Ising-nematic order

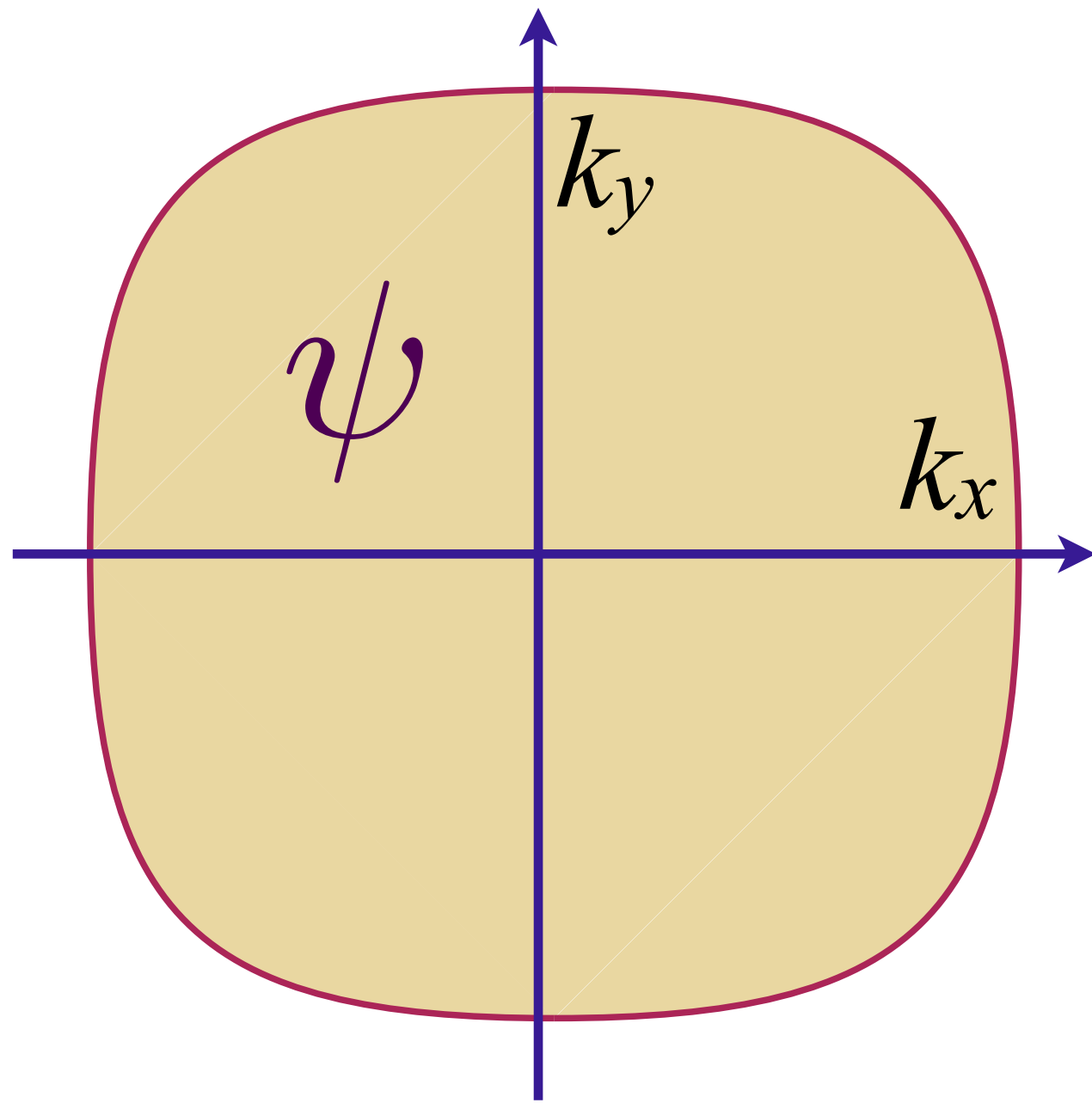
$$\frac{\phi^2}{J} + \psi^\dagger \psi \phi$$

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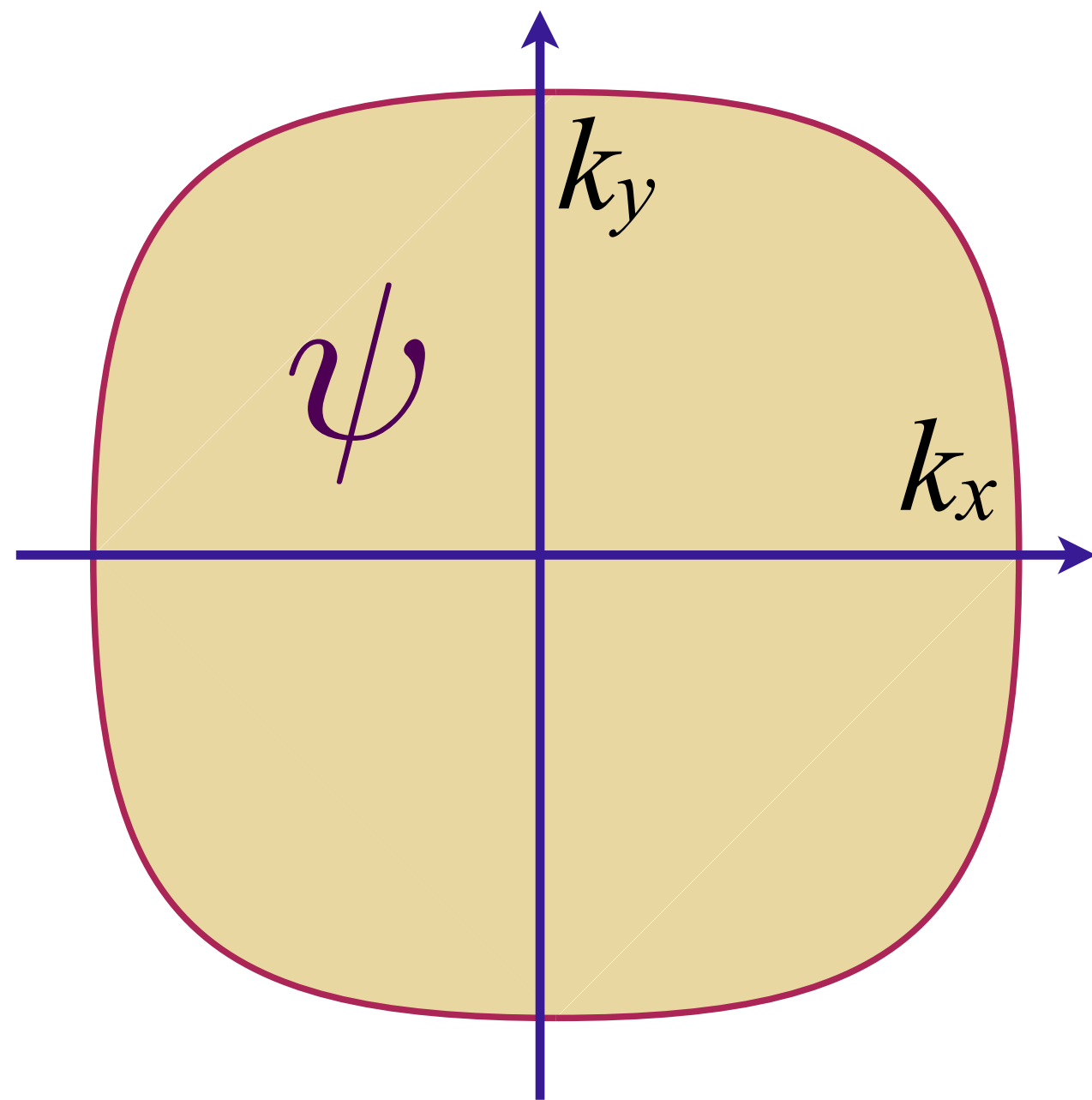


Eliashberg solution for electron (G) and boson (D) Green's functions at small ω :

$$\Sigma(\hat{\mathbf{k}}, i\omega) \sim -i \text{sgn}(\omega) |\omega|^{2/3}, \quad G(\mathbf{k}, i\omega) = \frac{1}{i\omega - \varepsilon(\mathbf{k}) - \Sigma(\hat{\mathbf{k}}, i\omega)}, \quad D(\mathbf{q}, i\Omega) = \frac{1}{\Omega^2 + q^2 + \gamma |\Omega|/q}$$

Fermi surface coupled to a critical boson

$$\mathcal{L}_\psi = \psi_{\mathbf{k}}^\dagger \left(\frac{\partial}{\partial \tau} + \varepsilon(\mathbf{k}) \right) \psi_{\mathbf{k}}$$



Transport—a perfect metal!

Conservation of momentum and fermion-boson drag imply:

$$\text{Re} [\sigma(\omega)] = D\delta(\omega) + \dots$$

a critical boson ϕ
e.g. Ising-nematic order

$$\frac{\phi^2}{J} + \psi^\dagger \psi \phi$$

S. A. Hartnoll, P. K. Kovtun, M. Muller, and S.S. PRB **76**, 144502 (2007)

D. L. Maslov, V. I. Yudson, and A. V. Chubukov PRL **106**, 106403 (2011)

S. A. Hartnoll, R. Mahajan, M. Punk, and S.S. PRB **89**, 155130 (2014)

A. Eberlein, I. Mandal, and S.S. PRB **94**, 045133 (2016)

Aavishkar Patel, Haoyu Guo, Ilya Esterlis, S.S. arXiv:2203.04990

Zhengyan Darius Shi, Hart Goldman, Dominic V. Else, T. Senthil

arXiv:2204.07585

Fermi surface coupled to a critical boson:

No spatial disorder

A non-Fermi liquid but NOT a strange metal

**Fermions
and
critical bosons
with spatial disorder**

Scaling Studies of Highly Disordered Spin- $\frac{1}{2}$ Antiferromagnetic Systems

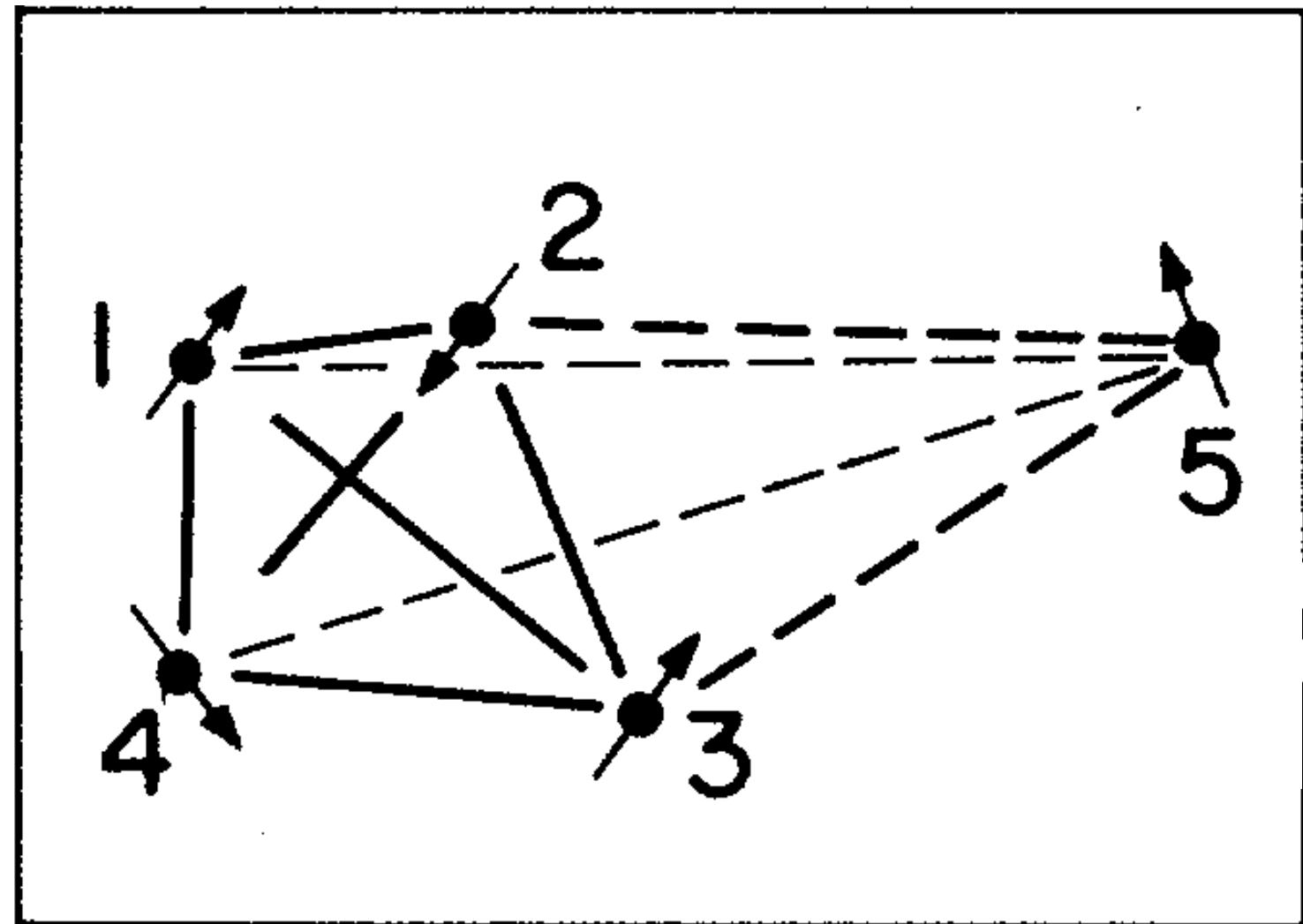
R. N. Bhatt and P. A. Lee

Bell Laboratories, Murray Hill, New Jersey 07974

(Received 6 May 1981; revised manuscript received 16 November 1981)

Physical Review Letters **48**, 344 (1982)

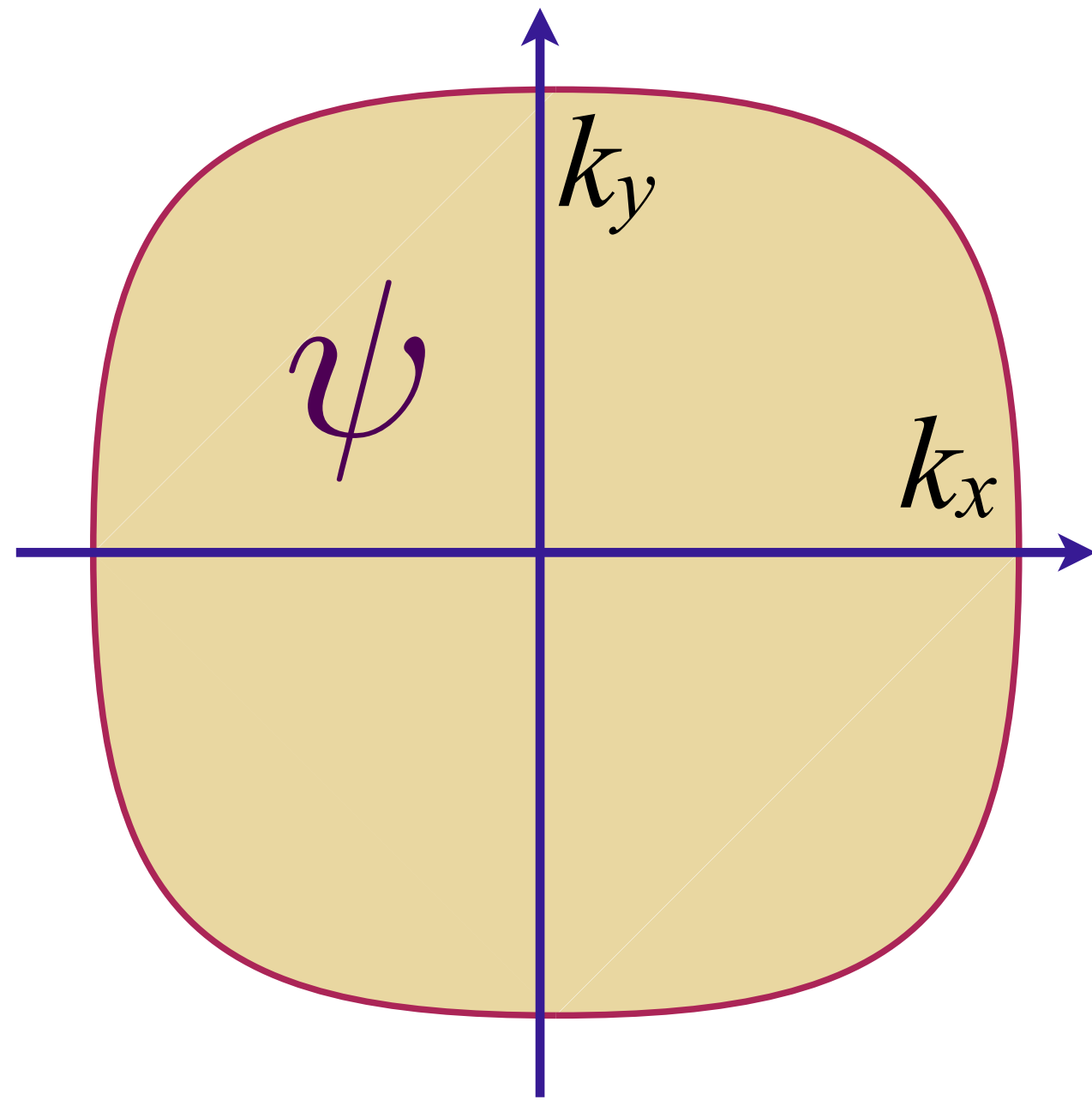
$$H = \frac{1}{2} \sum_{i \neq j} J(\vec{r}_i - \vec{r}_j) \vec{S}_i \cdot \vec{S}_j, \quad J(\vec{r}) \sim \exp(-2|\vec{r}|/a_0)$$



Strong disorder fixed point \Rightarrow Random singlet phase

Fermi surface coupled to a critical boson with disorder

$$\mathcal{L}_\psi = \psi_{\mathbf{k}}^\dagger \left(\frac{\partial}{\partial \tau} + \varepsilon(\mathbf{k}) \right) \psi_{\mathbf{k}}$$



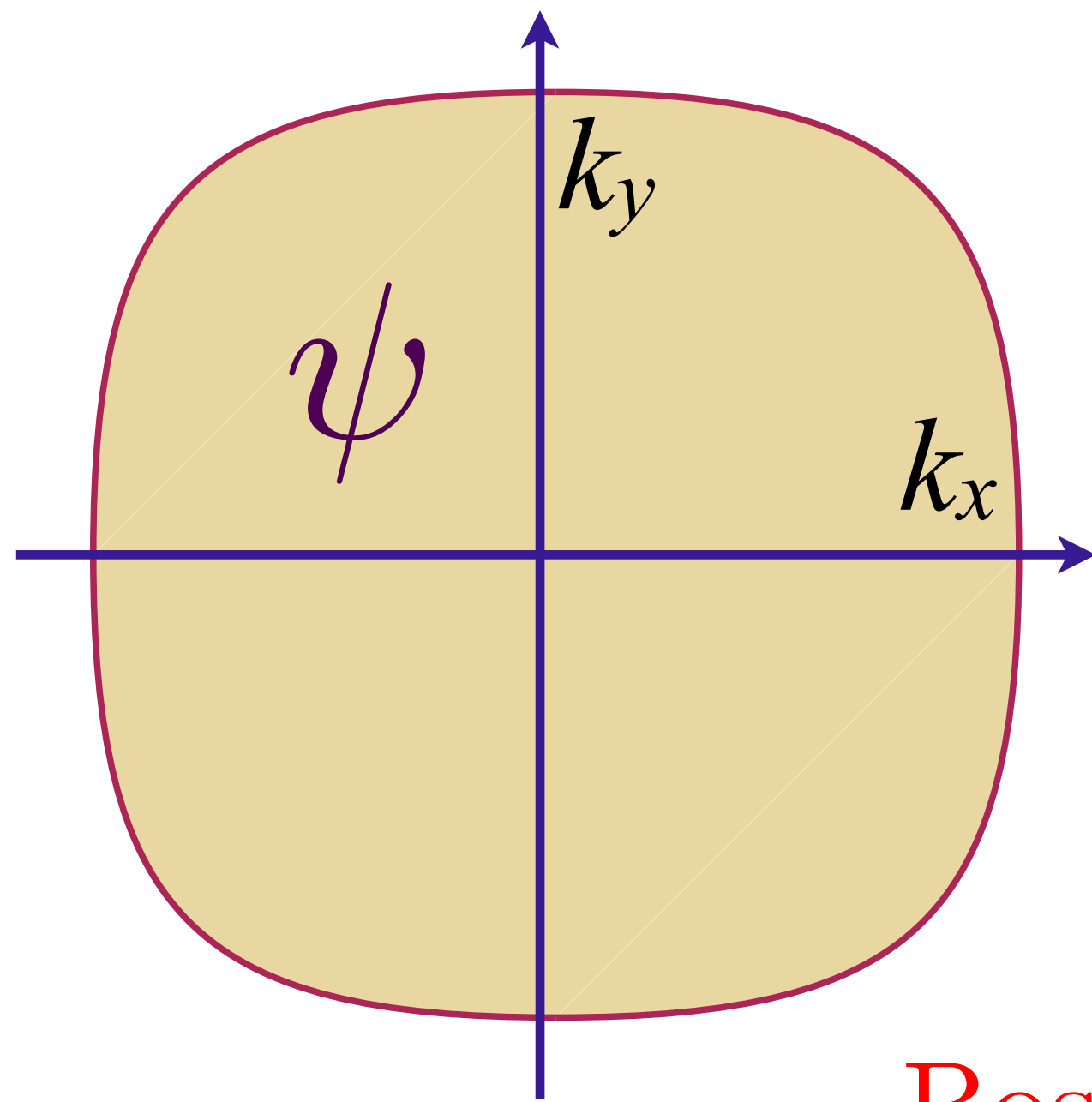
a critical boson ϕ
e.g. Ising-nematic order

$$\frac{\phi^2}{J} + \psi^\dagger \psi \phi$$
$$+ v(\mathbf{r}) \psi^\dagger \psi$$

Spatially random potential $v(r)$ with $\overline{v(r)} = 0$, $\overline{v(r)v(r')} = v^2 \delta(r - r')$

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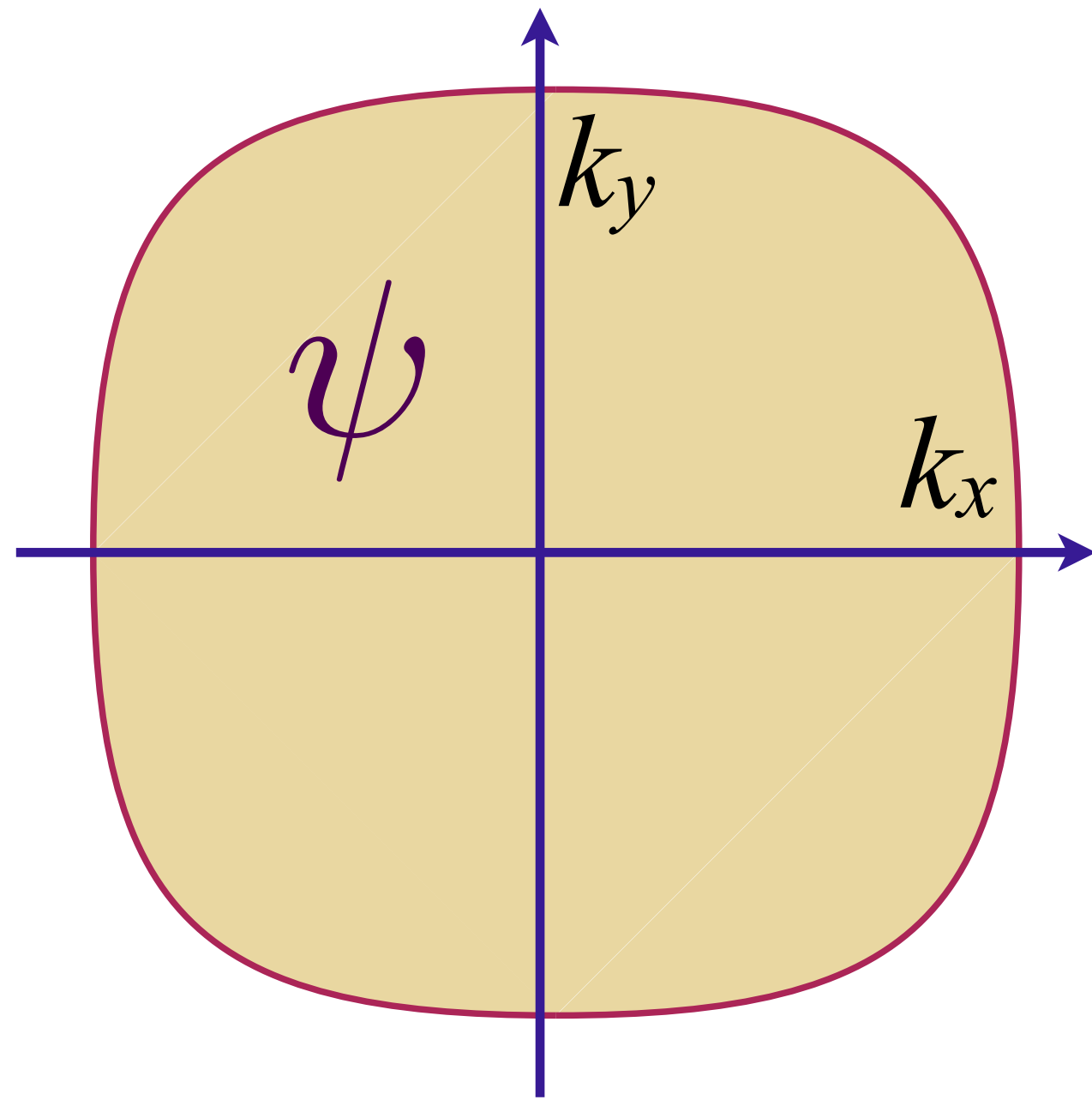
Boson Green's function: $D(q, i\Omega) \sim 1/(q^2 + \gamma|\Omega|)$

Fermion self energy: $\Sigma(i\omega) \sim -iv^2 \text{sgn}(\omega) - i \frac{g^2}{v^2} \omega \ln(1/|\omega|)$; $\frac{1}{\tau(\varepsilon)} \sim |\varepsilon|$

Marginal Fermi liquid self energy and $T \ln(1/T)$ specific heat

Fermi surface coupled to a critical boson with disorder

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$$\frac{\phi^2}{J} + \psi^\dagger \psi \phi$$
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But resistivity and optical conductivity are like a Fermi liquid,
with residual resistivity $\sim 1/v^2$.

Fermi surface coupled to a critical boson:

No spatial disorder

A non-Fermi liquid but NOT a strange metal

Fermi surface coupled to a critical boson:

Potential disorder v

A marginal Fermi liquid but NOT a strange metal

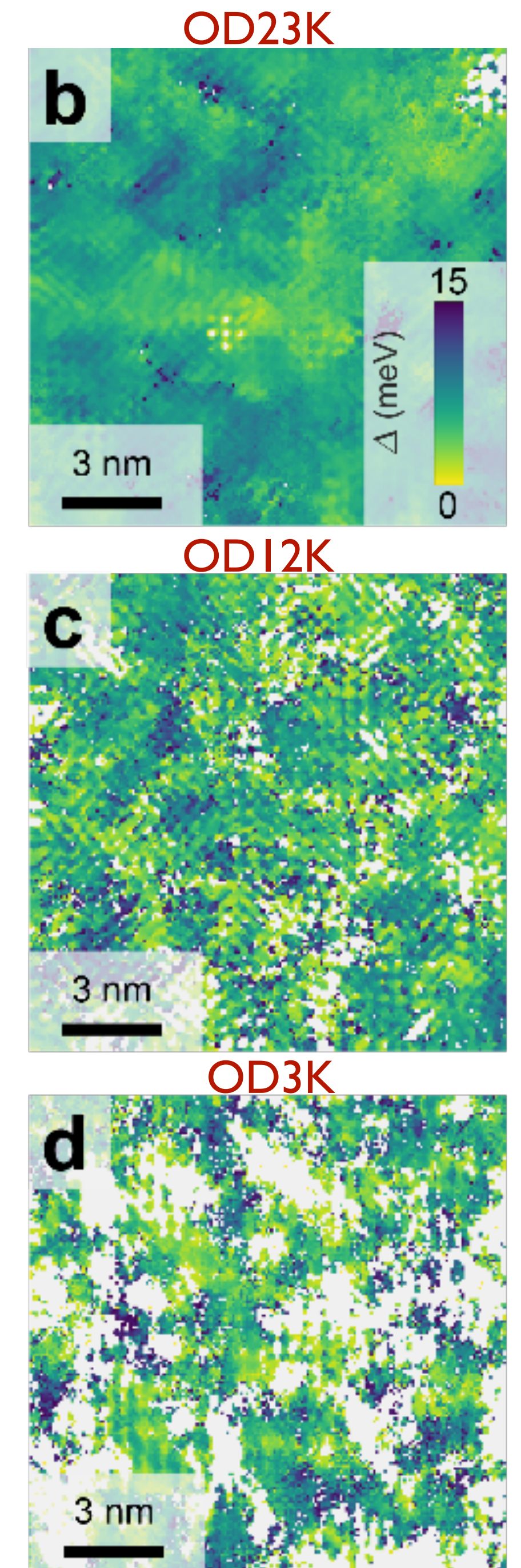
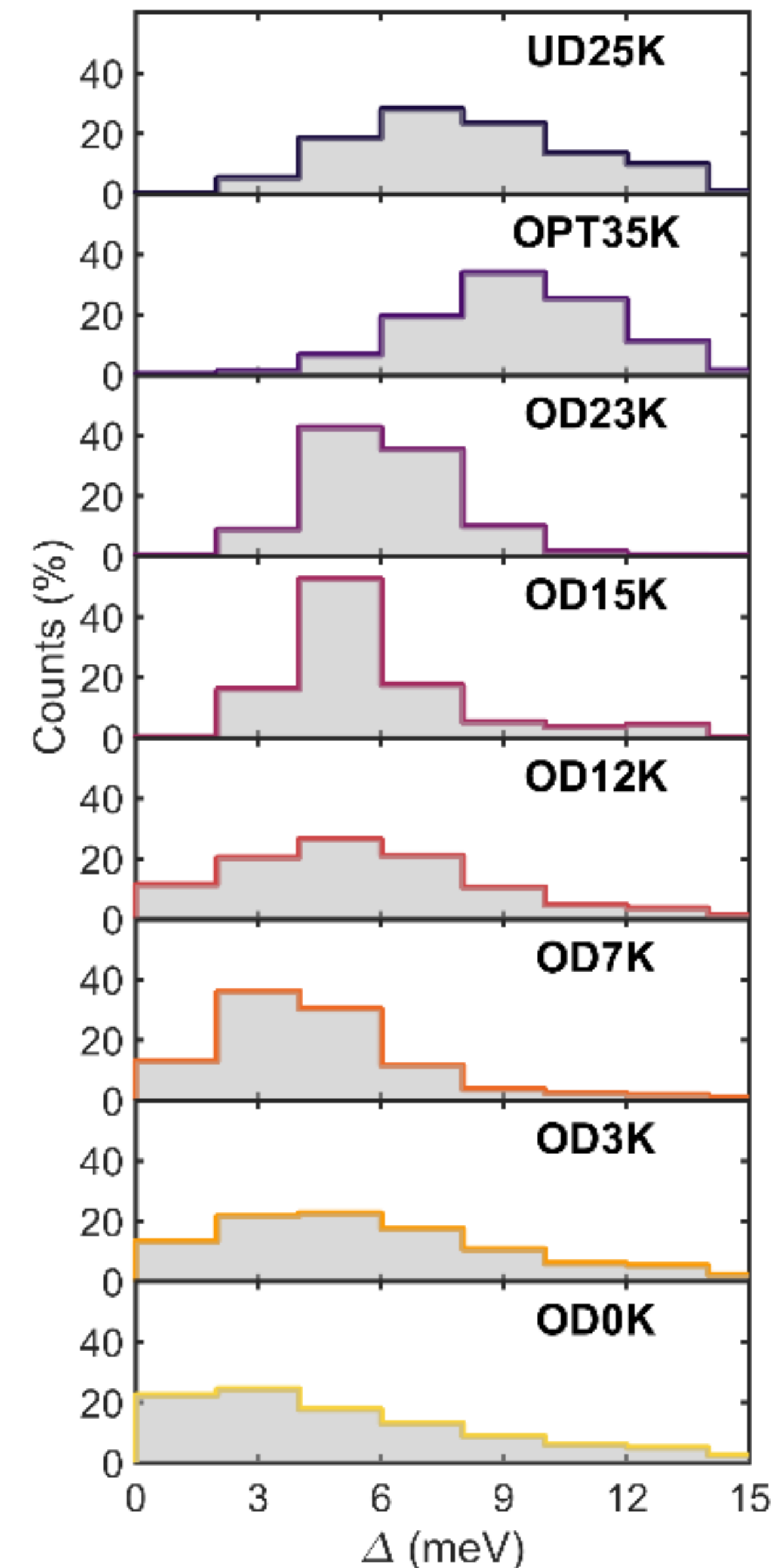
Spatially random interactions!

Puddle formation, persistent gaps, and non-mean-field breakdown of superconductivity in overdoped $(\text{Pb,Bi})_2\text{Sr}_2\text{CuO}_{6+\delta}$

Willem O. Tromp, Tjerk Benschop, Jian-Feng Ge, Irene Battisti, Koen M. Bastiaans, Damianos Chatzopoulos, Amber Vervloet, Steef Smit, Erik van Heumen, Mark S. Golden, Yinkai Huang, Takeshi Kondo, Yi Yin, Jennifer E. Hoffman, Miguel Antonio Sulangi, Jan Zaanen, Milan P. Allan

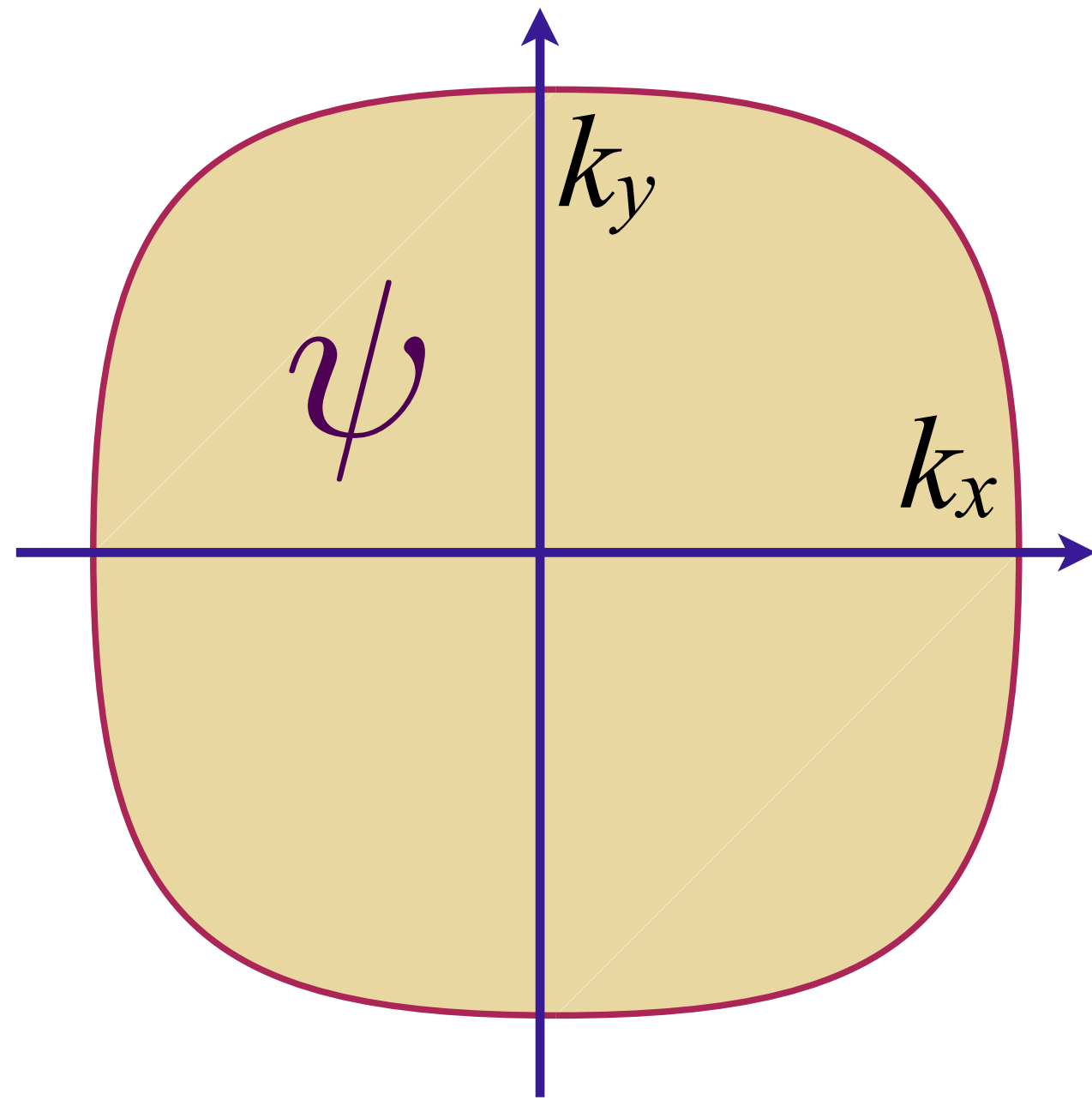
Our scanning tunneling spectroscopy measurements in the overdoped regime of the $(\text{Pb,Bi})_2\text{Sr}_2\text{CuO}_{6+\delta}$ high-temperature superconductor show the emergence of puddled superconductivity, featuring nanoscale superconducting islands in a metallic matrix

arXiv:2205.09740



Fermi surface coupled to a critical boson with disorder

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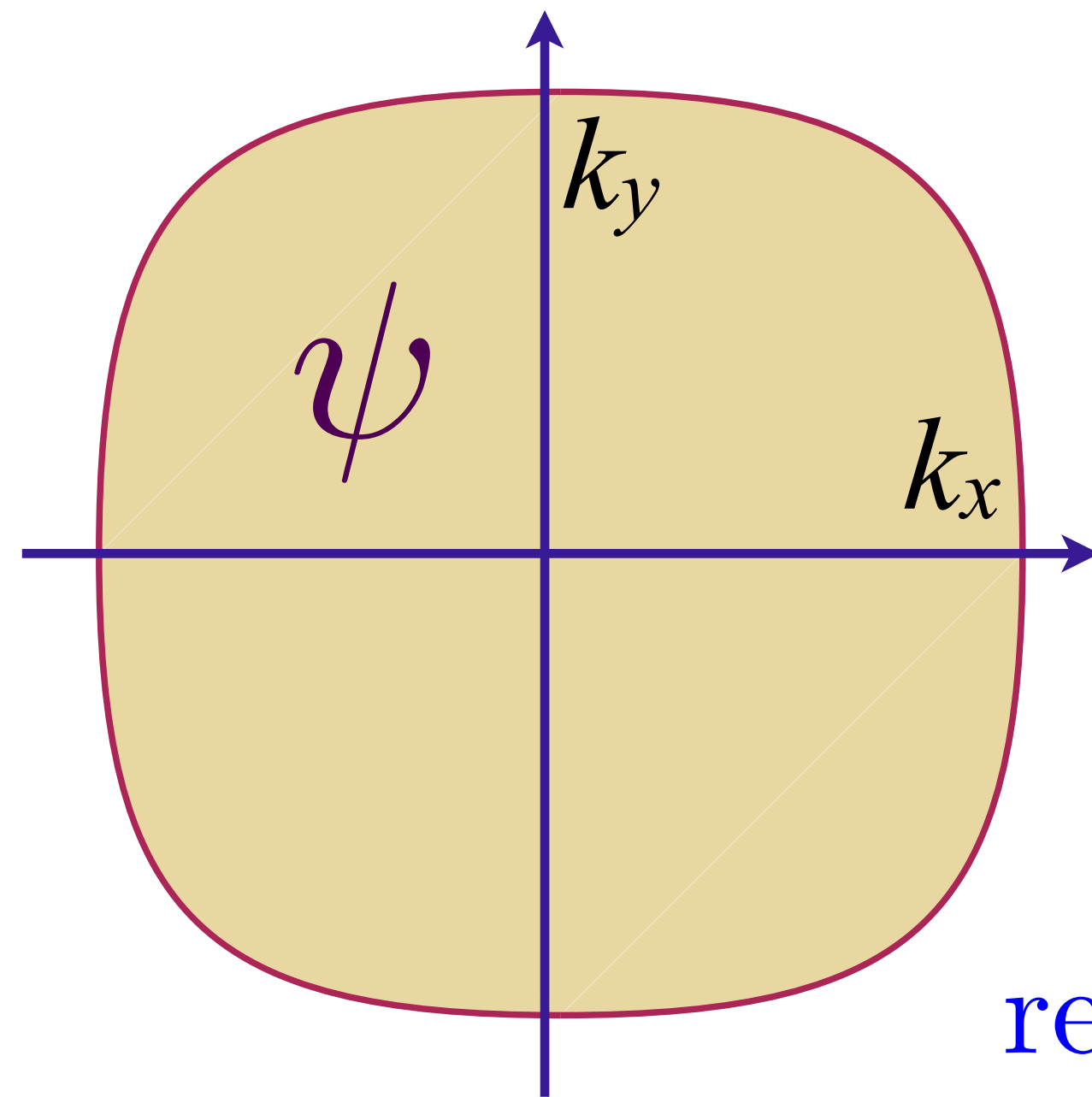
a critical boson ϕ
e.g. Ising-nematic order

$$\frac{\phi^2}{J + J'(\mathbf{r})} + \psi^\dagger \psi \phi + v(\mathbf{r}) \psi^\dagger \psi$$

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$$\phi^2 + [g + g'(\mathbf{r})] \psi^\dagger \psi \phi + v(\mathbf{r}) \psi^\dagger \psi$$

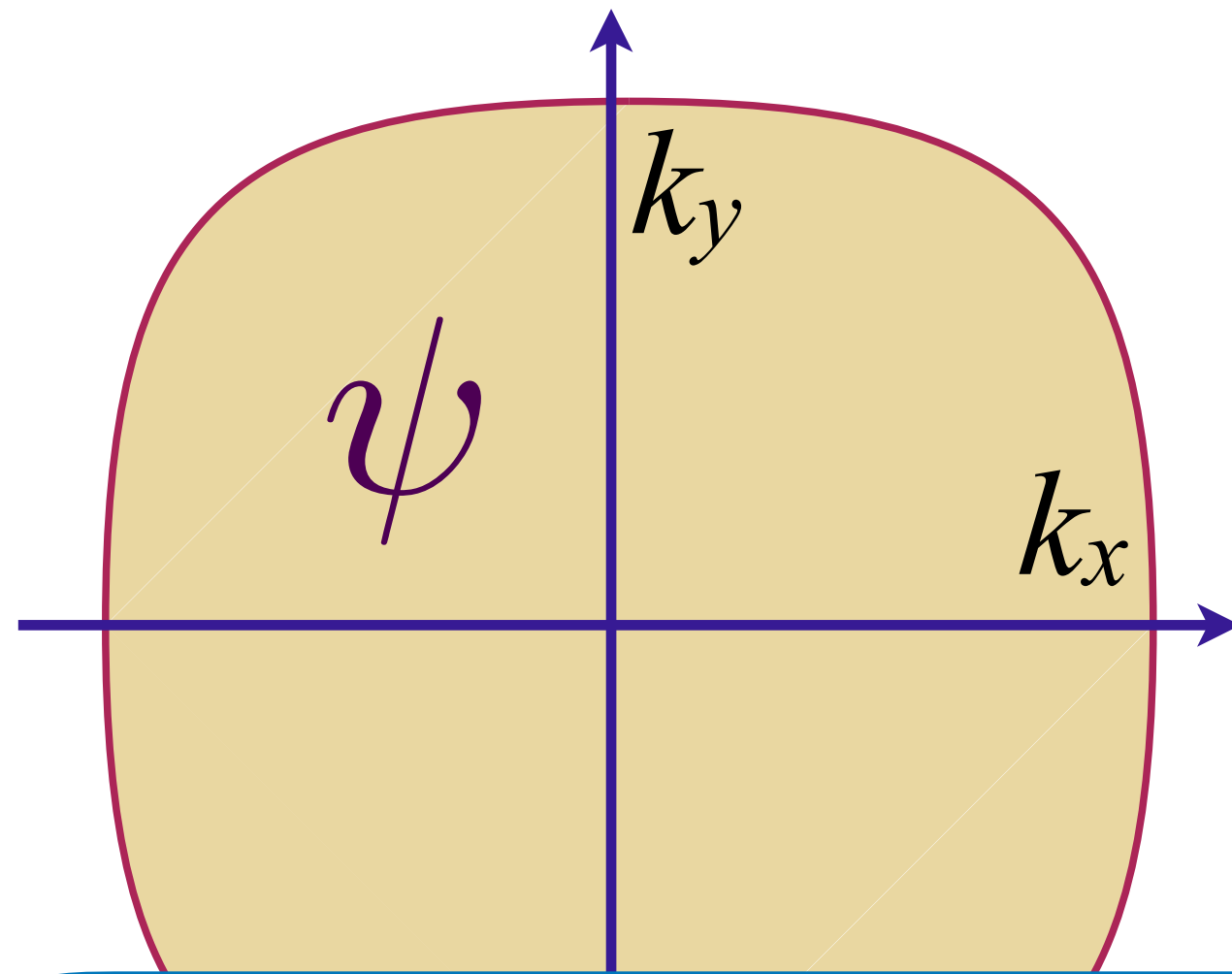
ϕ^2 “mass” disorder $J'(\mathbf{r})$ is strongly relevant;
 rescale ϕ to move disorder to the Yukawa coupling;
 can then be controlled in a SYK-like large N limit of ‘flavor’ indices,
 leading to a G - Σ - D - Π bi-local field theory.

Spatially random Yukawa coupling $g'(r)$ with $\overline{g'(r)} = 0$, $\overline{g'(r)g'(r')} = g'^2 \delta(r - r')$

Fermi surface coupled to a critical boson with disorder

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e.g. Ising-nematic order



$$\phi^2 + [g + g'(\mathbf{r})] \psi^\dagger \psi \phi$$

$$+ v(\mathbf{r}) \psi^\dagger \psi$$

Conductivity: $\sigma(\omega) \sim [1/\tau_{\text{trans}}(\omega) - i\omega m^*(\omega)/m]^{-1}$

$$\frac{1}{\tau_{\text{trans}}(\omega)} \sim v^2 + g'^2 |\omega| \quad ; \quad \frac{m^*(\omega)}{m} \sim \frac{2g'^2}{\pi} \ln(\Lambda/\omega)$$

Residual resistivity is determined by v^2 ; Linear-in- T resistivity determined by g'^2 ;
 Transport insensitive to g ; Marginal Fermi liquid self energy and $T \ln(1/T)$ specific heat.

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Potential disorder v

A marginal Fermi liquid but NOT a strange metal

Fermi surface coupled to a critical boson:

Interaction disorder g'

A marginal Fermi liquid AND a strange metal

The future

- Quantitative fitting of different experiments

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- More observables: magnetoresistance, cyclotron resonance, shot noise . . .

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- $1/N$ corrections
- Compute T_c to superconductivity